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## GYROKINETIC SIMULATIONS OF ENERGETIC PARTICLE DRIVEN TAE/EPM TRANSPORT EMBEDDED IN ITG/TEM TURBULENCE

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Energetic particle (EP) transport from local high-*n* toroidal Alfvén eigenmodes (TAEs) and energetic particle modes (EPMs) has been simulated with the gyrokinetic code GYRO [1]. Linear and nonlinear simulations have identified a parameter range where the TAE and EPM are unstable alongside the well-known ion-temperature-gradient (ITG) and trapped-electron-mode (TEM) instabilities. A new eigenvalue solver in GYRO facilitates this mode identification. States of nonlinearly saturated local TAE/EPM turbulent intensity are identified, showing a "soft" transport threshold for enhanced energetic particle transport against the TAE/EPM drive strength parameter  $[(n_{EP}/n_e)(T_{EP}/T_e)(R/L_{nEP})]$ . The very long-wavelength TAE/EPM transport is likely saturated by nonlinear interaction with ITG/TEM-driven zonal flows. These fixed gradient length nonlinearly saturated states are accessible over a relatively narrow range of TAE/EPM drive strength. Within this range, and in the local limit employed, TAE/EPM driven transport more closely resembles driftwave microturbulent transport than "stiff" ideal MHD transport with a clamped critical pressure gradient. At higher critical drive strength, nonlinear saturation fails (EP transport increases without limit and background transport decreases). Presumably saturation is obtained by (quasi-linear) relaxation of the EP pressure gradient to this critical gradient.

Previously, GYRO was used to study turbulent transport of fusion-produced alpha particles from short-wave, electrostatic ITG/TEM turbulence [2]. The new work here treats EP transport from long-wave, EP-driven, Alfvénic TAE and EPM modes. Linear simulations, including a new GYRO eigenvalue solver, have been used to track local growth rates and frequencies of these local high-n instabilities. Nonlinear simulations track the induced energy and particle transport in all species. All simulations are local, which is appropriate for sufficiently small  $\rho_* = \rho_s / a$  where  $\rho_{\rm s} = c_{\rm s} / \Omega_{\rm ci}$  is the ion sound gyroradius and *a* is the plasma radius. The short-wave ITG/TEM modes, whose driving gradients are kept in the simulation, are on the scale  $\rho_s$ . A radial density gradient in the sparse EP population drives TAE/EPM modes at much longer length scales - on the order of  $\rho_{\rm EP}$ , the EP gyroradius. Figure 1 shows the  $k_{\theta} = nq/r$  spectrum of rates.

The TAE and EPM are simultaneously destabilized primarily by the EP spatial pres-



Fig. 1: Linear frequency  $\omega$  and growth rate  $\gamma$  for leading local modes with q=2, s=1, r/a=0.5, R/a=3,  $T_i=T_e$ ,  $T_{EP}=100T_e$ ,  $n_{EP}=0.025n_e$ ,  $a/L_{Ti}=a/L_{Te}=3$ ,  $a/L_{ni}=a/L_{ne}=1$ ,  $a/L_{nEP}=4$ ,  $a/L_{TEP}=0$ . Inset shows same data over a greater range.

sure gradient. Reduction of the background ion and electron gradients with fixed magnetic geometry (circular *s*- $\alpha$  with  $\alpha$ =0) modestly lowers TAE/EPM growth rates. All length scales contribute to EP transport. However, as EP energy increases large perpendicular EP orbits partially average out ITG/TEM fluctuations, decreasing transport [2]. EP transport from long-wavelength TAE/EPM fluctuations shows no such reduction.

Nonlinear simulations of 40 interacting modes were performed over a range of EP drive strengths, represented by the relative EP density  $n_{\rm FP}/n_{\rm e}$ . A clear onset of TAE/EPM fluctuations is observed as  $n_{\rm EP}/n_{\rm e}$  increases (Fig. 2). At sufficiently weak  $n_{\rm EP}/n_{\rm e}$  drive, a state of finite saturated TAE/EPM intensity can be found in GYRO. The nonlinear saturation mechanism is interaction of finite- $k_{\theta}$  fluctuations with sheared  $k_{\theta}=0$ zonal flows. These zonal flows are driven primarily by ITG/TEM transport, and the system fails to saturate when ITG/TEM turbulence is reduced either by reduction of background gradients or exclusion of modes in the ITG/TEM range. Systems with stronger TAE/EPM drive  $(n_{\rm EP}/n_{\rm e} > 0.8\%)$  also fail to nonlinearly saturate for the parameters given. Proximity to the high-nideal MHD ballooning mode stability limit, and the related "sub-critical" limit in GYRO [3], is a likely component in this failure to saturate in this particular case.

In the range where fixed gradient length nonlinear saturation is obtained, the onset of TAE/EPM transport is not "stiff." Figure 3 shows a modest slope in the gyroBohm scaled EP diffusion coefficient with TAE/EPM drive strength  $(n_{\rm EP}/n_{\rm e} \text{ above } 0.5\%)$  corresponding to a stiffness of  $S = 1 + d \ln D_{\rm EP} / d \ln z = 1.5$ . Unlike the ideal MHD high-*n* ballooning mode onset, the high-*n* local EP-driven TAE/EPM threshold does not appear to impose a clamped EP critical pressure gradient. At a higher and critical drive strength  $[n_{\rm EP}/n_{\rm e}]$  above 0.8%] fixed gradient nonlinear saturation fails (EP transport increases without limit and background plasma transport actually decreases) and presumably the EP driving gradient relaxes to this critical EP pressure gradient.



Fig. 2: Time averaged fluctuation intensity per mode for  $n_{\rm EP}/n_{\rm e}=0.5\%$  (below TAE/EPM threshold, black) and  $n_{\rm EP}/n_{\rm e}=0.8\%$  (above TAE/EPM threshold, green). All other parameters the same as Fig. 1. Inset shows time traces of total finite- $k_{\theta}$  fluctuation intensity.



Fig. 3: GyroBohm scaled diffusion coefficient of energetic particle density  $D_{EP}/\chi_{gB}$  vs. TAE/EPM drive strength parameter (black curve). Low- $k_{\theta}$  linear growth rates shown for comparison (blue, right axis). Red dots and shaded areas had runaway at long times. Parameters same as in Fig. 1.

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