



## Effect of B-field Dependent Particle Drifts on ELM Behavior in the DIII-D Boundary Plasma,\*

M.E. Fenstermacher (LLNL)



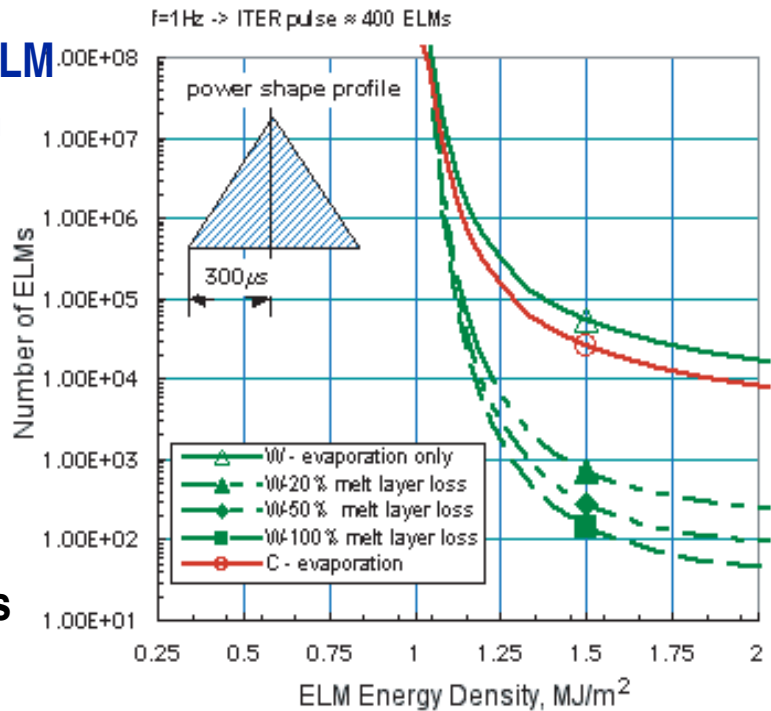
A.W. Leonard, T.W. Petrie, General Atomics,  
M. Groth, C.J. Lasnier, G.D. Porter - LLNL  
J. Boedo, D.S. Gray, E. Hollmann, UCSD  
M.R. Wade, ORNL, L. Zeng, UCLA  
J. Watkins, SNLA,

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# Motivation - Type I ELMs could limit PFC lifetime in burning plasma tokamaks - “Minimum Energy” ELMs at high $n_e$ might be tolerable

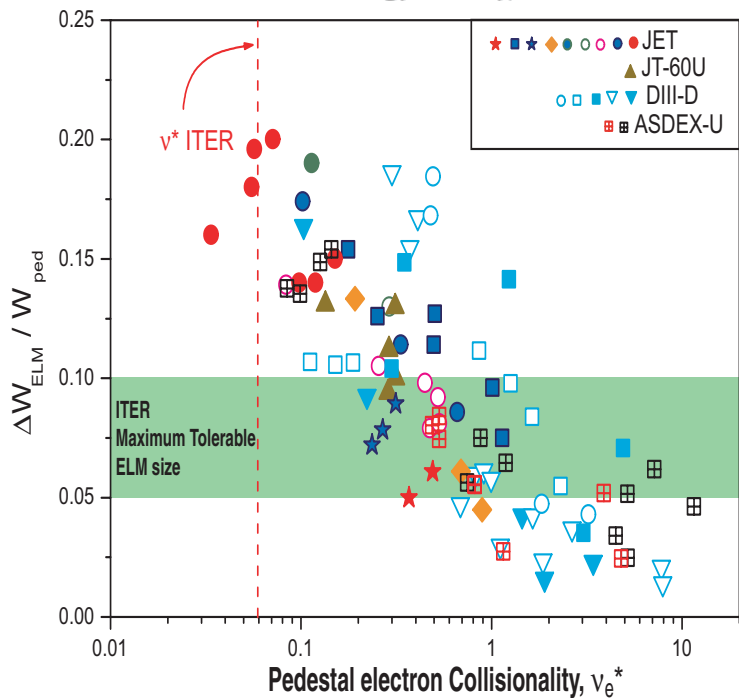
- To design future PFCs requires knowledge of ELM energy, spatial profile on surfaces and deposition time.

- Pedestal stability-key to reducing ELM sources
- SOL and divertor effects-key to reducing ELM fluxes on PFCs



- ELMs behavior in the boundary very different at low and high  $n_e$

- Midplane pedestal perturbation different
- Difference in divertor plasma response complex because pre-ELM divertor conditions strong function of density



# Summary I: SOL / divertor response to ELMs is a strong function of $n_e$ for LSN $\nabla B$ ↓

- ELM expelled pedestal particles appear far out (~ 4 cm) in SOL;  $T_e^{\text{SOL}}$  not perturbed even with  $\Delta T_e^{\text{ped}}$  at low  $n_e$ .
- Pedestal refueled by multi-step charge exchange neutrals; fast response consistent with local neutral source during ELM (main chamber surfaces)
- At low  $n_e$ : inner leg burns-through during ELM; large heat flux observed
- At high  $n_e$ : outer leg during ELM:
  - Carbon radiation burns-through to near target
  - Large particle flux increase
  - Target electron density and  $D_\alpha$  drop
  - No heat flux observed
- Rapid rise in midplane and divertor  $D_\alpha$ , and target  $j_{\text{sat}}$  at thermal energy loss
- SOL parallel pulse propagation times consistent with ion sound speeds at moderate - high density
  - Inner  $D_\alpha$  delayed ~ 250  $\mu\text{sec}$  after outer  $D_\alpha$
  - Midplane to divertor radiation pulse propagation times ~ 100  $\mu\text{sec}$

# Summary II: SOL / divertor response to ELMs is a strong function of $B_T$ direction

- $D_\alpha$  and  $P_{\text{rad}}$  Timing during ELMs:
  - LSN  $\nabla B \downarrow$  :
    - **Low  $n_e$** : Delay of inner vs. outer divertor  $D_\alpha$  reduced below ion convection times in SOL,  $P_{\text{rad}}$  delay negative (inner occurs before outer)
    - **High  $n_e$** : Delays consistent with ion convection timescale
  - LSN  $\nabla B \uparrow$  :
    - **High and low  $n_e$** : Delay below ion convection timescale
- Heat Flux during ELMs:
  - LSN  $\nabla B \downarrow$  **High  $\delta$**  : Peak inner / outer heat flux asymmetry  $\leq 2x$  at **low  $n_e$** , even larger at **high  $n_e$** 
    - Outer ELM peak heat flux reduced with  $n_e$
  - LSN  $\nabla B \downarrow$  **Low  $\delta$**  : Peak inner / outer heat flux asymmetry  $\sim 2x$  at **low  $n_e$** , drops to  $\sim 1.5$  at **high  $n_e$**
  - LSN  $\nabla B \uparrow$  **Low  $\delta$**  : Peak inner / outer ratio  $\sim 2$  independent of density
  - Surface layer effects may play a role in measurement

# Summary III: Initial UEDGE ELM modeling with drifts shows features of $B_T$ dependence

- **Model assumptions guided by measurements**
  - Midplane instability and particle loss appear for 200-500  $\mu\text{sec}$  before pedestal thermal energy loss
- **Model Verification and Fluid Simulations**
  - ELM energy transport by parallel ion convection at ion sound speeds verified by measurements
  - Initial UEDGE simulations show characteristics of ELM propagation at ion sound speed
    - Delays of inner  $D_\alpha$  from outer  $D_\alpha$  timing
    - Slower  $D_\alpha$  rise time in inner vs outer divertor
- **Some of  $B_T$  dependence consistent with each of two models:**
  - Changing  $E_x B$  produces vastly different pre-ELM divertor conditions  $\implies$  ELM response is different even though  $E_x B$  and other particle drifts not playing a role during ELM event
  - $E_x B$  particle drift play a strong role during ELM due to large  $T_e$  gradients (E-field) created by ELM perturbation

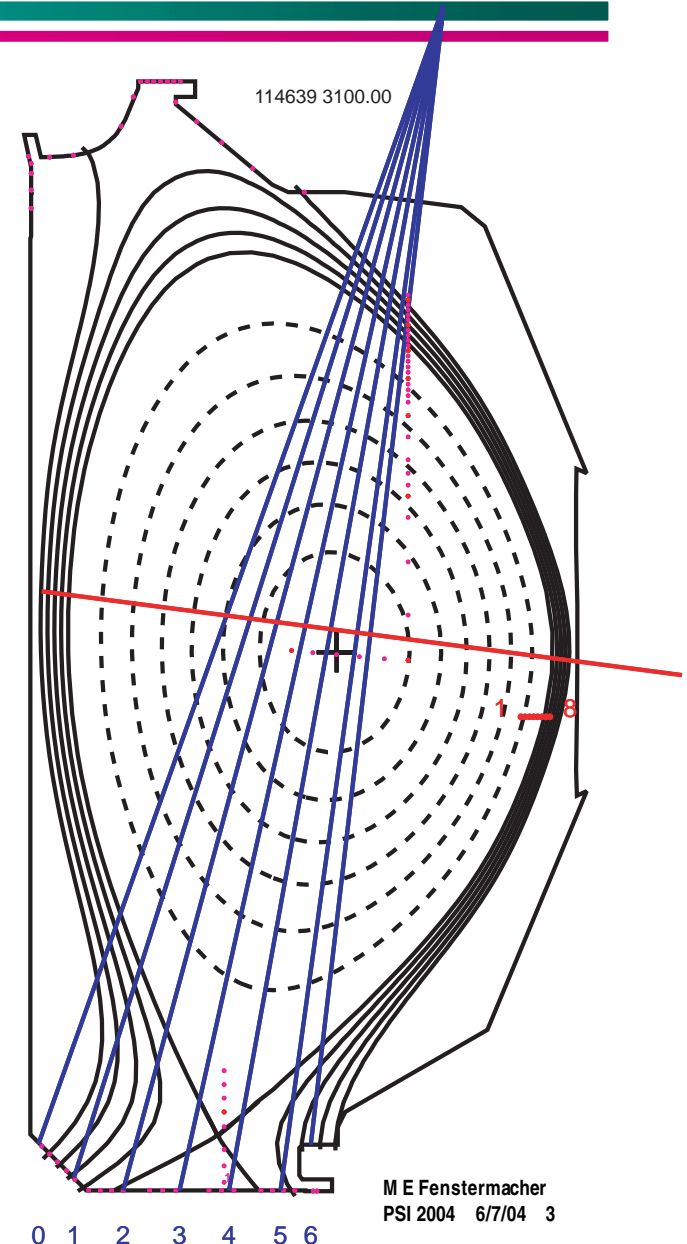
# Configuration and Diagnostics

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# DIII-D fast diagnostics used in this poster cover both the outer midplane and lower divertor

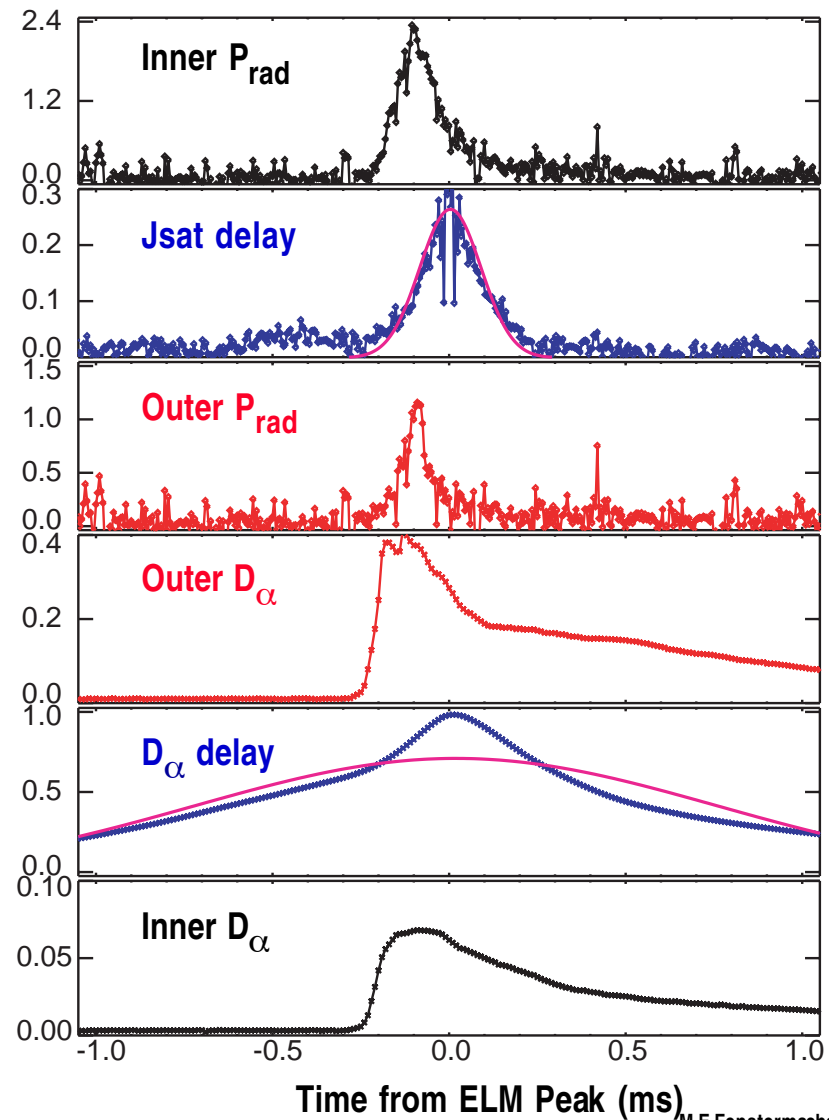
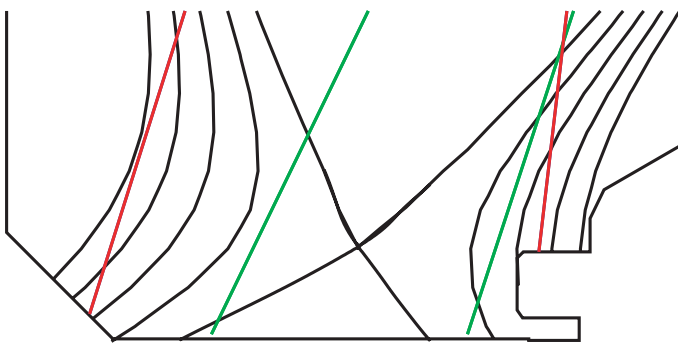
Parameter	Fast Diagnostic	Rate / Integration time
SOL $n_e$ , $T_e$ profiles	Reciprocating probe	$\leq 1000$ kHz
Pedestal $n_e$ , $T_e$	Thomson scattering	1 ns @ 6 ms
Midplane $D_\alpha$	Filterscopes array	$\leq 100$ kHz
Midplane inner SOL line radiation	Gated, intensified camera	20 us @ 17 ms
$n_e^{\text{ped}}$ gradient	Reflectometry	$\leq 10$ kHz
Total radiated power	Bolometer array	$\leq 500$ kHz
Divertor line radiation	Gated, intensified camera	20 us @ 17 ms
Target heat flux	IRTV line scan	$\leq 9$ kHz
Target ion flux	Target probes	$\leq 100$ kHz
Toroidal target current	Tile current array	$\leq 200$ kHz
Calibrated divertor line radiation	Filterscopes array	$\leq 100$ kHz
Divertor line density	Interferometer	$\leq 50$ kHz
Edge ion temperature	CER	$\leq 2$ kHz



# Cross correlation analysis finds delay of ELM response in one signal compared with another.

- Cross correlation of signals applied in  $\sim 8$  ms window centered on the ELM event at midplane
- Delays of peak response dominates over delay of initial response.

Geometry of divertor  $D_a$  and  $P_{rad}$  Chords





# Average ELM behavior in $\nabla B \downarrow$ vs $\nabla B \uparrow$ shows changes in inner vs outer asymmetry but similar timing

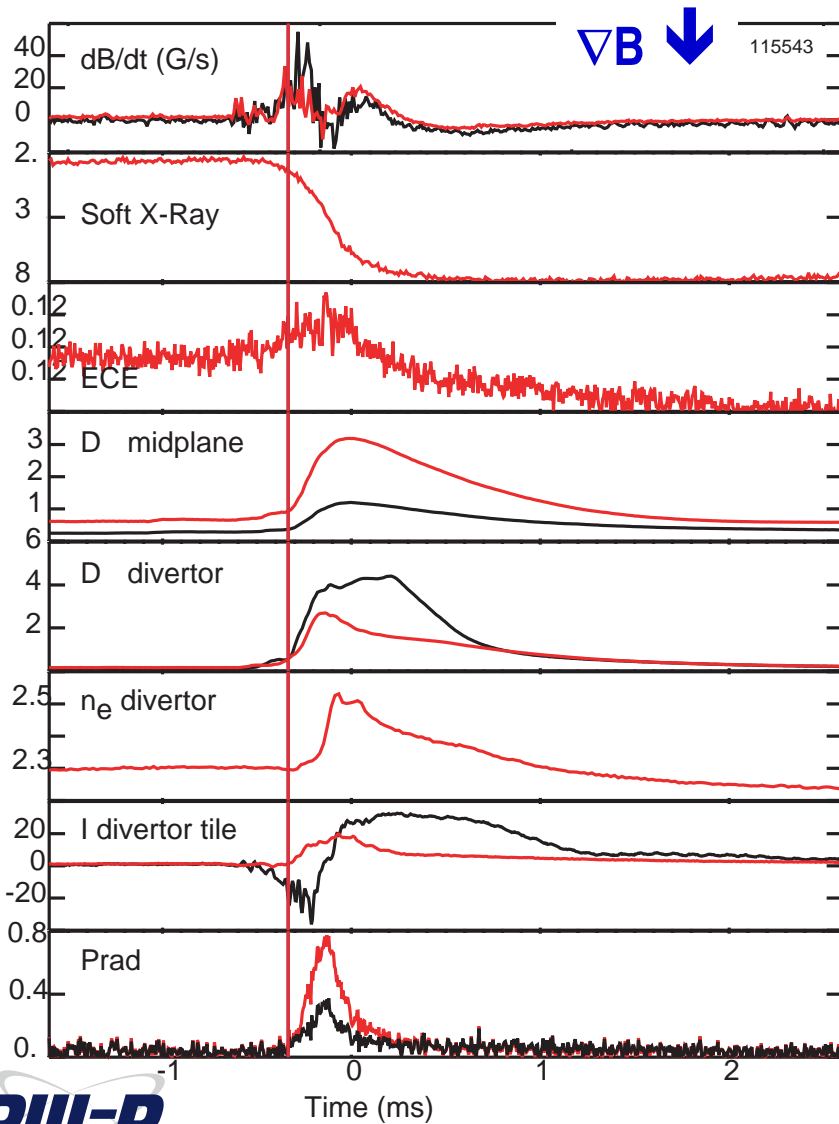


Fig. 1

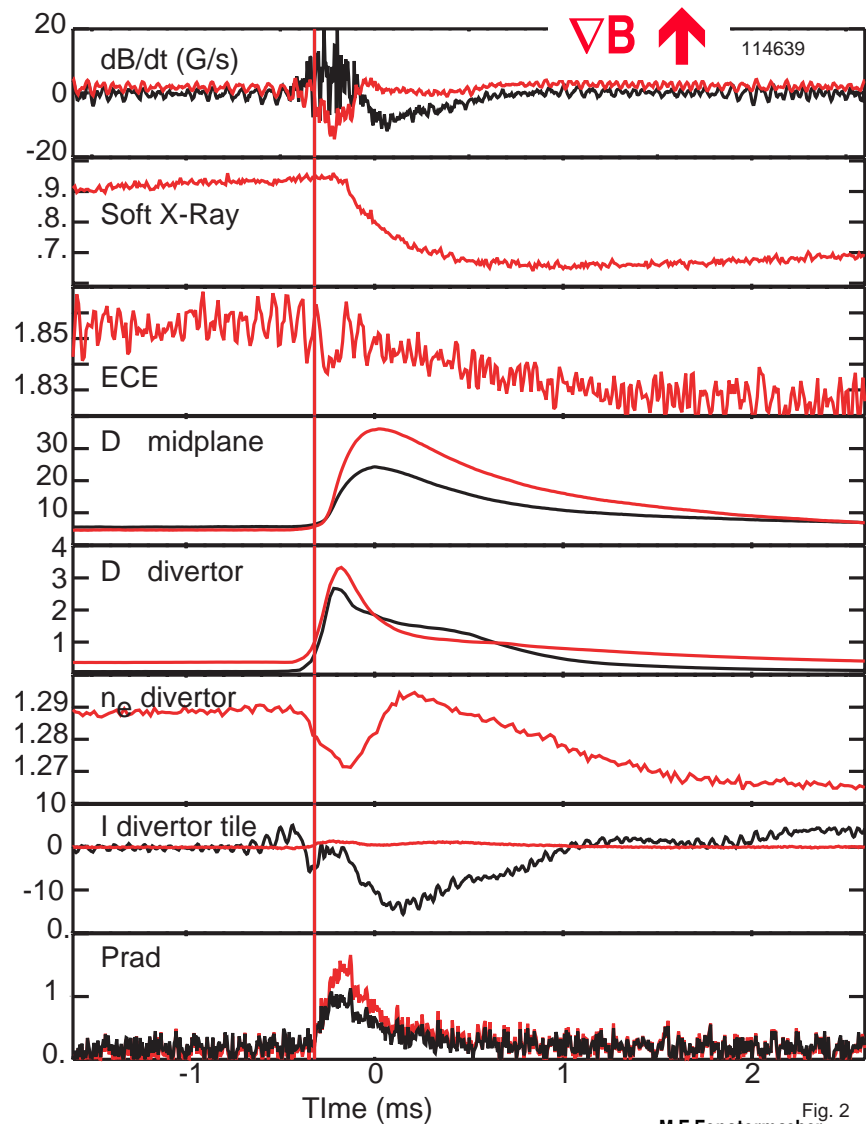


Fig. 2

# Background: LSN $\nabla$ B $\downarrow$ from 2002

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# Conclusions: Model of SOL ELM propagation by ion convection supported by some, but not all, of the ELM data

- **Model says:**
  - Deposition profile should be set by perpendicular vs. parallel transport in SOL
  - Deposition time set by  $L_{||} / C_s$  for ELM expelled ions
  - ELM energy may be limited if ELM duration < ion transit time to targets
- **Model supported by data:**
  - Density dependence of inner vs. outer target delays
  - $\Delta T_e^{\text{ped}}$  delay until  $\Delta t_{\text{ELM}} \sim L_{||} / C_s$
  - Divertor density rise higher than  $n_e^{\text{ped}}$  due to release of trapped neutrals
- **Model not supported by data:**
  - Some inner vs. outer SOL delays backwards ( eg.  $P_{\text{rad}}$ ,  $J_{\text{sat}}$  )
  - Outer target heat flux width not wide enough to be consistent with observed midplane density perturbation in the far SOL, even narrower on the inside
  - Fast  $T_e^{\text{ped}}$  drop in low density case - more like reconnection
- **Comparison of ELM propagation with ion  $B_x \nabla B$  drift into vs. out of divertor should increase understanding of ELM propagation physics.**

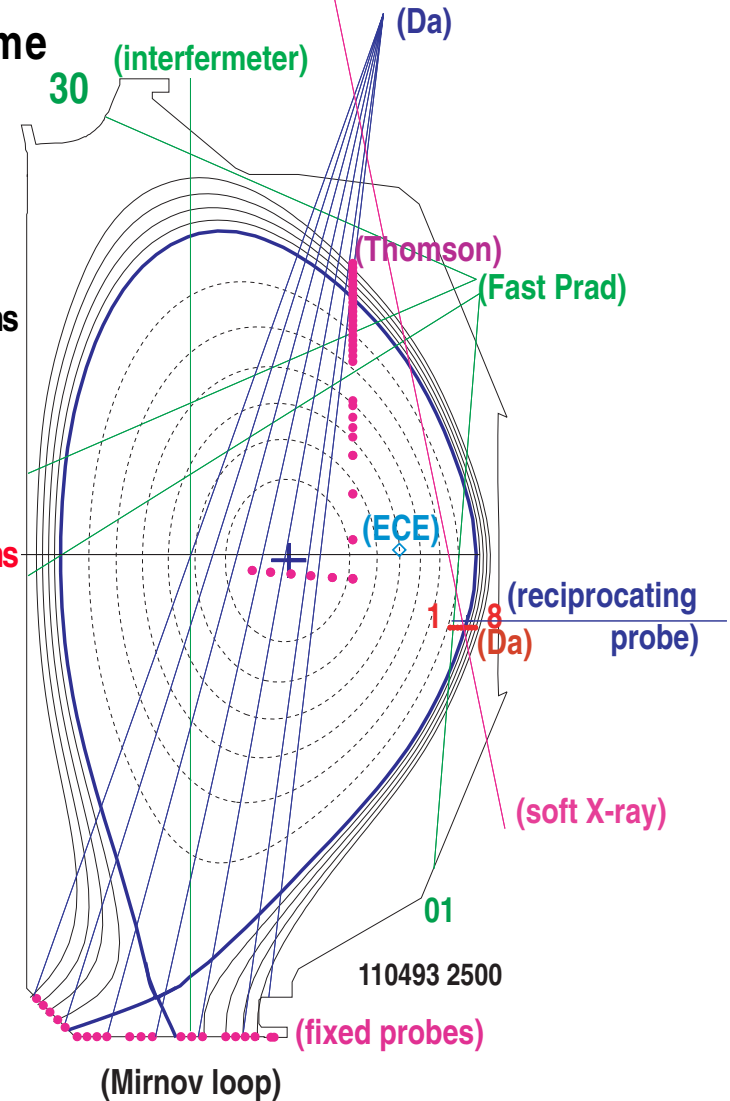
# Backup Slides

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# Simple model of ELM particle and energy transport in the SOL and divertor supported by calculation results\*

- Instability flattens density and temperature profiles (electrons and ions) at the outer midplane separatrix
- Fast electrons on field lines connected to targets go to targets in electron transit time (~ several  $\mu\text{sec}$ )
  - Sheath potential raised and electron conduction gets cut-off
  - $T_e$  in SOL equilibrated somewhat
- Local ions in sheath strike target at high energy - take out some fraction of ELM electron energy (~ 10  $\mu\text{sec}$ )
- ELM expelled ions transit to elevated sheath at ion sound speed ( $T_i^{\text{ped}}$ ) ~ several 100  $\mu\text{sec}$ 
  - ELM ions falling through sheath remove ELM electron and ion energy
- Neutrals from increased recycling of ELM ions dissipate in recycling time scale (~ several ms).

\* A. Bergmann 2002 submitted to NF  
D. Tskhakaya PSI02 submitted to JNM  
T. Rognien PSI02 submitted to JNM

# Complicating effects may be important in SOL / divertor ELM transport

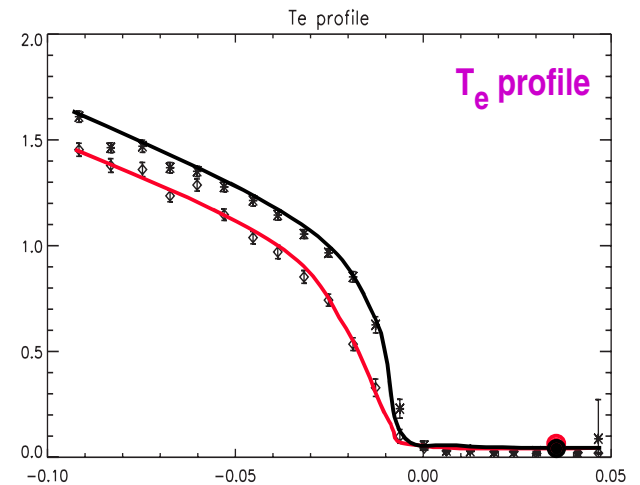
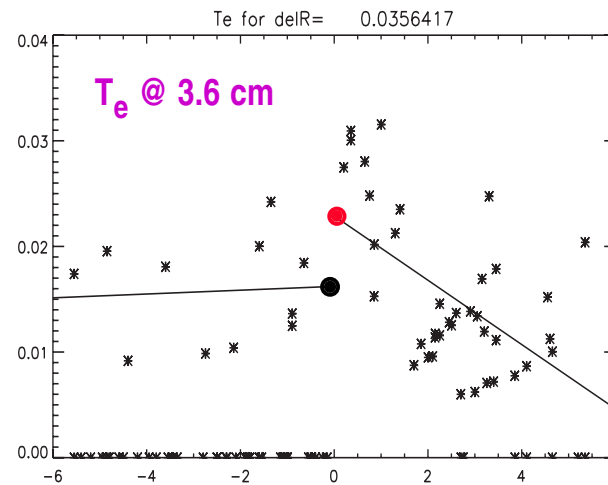
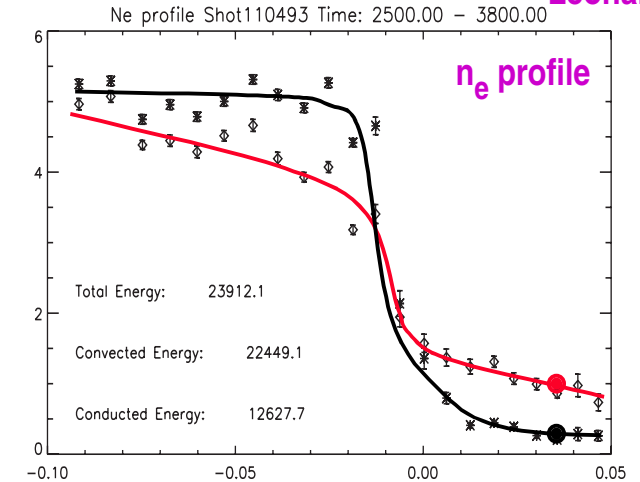
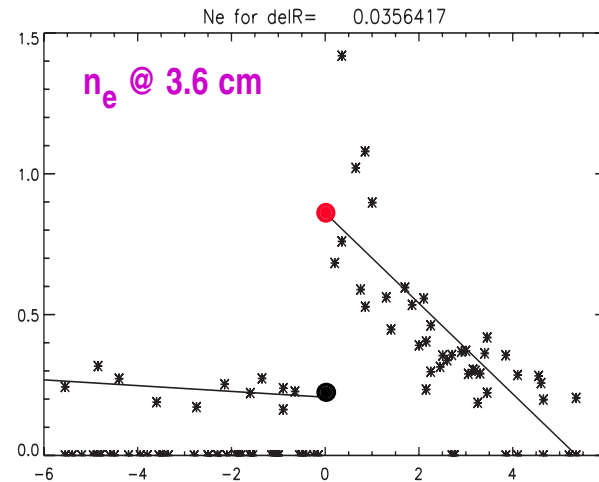
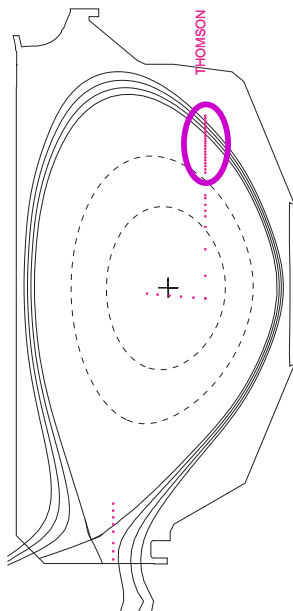
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- More ELM electron energy may get to targets on short time scale if:
  - Secondary electron emission at targets reduces sheath build-up
  - High energy ions striking targets liberate trapped neutrals increasing local ion source
- Perpendicular transport in upper SOL may reduce ions available to carry ELM energy to targets
- Impurity release by fast ion physical sputtering on targets produces radiation
- Loss of pedestal thermal energy ( $\Delta T_e^{\text{ped}}$ ) may require instability duration  $>$  ion transit time to targets
  - $\Delta T_e^{\text{ped}}$  may not occur until instability has been growing for an ion transit time
  - If ion transit time is long,  $\Delta T_e^{\text{ped}}$  may not occur at all

# Low $n_e$ ELMs: Thomson profiles show particles lost from pedestal appear in the far SOL; pedestal $\Delta T_e$ not seen in SOL

Leonard

- Linear extrapolation to ELM time may underestimate perturbation in the SOL



Time,  $t - t_{ELM}$  (ms)

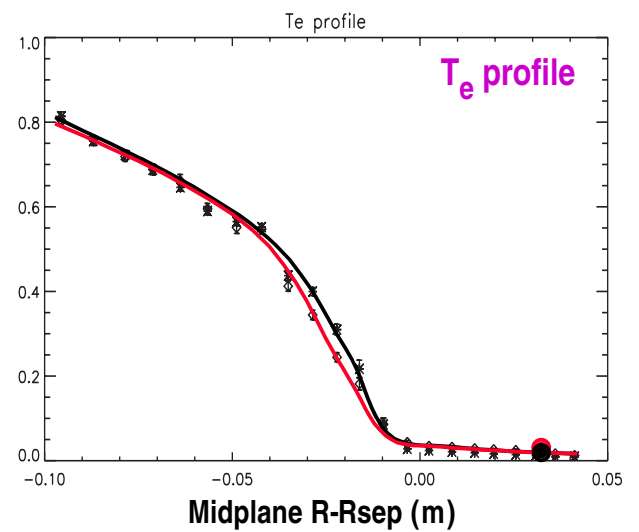
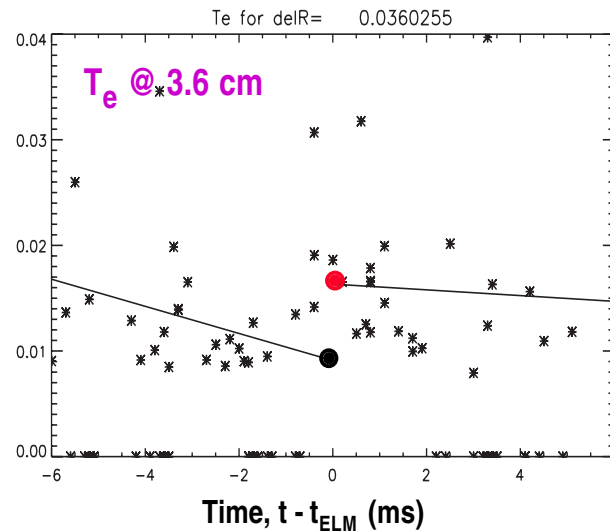
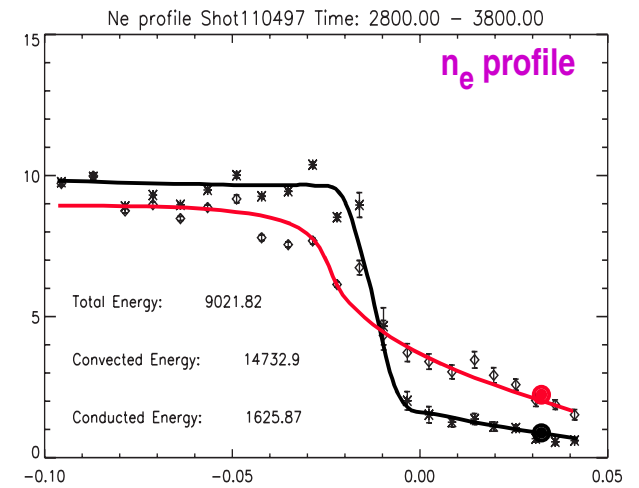
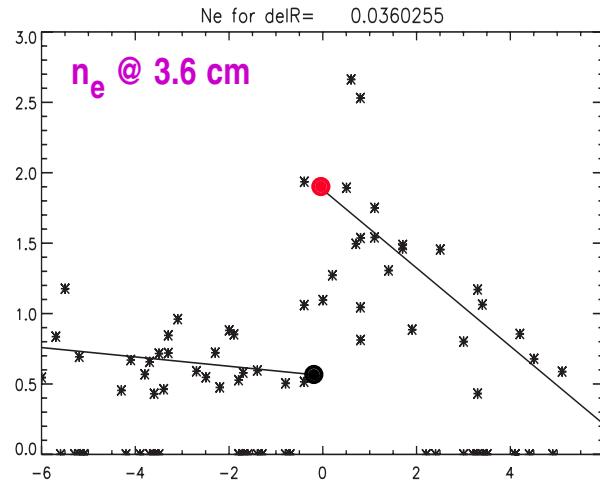
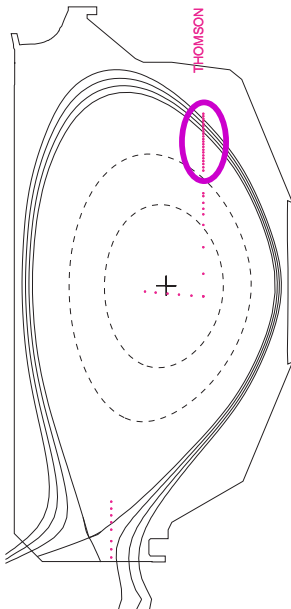
Midplane R-Rsep (m)



# High $n_e$ ELMs: Particles seen far out in SOL at midplane; pedestal $\Delta T_e$ very small

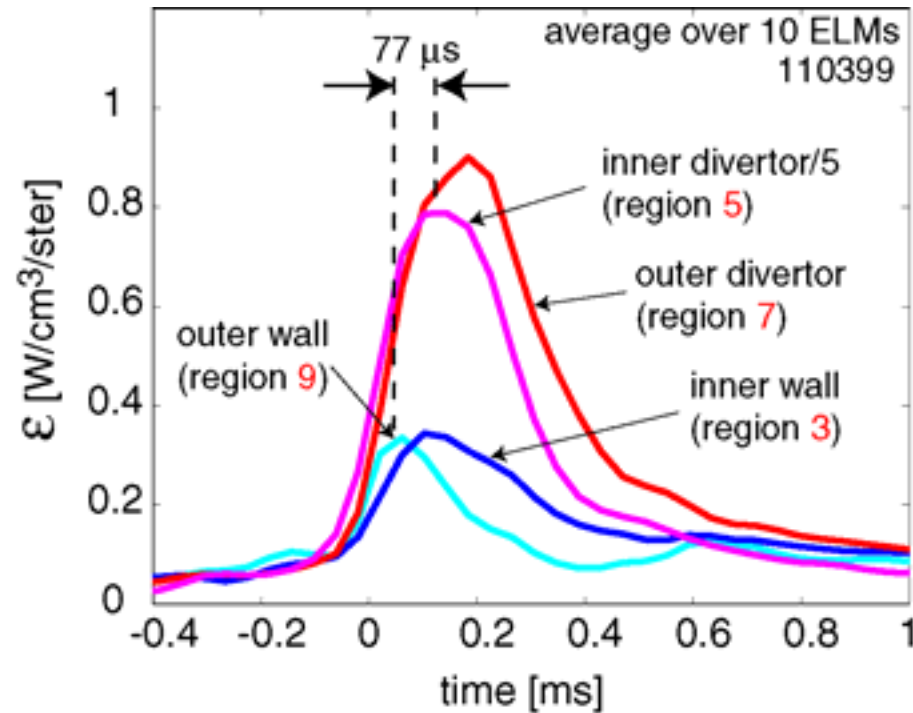
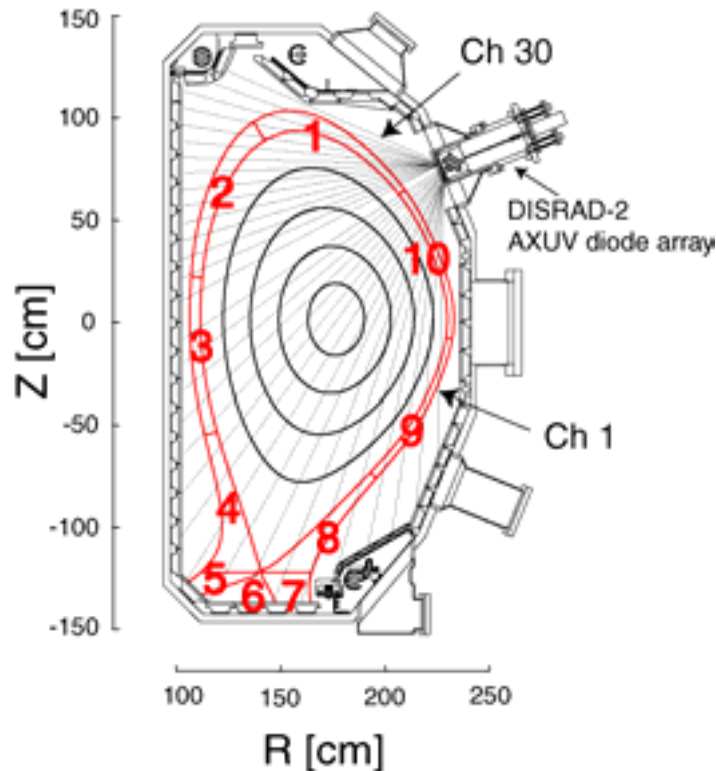
Leonard

- Greater scatter in high  $n_e$  data - higher SOL turbulence levels



# Low $n_e$ ELMs: fast bolometer chords show propagation of pulse around SOL to divertors

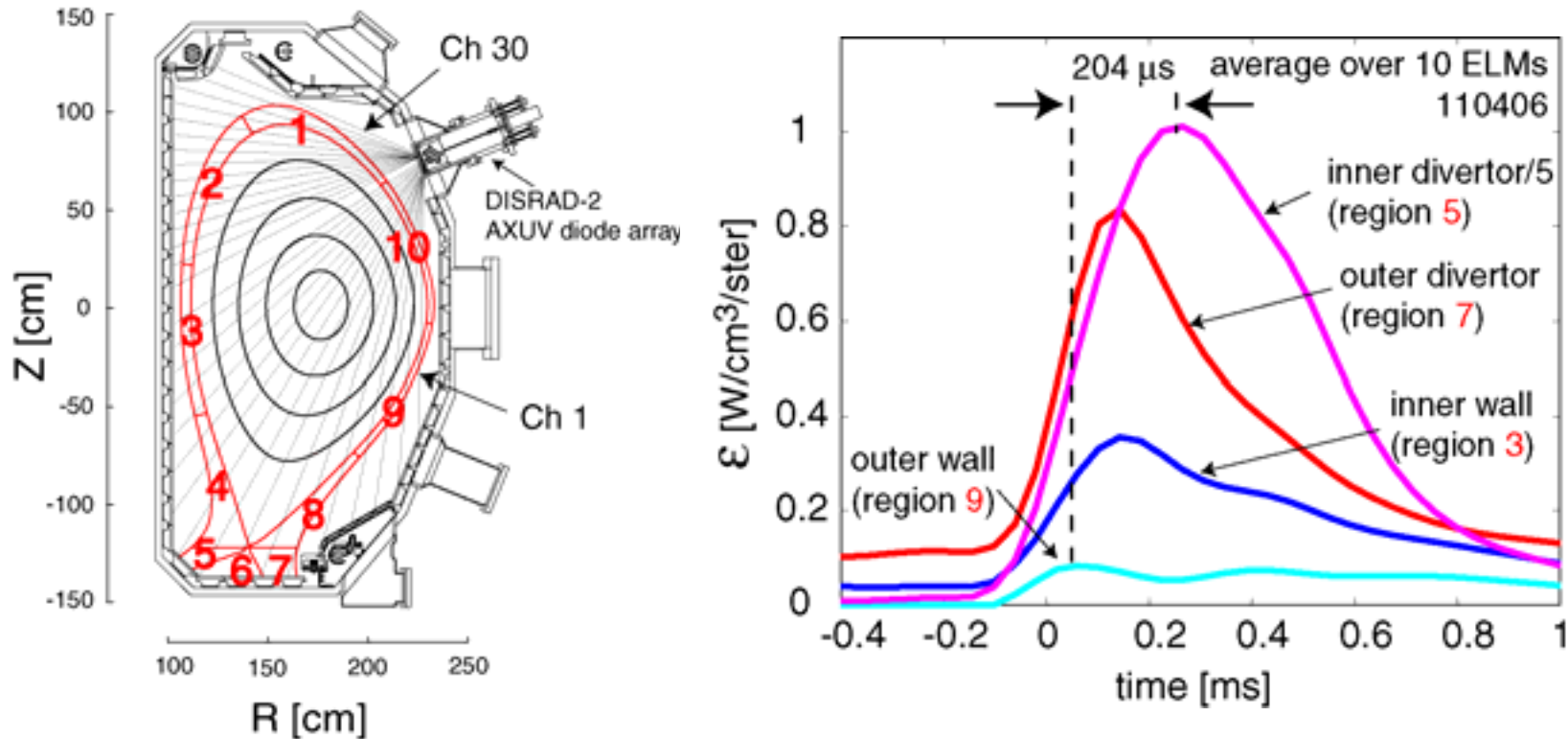
E. Hollmann



- Delays are consistent with ion transit time (outer midplane to inner strike point ~ 100  $\mu$ s) not electron conduction time.

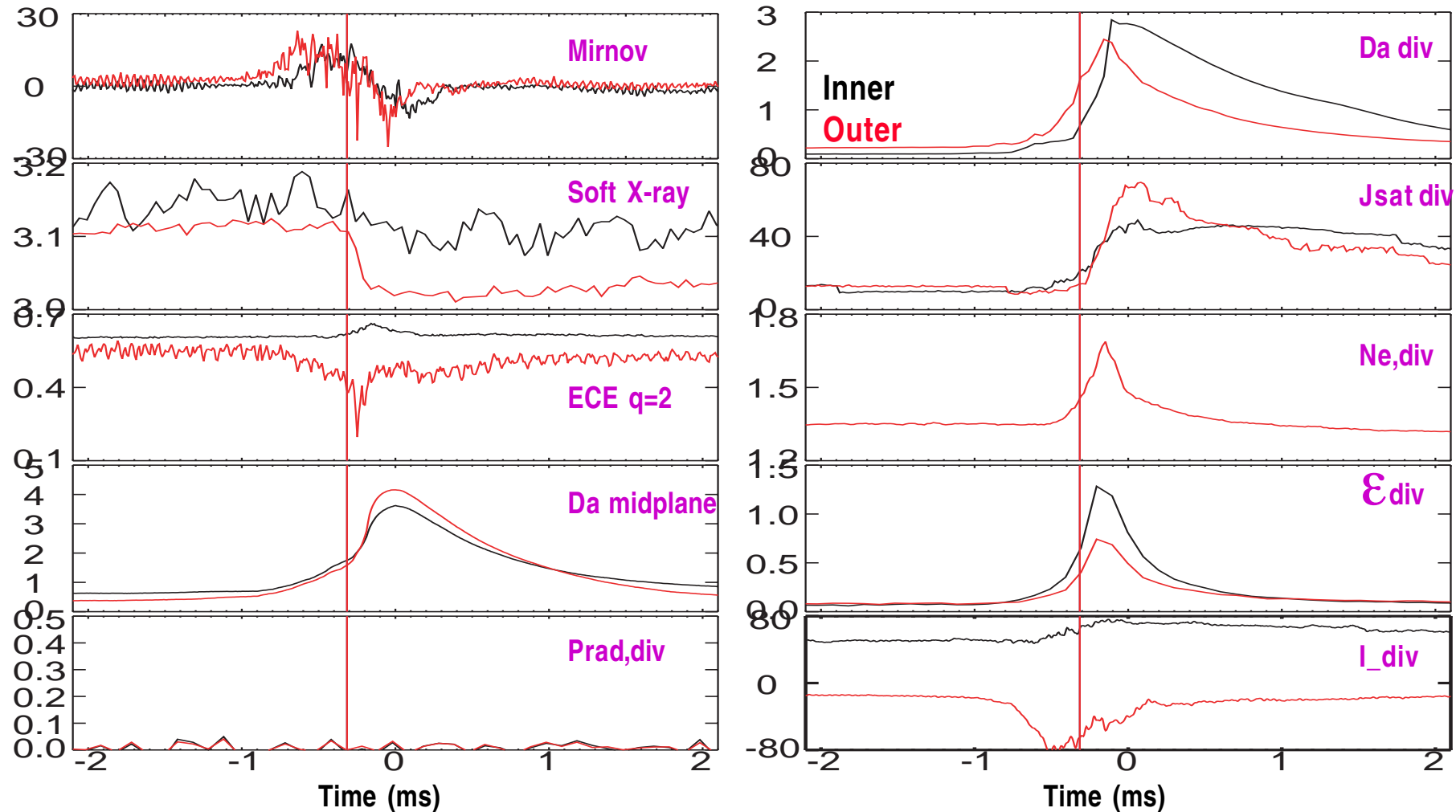
# High $n_e$ ELMs: Fast bolometer chords show propagation of pulse around SOL to divertors

E. Hollmann



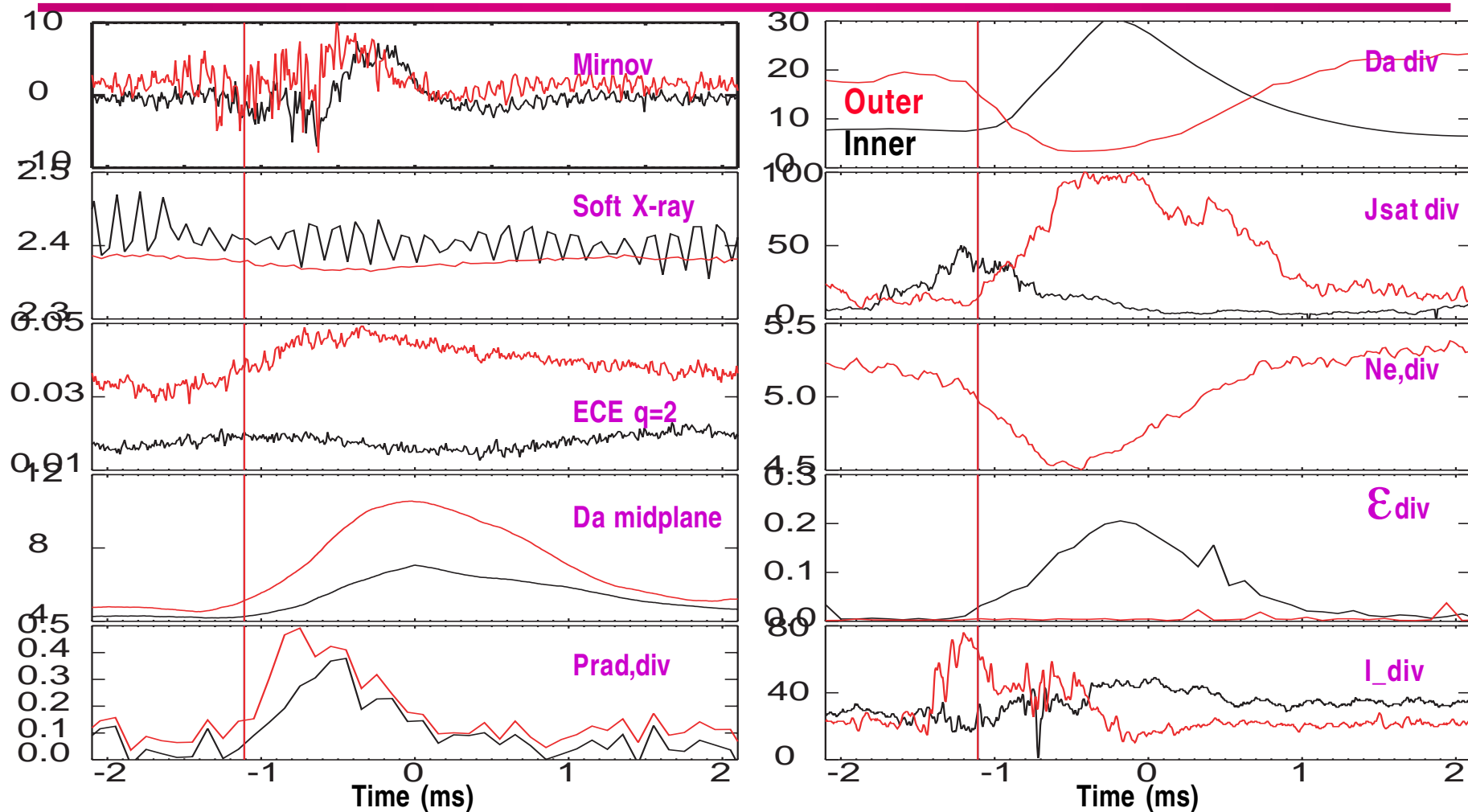
- Delays are consistent with ion transit time (outer midplane to inner strike point ~ 100  $\mu$ s) not electron conduction time.

# Low $n_e$ ELMs: Multi-diag. timing shows evidence of ELM particle transport from pedestal before thermal energy loss



- Two phases to ELM build-up: particle loss followed by rapid thermal energy loss

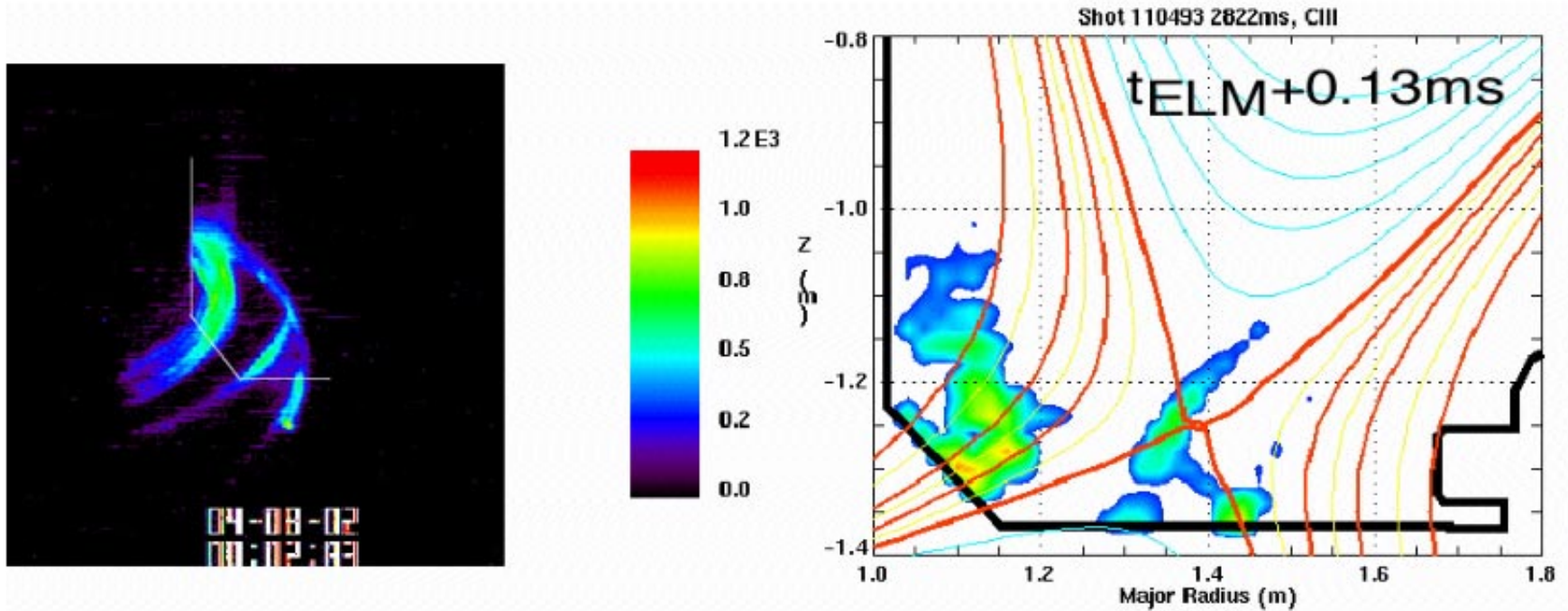
# High $n_e$ ELMs: Multi-diag. timing shows completely different behavior of outer divertor $n_e$ and heat flux vs. low $n_e$ ELMs



- No rapid rising phase observed consistent with lack of thermal energy loss
- No measurable outer target heat flux - still unexplained

# Low $n_e$ ELMs: Gated divertor TV shows burn-through of inner divertor leg: CIII moves from X-point to inner strikepoint

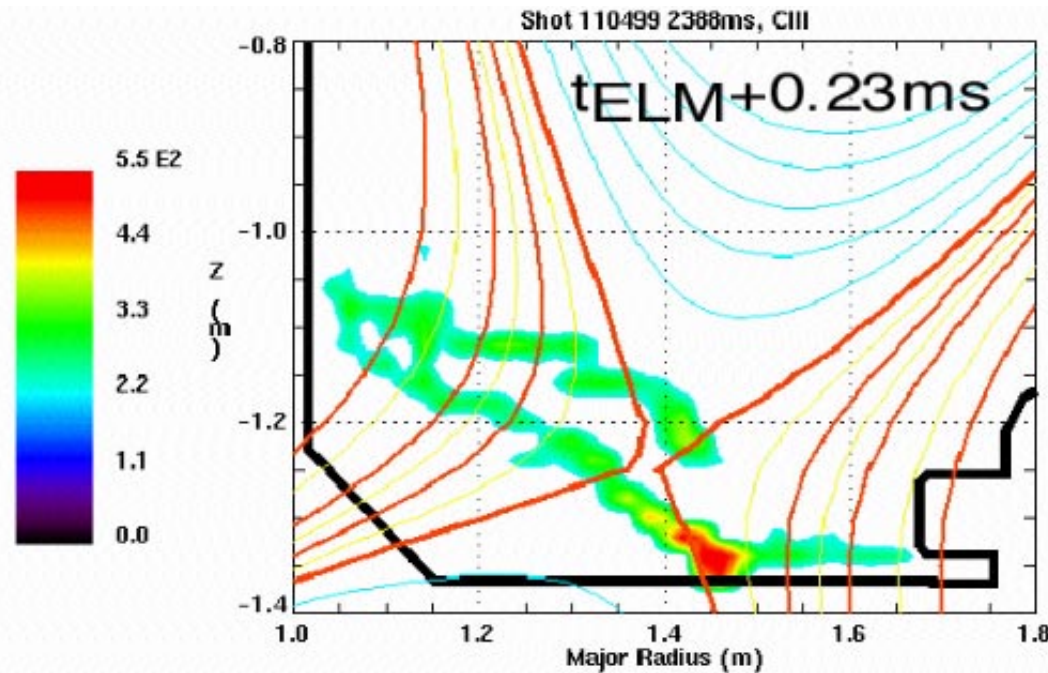
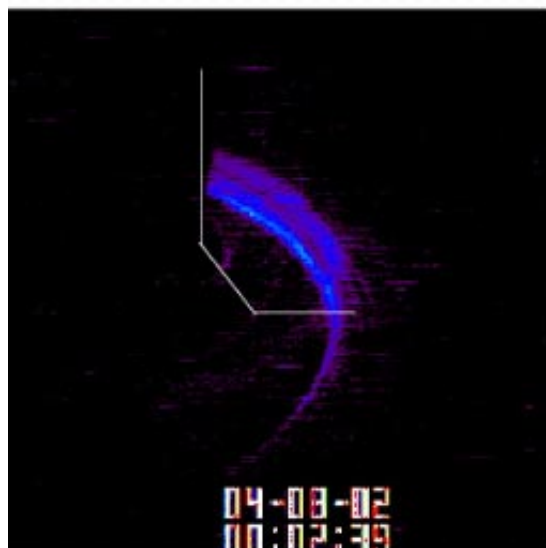
Groth PSI02



- Burn-through occurs between 10 and 130  $\mu$ sec after the ELM start

# High $n_e$ ELMs: Gated divertor TV shows burn-through of outer divertor leg: CIII moves from X-point to outer strikepoint

Groth PSI02



- Radiation increase near X-point occurs between 80 and 110  $\mu$ sec after ELM start
- Burn-through occurs between 110 and 230  $\mu$ sec after the ELM start

# Comparison of $\nabla B \downarrow$ vs. $\nabla B \uparrow$ Effects

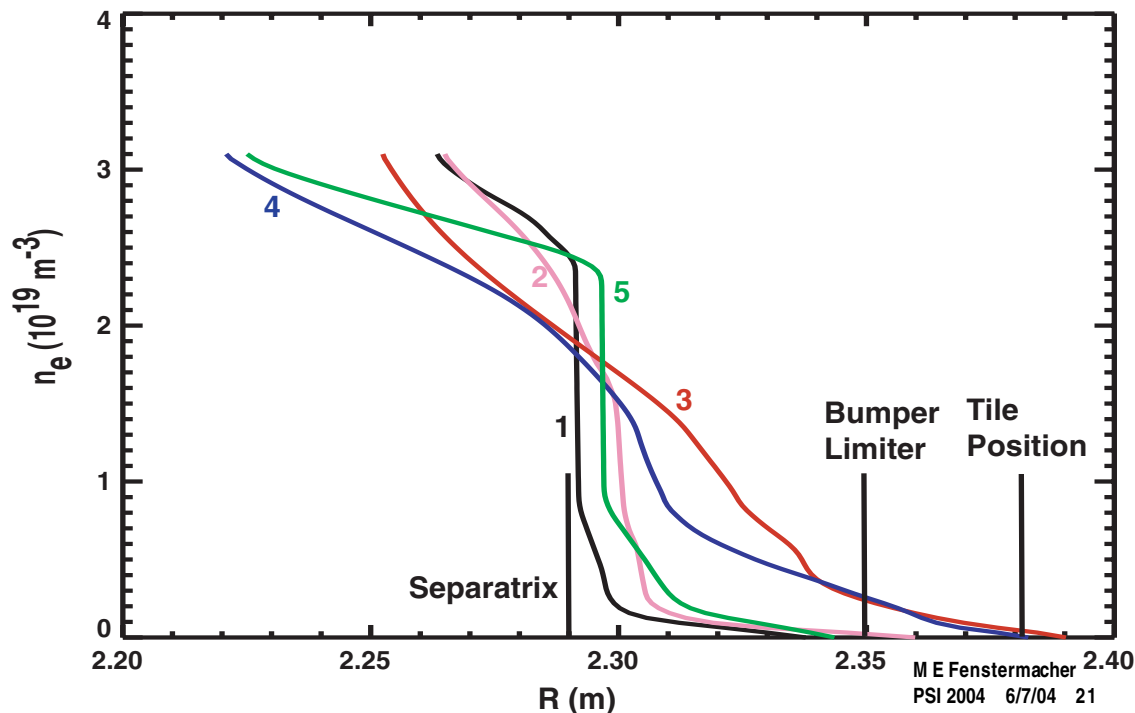
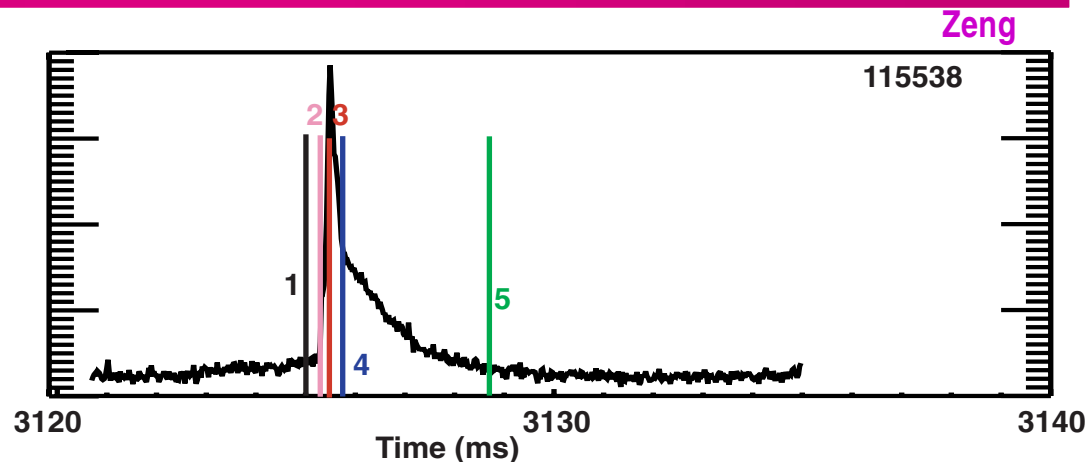
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# LSN $\nabla B \downarrow$ Low $\delta$ ELMs: Particle perturbation seen much farther out in midplane SOL than $\Delta T_e$ , especially at low $n_e$

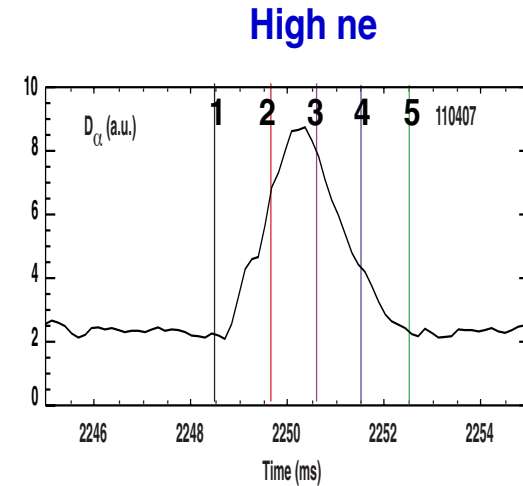
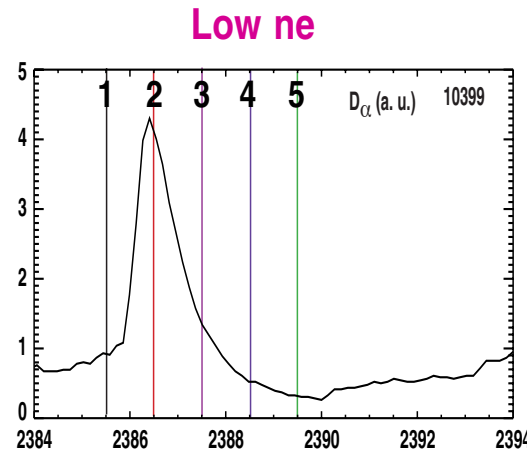
- Reflectometer shows reduction of pedestal density [curves 1,2,3,4]
- Density lost from pedestal appears in SOL at limiters [curves 3,4]
- Recovery of pre-ELM density profile takes  $\sim 3$  ms [curves 4, 5]
- During ELM  $n_e \sim 10^{19} \text{m}^{-3}$  at 3.5 cm ( $3 \lambda_{ne}^{\text{pre-ELM}}$ ) from pre-ELM separatrix [curve 3]



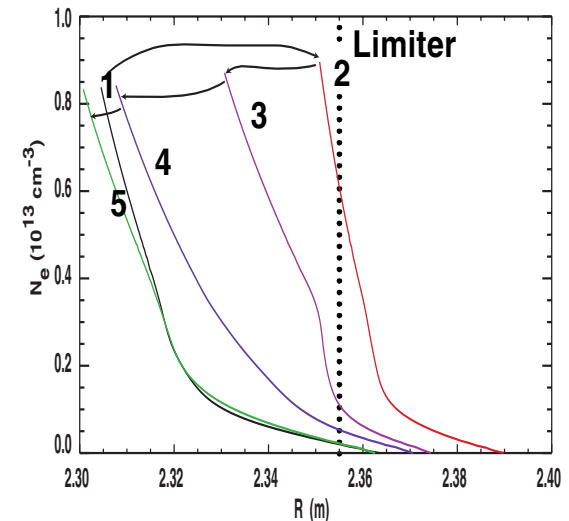
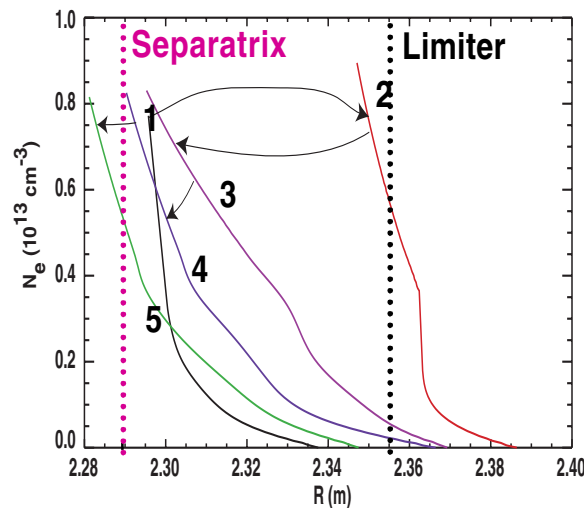
# LSN $\nabla B \downarrow$ High $\delta$ ELMs: Particle perturbation seen much farther out in midplane SOL than $\Delta T_e$ for both low and High $n_e$

Zeng, Leonard

- Reflectometer shows motion of density out to the limiter region (5cm from separatrix) in  $\sim 500 \mu\text{sec}$  [curves 1  $\rightarrow$  2]

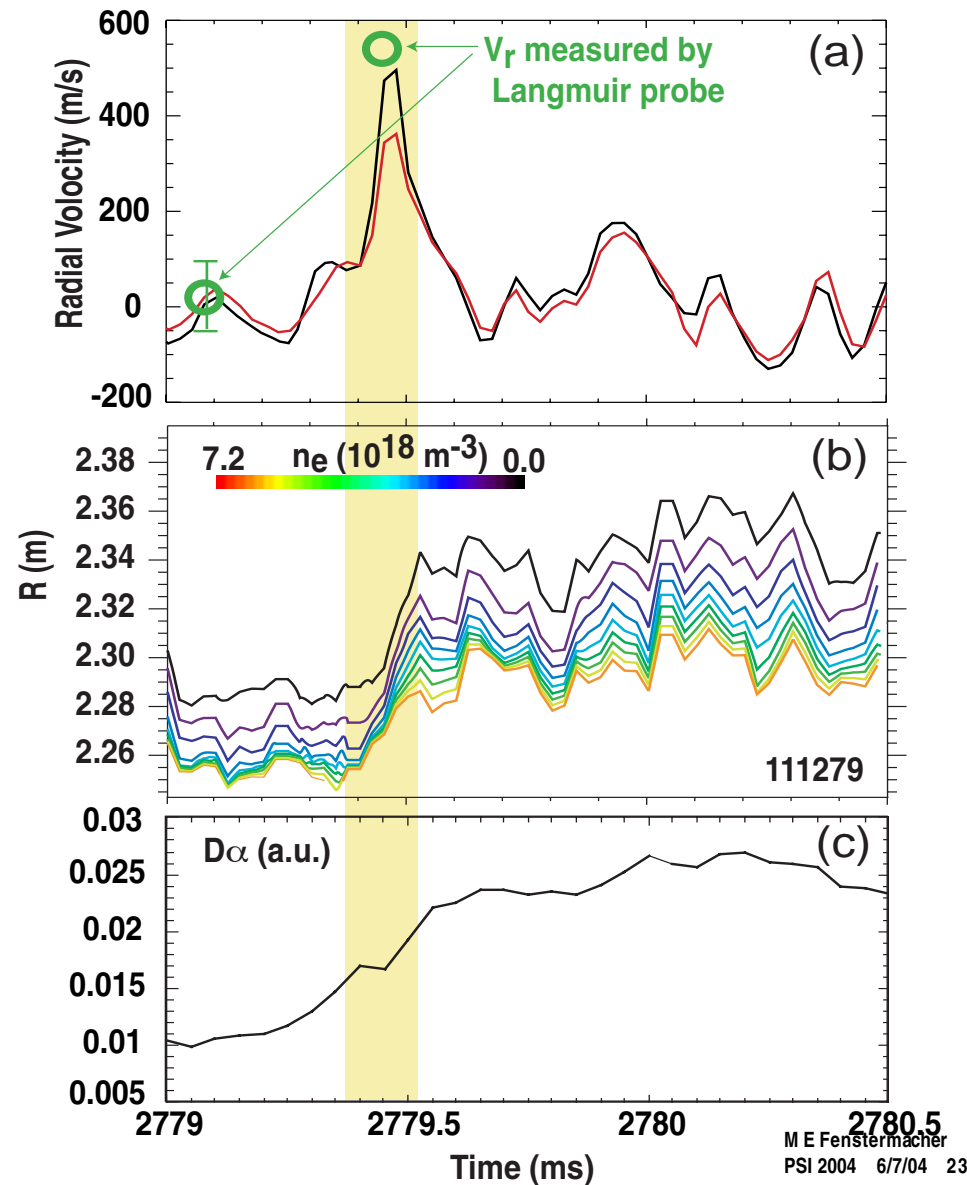


- Recovery of pre-ELM density profile takes  $\sim 1.5 \text{ ms}$  [curves 2  $\rightarrow$  5]



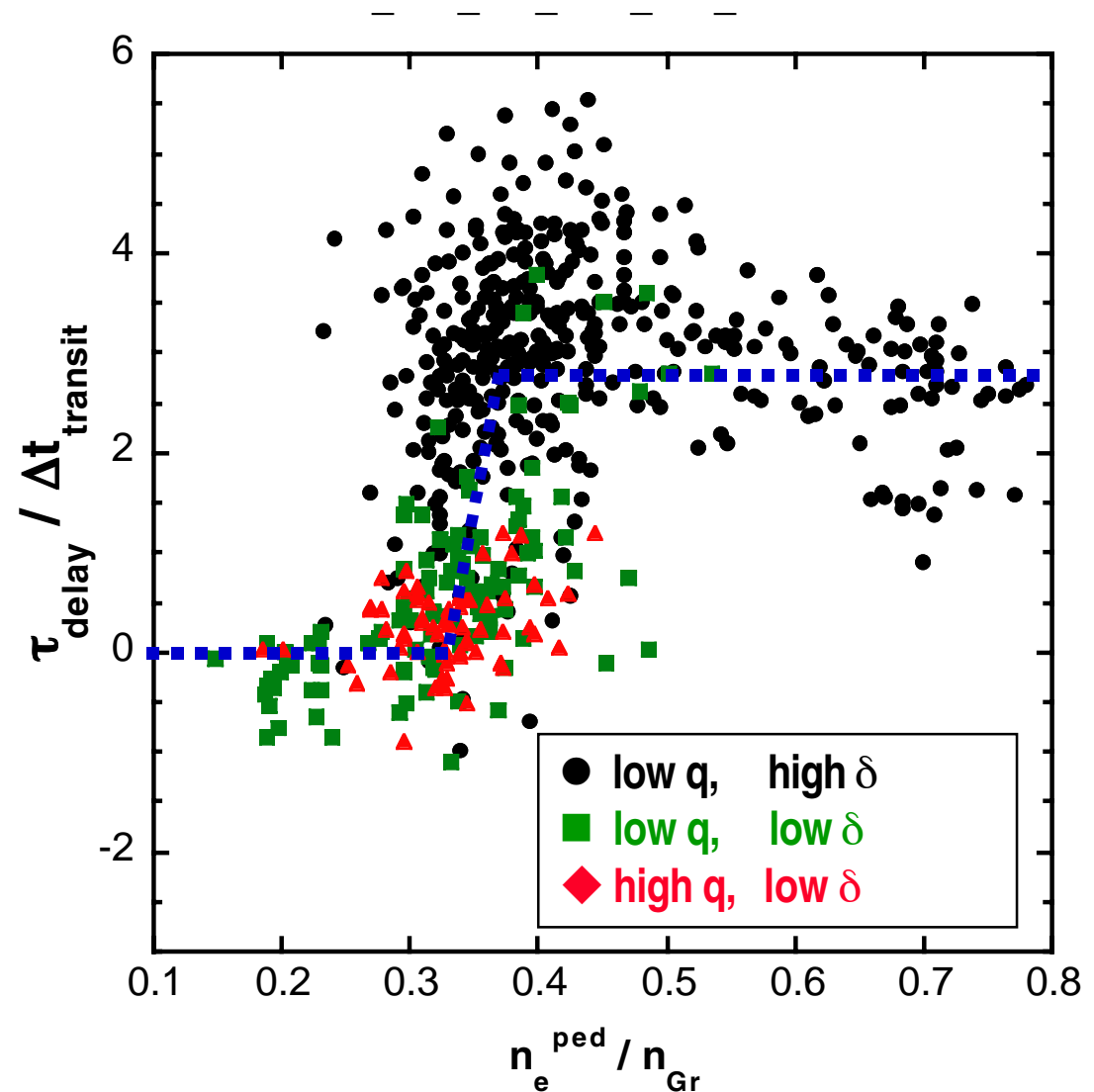
# Low $n_e$ High $\delta$ ELMs: Radial velocity $\sim 600$ m/s from reflectometer agrees with ExB velocity from probes Zeng, Boedo

- Reflectometer data to 40 kHz shows radial velocity of 500 m/s for  $5 \times 10^{19} \text{ m}^{-3}$  surface at midplane
- Probe measurement of “density blobs” shows ExB radial velocity of 550 m/s



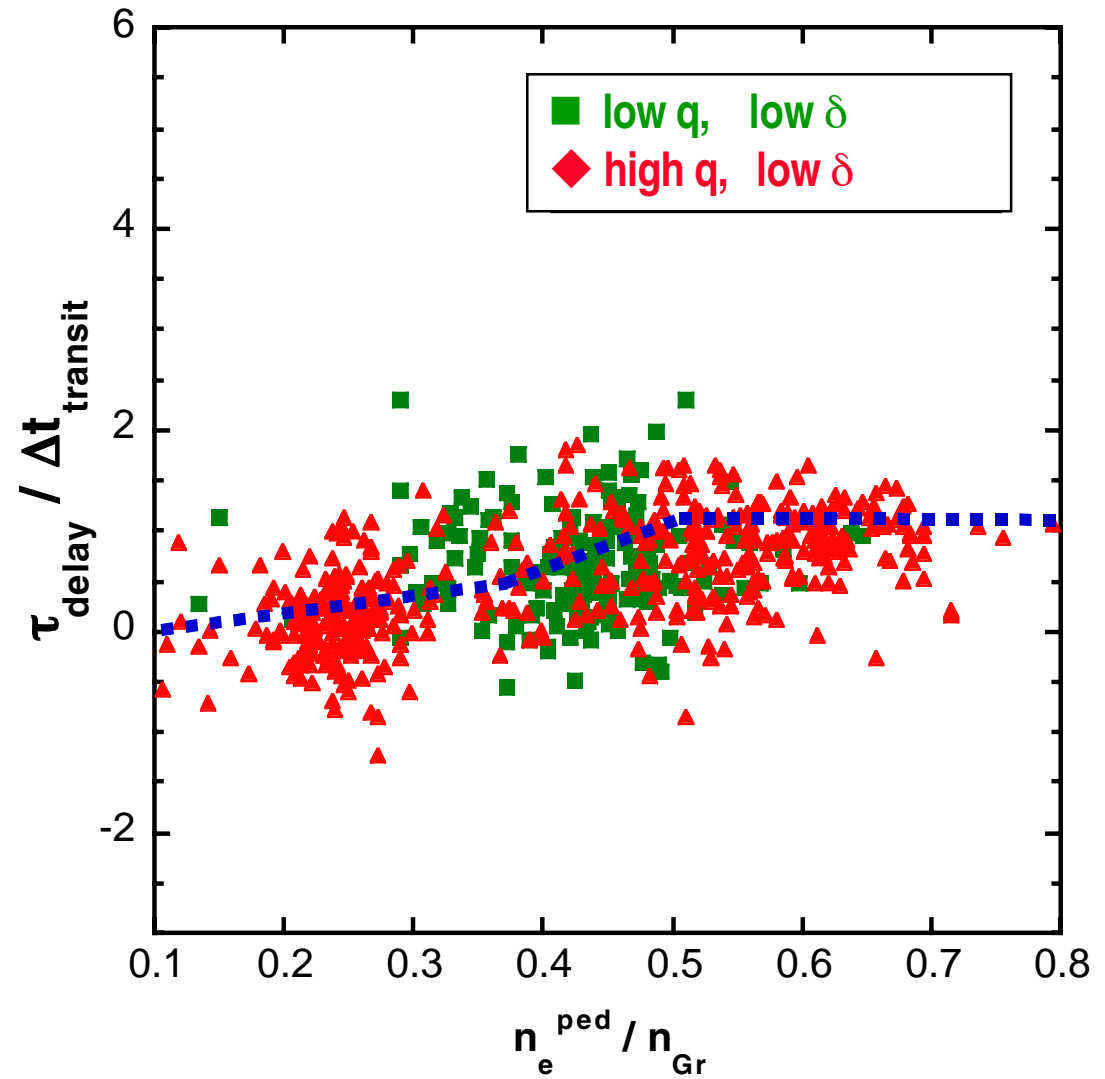
# LSN $\nabla B \downarrow$ ELMs: Delay of inner vs outer $D_\alpha$ about 3x the difference in ion transit times from midplane to targets.

- Ion transport assumed at sound speed evaluated at pedestal  $T_e$
- Scatter increases and delay time drops to small value at very low density
  - Evidence of fast electron effects ?
  - Evidence of change in character of ELM from ballooning to peeling dominated ?



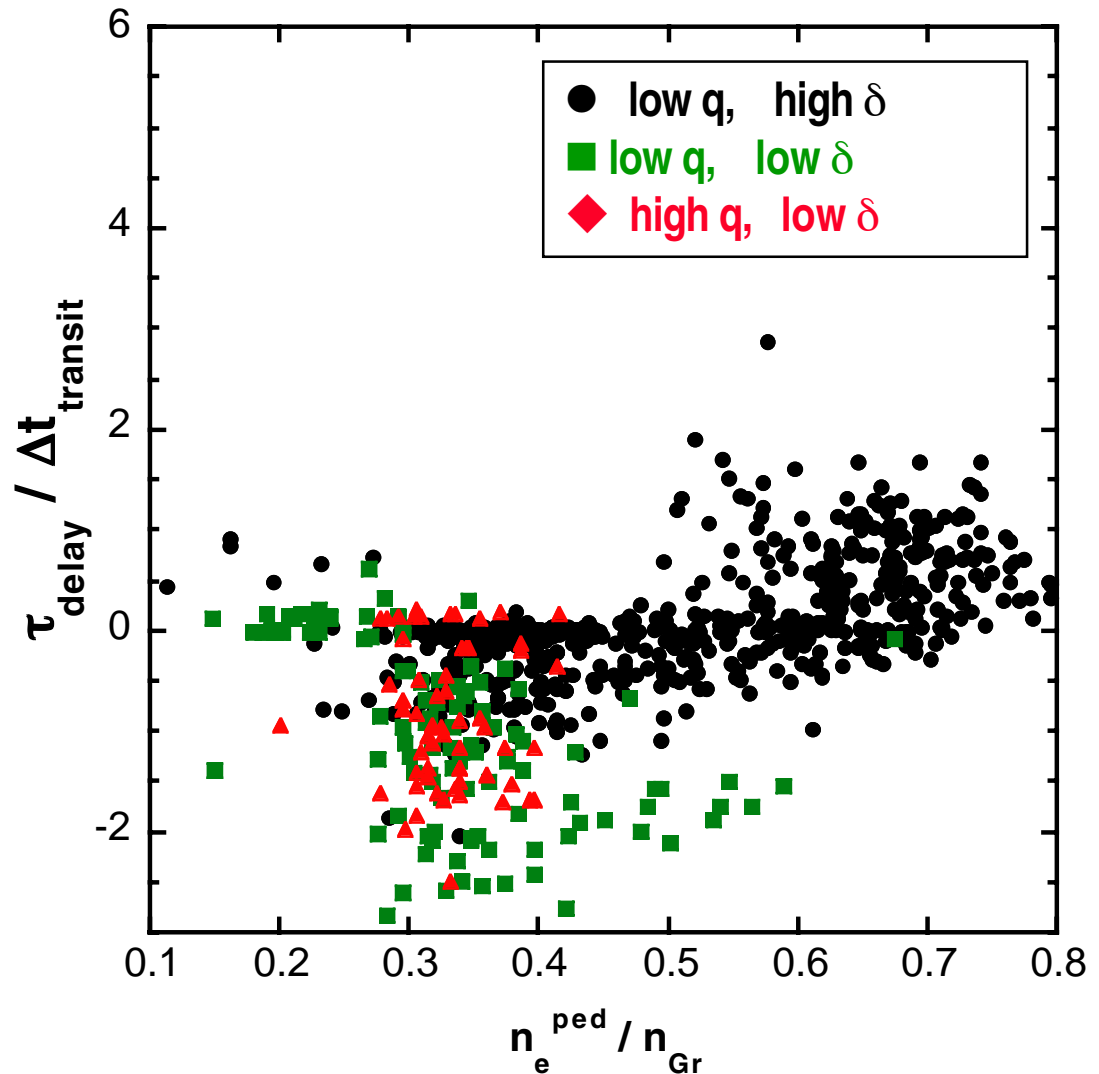
# LSN $\nabla B \uparrow$ ELMs: Delay of inner vs outer $D_\alpha$ with $\nabla B \uparrow$ much smaller than in $\nabla B \downarrow$ case.

- With  $\nabla B$  out of divertor inner leg plasma conditions similar to outer leg
- Dependence of delays on  $\nabla B$  direction may be due to:
  - Difference in pre-ELM divertor conditions ?
  - Role of ExB drifts during ELM event ?



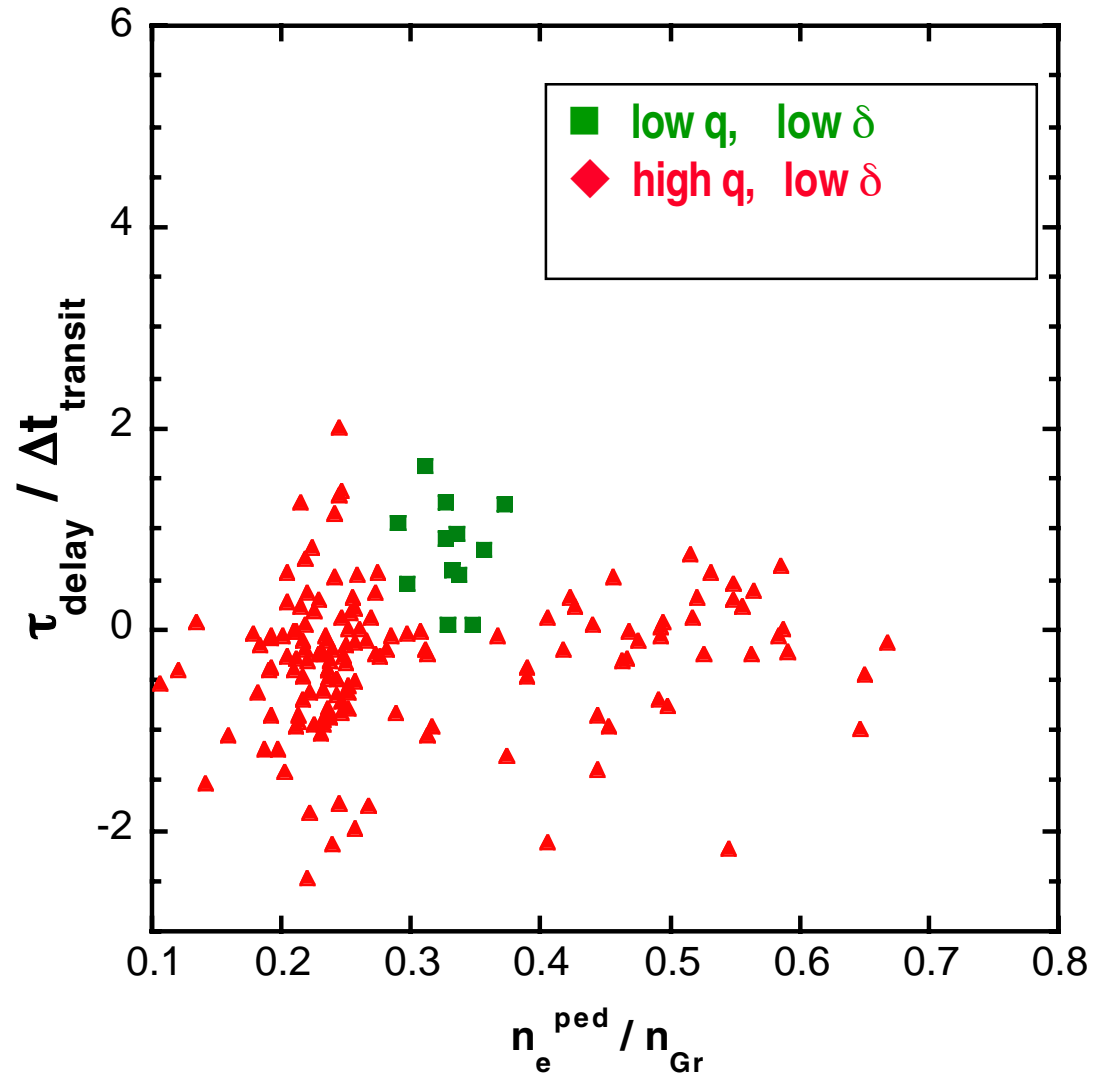
# LSN $\nabla B \downarrow$ ELMs: Delay of inner vs outer $P_{\text{rad}}$ less than $D_{\alpha}$ delay

- At high  $n_e$  the delay is 2x smaller than in  $D_{\alpha}$
- At low  $n_e$ , ELM  $P_{\text{rad}}$  (inner) before ELM  $P_{\text{rad}}$  (outer)
  - This also seen in analysis of spatial zones by Hollman (2002)
  - Fast electron pulse burns through detached inner divertor ?



# LSN $\nabla B \uparrow$ ELMs: Delay of inner vs outer $P_{\text{rad}}$ similar to delays of $D_{\alpha}$

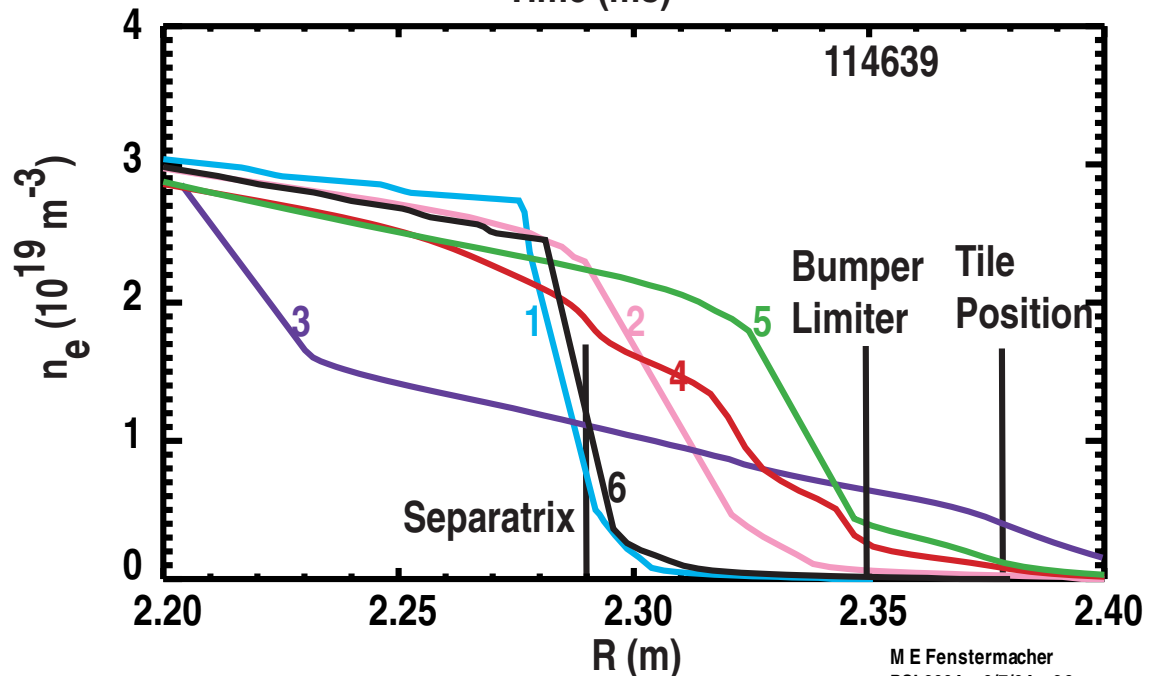
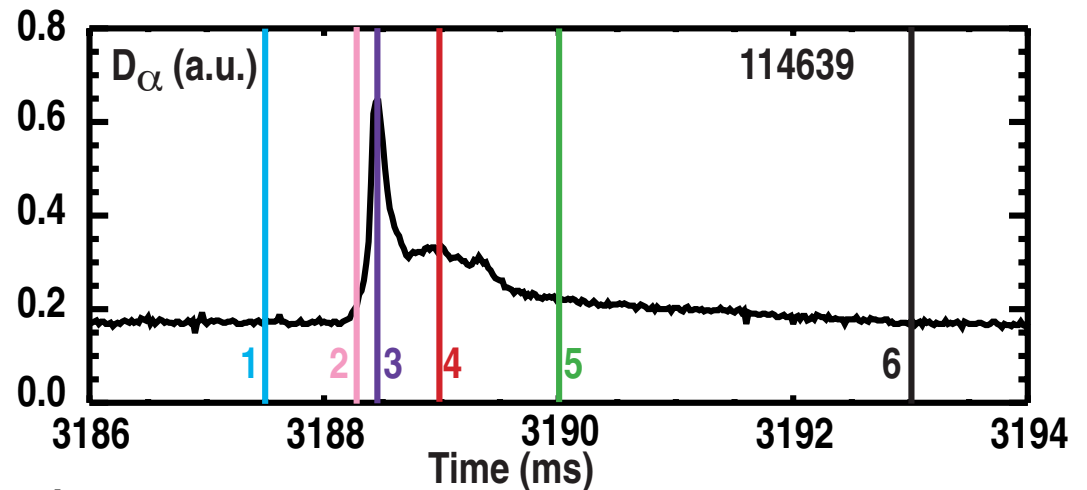
- Data set limited to high  $q$  shots because outer  $P_{\text{rad}}$  saturated at low  $q$
- Small delay (albeit with large scatter)
  - No clear variation with density



# LSN $\nabla B \uparrow$ Low $\delta$ ELMs: $\Delta n_e$ seen much farther out in midplane SOL than $\Delta T_e$ , especially at low $n_e$

Zeng

- Reflectometer shows reduction of pedestal density [curves 2 -->3]
- Density lost from pedestal appears far out in SOL;  $n_e \sim 10^{19} \text{m}^{-3}$  at 4.5 cm ( $4 \lambda_{ne}^{\text{pre-ELM}}$ ) from pre-ELM separatrix [curve 3]
- Recovery of pre-ELM density profile takes  $> 1.5$  ms [curves 4, 5, 6]
  - Intermediate recovery stage with “pedestal” in the SOL (curve 5)
  - Full recovery after  $\sim 5$  ms (curve 6)

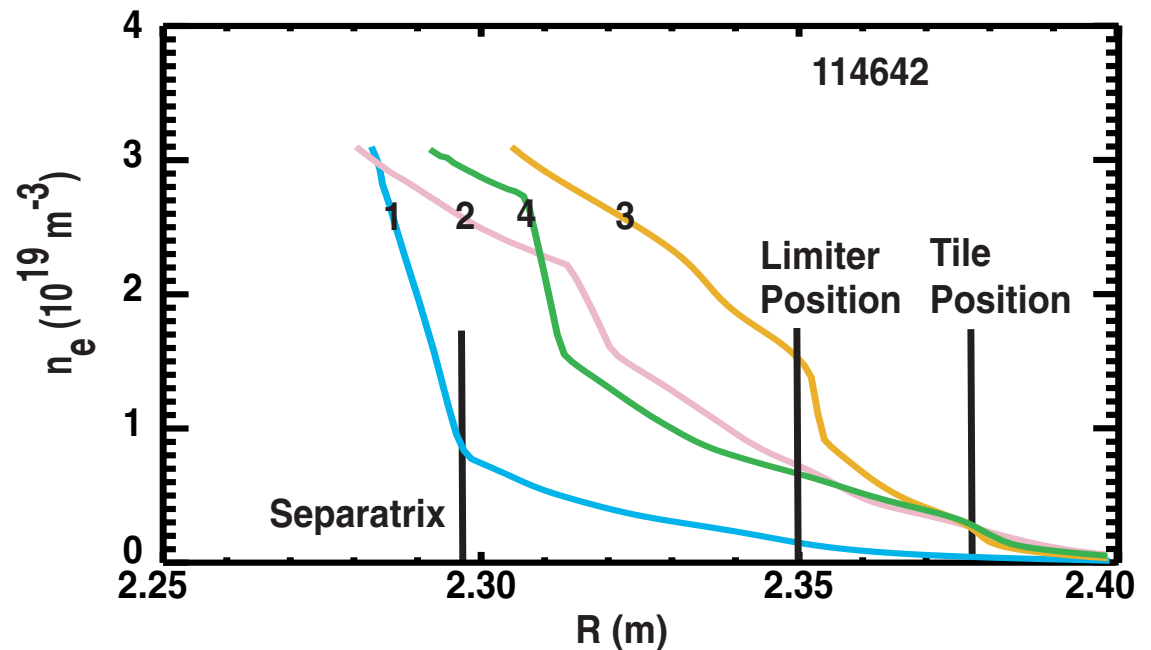
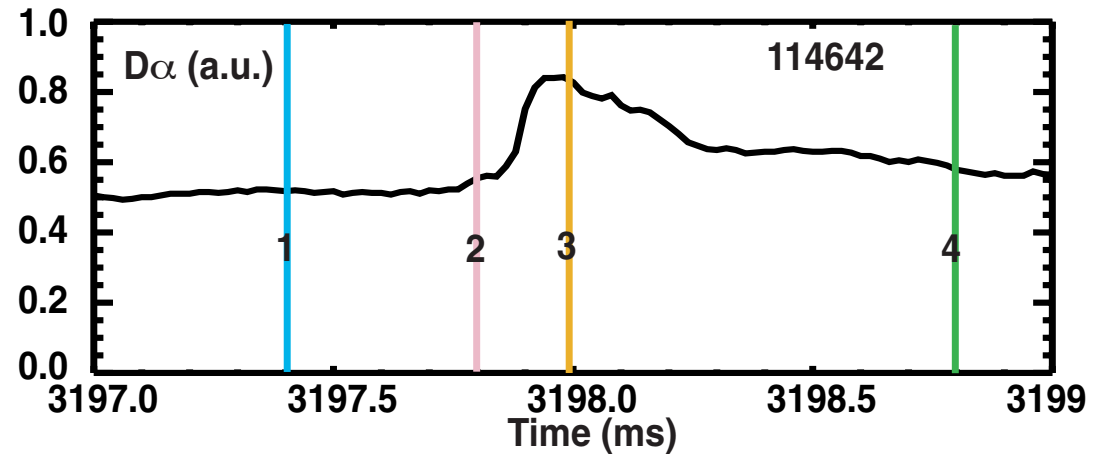




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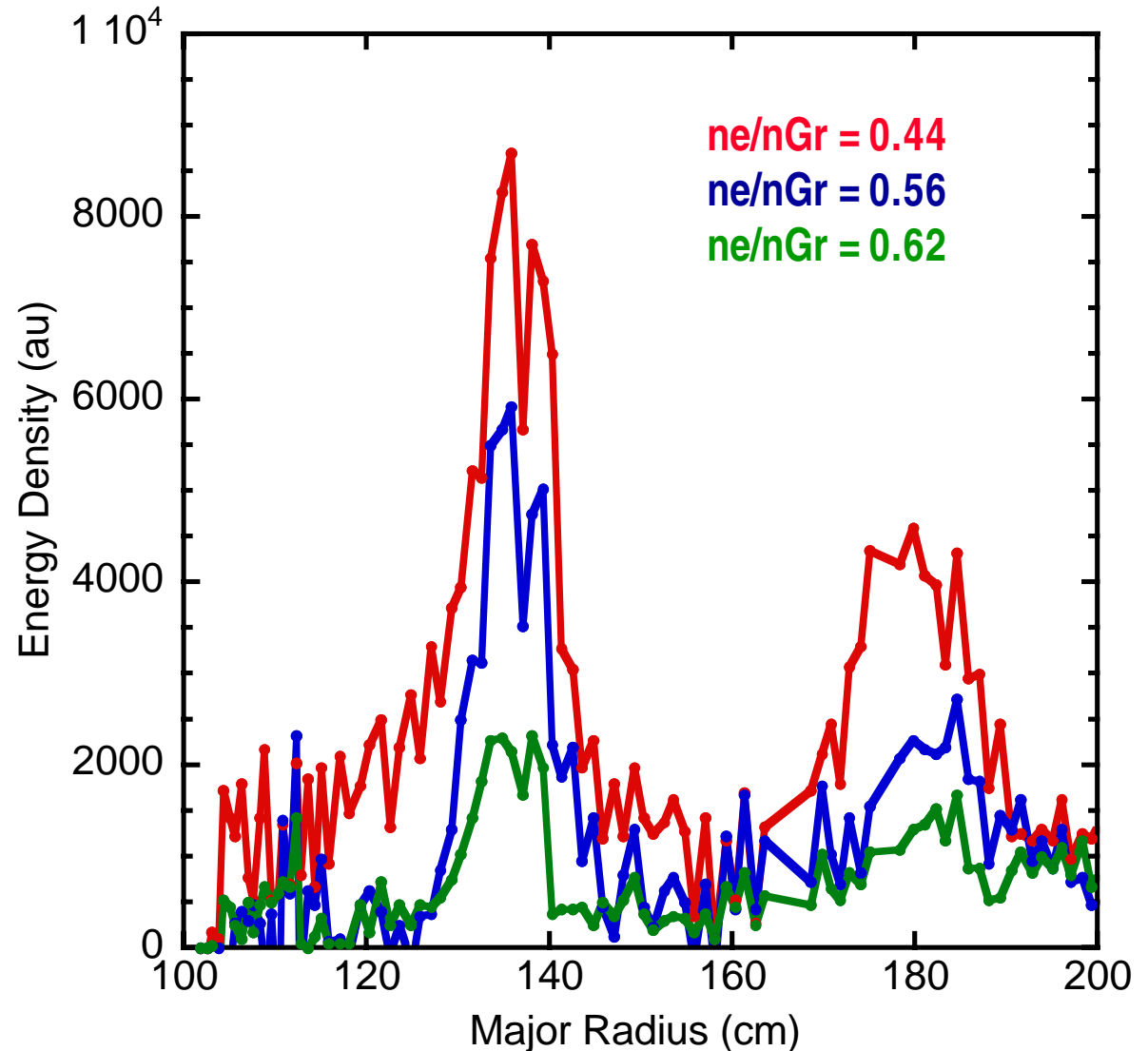
Zeng,

- Reflectometer shows particles ejected into SOL [curves 1, 2, 3]
- Density profile modified before large  $D\alpha$  rise [curve 2]
- Far SOL density rise to  $n_e \sim 10^{19} \text{m}^{-3}$  at 6 cm ( $5 \lambda_{ne}^{\text{pre-ELM}}$ ) from pre-ELM separatrix [curve 3]
- Recovery of pre-ELM density profile takes  $\gg 1$  ms [curves 3, 4]



# LSN $\nabla B \downarrow$ ELMs: Inner / outer target energy density asymmetry during ELMs decreases slightly with $n_e$

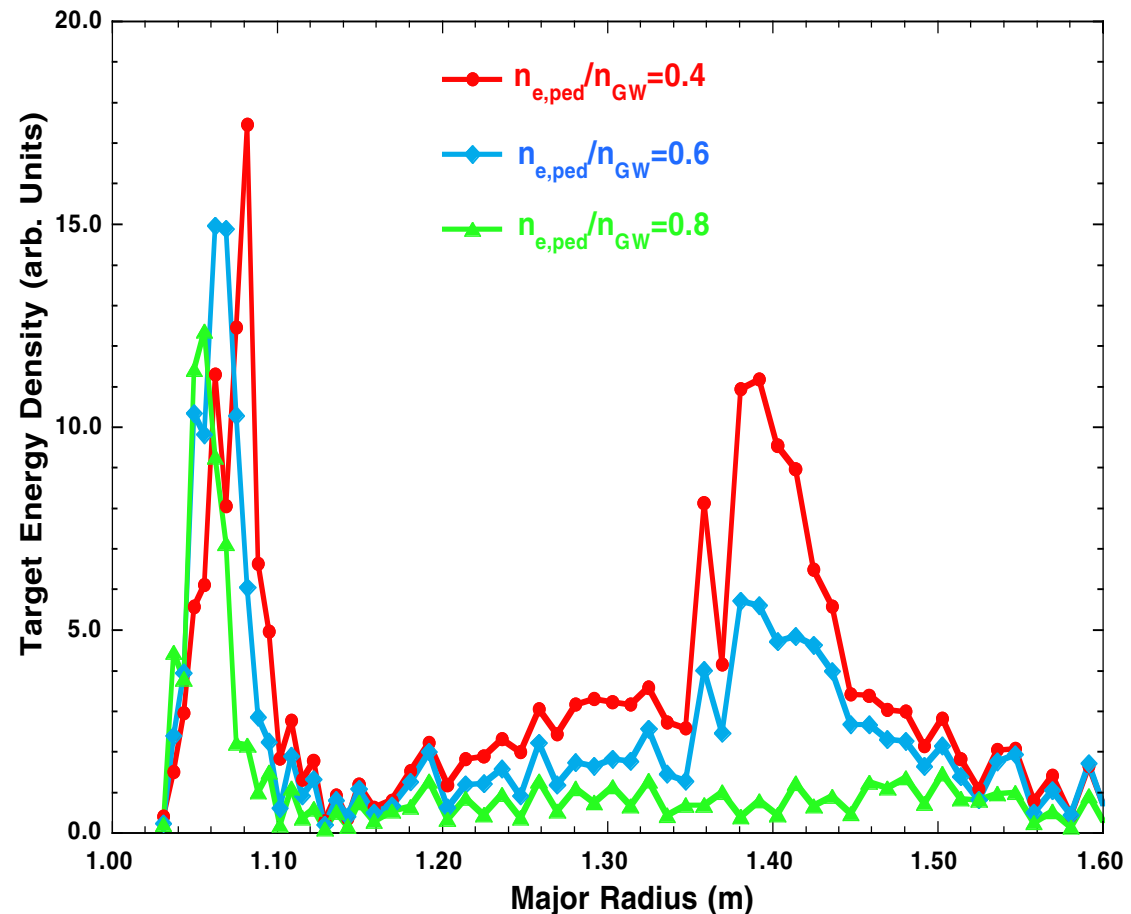
- Inner / outer peak energy density ratio  $\sim 2$  at low  $n_e/n_G \sim 0.4$ , ratio decreases to 1.5 at higher density,  $n_e/n_G > 0.6$
- Profiles averaged over 10 - 20 ELMs.
- Surface layer effects may be playing an important role in these results.



# LSN $\nabla B \downarrow$ ELMs: Inner / outer target heat flux asymmetry during ELMs increases with $n_e$ ; profile broadens $\leq$ factor of 2

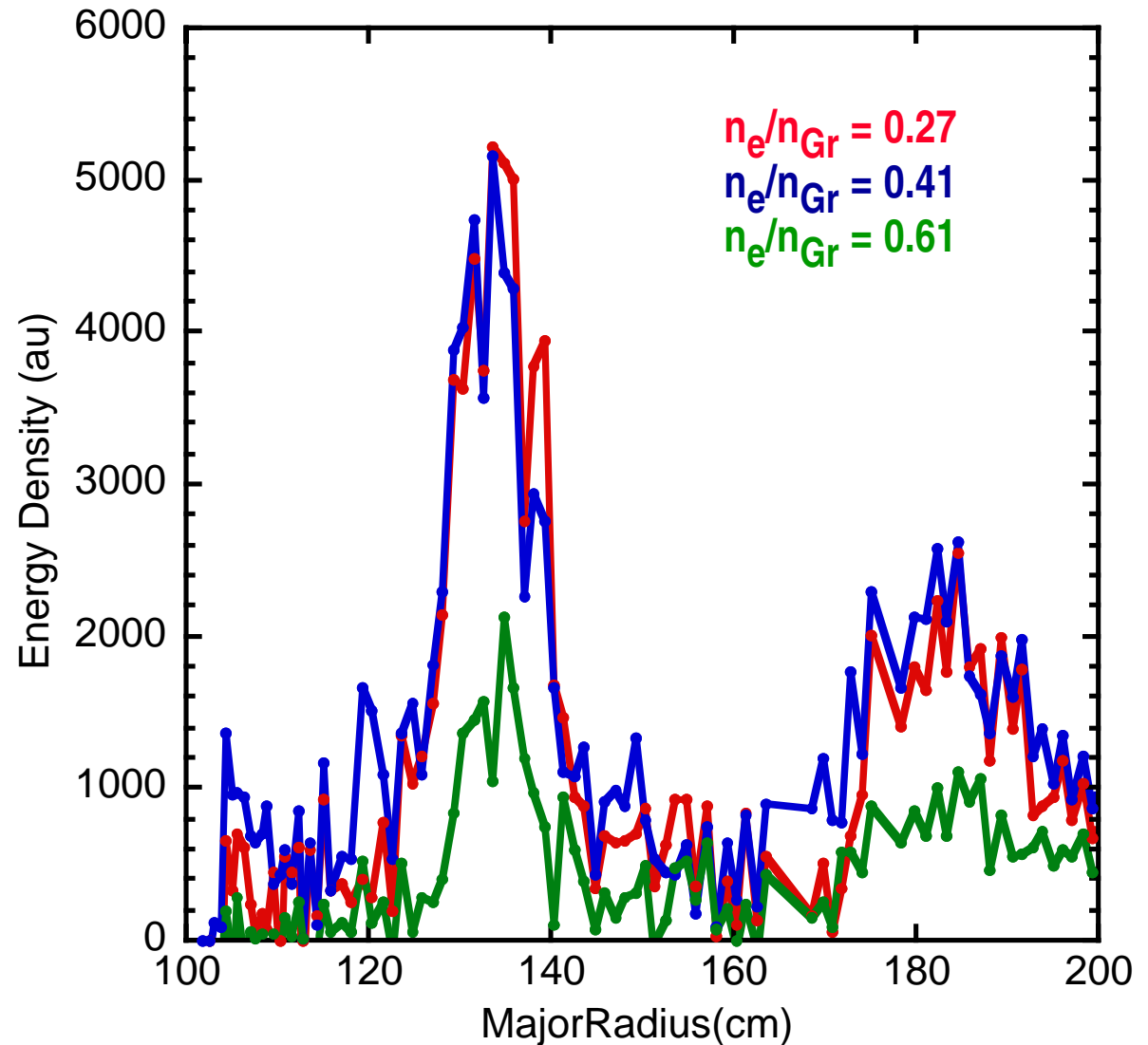
Lasnier, Leonard PSI02 Invited

- Outer target heat flux drops to near zero at high density
- Peak of inner heat flux profile moves away from SP
- Inner / outer energy ratio  $\sim 2$  from previous experiments - still working on present calibrations
- ELMs broader than time averaged by 2x on outer leg but narrower by up to 1.5x on inner leg.



# LSN $\nabla B$ $\uparrow$ ELMs: Inner / outer target energy density asymmetry during ELMs nearly constant with $n_e$

- Little change in profiles from low to moderate density,  $0.27 < n_e/n_{Gr} < 0.4$
- In / out asymmetry  $\sim 2.0$  independent of density
- Profiles averaged over 10 - 20 ELMs.
- Surface layer effects may be playing an important role in these results.



# Comments: SOL/divertor ELM behavior depends on both $n_e$ and $B_T$ -dependent particle drifts

- Delays of inner vs outer  $D_\alpha$  and  $P_{\text{rad}} = f(n_e, B_T)$ 
  - ELM poloidal character may change with  $n_e$
  - Fast electron effects may dominate at low  $n_e$ ; ion convection at high  $n_e$
  - Difference in pre-ELM divertor conditions with  $B_T$  may play a role
- Pedestal particles ejected far into midplane SOL,  $3 - 5 \lambda_{ne}^{\text{pre-ELM}}$ , independent of  $n_e, B_T$ 
  - Ejected  $T_i$  (and heat flux) at main chamber wall not known
  - Ejected  $T_e$  falls rapidly with radius in SOL

(see also Zeng O-29 Rudakov O-24 , Boedo P2-5)
- Asymmetry of peak energy density weak  $f(n_e, B_T)$ 
  - Asymmetry decreases slightly with  $n_e$  for  $\nabla B \downarrow$ ; nearly constant for  $\nabla B \uparrow$
  - May be contaminated by surface layer effects

# UEDGE SS and ELM Modeling - $\nabla B$ ↓

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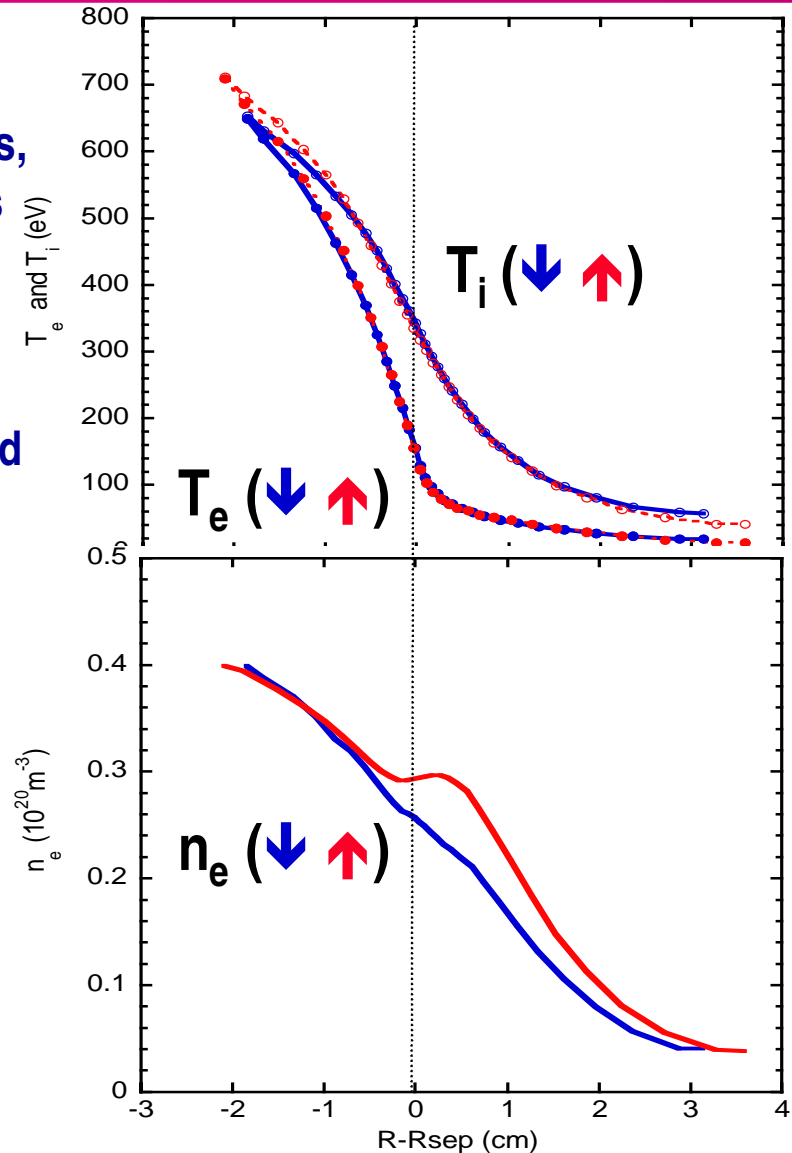
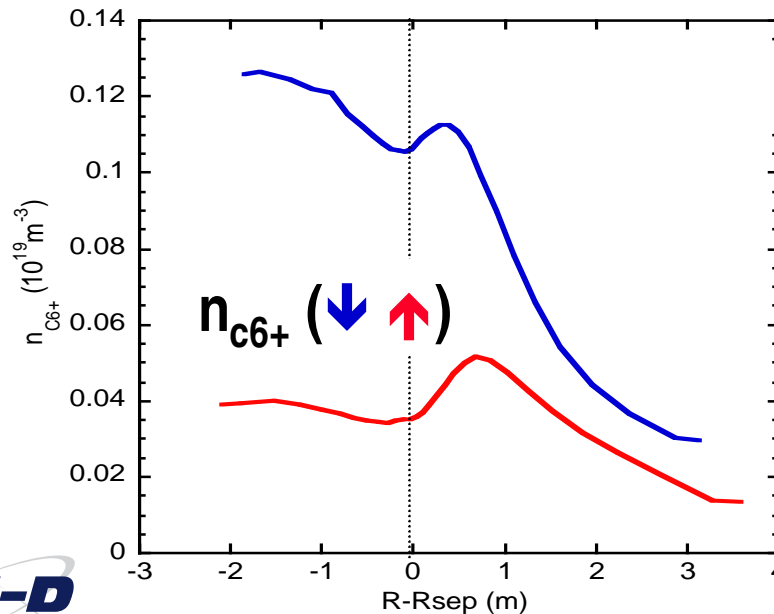
# UEDGE SS and ELM Modeling - $\nabla B$

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# UEDGE simulations of pre-ELM $\nabla B \downarrow$ vs $\nabla B \uparrow$ cases show similar midplane $T_e, T_i$ but changes in density profiles

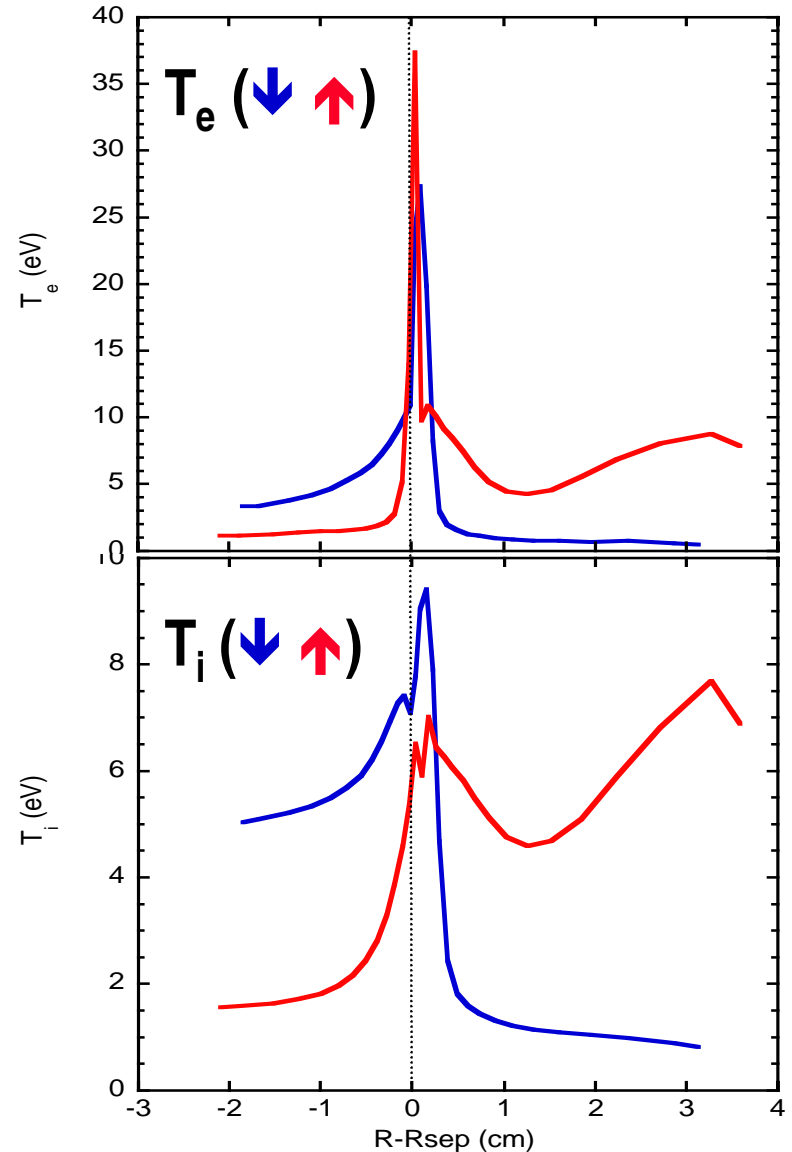
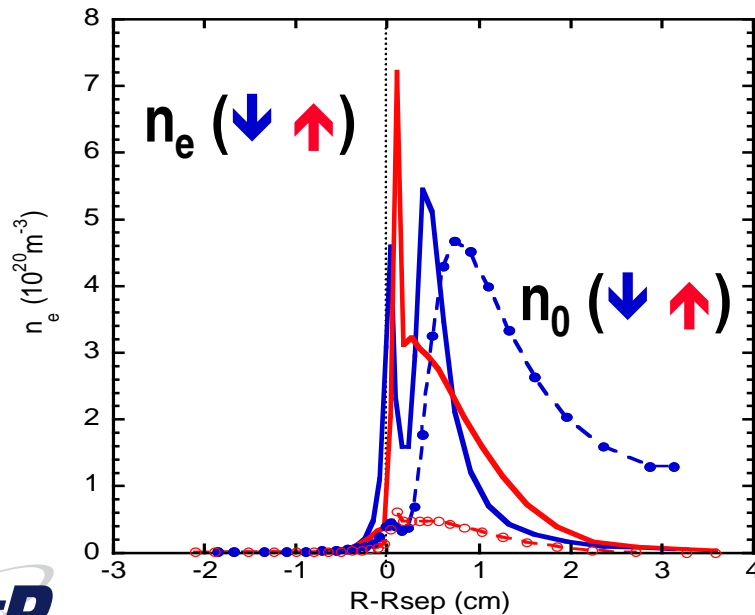
- $D = 0.09 \text{ m}^2/\text{s}$ ,  $\chi_e = \chi_i = 0.35 \text{ m}^2/\text{s}$ , 5% pumping at walls, unity recycling at targets, 6 carbon species, 40% of theoretical drifts
- Midplane  $T_e$  and  $T_i$  profiles independent of  $\nabla B$  direction
- Elevated density in midplane SOL, reduced C6+ density in midplane pedestal





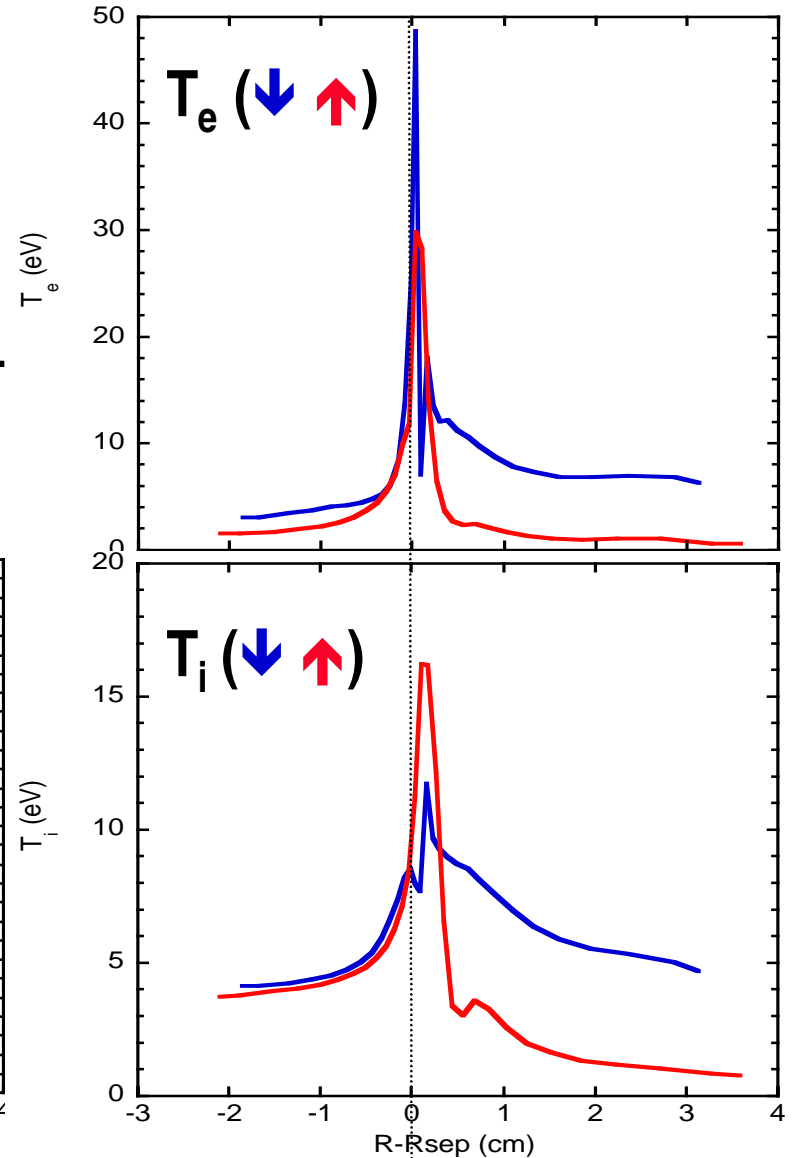
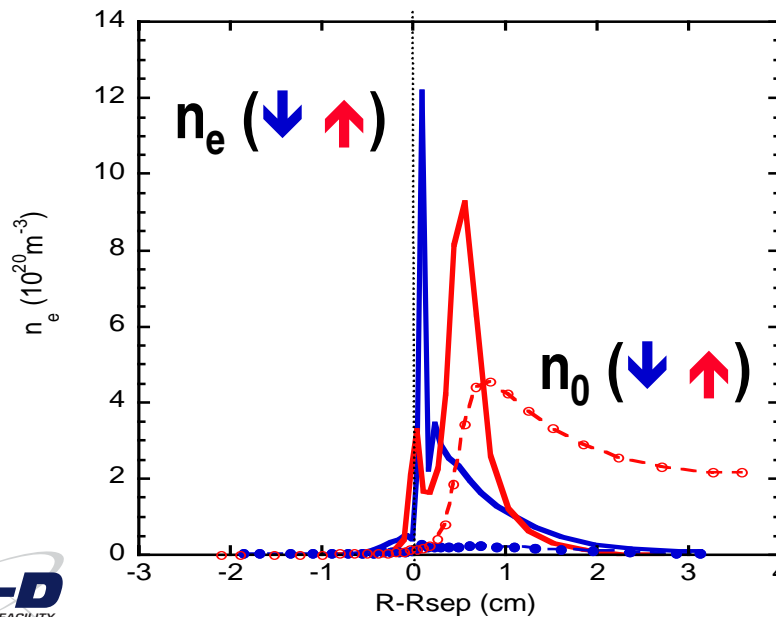
# UEDGE simulations of pre-ELM $\nabla B \downarrow$ vs $\nabla B \uparrow$ cases show completely different inner divertor conditions.

- $\nabla B \downarrow$  inner divertor - Detached
  - $n_e$  and  $n_0$  high to far SOL
  - $T_e$  and  $T_i$  below 2 eV except very close to ISP
- $\nabla B \uparrow$  inner divertor - Attached
  - $n_e$  2x lower and  $n_0$  10x lower in SOL
  - $T_e$  and  $T_i = 4 - 8$  eV throughout SOL



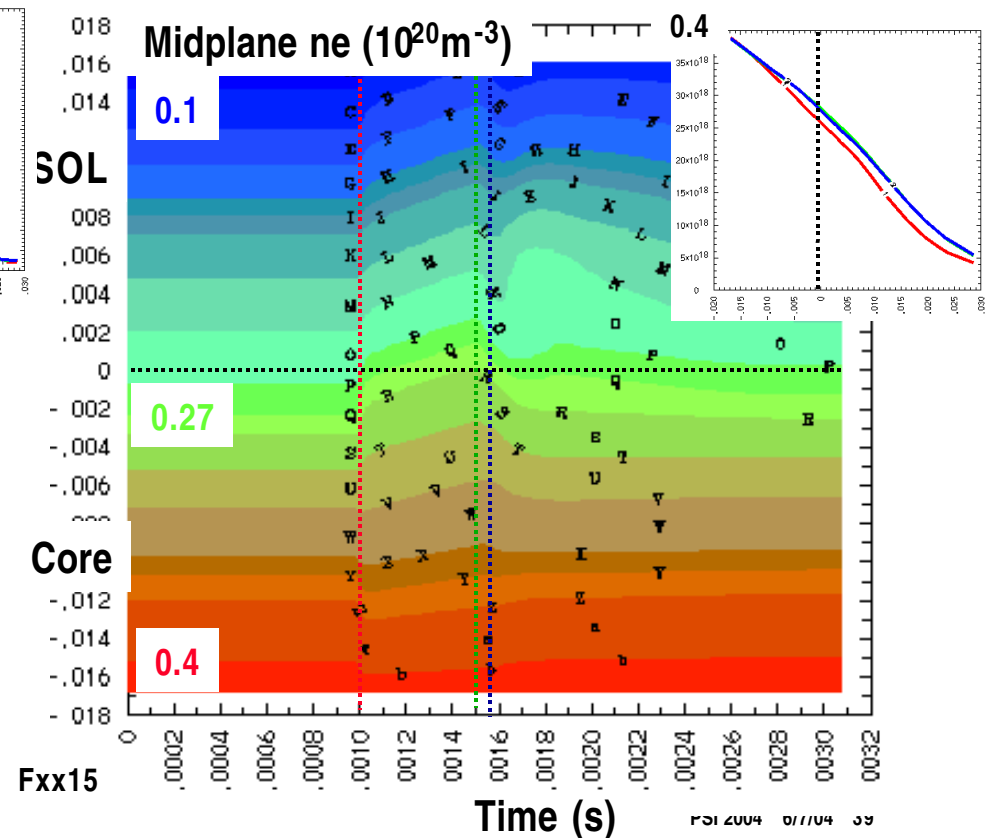
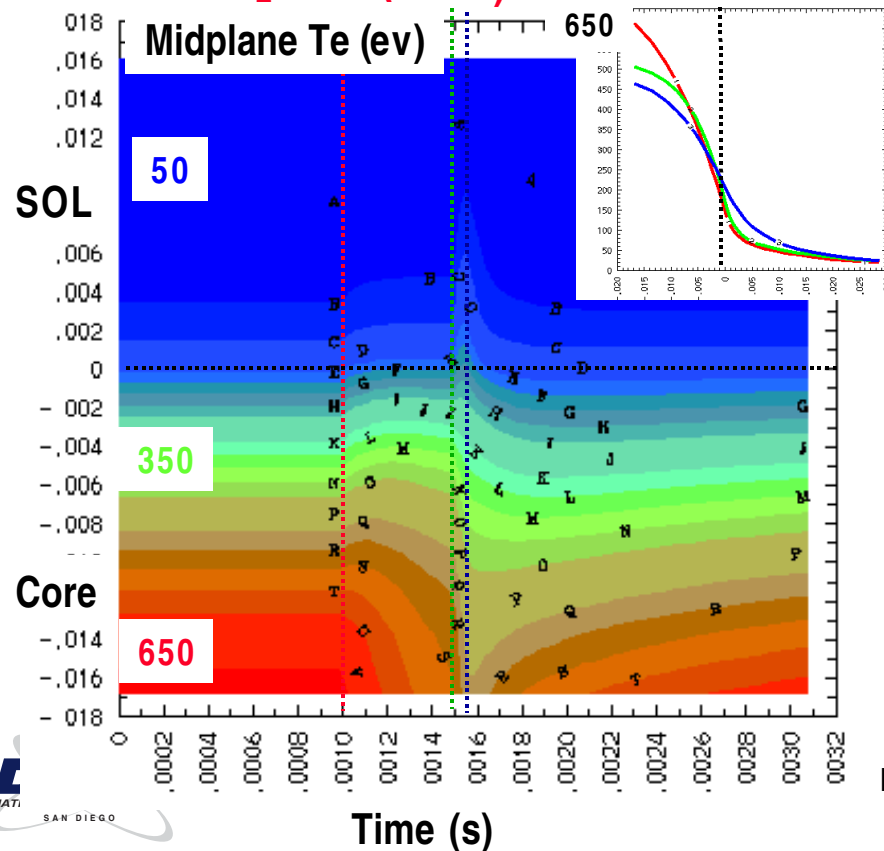
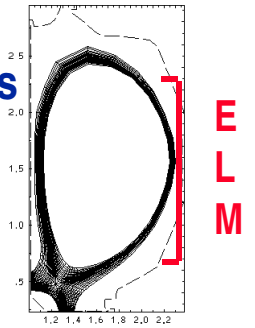
# UEDGE simulations of pre-ELM $\nabla B \downarrow$ vs $\nabla B \uparrow$ cases show very different outer divertor SOL plasma

- $\nabla B \downarrow$  outer divertor - Attached far SOL
  - $n_e$  and  $n_0$  low to far SOL
  - $T_e$  and  $T_i = 5-8$  eV into far SOL
- $\nabla B \uparrow$  outer divertor - Detached far SOL
  - $n_e$  5x higher and  $n_0$  10x higher in SOL
  - $T_e$  and  $T_i < 2$  eV throughout SOL



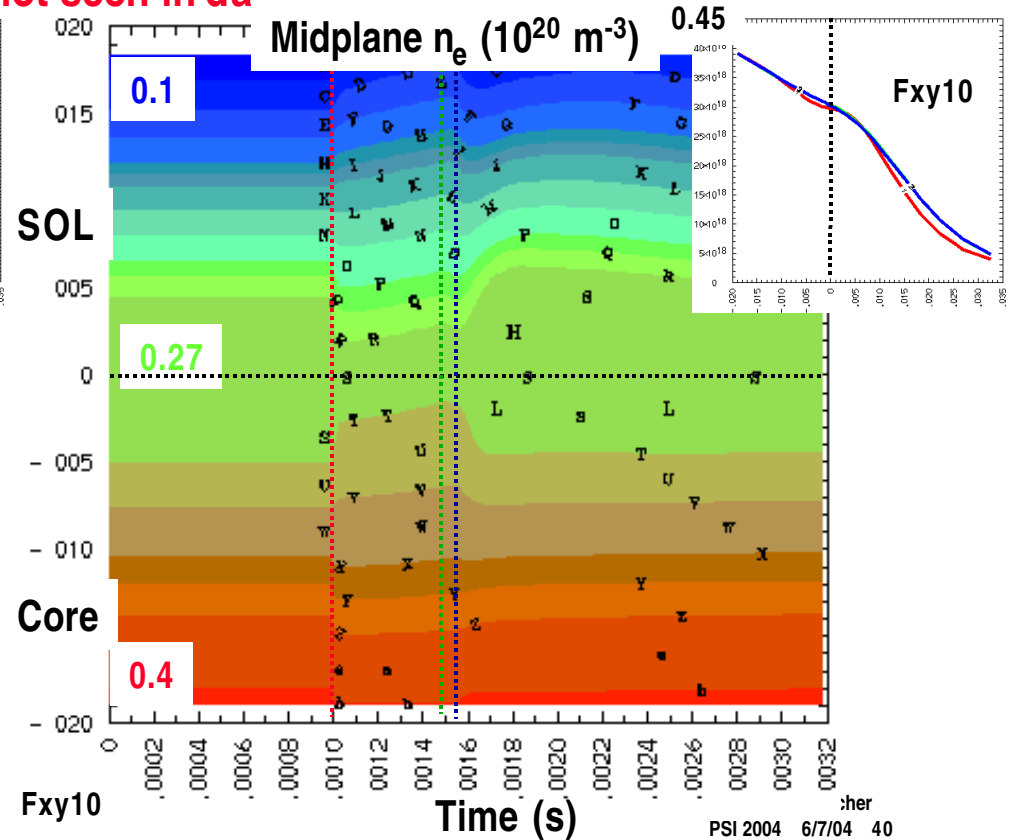
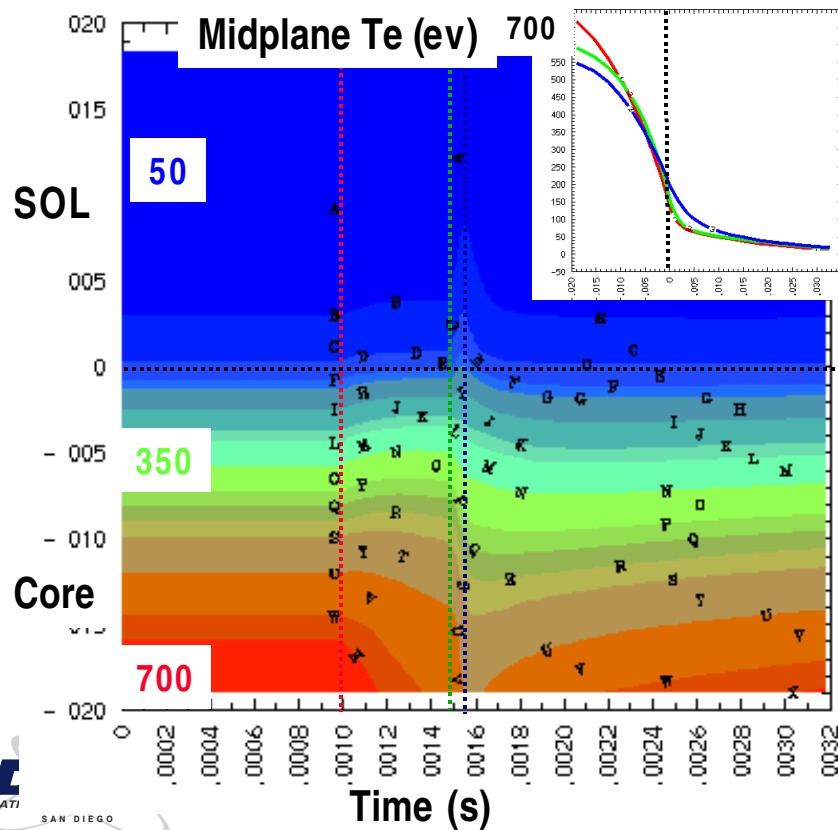
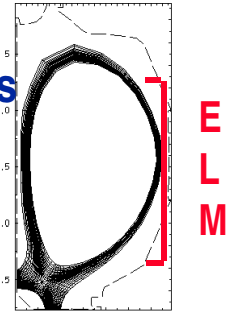
# UEDGE ELM simulation shows pedestal behavior similar to low $n_e$ case

- Exponential radial and Gaussian poloidal perturbation near midplane
- At 1.0 ms, increase  $D_{\perp}$  by 10x for 500  $\mu$ s, then add increase of  $\chi$  by 10x for 50  $\mu$ s
- Relaxation phase with transport coefficients from between-ELM solution
- Pedestal  $T_e$  loss without substantial SOL increase - similar to data
- Pedestal  $n_e$  loss (small) with  $n_e$  increase into far SOL - similar to data



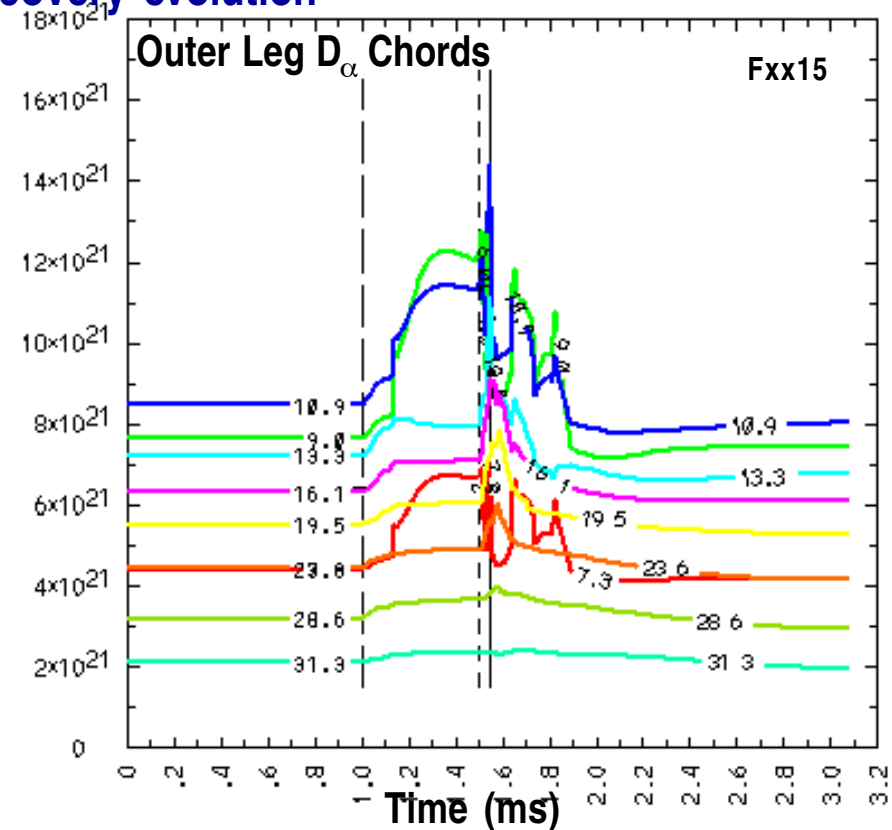
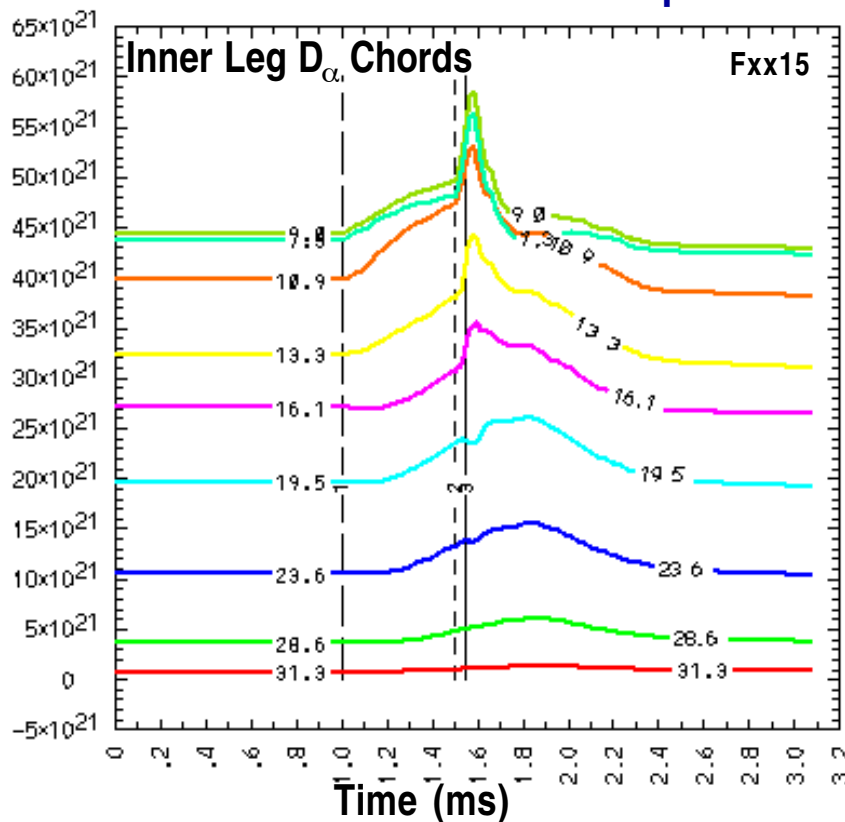
# UEDGE ELM simulation with $\nabla B \uparrow$ shows weak perturbation of midplane profiles

- Exponential radial and Gaussian poloidal perturbation near midplane
- At 1.0 ms, increase  $D_{\perp}$  by 10x for 500  $\mu\text{s}$ , then add increase of  $\chi$  by 10x for 50  $\mu\text{s}$
- Relaxation phase with transport coefficients from between-ELM solution
- Almost no SOL  $T_e$  perturbation - similar to data
- SOL density bump flattens during ELM - not seen in data



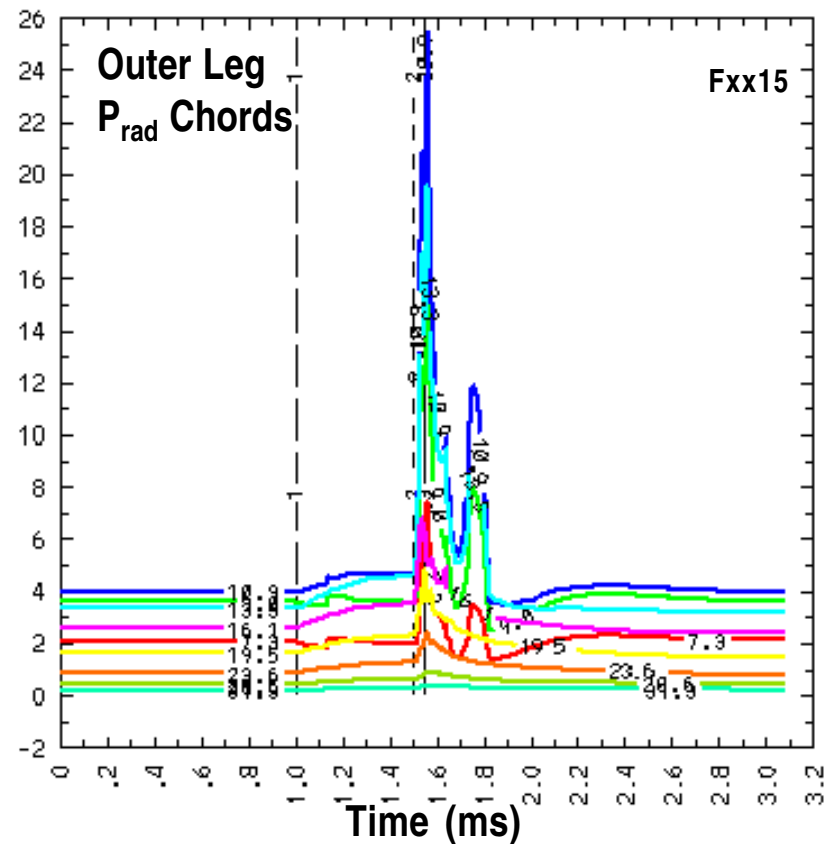
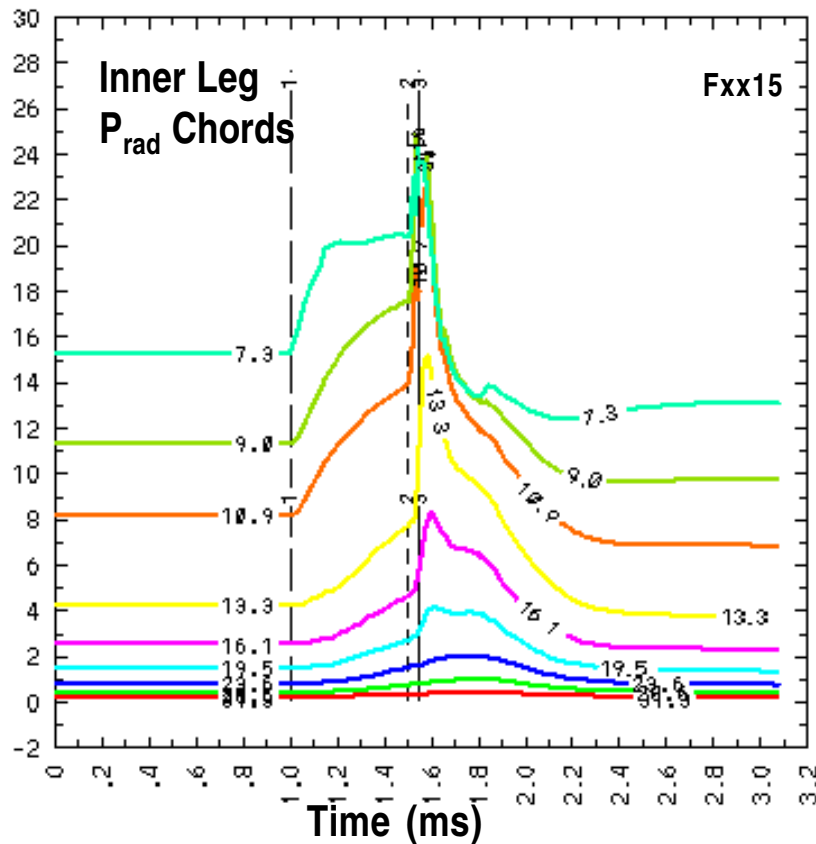
# $D_{\alpha}$ chord integrals vs. time from UEDGE solution simulate filterscope signals

- Inner Divertor:
  - Initial slow  $D_{\alpha}$  rise at  $D_{\perp}$  increase
  - Fast  $D_{\alpha}$  rise at  $\chi$  increase
  - Long slow (several ms) recovery on recycling timescale
- Outer Divertor:
  - Similar response to  $D_{\perp}$  and  $\chi$  increases
  - More complicated recovery evolution



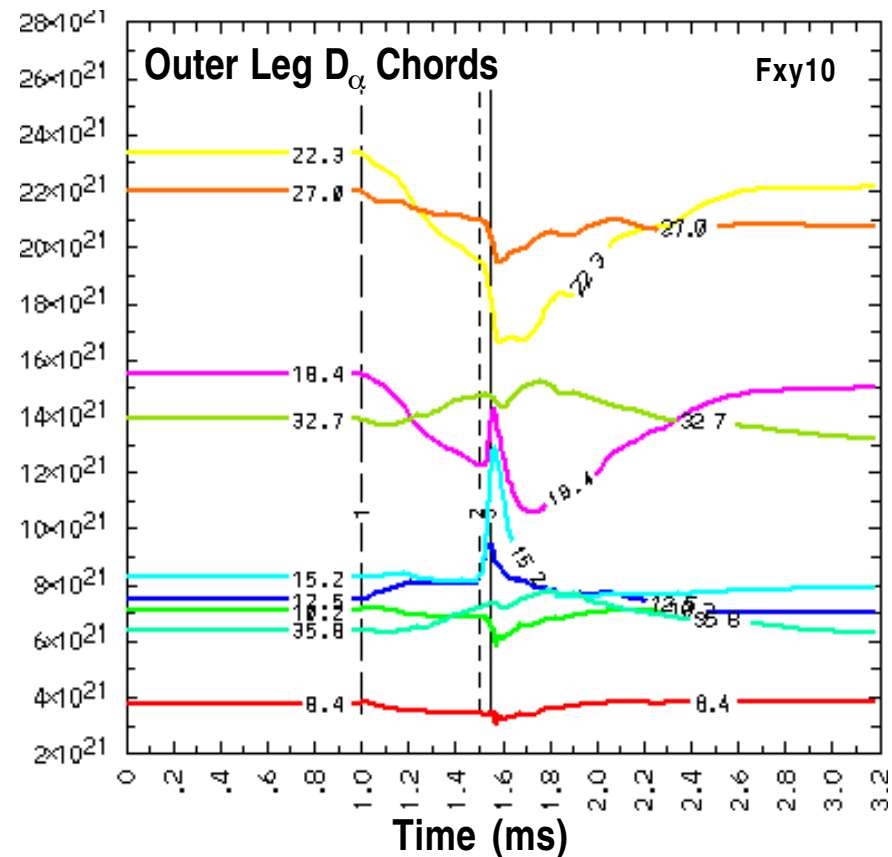
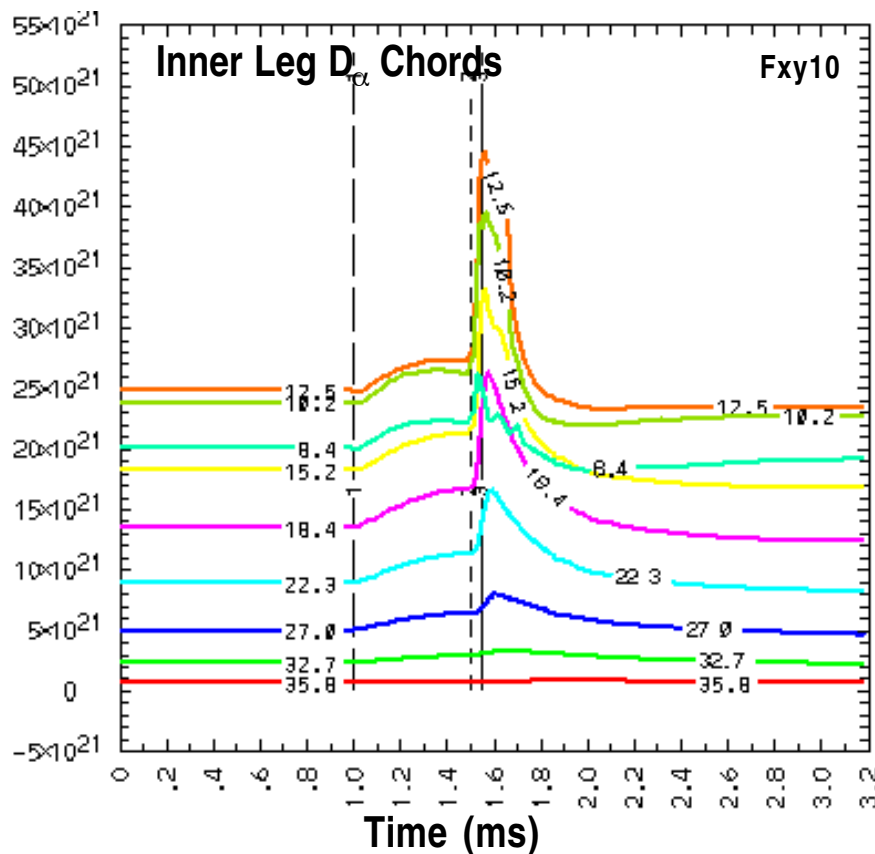
# $P_{\text{rad}}$ chord integrals vs. time from UEDGE solution simulate DISRAD-II signals

- Inner Divertor: - Sharper rise at  $D_{\perp}$  change than in  $D_{\alpha}$   
- More rapid recovery than in  $D_{\alpha}$
- Outer Divertor - Relative response to  $\chi$  change much larger than for  $D_{\alpha}$



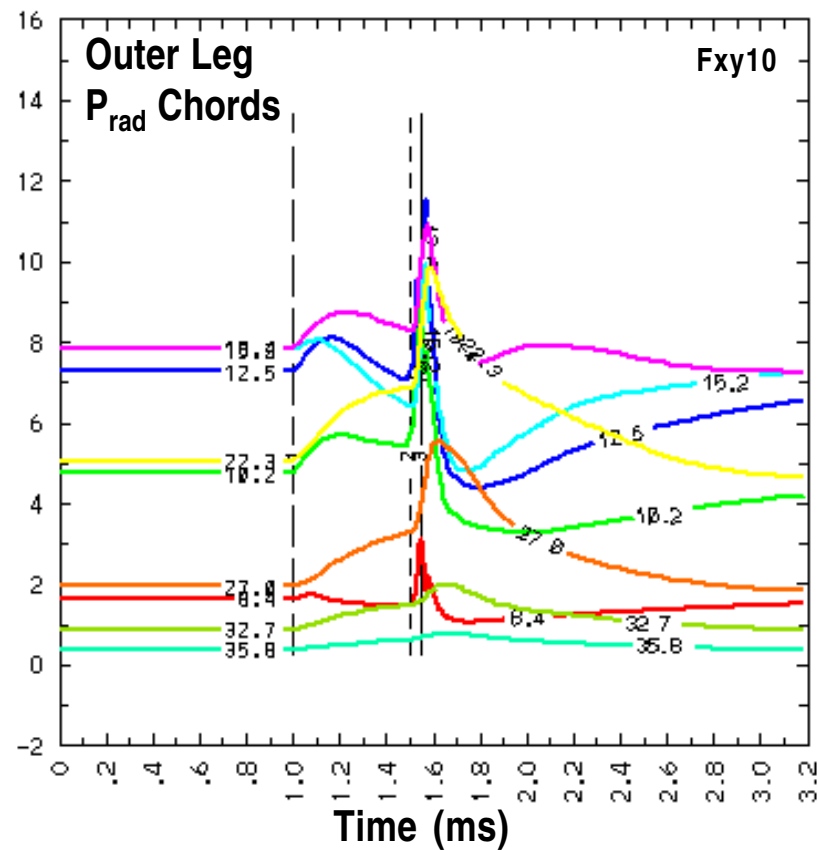
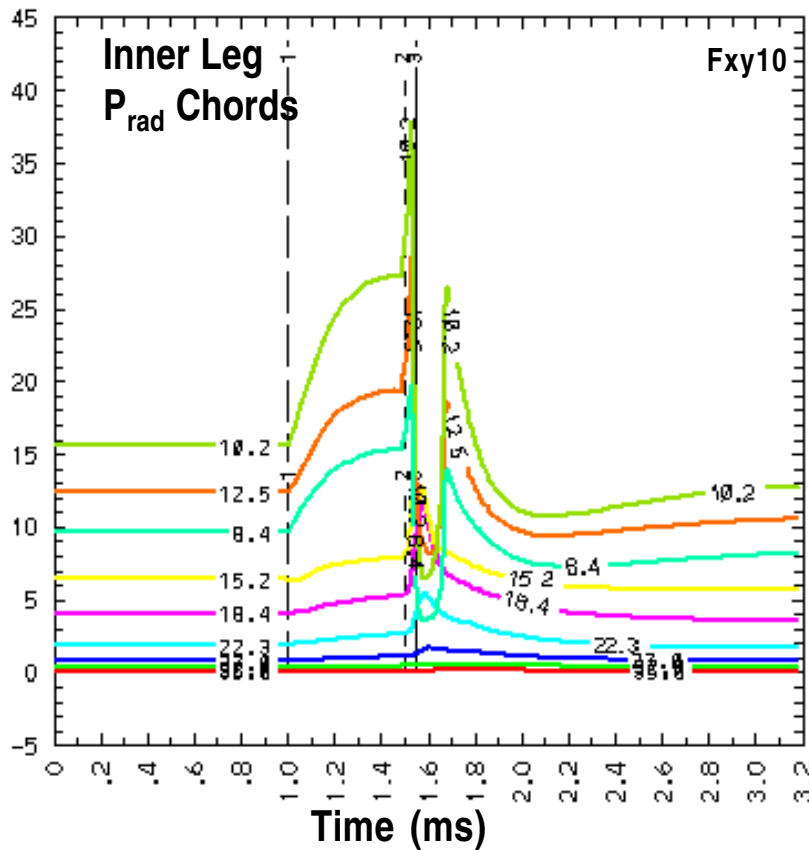
# $D_\alpha$ chord integrals vs. time from UEDGE solution with $\nabla B \uparrow$ simulate filterscope signals

- **Inner Divertor:**
  - Response to  $D_\perp$  change similar in  $\nabla B \downarrow$  and  $\nabla B \uparrow$
  - Response to  $\chi$  change is larger in  $\nabla B \uparrow$  than in  $\nabla B \downarrow$
- **Outer Divertor** - Positive and negative response to both  $D_\perp$  and  $\chi$  changes - **Unexplained**



# $P_{\text{rad}}$ chord integrals vs. time from UEDGE solution with $\nabla B \uparrow$ simulate DISRAD-II signals

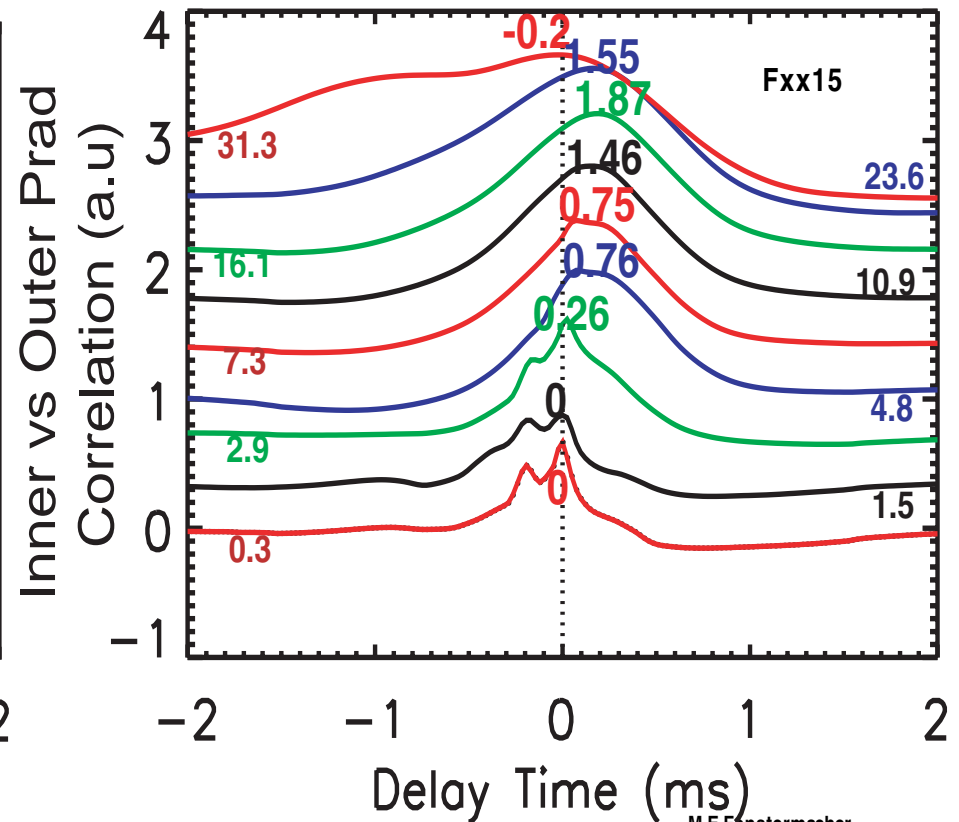
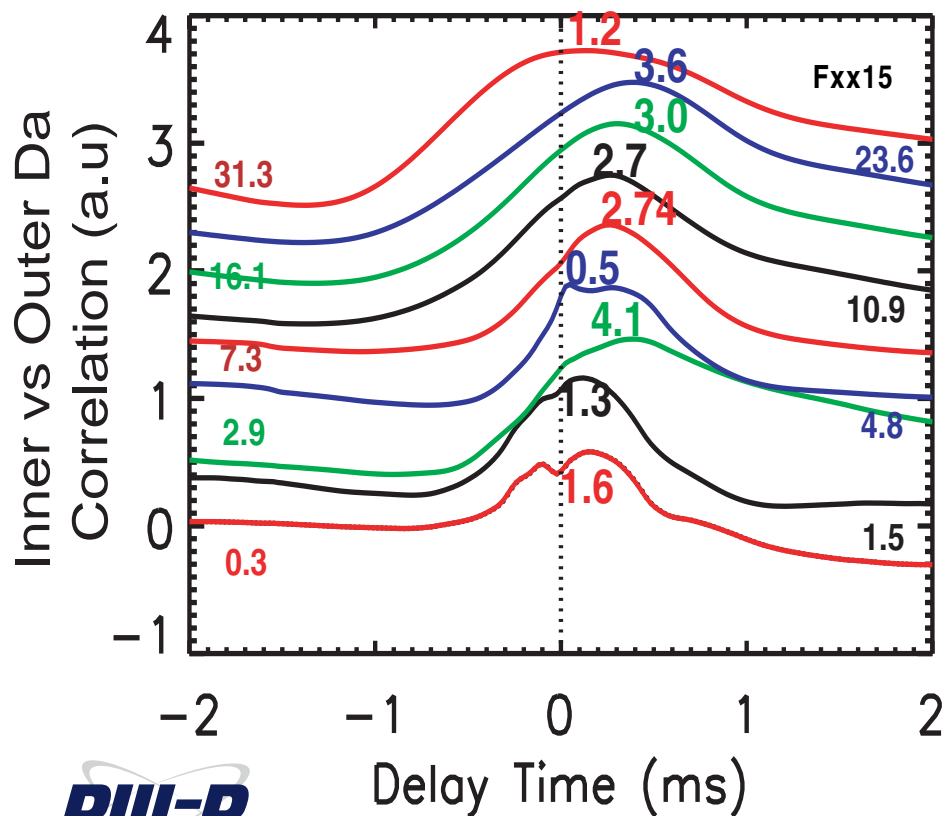
- **Inner Divertor:** - Initial response to  $D_{\perp}$  and  $\chi$  change similar in  $\nabla B \downarrow$  and  $\nabla B \uparrow$   
- Recovery phase more complicated in  $\nabla B \uparrow$  than in  $\nabla B \downarrow$
- **Outer Divertor** - Response to  $\chi$  much less in  $\nabla B \uparrow$  than in  $\nabla B \downarrow$





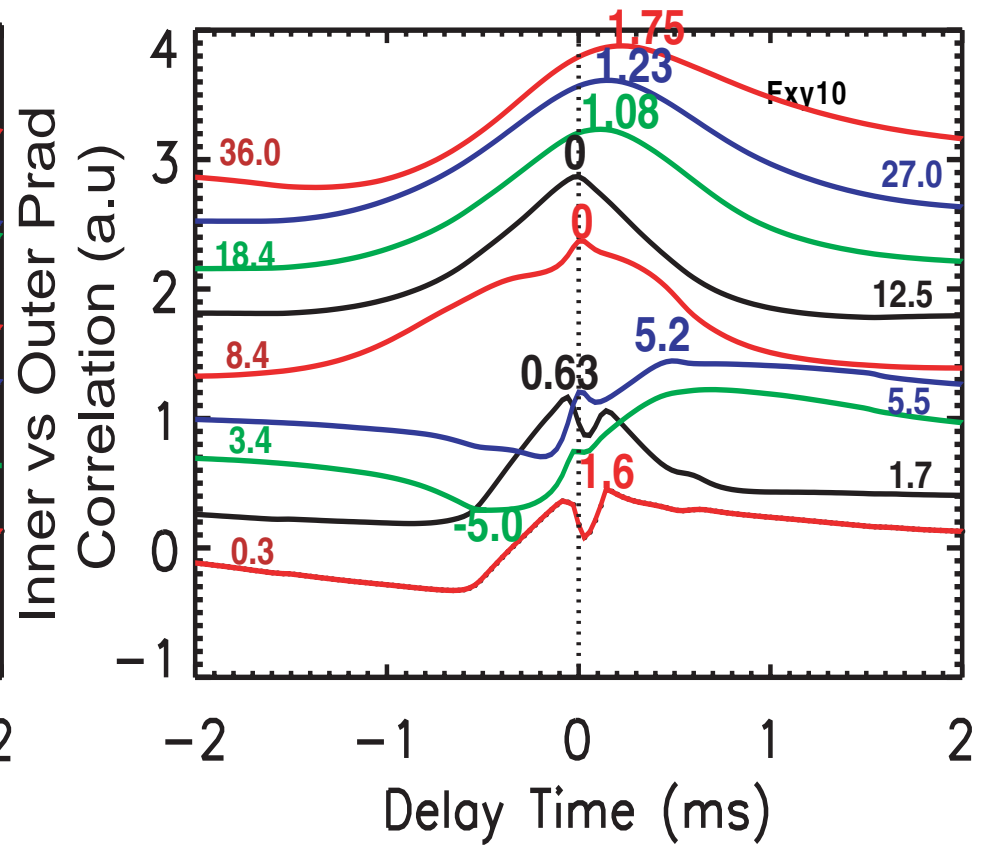
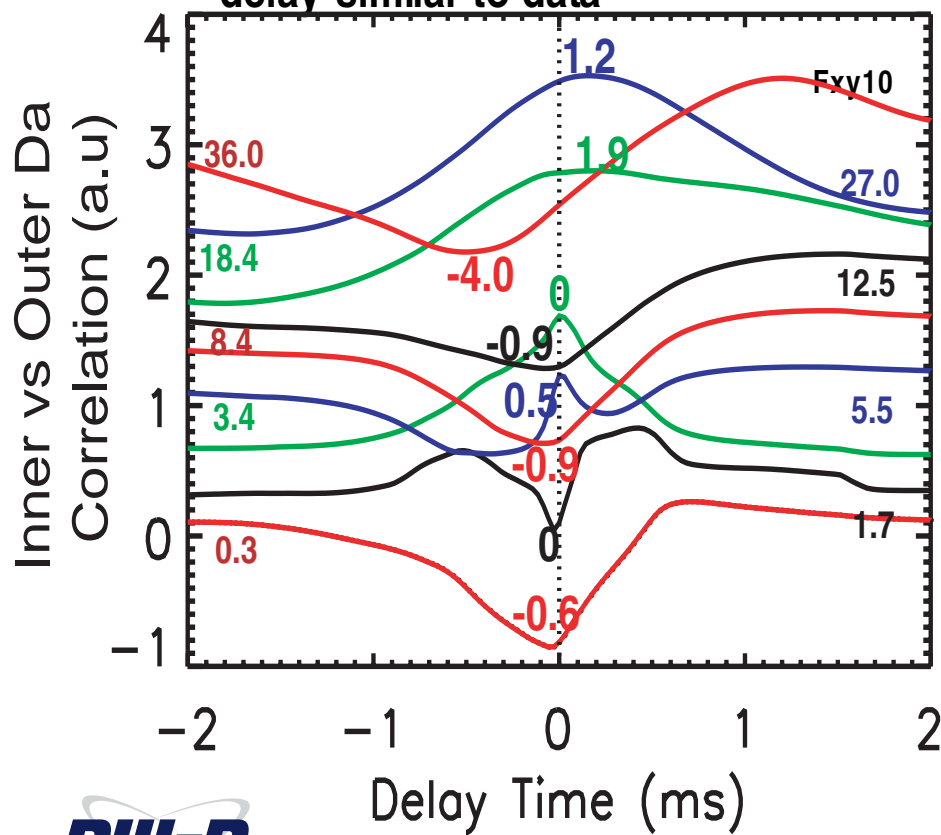
# Inner vs Outer correlation of UEDGE simulated Da and Prad signals show features similar to data correlations

- Correlation of inner vs. outer divertor synthetic DISRAD-II signals yields predictions of delays similar to observations
  - Normalized  $D_{\alpha}$  delay in the range [0.5 - 3.6] similar to data at  $n_e/n_{Gr} \sim 0.4$
  - Normalized  $P_{rad}$  in the range [0 - 1.9] - less delay than in  $D_{\alpha}$  as seen in the data



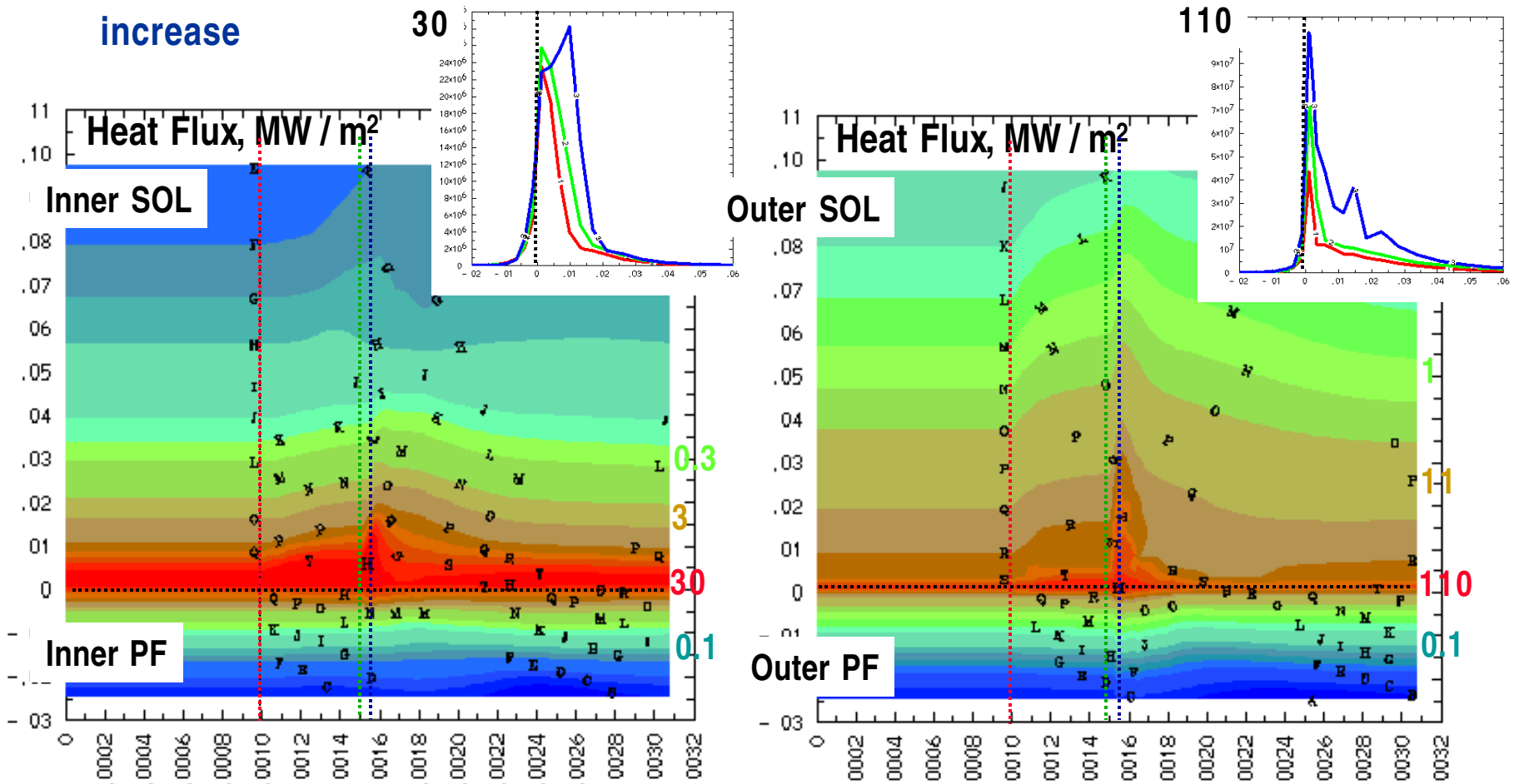
# Inner vs outer correlation of Prad signals with $\nabla B \uparrow$ show features similar to data correlations

- Correlation of inner vs. outer divertor synthetic  $D_\alpha$  and  $P_{rad}$  signals yields
  - Normalized delay of  $D_\alpha$  in the range [-4.0 - +1.9]: similar timing inversion occurs in data at  $n_e/n_G \sim 0.5$
  - Normalized delay in  $P_{rad}$  in the range [-5.0 - +1.8], However most radii have small delay similar to data



# Simulated inner and outer target heat fluxes broaden at most by a factor of 2 during ELM

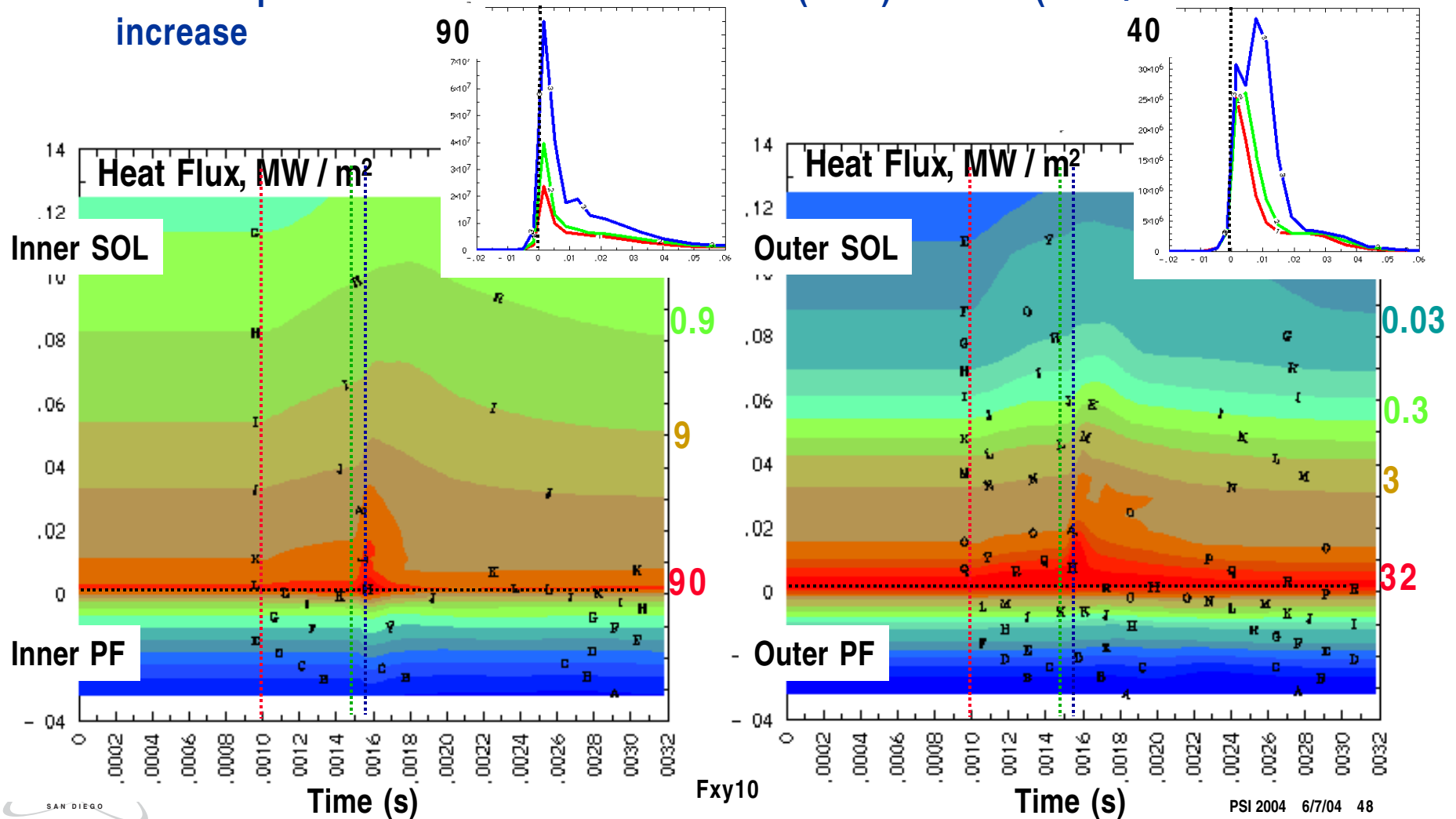
- Heat flux broadens by factors of 1.5 x (inner) and 1.2 (outer) during  $D_{\perp}$  increase
- Heat flux profile broadening increases to 2.0x (inner) and 1.8 (outer) by end of  $\chi$  increase



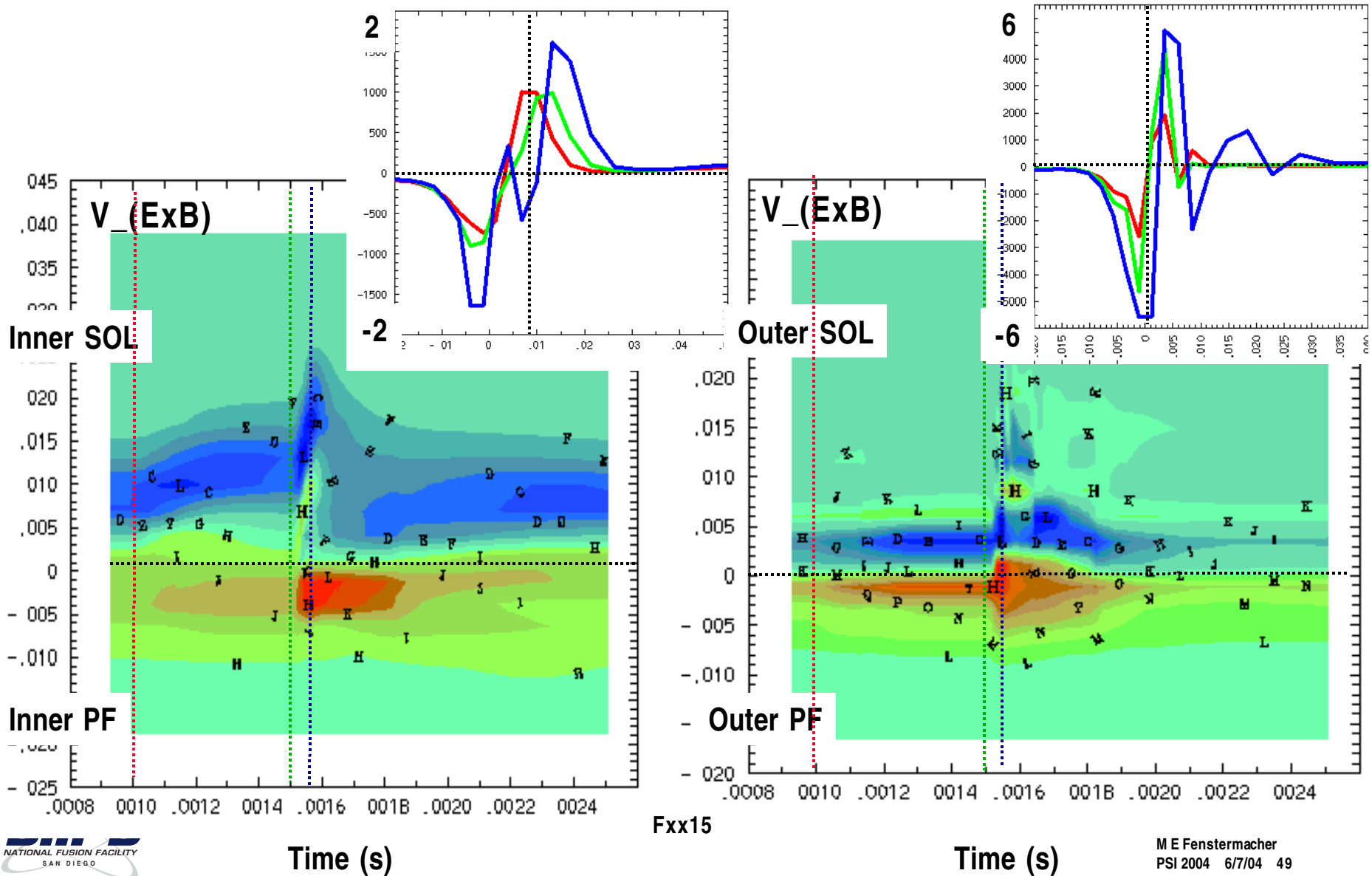
Fxx15

# With $\nabla B \uparrow$ broadening of heat flux during ELM less on inner and greater on outer target than with $\nabla B \downarrow$

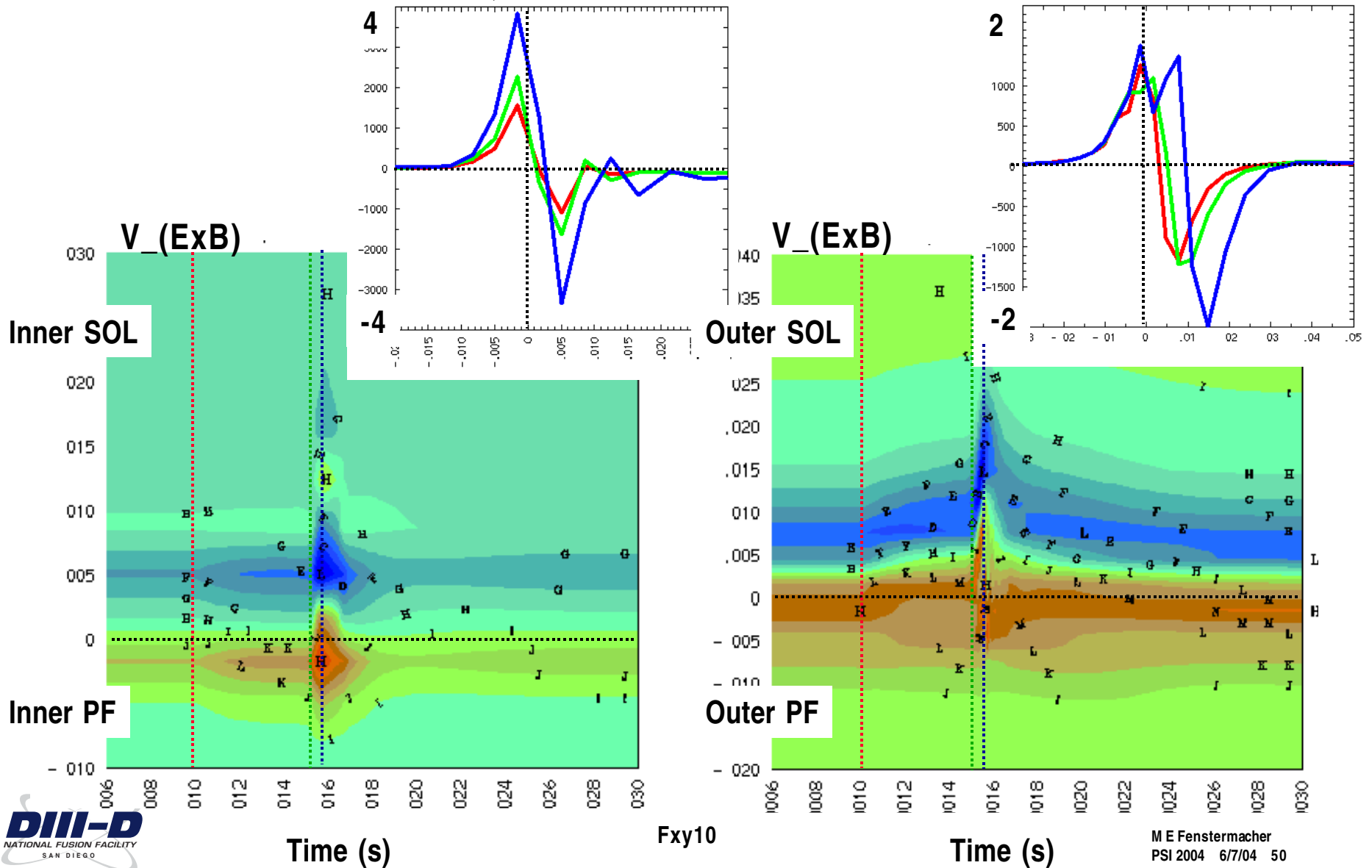
- Heat flux broadens by factors of 1.1 x (inner) and 1.5 (outer) during  $D_{\perp}$  increase
- Heat flux profile broadening increases to 1.5x (inner) and 2.2x (outer) by end of X increase



# With $\nabla B \downarrow$ private flux region poloidal ExB velocity increases 2x at ELM crash

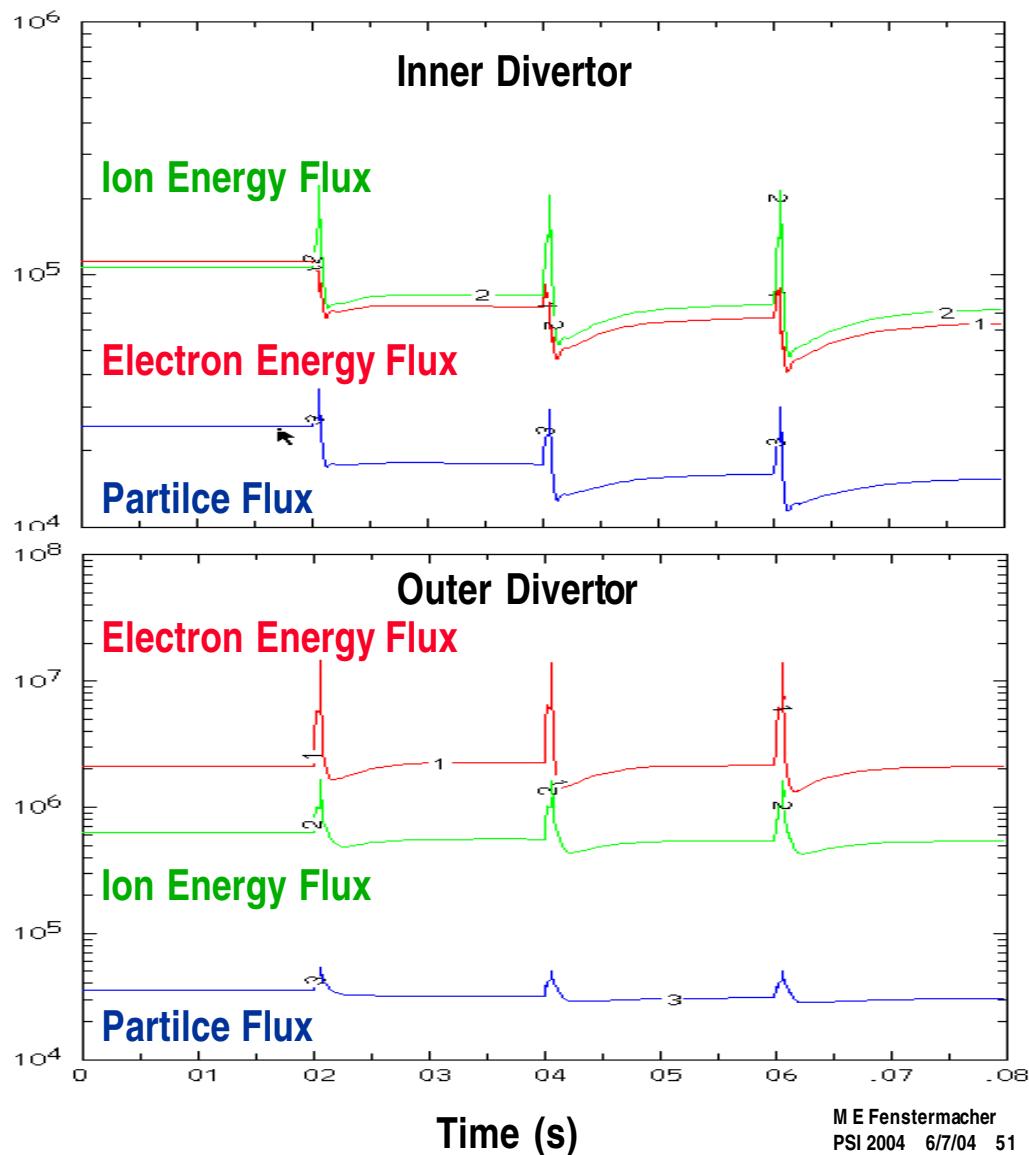


With  $\nabla B \uparrow$  change to PF poloidal  $v_{ExB}$  larger at inner target and smaller at outer target than with  $\nabla B \downarrow$



# Simulation with multiple ELMs shows slow relaxation to new parameter regime between ELMs

- Divertor particle and energy fluxes between ELMs are still evolving after 3 ELMs
  - Effect stronger on inner divertor
  - Indicates long time-scale effects (carbon, neutrals) still responding to ELMs
- Future single ELM simulations should start from “ELMing equilibrium” not steady state between-ELM solution



## Summary: SOL/divertor ELM behavior depends on both density and B-field dependent particle drifts

- Normalized delays of inner vs outer  $D_{\alpha}$  and  $P_{\text{rad}}$  depend on  $n_e$ 
  - Observations
    - Stronger  $n_e$  dependence in normal drifts direction
    - Delay greater and recovery longer for  $D_{\alpha}$  than for  $P_{\text{rad}}$
  - Possible Explanations
    - ELM poloidal character may change with  $n_e$
    - Fast electron effects may dominate at low  $n_e$ ; ion convection at high  $n_e$
- Normalized delays of inner vs outer  $D_{\alpha}$  and  $P_{\text{rad}}$  change with B-dependent drifts
  - Delays much less in reversed drifts case
  - Differences in pre-ELM divertor conditions with  $B_T$  play a role
  - Different response of  $E_r$  to ELM in normal and reversed drifts cases may affect  $E \times B$  drifts during ELM evolution
- Pedestal particles ejected far into SOL independent of  $n_e$  or drifts direction



## **Summary:** UEDGE ELM simulations including drifts show evolution and B-field dependence similar to data

- **Model of ELM as  $D_{\perp}$  and  $\chi$  increases supported by similarity of calculated and measured ELM evolution**
  - Initial response of simulated  $D_{\alpha}$  and  $P_{\text{rad}}$  to  $D_{\perp}$  increase and larger response to  $\chi$  increase similar to measured ELM signals
  - Pedestal density and temperature drops with SOL  $n_e$  increase and unchanged SOL  $T_e$  similar to data from low  $n_e$  plasmas
  - As in the data, simulated delays larger for  $D_{\alpha}$  than for  $P_{\text{rad}}$  in normal drifts case; small  $P_{\text{rad}}$  delays and positive/negative delay in  $D_{\alpha}$  for reversed drifts case
  - Simulated  $Q_{\text{div}}$  broadens  $\sim 2x$  at ELM crash in normal drifts case, broadening less in reversed drifts case.
- **UEDGE cases with normal and reversed drifts shows B-field dependent features seen in data**
  - Delays in  $D_{\alpha}$  and  $P_{\text{rad}}$  less in reversed drifts simulations consistent with measurements
  - ELM perturbation of divertor  $E_r$  and poloidal particle drifts may contribute to divertor ELM response