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The first systematic investigation of core electron thermal transport and the role of local ITG/TEM/ETG-scale core turbulence is performed in high temperature, low collisionality DIII-D H-mode plasmas. Core ITG/TEM-scale turbulence is substantially reduced/suppressed by $E \times B$ shear promptly after the L-H transition, resulting in reduced electron thermal transport across the entire minor radius. Nonlinear gyrokinetic (GYRO [1]) simulations indicate that a significant portion of the remaining H-mode electron heat flux results directly from residual short-scale TEM/ETG turbulence. The studies are performed at ITER-relevant collisionality ($\nu_e^* \sim 0.05$, $r/a \leq 0.6$) and are important since ITER plasmas will be electron heat-dominated due to α -particle heating. Taking advantage of the unique set of DIII-D turbulence [2,3] and profile diagnostics, experimentally determined H-mode electron transport fluxes and turbulence spectra are directly contrasted for the first time with nonlinear gyrokinetic simulation results.

Figures 1(a-c) show reduced ITG-scale \tilde{n} and \tilde{T}_e fluctuations and intermediate-scale core density fluctuations measured by Doppler backscattering (DBS) across the L-H transition (indicated by a vertical dashed line), and a comparison of the (normalized) maximum linear instability growth rate in the ITG/TEM range ($k_\theta \rho_s \leq 3$) with the $E \times B$ quench rate $\alpha_E \langle \omega_{E \times B} \rangle a/c_s$ evaluated from charge exchange recombination (CER) data. Here, $\langle \omega_{E \times B} \rangle a/c_s$ is the normalized Waltz-Miller shearing rate,

$\alpha_E \sim 0.3(\kappa)^{1/2}$ is the quench factor determined empirically by systematic variations of the shearing rate in GYRO simulations [4], and κ is the local plasma elongation. Linear instability growth rates are calculated using the trapped gyro-Landau fluid code (TGLF [5]). Around the time of the L-H transition [indicated by the drop in D_α line intensity, Fig. 1(e)], the shearing rate exceeds the maximum linear ITG/TEM growth rate in the core plasma (shown for $r/a=0.4$). In contrast, the ETG growth rate (not shown) always exceeds the shearing rate by a factor of 3-10. The electron temperature starts increasing first near the edge (pedestal formation) but also within 5-10 ms in the core [Figs. 1(d,g)]. Time-dependent TRANSP transport analysis shows that the electron diffusivity is significantly reduced across the minor radius [Fig. 1(f)], implying globally reduced electron thermal transport.

Figures 2(a,c) show wavenumber spectra of core density turbulence, measured by DBS. In L-mode, spectra are best described by an exponential dependence on $k_\theta \rho_s$ as observed previously in the TORE SUPRA tokamak. In H-mode, density fluctuation levels are reduced at low and intermediate wavenumber (ITG/TEM scale), and the spectra are flatter in this spectral region. At higher wavenumbers ($k_\theta \rho_s > 4$) fluctuation levels continue to decay exponentially. Comparing

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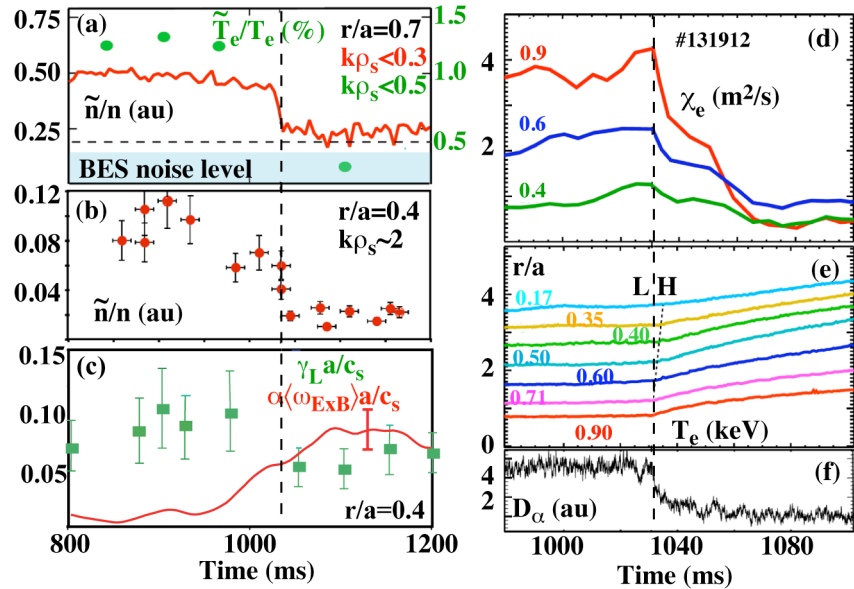


Fig. 1. (a) ITG-scale density and electron temperature fluctuations ($r/a=0.7$); dashed line: CECE detection limit; (b) density fluctuation level for $k_\theta \rho_s \sim 2$; (c) comparison of maximum normalized linear ITG/TEM growth rate ($k_\theta \rho_s < 3$, from TGLF) with the $E \times B$ quench rate ($r/a=0.4$); (d) electron thermal diffusivity from TRANSP for three normalized radii, (e) time evolution of local electron temperature; (f) D_α signal.

linear instability growth rates from TGLF with Miller-Waltz quench rates [Figs. 2(b,d)] we find that the wavenumber range of reduced large/intermediate-scale turbulence is well predicted by the wavenumber range where the quench rate $\alpha_E < \omega_{E \times B} > a/c_s$ exceeds the normalized linear growth rate.

Figure 3 shows the wavenumber spectrum of density fluctuations predicted by a preliminary non-linear GYRO simulation with 40 radial and 40 poloidal modes, spanning $1 \leq k_\theta \rho_s \leq 40$. The spectrum labeled $k_r=0$ takes into account the instrumental DBS sensitivity which intrinsically detects modes with $k_r=0$ [2]. Multi-scale simulations with improved low- k resolution are underway to more adequately simulate the experimental results. However, initial flux estimates show that turbulence in the range $k_\theta \rho_s > 1$ accounts for roughly 50% of the local electron thermal flux in H-mode ($r/a=0.6$).

In summary, we have observed significantly reduced large/intermediate-scale turbulence and reduced electron thermal flux in the core of low collisionality H-mode plasmas. Initial GYRO simulations indicate that TEM/ETG-scale turbulence accounts for a substantial part of the residual H-mode electron thermal flux. Measured wavenumber spectra show qualitative agreement with initial nonlinear gyrokinetic calculations which indicate a somewhat flat fluctuation spectrum in the ITG/TEM spectral range.

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- [1] J. Candy and R.E. Waltz, *J. Comput. Phys.* **186**, 545 (2003).
- [2] L. Schmitz, *et al.*, *Rev. Sci Instrum.* **79**, 10F113 (2008)
- [3] A.E. White, *et al.*, *Rev. Sci Instrum.* **79**, 103505 (2008)
- [4] J.E. Kinsey, *et al.*, *Phys. Plasmas* **15**, 055908 (2008).
- [5] G.M. Staebler, *et al.*, *Phys. Plasmas* **14**, 055909 (2007).

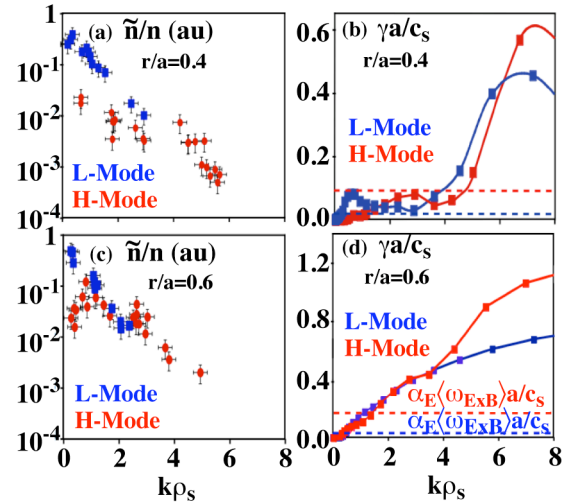


Fig. 2. Density fluctuation wavenumber spectra (L/H-mode) at $r/a=0.4$ (a) and $r/a=0.6$ (c); corresponding normalized linear growth rates and $E \times B$ quench rates (b,d).

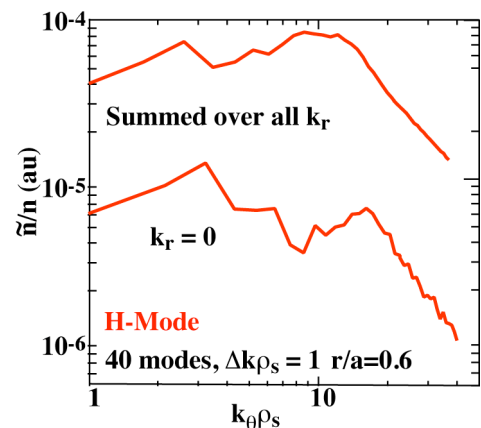


Fig. 3. Density fluctuation spectra from GYRO; the spectrum labeled $k_r=0$ takes into account the DBS detection sensitivity.