## Design Aid Diagram for FRC Fusion-research Devices: Research paths defined by dimensionless variables

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## I. FRC PARAMETERS AND EQUATIONS

To explore FRC reactor physics properly, experimental devices should be designed to have their pertinent dimensionless parameters approaching those of their target reactor. Among the dimensionless parameters considered most relevant are  $\nu^*$ , collisionality – to address transport issues – , and s, a measure of the degree of MHD-like vs. kinetic-like behavior – to address stability questions.  $\nu^*$  is defined by the ratio of the machine size,  $\kappa r_s$ , to the Spitzer-scattering mean-freepath (mfp),  $\lambda_S \equiv \lambda_C/4ln\Lambda$ , where  $\lambda_C = 90^\circ$  Coulomb-scattering mfp,  $\kappa =$  FRC elongation, and  $r_s$ = separatrix radius. The parameter s is approximated as the number of ion gyro-radii,  $\rho_i$ , between the FRC O-point radius,  $r_o$ , and  $r_s$ :  $s \sim 0.3r_s/\rho_i$ . Both s and  $\nu^*$  are determined by relatively few basic machine parameters, e.g., primarily external magnetic field,  $B_e$ , and  $r_s$ , with the plasma species temperatures,  $T_{e,i}$ , remaining the only free parameters. Plasma temperatures would be set in a reactor by the desired fusion power density.

The equations used to find s and  $\nu^*$  are (in the commonly used mixed units,  $n \ (cm^{-3}), r_s \ (cm), B_e \ (G), T \ (eV)$  and  $m_i$  in amu):

$$\langle \beta \rangle = 4.03 \times 10^{-11} \Sigma \langle nT \rangle / B_e^2 \tag{1}$$

$$\phi \sim \pi r_s^2 B_e/4 \tag{2}$$

$$s = 3 \times 10^{-3} r_s B_e / (m_i T_i)^{1/2} \tag{3}$$

$$\nu^* = 2.5 \times 10^{-14} n \kappa r_s / T^2 \tag{4}$$

Equation (1) is used to eliminate density from the other equations. Equation (2) assumes reasonable plasma profiles. What remains are equations that contain  $\kappa$ ,  $B_e$ ,  $r_s$ ,  $\langle\beta\rangle$ ,  $m_i$ , and T. In the following we assume singly charged ions,  $m_i = 1$ ,  $\kappa = 5$ ,  $\langle\beta\rangle = 0.7$ ,  $T_e = T_i$  and  $n_e = n_i$ . These choices have little impact on the overall picture, due to the limited ranges these parameters may have and the log-log scales on the diagrams.

When considering a reactor, an important issue is the confinement of fusion products. Ignited operation requires, at a minimum, that the gyro-radii of the fusion products,  $\rho_{fp}$ , be considerably small than  $r_s$ . For  $r_s/\rho_{fp} \ge 6$ , the minimum required magnetic flux for confining 3.5 MeV  $\alpha$  particles is  $\phi_m \ge 2.8 \times 10^5 E/B_e \sim 0.7$  Wb at  $B_e = 15$  kG, with E the  $\alpha$ -particle energy in eV and

 $\phi$  in  $G \ cm^2$ . (1 Wb = 10<sup>8</sup>  $G \ cm^2$ .) Driven reactors, ones that do not rely on direct plasma heating by fusion products, do not have this constraint.

Ignited D-T reactors are estimated have  $r_s \sim 2.5$  m,  $B_e \sim 15$  kG, and temperature and density near 10 keV and  $4 \times 10^{14}$  cm<sup>-3</sup>, respectively, yielding  $s \sim 30$  and  $\nu^* \sim 10^{-2}$ . A driven,  $D - He_3$ reactor would require about 10× higher temperature but could have a proportionally smaller  $r_s$ ,  $B_e \sim 50$  kG, resulting in  $s \sim 8$  and  $\nu^* \sim 10^{-5}$ . An ignited  $p - B_{11}$  reactor, if possible, would have  $r_s \sim 2.5$  m and  $B_e \sim 150$  kG, with  $s \sim 40$  and  $\nu^* \sim 10^{-5}$ .

The final parameter that describes the FRC reactor system is the energy confinement time, embodied in the diffusivity, D. For classical collisional transport,  $D = \rho_i^2 \nu_S = r_s^2 / \tau_E$ , with  $\nu_S$  the Spitzer scattering frequency. D can be rewritten in terms of s and  $\nu^*$ ,

$$D = 0.1 \frac{r_s v_t}{\kappa} \frac{\nu^*}{s^2},\tag{5}$$

where  $v_t$  is the ion thermal velocity. The resulting classical D's are 0.6  $m^2/s$  for the ignited, 10keV, D-T, 2.5-m device and 0.003  $m^2/s$  for the driven, 100-keV,  $D - He_3$ , 25-cm device. Both the ignited D-T and driven  $D - He_3$  FRC reactors described above require a (thermal) diffusivity less than a factor of 10 above classical. Non-classical mechanisms could increase transport and D. At their expected  $\nu^*$  and s values, ignited reactors face extreme stability issues while driven reactors would face challenging confinement requirements. Equation (5) indicates that confinement and stability issues are not strictly independent.

Other dimensionless parameters should be considered in the design of next-step FRC research devices. For example, particularly important to pulse-length requirements and flux sustainment is the Lundquist number.

## II. THE DESIGN AID DIAGRAM (DAD)

Figure 1 shows iso-contours of the parameters  $\phi$ ,  $\nu^*T^3$ , and  $sT^{1/2}$  plotted in the  $r_s - B_e$  plane. Also plotted are points representing four operating FRCs (TCSU, MRX-CT, PFRC and PHD) and proposed new ones, including that of the Tri-Alpha Corporation, a  $p - B_{11}$ -fueled device. Research devices intended to explore stability or transport separately would have a wide area of the  $r_s - B_e$ plane available, depending on the temperature achievable. However, if simultaneous reactor-like values for s and  $\nu^*$  were necessary for more realistic experiments, then temperature is no longer a free parameter and a constraint is found from equations (1)-(4):

$$B^{4/5}r_s = 243(\frac{s^6}{\nu^*})^{1/5}(m^3\kappa\langle\beta\rangle)^{1/5},\tag{6}$$

The value of  $(s^6/\nu^*)^{1/5}$  is about a factor of 1.5 higher for ignited operation than for driven. Isocontours are shown for both the ignited and driven base cases and approximately bisect figure diagonally. Clearly, high ion temperatures are necessary to achieve reactor-like s and  $\nu^*$  values simultaneously. An important conclusion is that tests at reactor-relevant s and  $\nu^*$  can be performed in devices with  $B_e$  and  $r_s$  3-5 times smaller than required for reactors or to contain  $\alpha$  particles.

FRC research programs must attack fundamental issues and simultaneously improve plasma parameters by many orders of magnitude. Should research explore energy confinement or stability separately? Is it proper to investigate D at low s or stability at high  $\nu^*$ ? Should FRC research programs make large steps in size and power to a next-step device? What role does the privately financed Tri-alpha device play in DOE planning?



FIG. 1: Design aid diagram showing contours of constant  $\phi$ ,  $sT^{1/2}$ , and  $log(\nu^*T^3)$  in the  $r_s - B_e$  plane. Positions of current and proposed FRC devices are shown. Note, for calculating s, T is the ion temperature. Also shown are iso-contours of constant  $(s^6/\nu^*)^{1/5}$  for two generic FRC reactors, ignited D-T and driven D-He<sub>3</sub>, described in the text.

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