
Toroidal Rotation and 3D Nonlinear Dynamics in the Peeling-Ballooning Theory of ELMs

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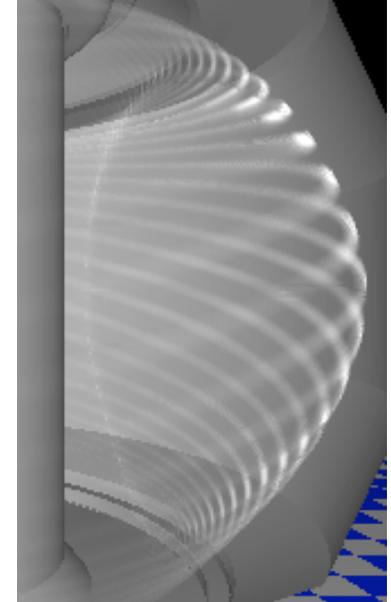
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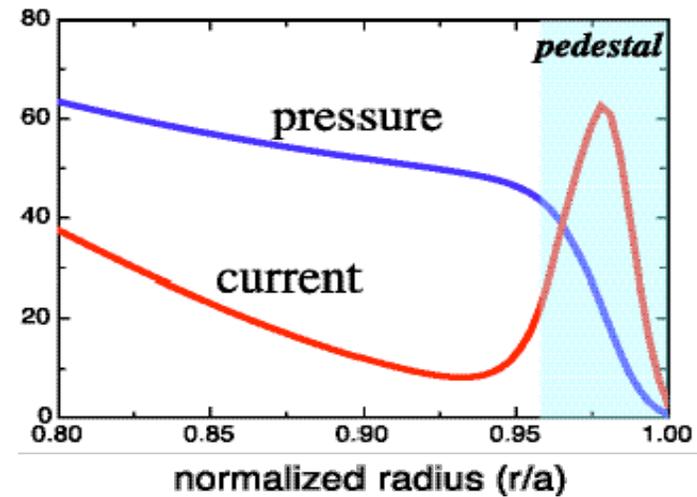
APS Division of Plasma Physics Meeting

Savannah, GA, 17 November 2004

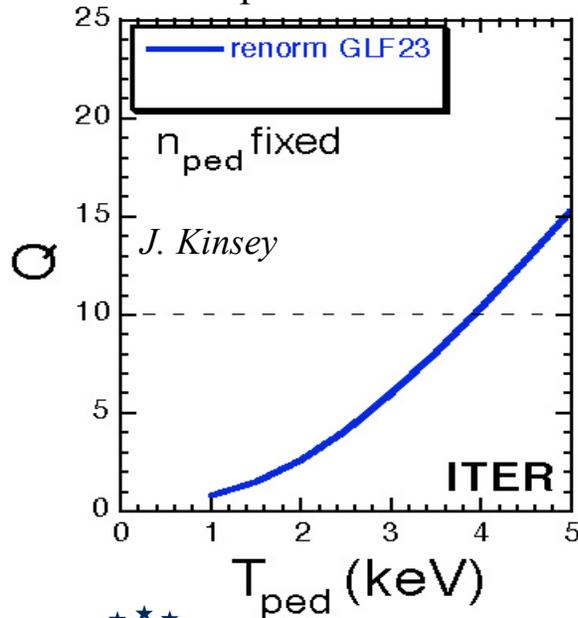


Motivation and Background

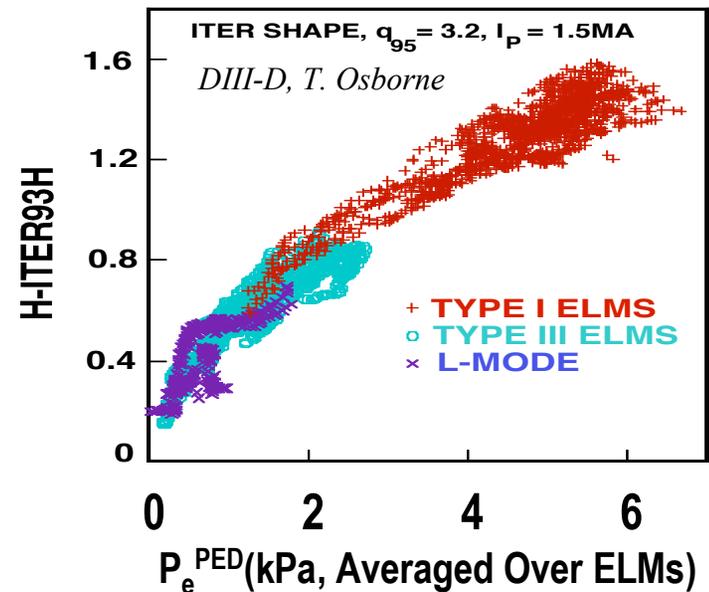
- ELMs and the edge pedestal are key fusion plasma issues
 - “Pedestal Height” controls core confinement and therefore fusion performance (Q)
 - ELM heat pulses impact plasma facing materials



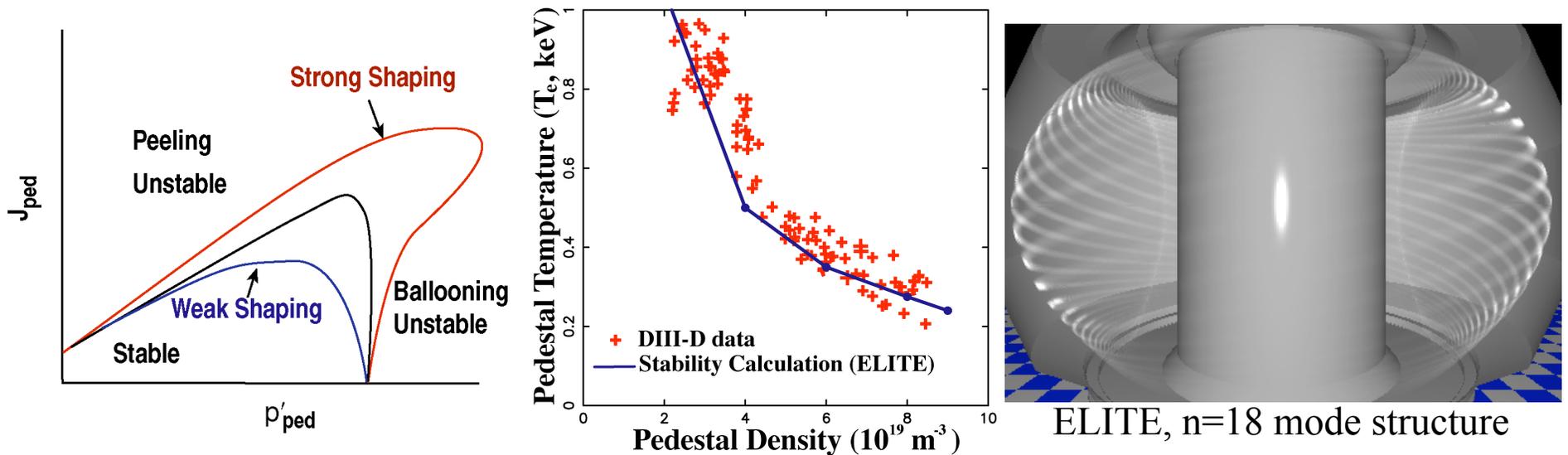
Predicted Impact of Pedestal Height



Observed Impact of Pedestal Height



Background: Extending the Peeling-Ballooning Model



- Peeling-Ballooning Model of ELMs - significant successes
 - ELMs caused by intermediate wavelength ($n \sim 3-30$) MHD instabilities
 - Both current and pressure gradient driven
 - Complex dependencies on v_* , shape etc due to bootstrap current and “2nd stability”
 - Successful comparisons to experiment both directly and in database studies
- Need to understand sources and transport to get profile shapes (“pedestal width”)
- Rotation and non-ideal effects to precisely characterize P-B limits, nonlinear dynamics for ELM size and heat and particle loading on material surfaces

Outline

- Toroidal Flow Shear
 - How toroidal rotation complicates ballooning theory (1D \Rightarrow 2D)
 - Eigenvalue formulation
 - Impact on peeling-ballooning modes in the tokamak edge region
- Nonlinear ELM Simulations
 - General challenges
 - 2 fluid reduced Braginskii (BOUT) simulation results
 - Expected peeling-ballooning characteristics in linear phase
 - Explosive, radially propagating filaments in nonlinear phase
 - Comparison to Observations
 - Proposals for dynamics of the full ELM crash
- Summary and Future Work

Ballooning mode theory with rotation

Static:

For large n and solutions $\sim e^{\gamma t}$, derive the ballooning equation:

$$L\left(\frac{\partial}{\partial \theta}, q'(\theta - \theta_0); \gamma(\theta_0)\right) \hat{\xi} = 0 \quad \text{A 2nd order ODE, 1D eigenvalue problem}$$

Higher order theory \Rightarrow **choose θ_0 to maximize $\gamma(\theta_0)$**

With sheared toroidal flow: $\mathbf{v} = R^2 \Omega(\psi) \nabla \phi$ $R\Omega / C_s \sim n^{-1} \ll 1$ $\frac{1}{q'} \frac{\partial(R\Omega / C_s)}{\partial \psi} \sim 1$

Using a time dependent eikonal approach

$$L\left(\frac{\partial}{\partial \theta}, q'(\theta - \theta_0 + \Omega' t / q'); \frac{\partial}{\partial t}\right) \hat{\xi} = 0 \quad \text{A 2D initial value problem}$$

Cooper, PPCF 30, 1805 (1988)

Low flow shear, separable solution \Rightarrow **average of $\gamma(\theta_0)$ over θ_0**

$$\gamma = \frac{1}{2\pi} \oint \gamma(\theta_0) d\theta_0 \quad \text{Waelbroeck and Chen Phys Fluids B3 601 (1991)}$$

- There is a **discontinuity in the theory**, which we would like to understand
- Suggests that **flow shear could in principle have a big effect on ballooning modes**

Flow shear and the Eigenmode Formalism

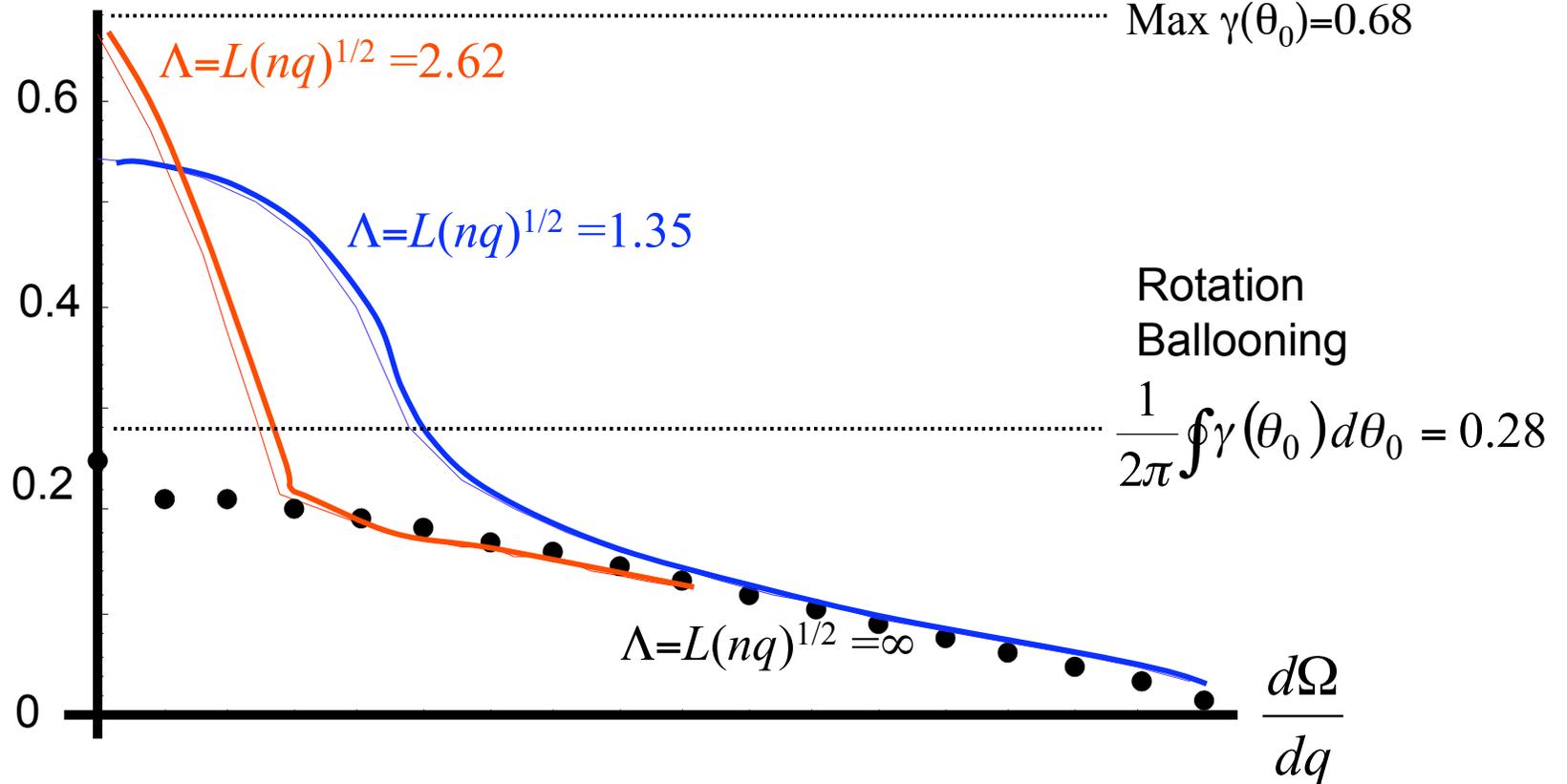
- Would like to develop an eigenmode formalism for the effect of flow shear on ballooning modes:
 - Smoothly connect to the conventional ballooning modes as $\Omega' \rightarrow 0$ and understand this 'discontinuity'
 - Calculate the radial eigenmode structure
 - Provides an eigenmode frequency
 - Enables consideration of finite n corrections
 - Permits flow shear to be incorporated into ELITE (an eigenmode code)
 - Evaluate impact on P-B modes in experimental equilibria
- Eigenmode formalism derived and implemented

Including finite- n via eigenmode formulation resolves small rotation discontinuity

A. Webster et al PRL 92 (2004) 165004

s - α geometry: $s=1.0$ $\alpha=1.7$

$\text{Re}(\gamma)$



Discontinuity resolved, transition from static ballooning slows with decreasing n

ELITE is a Highly Efficient MHD Stability Code for $n > \sim 5$

ELITE is a 2D eigenvalue code, based on ideal MHD (amenable to extensions):

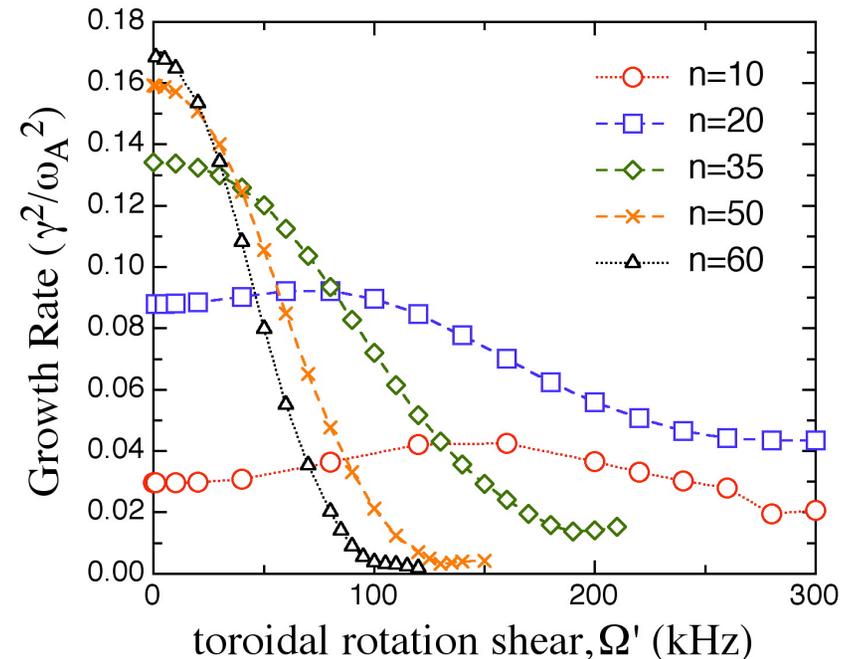
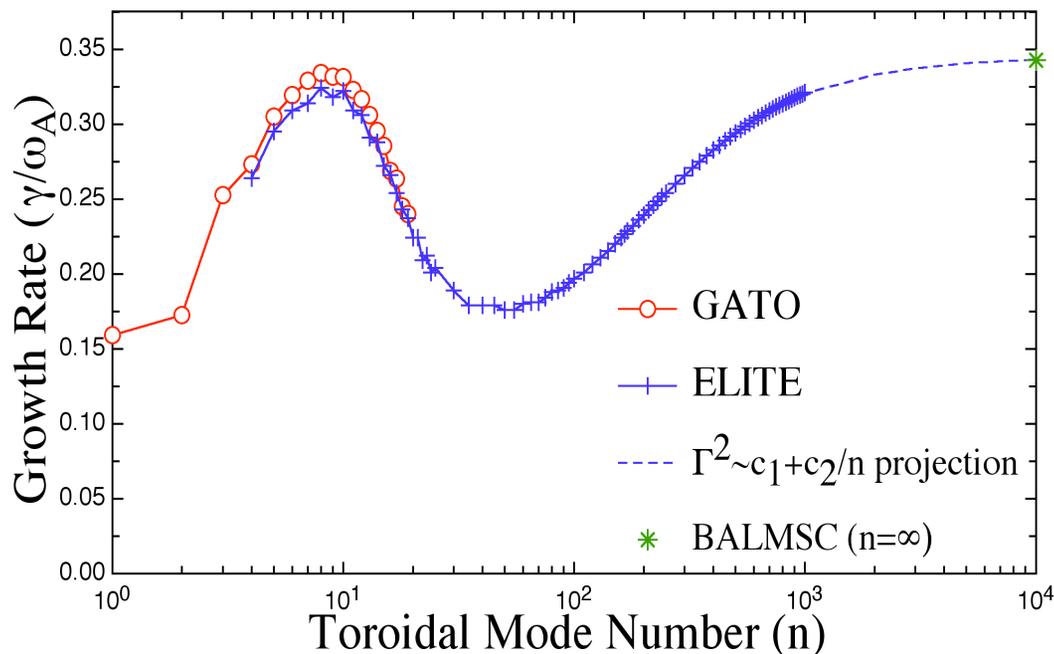
-Generalization of ballooning theory:

- 1) incorporate surface terms which drive peeling modes
- 2) retain first two orders in $1/n$ (treats intermediate $n > \sim 5$)

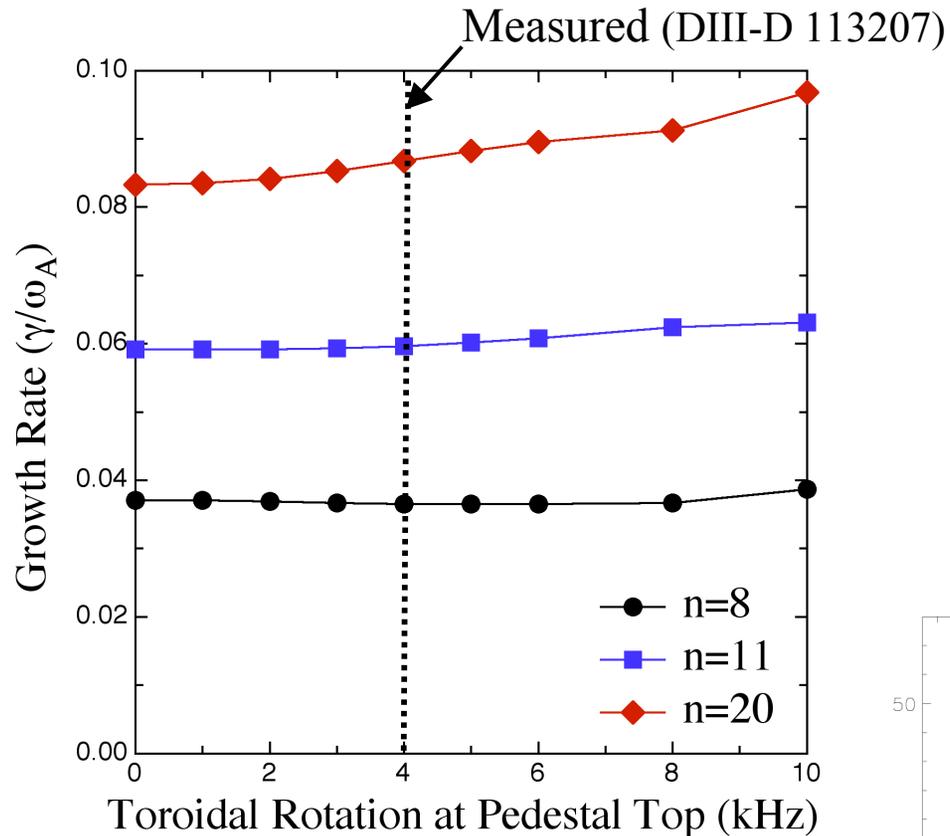
-Makes use of poloidal harmonic localization for efficiency

-Successfully benchmarked against GATO, MISHKA, MARS, BAL-MSC

-Code extended to include leading order ($n\Omega \sim 1$, $\Omega' \sim 1$) sheared toroidal flow and compression - results qualitatively similar to $s-\alpha$



Flow Shear Effect on Growth Rates is Modest in Standard ELMing Discharges, Mode Structure Does Change

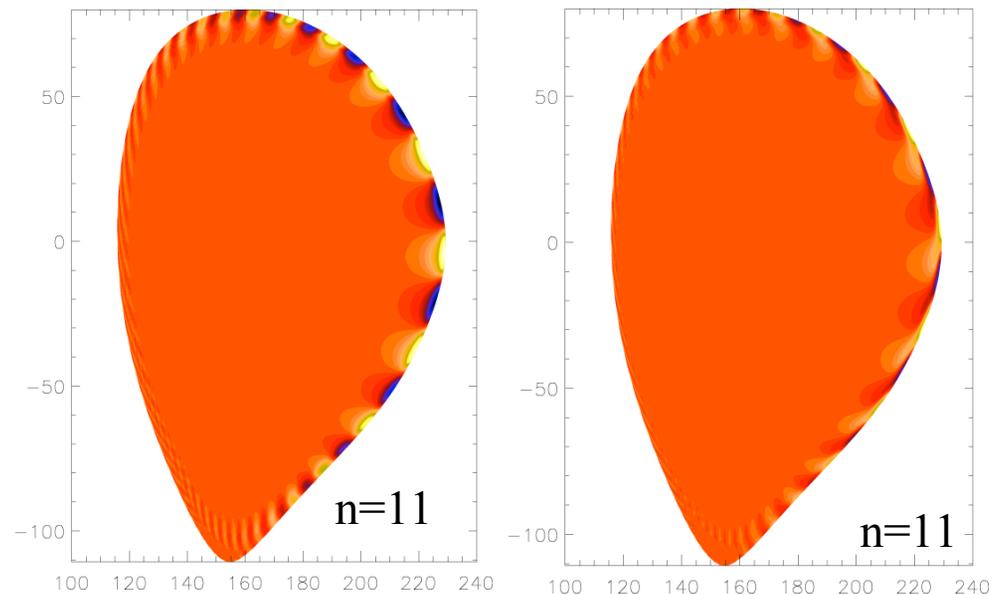


Rotation Shear on P-B Modes:

- Stabilizing near marginality
- Finite n and large γ dramatically reduce effect
- Does not measurably change expected ELM onset time in typical ELMing discharges

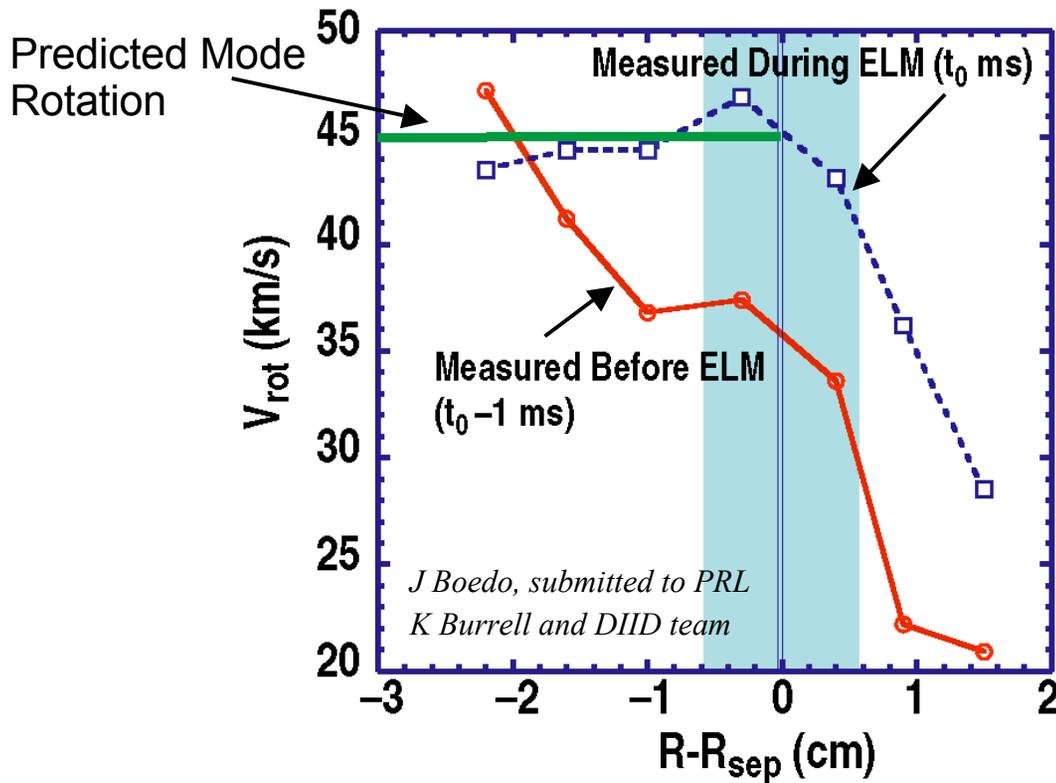
$\Omega_{\text{ped}}=0$

$\Omega_{\text{ped}}=10$ kHz

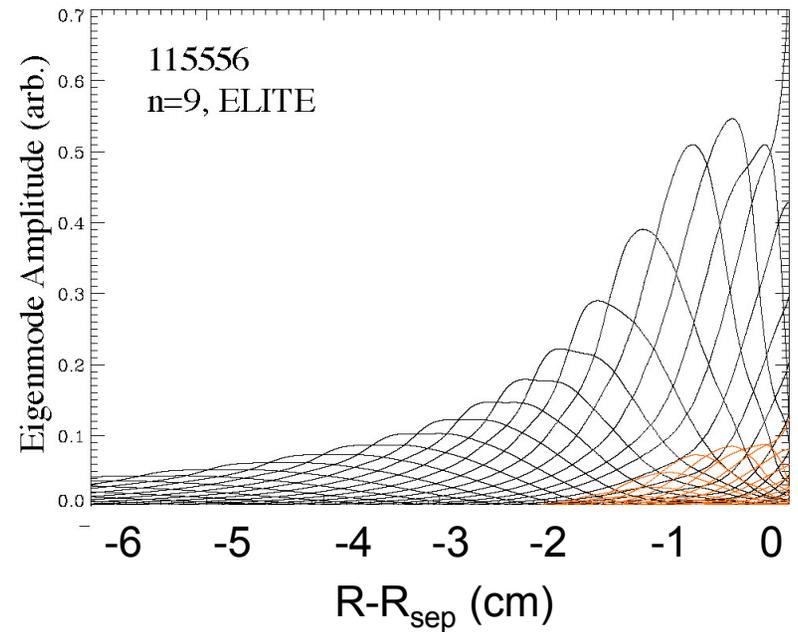


- Mode structure strongly altered
 - Narrowing and phase changes
 - May impact dynamics, ELM size

Calculated Mode Rotation Agrees with Observation during ELM



Calculated Structure of Most Unstable Mode



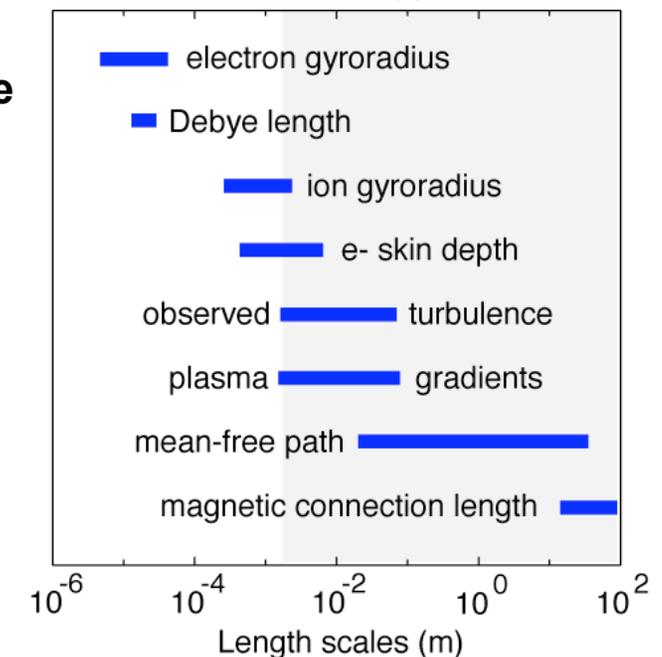
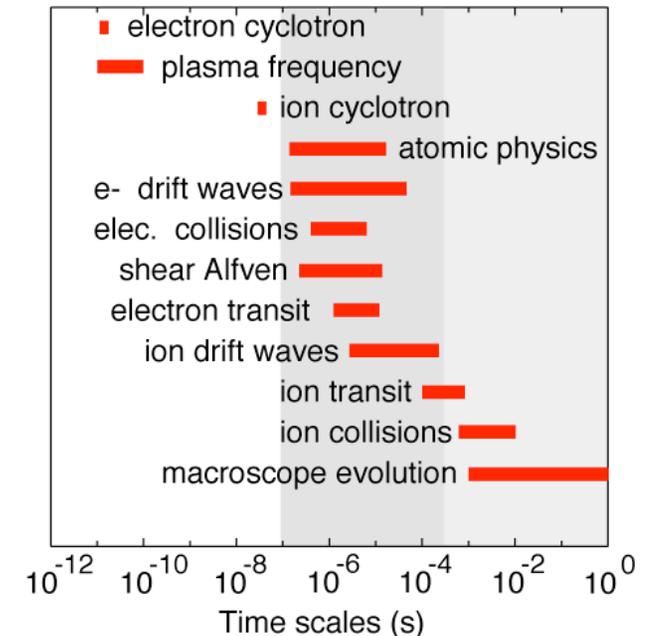
- Measured rotation profile strongly sheared just before the ELM, becomes \sim flat at \sim **45km/s** across pedestal region at ELM onset
- Study with ELITE finds peeling-ballooning unstable just before ELM - most unstable mode ($\max \gamma/\omega_*$) is $n=9$
- Calculated frequency for this $n=9$ mode is $\omega/\omega_A=0.0082$, $\mathbf{V_{rot}=45km/s}$
- Suggests “locking” of pedestal region to the mode during initial phase of ELM crash \Rightarrow edge barrier collapse

Summary of Toroidal Flow Shear Impact on P-B Modes

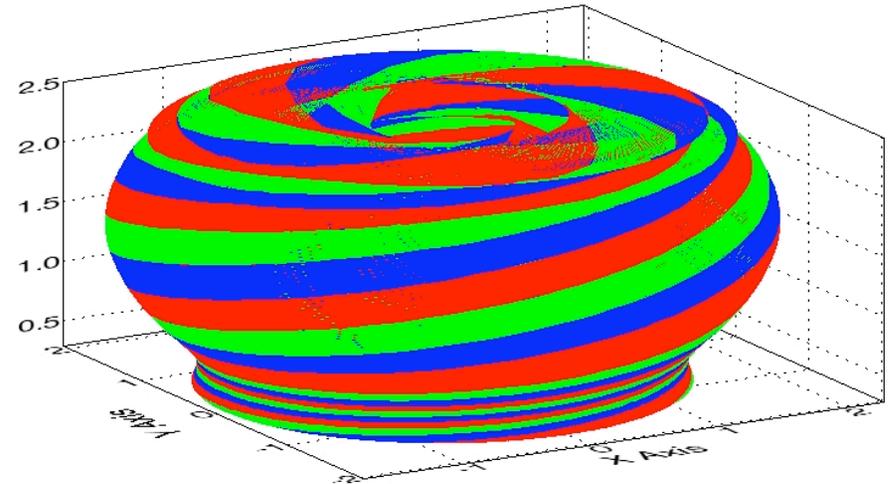
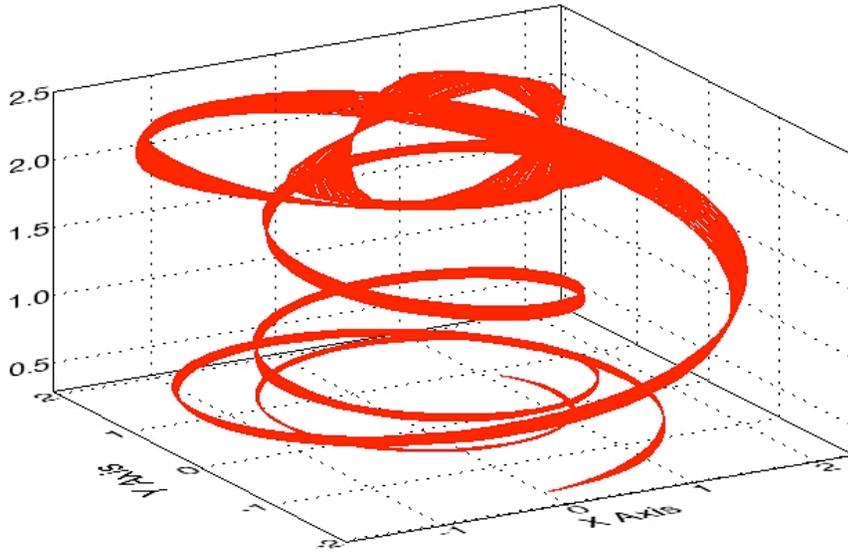
- Toroidal flow shear generally stabilizing at high n , effect reduced with decreasing n
- For experimental profiles:
 - Stabilization near marginal point, weak effect on growth rate away from marginal point (except for high n)
 - Slightly delay ELM onset time, and reduce most unstable n value
 - Effect stronger at low s (high Ω'/q'), e.g. where local shear is reduced by high bootstrap current (low v^* , high pedestal). May play a significant role in QH and “grassy ELM” regimes
- Substantial radial narrowing of eigenmode
- Mode eigenfrequency matches plasma Ω near top of pedestal
 - Observations suggest “locking” of bulk rotation during early ELM crash
- **Both of the above effects can have important impacts on the dynamics of the ELM crash**

Nonlinear Edge/Pedestal Simulations

- Many challenges for nonlinear simulations of the edge region
 - Broad range of overlapping scales and physics (L-H transition, sources and transport, ELMs, density limit..)
 - Many techniques used to simplify core simulations not applicable in edge
 - Long term goal is to unite full set of physics into massive scale simulations
- Here we focus on the fast timescales of the ELM crash event itself
 - Goal is to **understand physics determining ELM size and heat deposition**
 - Initialize with P-B unstable equilibria, evolve dynamics on fast timescales
- Reduced Braginskii 2 fluid simulations with the 3D BOUT code [X Q Xu et al Nucl Fus 42 21 2002]



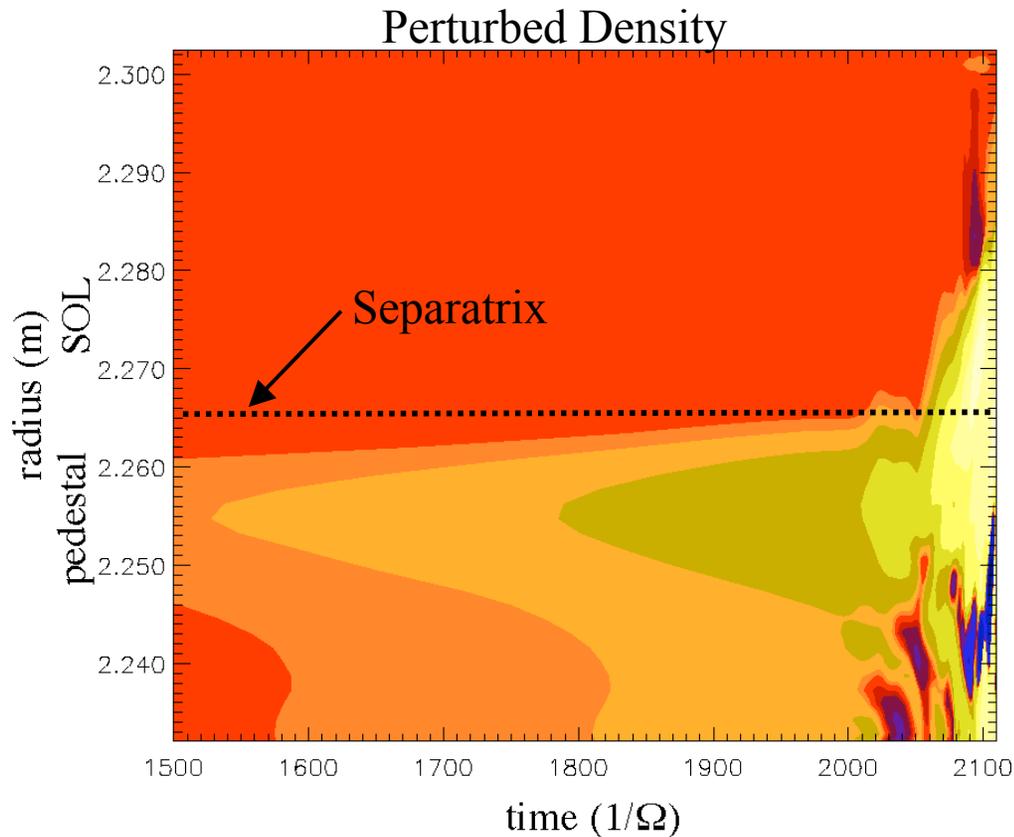
BOUT Simulation Geometry



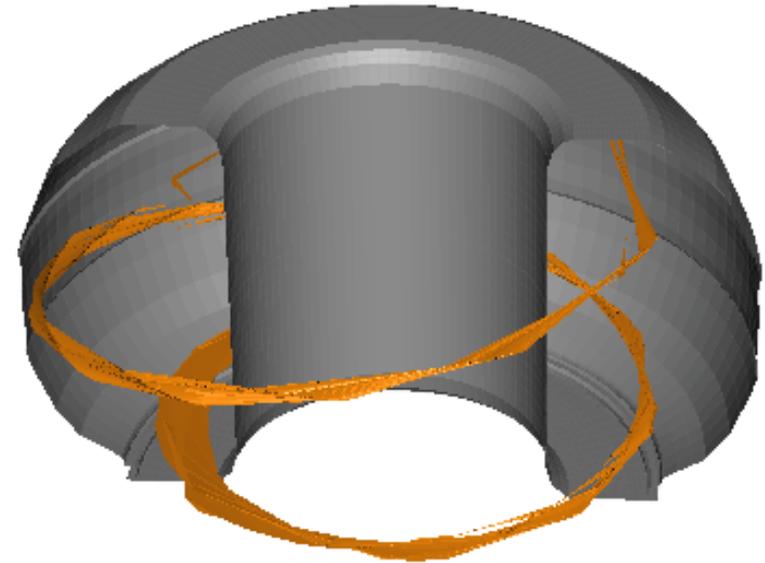
- BOUT incorporates 2 fluid/diamagnetic physics and uses field line following coordinates
 - Bundle of lines (left) wraps around 2π poloidally
 - A group of such bundles (right) spans the flux surface
 - For ELM simulations, generally go 1/5 (or 1/2) of the way around the torus, ie treat $\Delta n=5$ (or $\Delta n=2$), $n=0,5,10\dots\sim 160$, $0.9 < \Psi < 1.1$ both closed and open flux surfaces
 - Equilibrium current (kink term) added for ELM studies



Fast ELM-like Burst Seen in BOUT Simulations



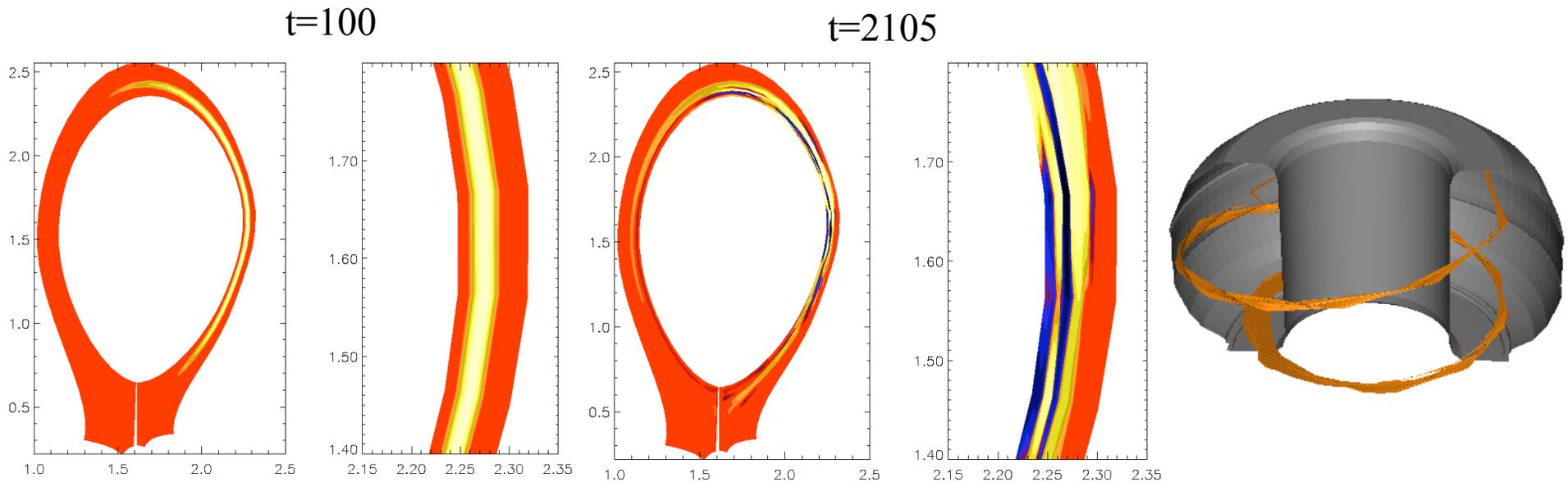
$t=2106$, surface of constant δn



- High density (small ELM), DIII-D LSN case, $0.9 < \psi < 1.1$
- Initial linear growth phase, then fast radial burst begins at $t \sim 2000$, can see positive density (light) moving into SOL and negative density perturbations near pedestal top
- Radial burst has filamentary structure, extended along B field



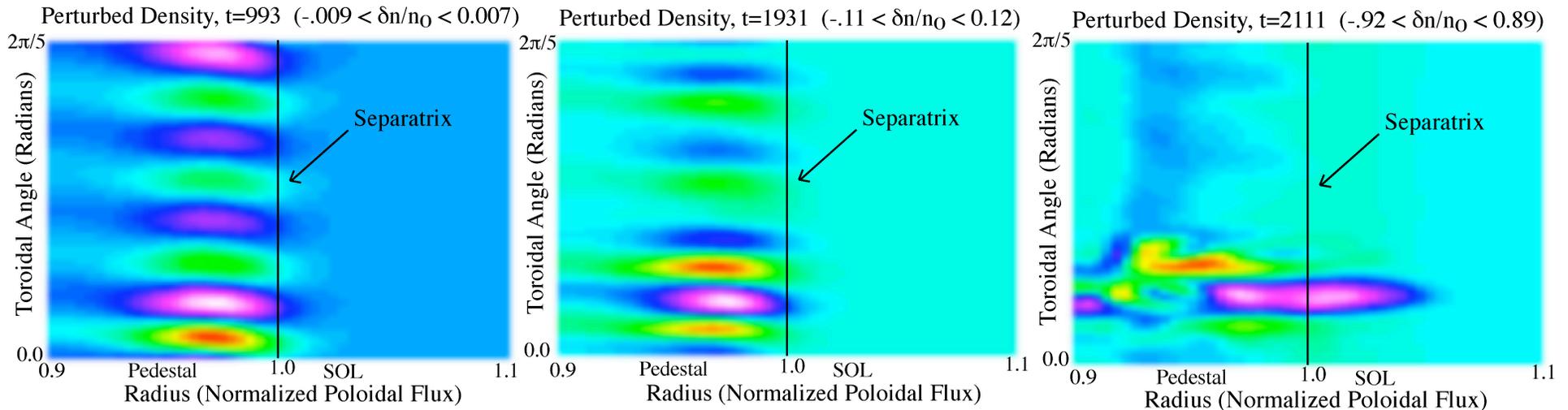
Expected Peeling-Ballooning Character in Early Phase



- Plots show projections of bundles of field lines onto the RZ plane - field lines extend into and out of page (radial vs parallel)
- Linear phase: Mode has ~expected characteristics of linear mode, radial and poloidal extent, $n \sim 20$, $\gamma/\omega_A \sim 0.15$
 - Reducing gradients slightly stabilizes the mode- abrupt onset near P-B boundary
- Fast Burst: Filaments extended along the field, but irregular



Fast ELM Burst Shows Toroidal Localization, “finger”

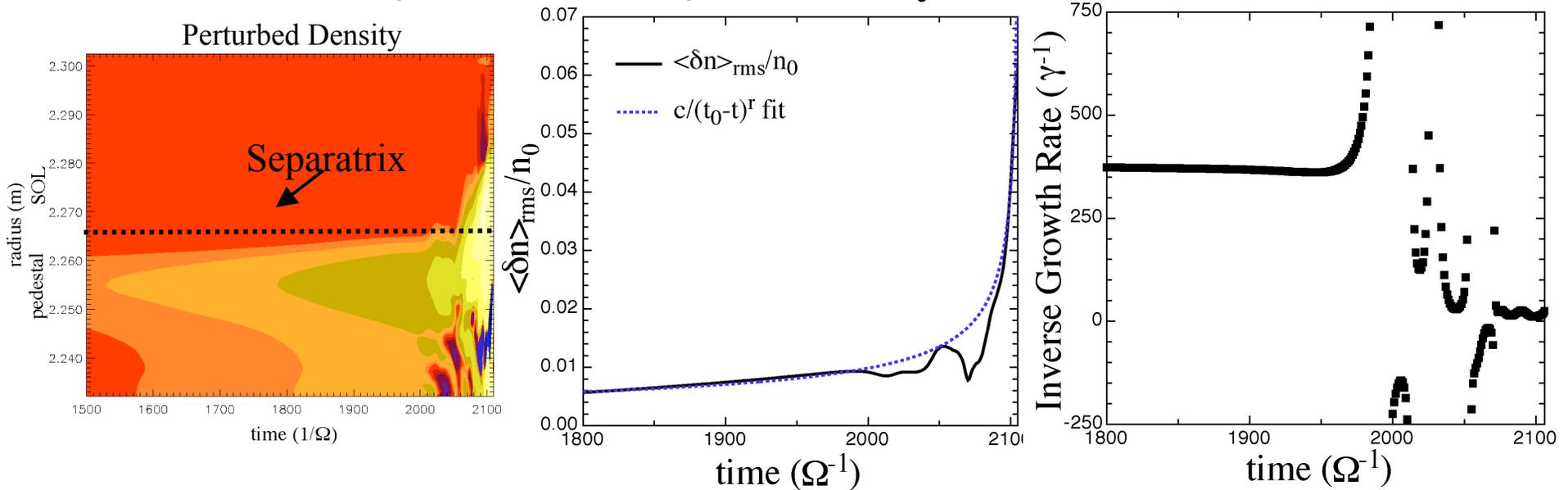


- R, ϕ plots on outer midplane
- Linear phase, $n=20$. Burst occurs asymmetrically at a particular toroidal location.
- Burst location is point of maximum resonance between dominant linear mode ($n=20$) and dominant nonlinearly driven beat wave
- “Finger” is an extended filament along the field, which propagates rapidly into the open field line region



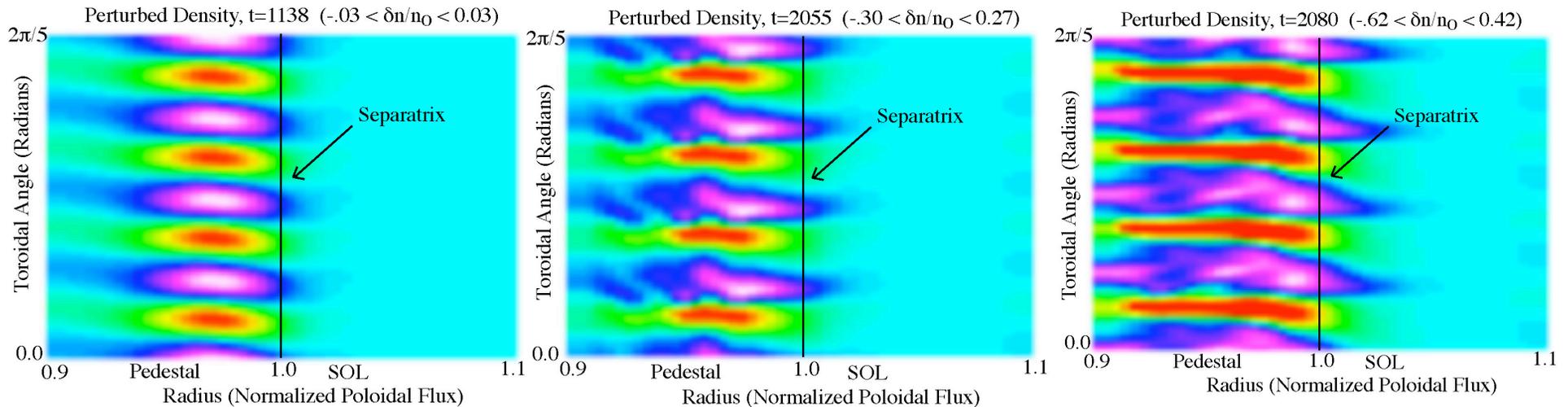
Similarities to Nonlinear Ballooning Theory

- Nonlinear ideal ballooning theory [Cowley and Wilson PRL **92** (2004) 175006] predicts explosive growth of a number of filaments
 - Nonlinear terms weaken field-line bending, accelerating mode growth
 - In nonlinear regime, perturbation grows like $\sim 1/(t_0-t)^r$



- Perturbed density in nonlinear simulations grows like $\sim 1/(t_0-t)^{0.5}$ (theory $r \sim 1.1$)
- Growth rate increases with time, increases rapidly during burst
 - Significant complexity, characteristic lull in growth rate prior to radial burst
 - Possible association with symmetry-breaking event when $\delta n \sim n_{0\text{loc}}$

One Filament or Many?

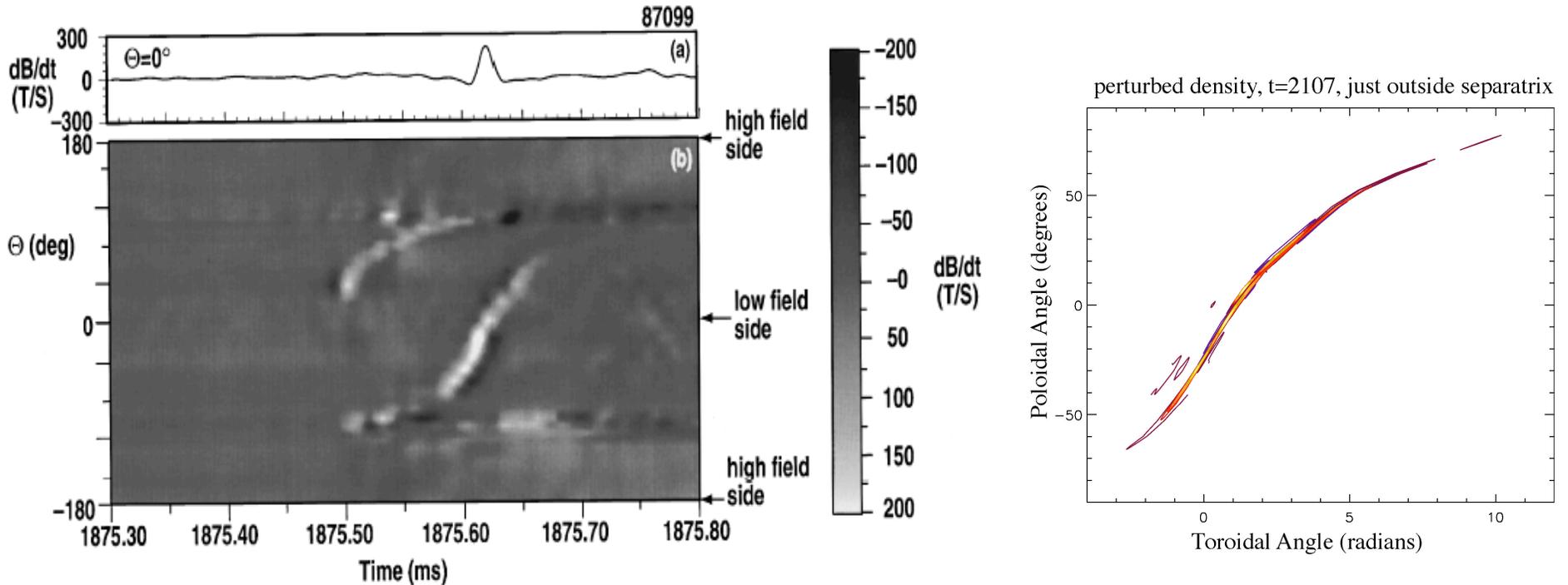


- Same case, initialized with a pure $n=20$ mode
 - Largely eliminates nearest neighbor coupling to generate beat wave
- Remains dominated by harmonics ($n=0,20,40,60,80$) well into nonlinear phase
- Burst occurs fairly symmetrically, multiple propagating filaments
- Evidence of secondary instability breaking up filaments

Both single and multiple filament cases are possible. Dependence on flatness of γ spectrum and rate at which profiles are driven across the marginal point.



Filaments Observed During ELMs



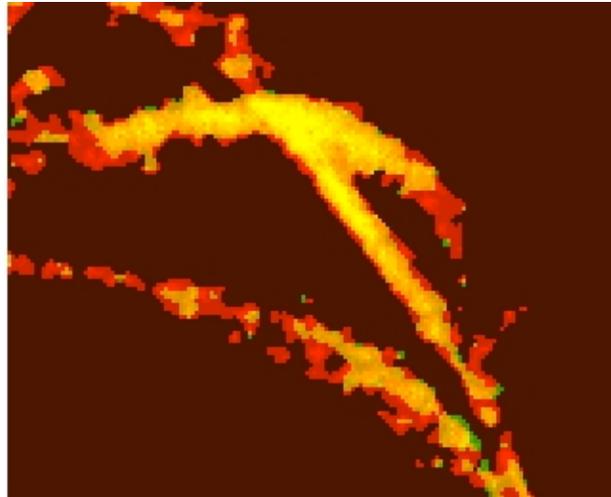
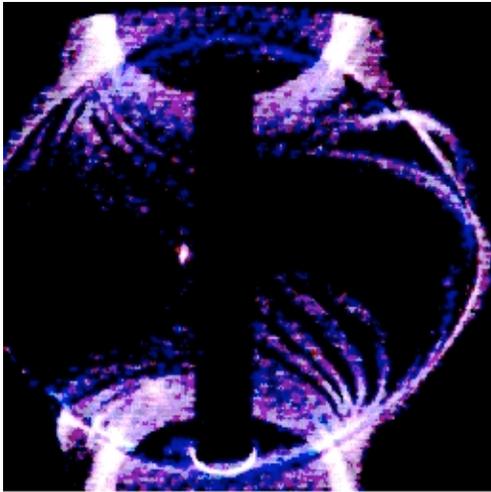
DIII-D Observation [E Strait, Phys Plas 1997]

3D Simulation

- Filament observed in fast magnetics during ELM (left)
- Finger-like structure from simulation (right) is extended along the magnetic field
- Qualitatively similar (rotation rate consistent with toroidal extent)



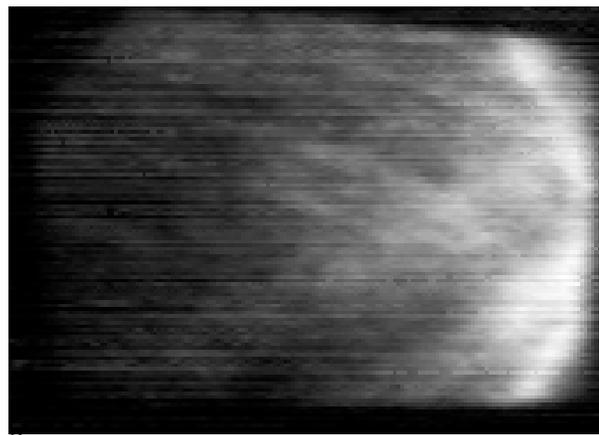
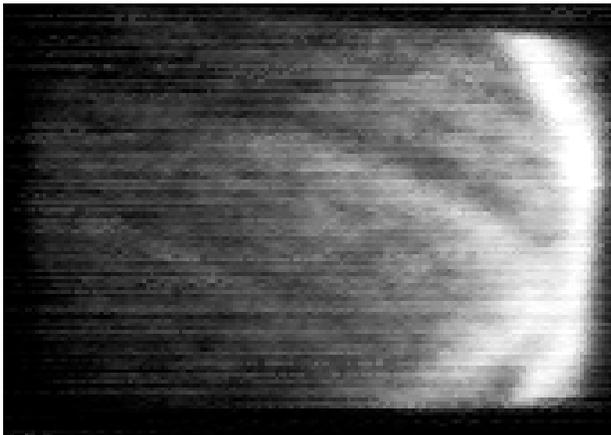
Fast ELM Observations: Multiple Machines



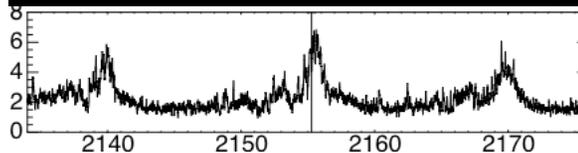
- $n=10$ structure on outboard side
- Filaments moving radially outward

A. Kirk, MAST, PRL 92 (2004) 245002-1

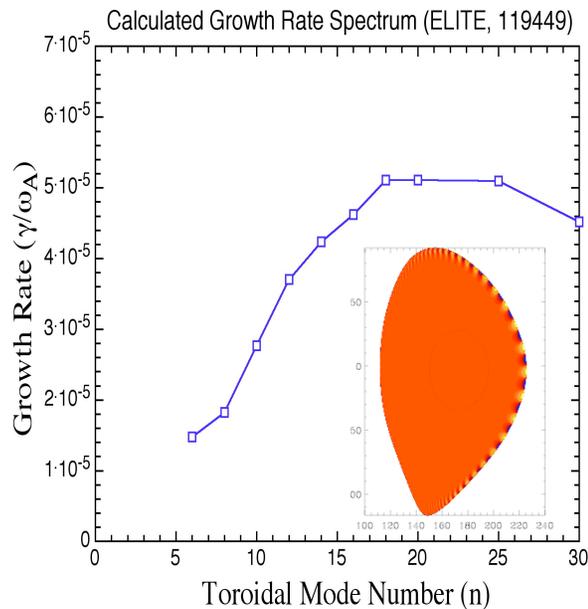
M. Fenstermacher, DIII-D, IAEA 2004



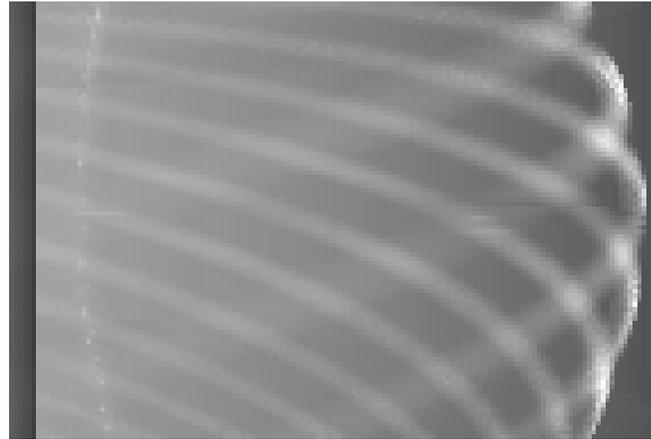
- CIII images from fast camera on DIII-D
- $n \sim 18$ inferred from filament spacing



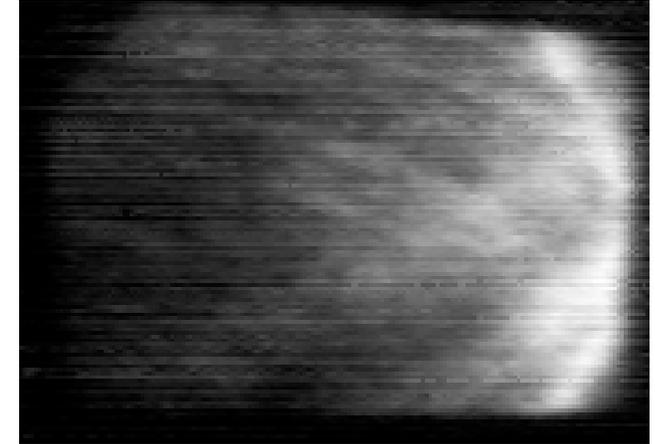
Fast CIII Images on DIII-D show filamentary structure



ELITE, n=18

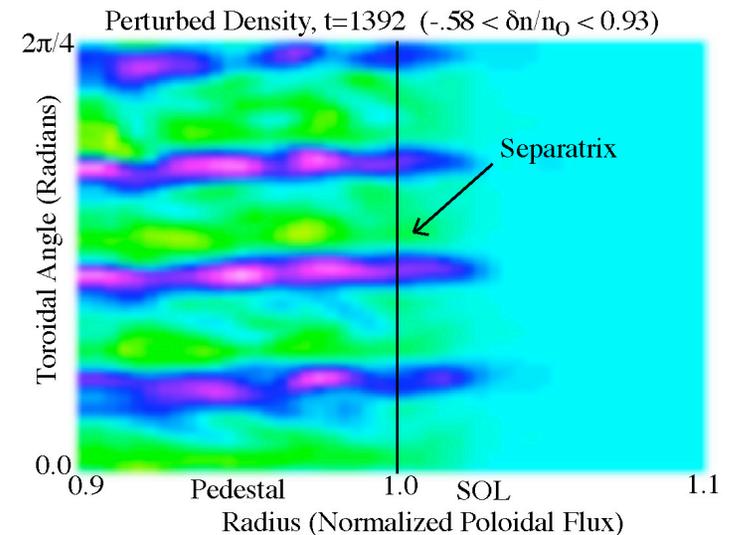


Fast CIII Image, DIII-D 119449
M. Fenstermacher, DIII-D/LLNL



- **ELITE linear P-B calculation on kinetic equilibrium shows peak $15 \leq n \leq 25$; mode in this range predicted to be first to go unstable**
- **Calculated structure of $n = 18$ mode similar to images**
 - Poloidal structure similar to outer midplane SOL structure in images
 - 3D structure has similar m/n structure seen in images
- **Nonlinear simulations show symmetric structure in early phase, extended uneven filaments later**

BOUT, nonlinear burst phase



Proposals for ELM Particle and Energy Losses

- Radially propagating filaments (one or many) carry only a small fraction of energy lost during an ELM
- Two possible mechanisms for the full ELM loss:
 - 1) **Conduits:** Heat and particles flow along filaments while ends remain connected to hot core. Fast diffusion/secondary instabilities allow flow across filament to open flux SOL plasma.
 - 2) **Barrier Collapse:** Radial eruption of filament (with fixed eigenfrequency) damps sheared rotation as it moves outward, collapses sheared E_r and edge transport barrier. Temporary return to L-mode-like transport+. Reduced gradients re-stabilize mode, allowing shear and pedestal to be re-established. [E_r well collapse during ELM observed on DIII-D, see Wade CI2A.002]
- Possible that both mechanisms are active. Collisional restriction of heat flow along filaments may explain transition to convective ELMs at high collisionality.

Summary

- Peeling-ballooning model has achieved a degree of success in explaining pedestal constraints, ELM onset and a number of ELM characteristics
 - Extend to include rotation and nonlinear, non-ideal dynamics
- Toroidal rotation shear included in ELITE
 - Discontinuity in previous studies removed via eigenmode formulation
 - Small effect on predicted ELM onset, but significant modification of mode structure
 - Real frequency of mode matches plasma rotation near center of mode
 - Encouraging comparisons with fast CER observations, suggests mode damps flow shear
- 3D nonlinear ELM simulations carried out with BOUT
 - Early structure and growth similar to expectations from linear Peeling-Ballooning
 - Radially propagating filamentary structures, grow explosively
 - One or many filaments possible (dependence on spectral shape and heating rate; single filament due to resonance of lin & nonlin modes)
 - similar to observations (eg MAST and DIII-D), and nonlinear ballooning theory
 - Filaments acting as conduits and collapse of the edge barrier provide possible mechanisms for full ELM particle and energy losses

Future Work

- Extend duration of existing simulations, test proposals for ELM losses, compare to expt
- Move on to larger problems:
 - 1) Toroidal scales – For some types of ELMs, need full torus ($n=1$ to $\sim\rho_j$)
 - 2) Radial scales – extend to wall and further into core
 - 3) Time scale – Include sources and drive pedestal slowly across P-B boundary
- Scale overlap and close coupling with pedestal formation (L-H) physics, inter-ELM transport and source (including atomic) physics
- Need optimal formulations (collisionless), efficient numerics and large computational resources

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