

Two sets of MHD simulations are presented in which resonant magnetic perturbation (RMP) vacuum fields are imposed on DIII-D equilibria

Three simulations with DIII-D shot 113317

(NIMROD code is used for all simulations)

Four simulations with DIII-D shot 119690

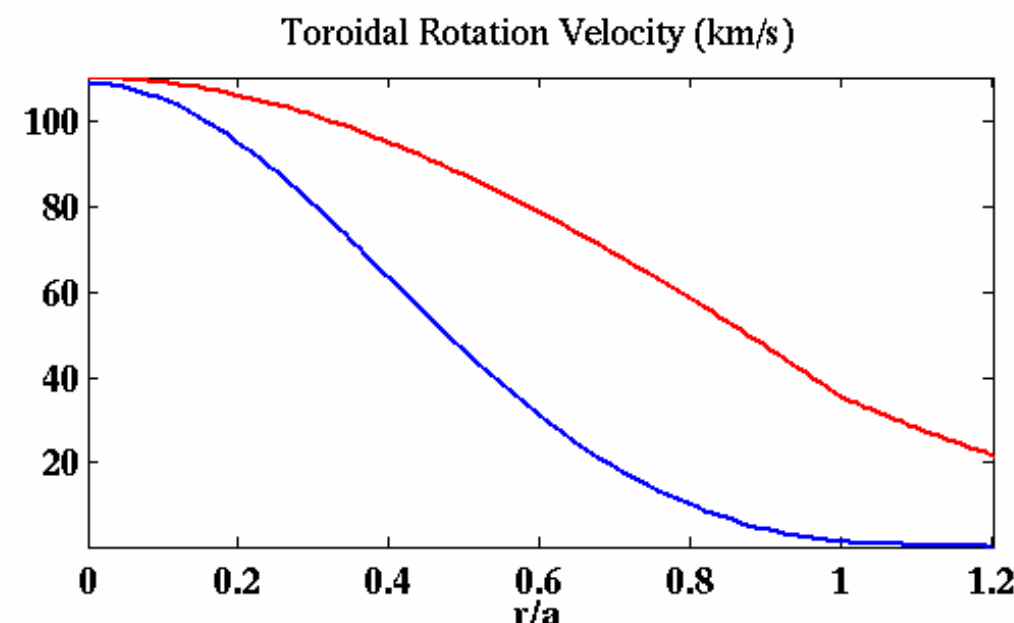
- Not an RMP experimental discharge • Measured density and temperature profiles
- Three artificially imposed rotation profile • 3 kA even parity I-coil currents
- Measured error fields + 1 kA C-coil currents • $S (= \tau_{L/R} / \tau_A) = 3 \times 10^5$

• Three simulations with:

1) No rotation

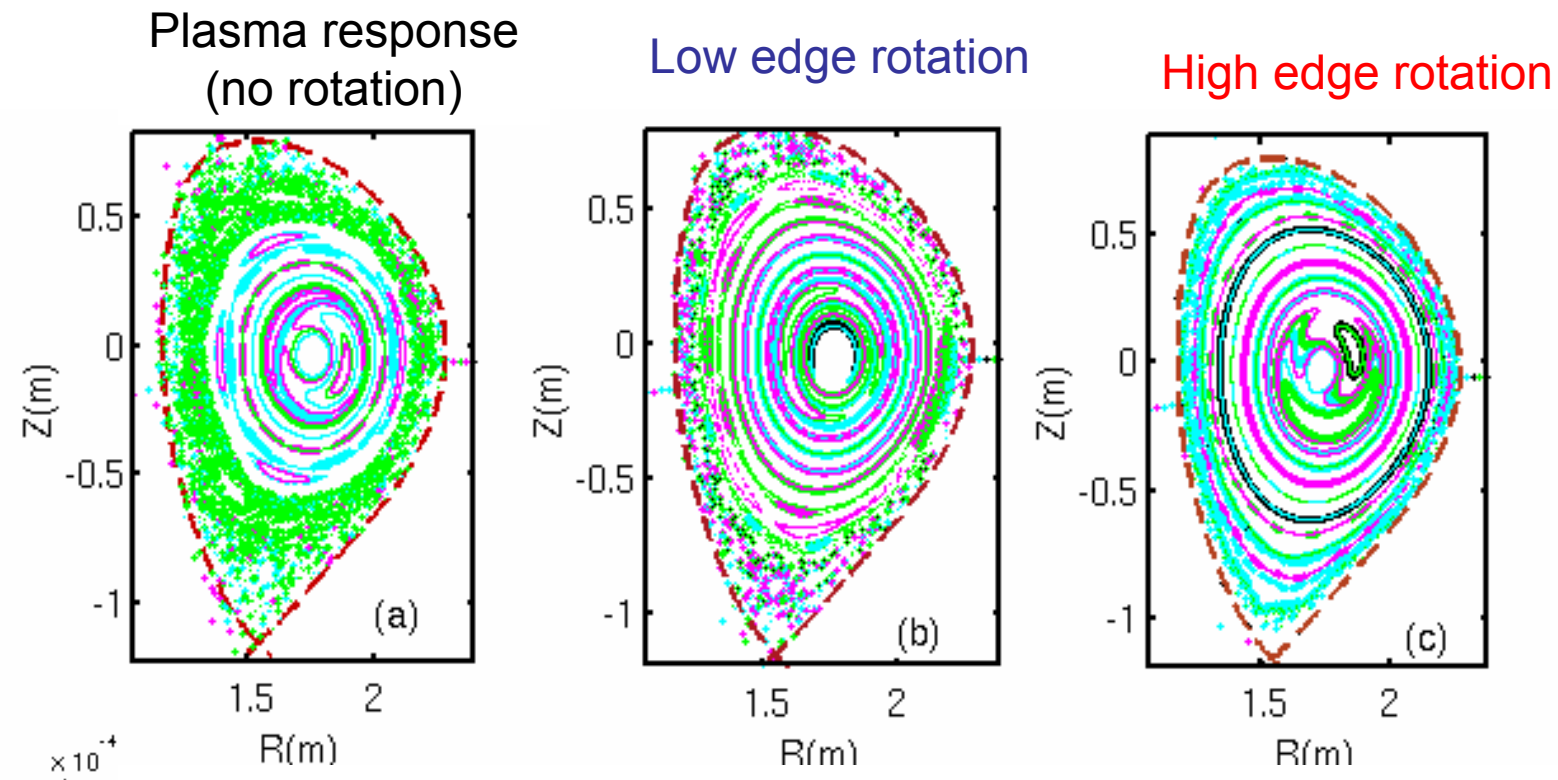
2) Low edge rotation

3) High edge rotation



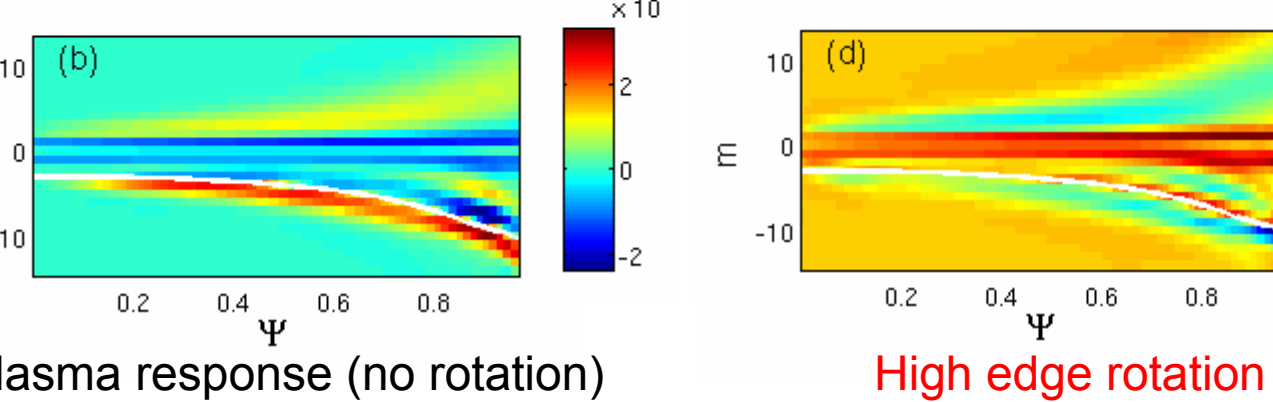
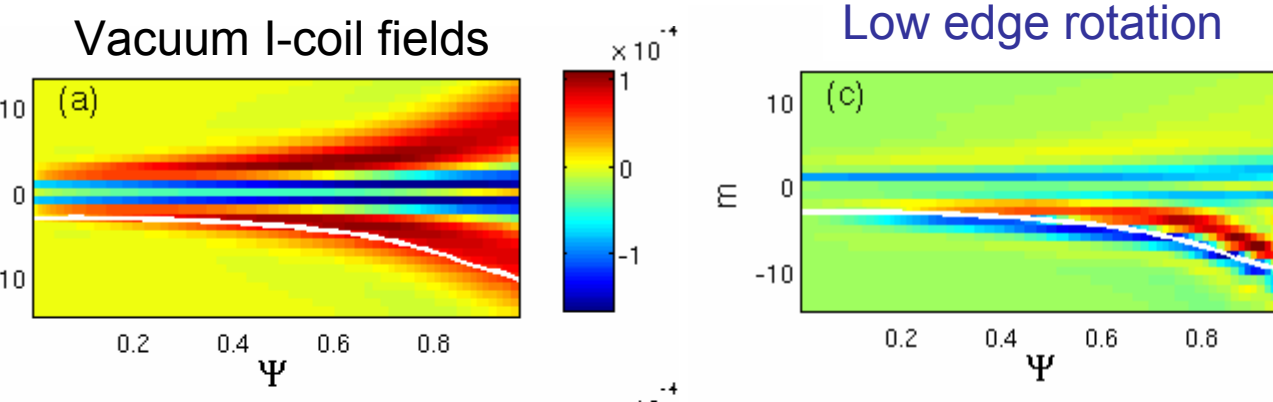
• Effects of rotational screening on applied RMP fields are investigated in real geometry

Each simulation is evolved for several reconnection times until a steady state magnetic field configuration is reached

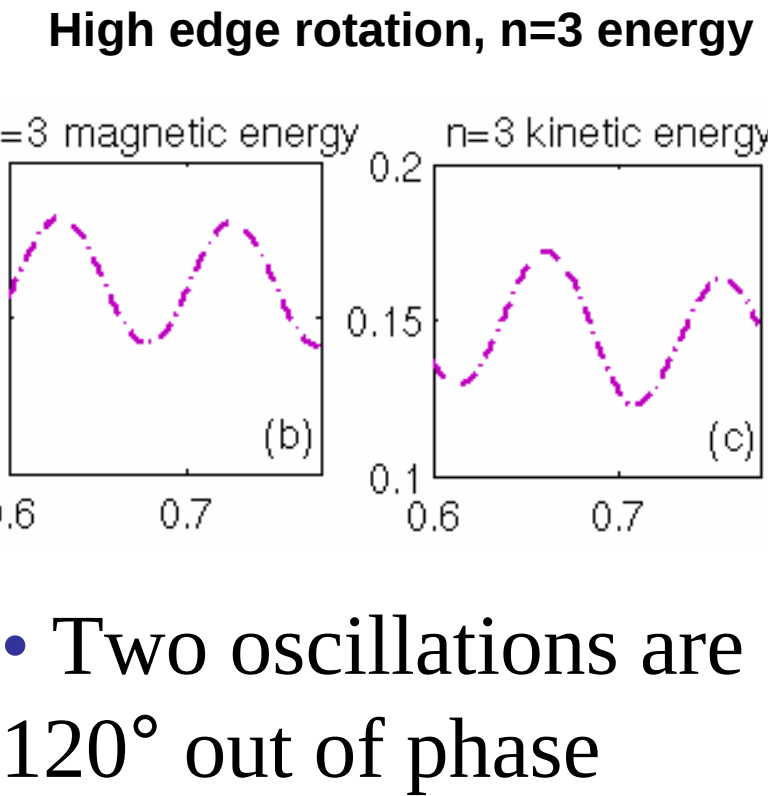
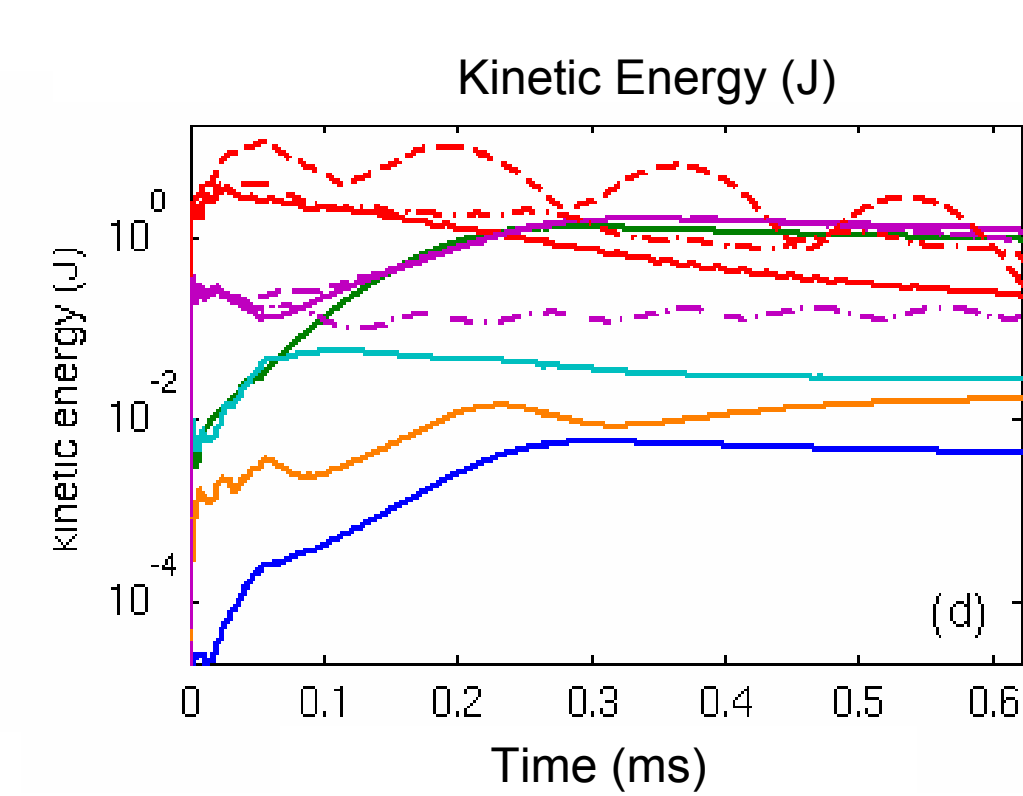
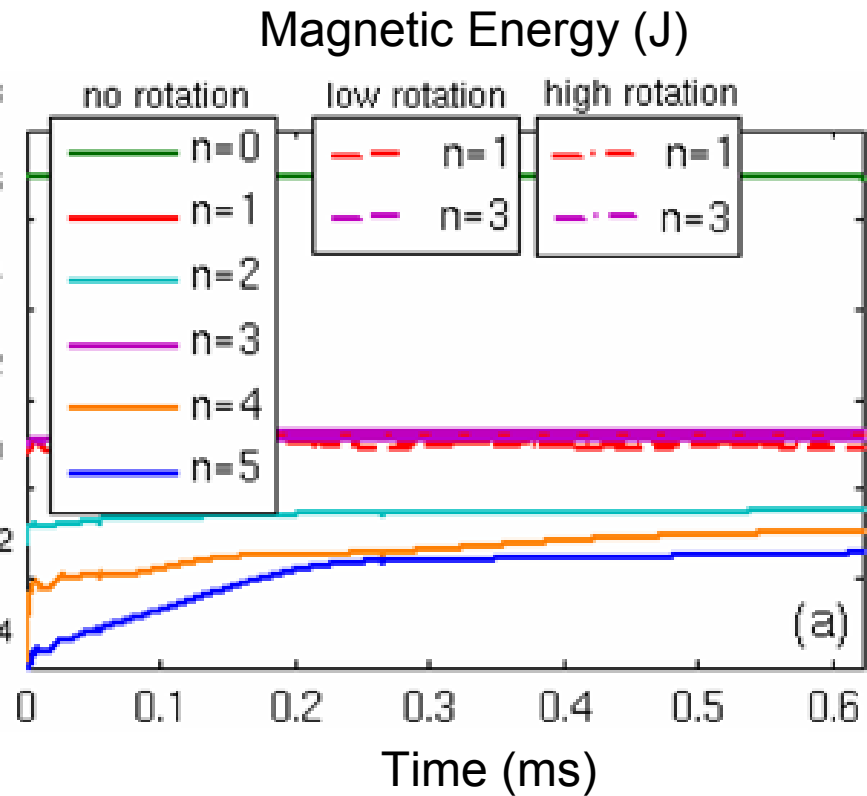


• Vacuum fields superposed on the equilibrium is a non-equilibrium initial condition

Normal component of the n=3 B-fields (T)
White line on each plot is the resonant line $m=3q$



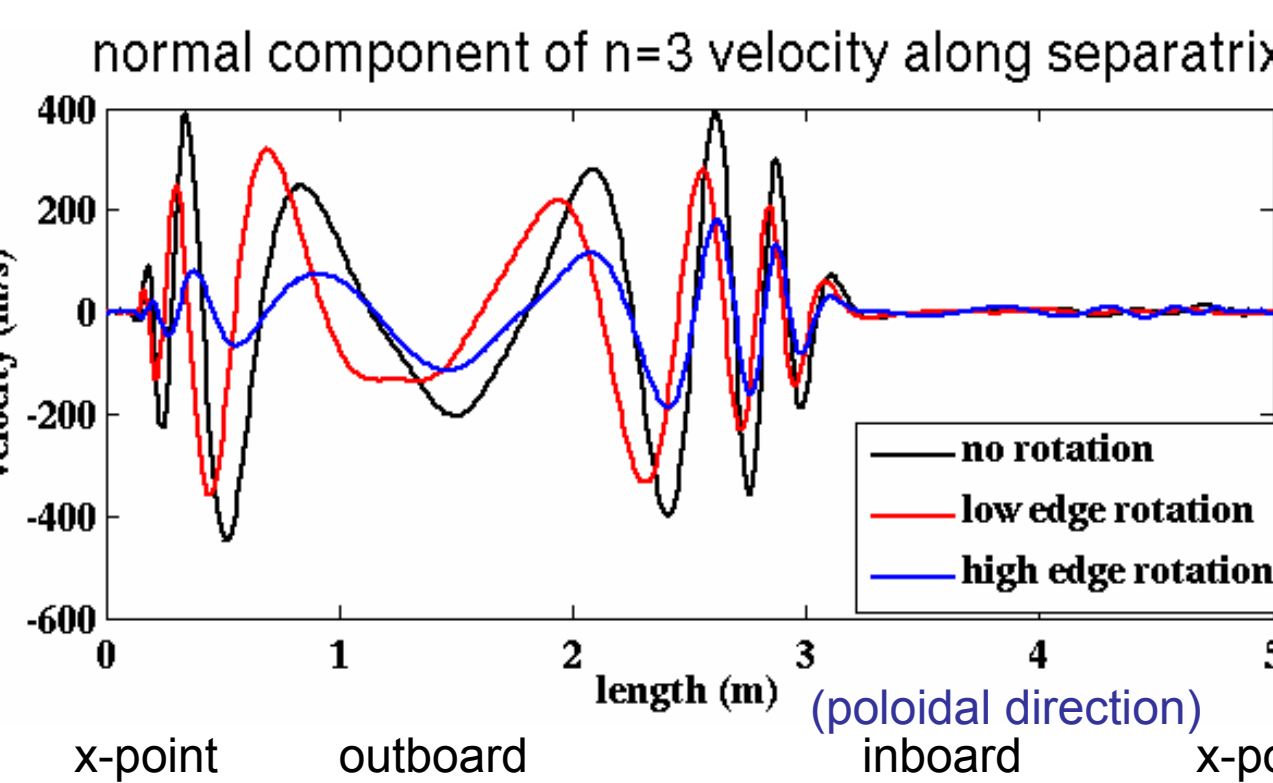
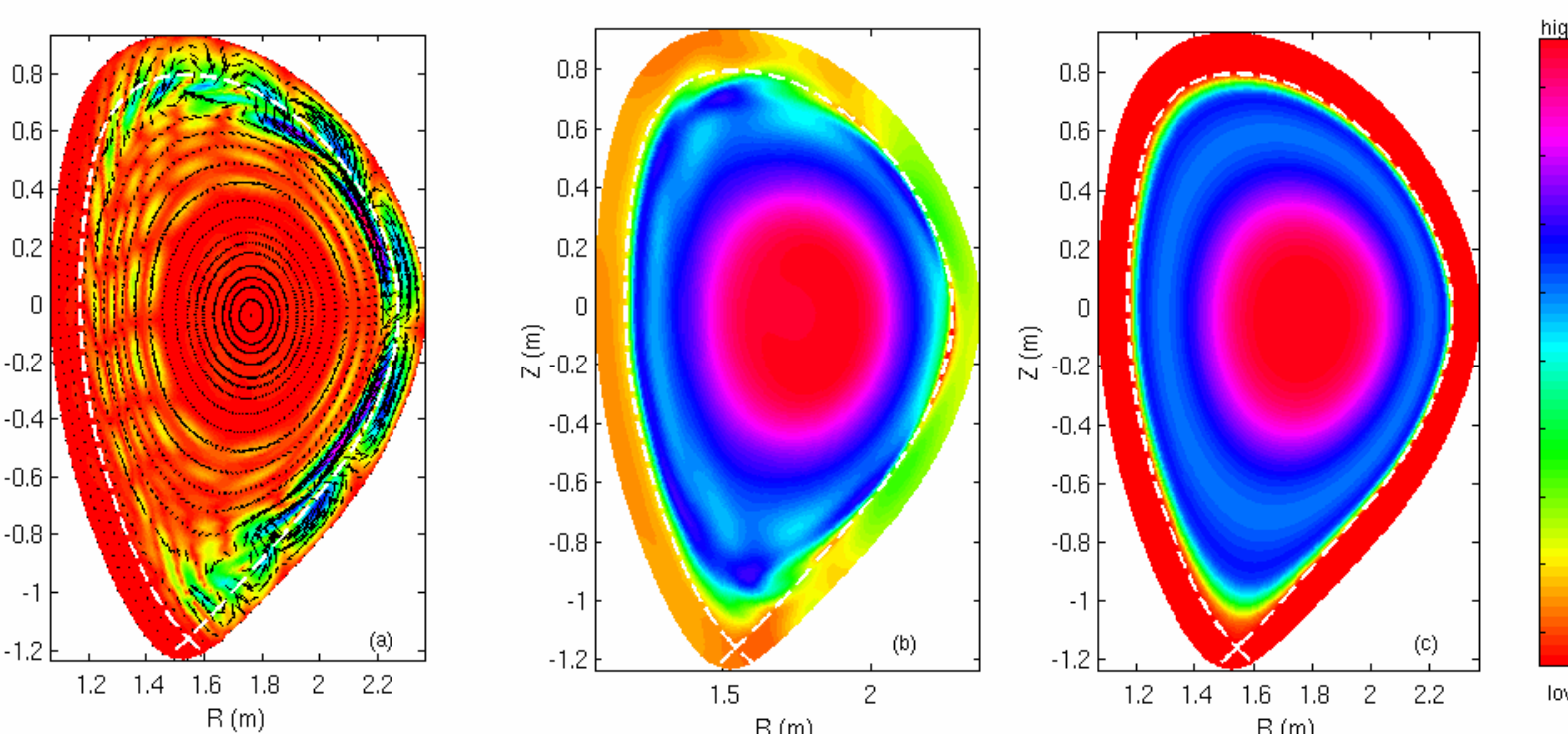
- Growth of n=3 kinetic energy observed in no rotation and low edge rotation simulations, but no growth for higher edge rotation



- 10 kHz n=3 oscillation for high edge rotation, and 5.8 kHz n=1 oscillation for low rotation are observed in the mode energy time history

- Density evolution depends on edge rotation velocity:

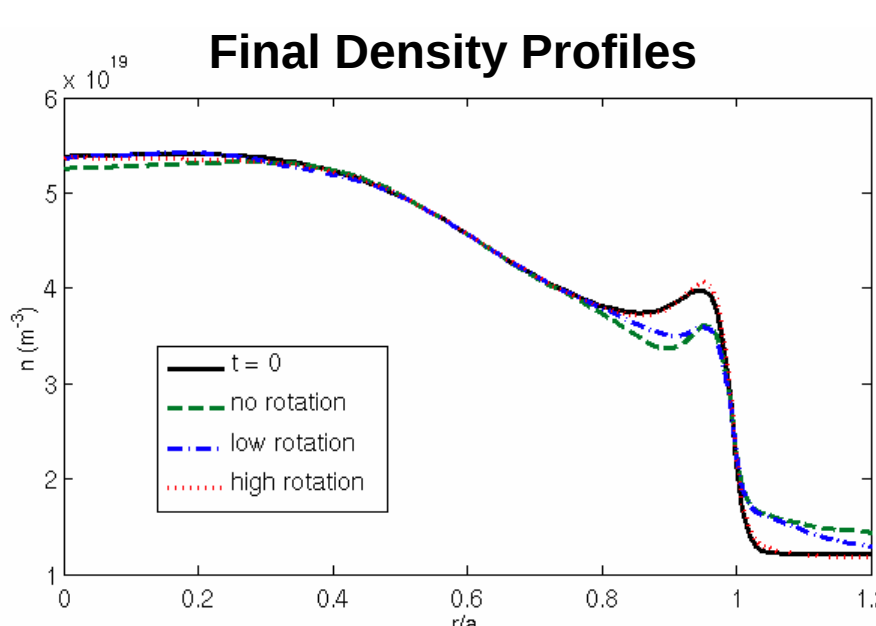
Poloidal Velocity Density (no rotation) Density (high rotation)



- Radial velocity at the separatrix is reduced as rotation is increased

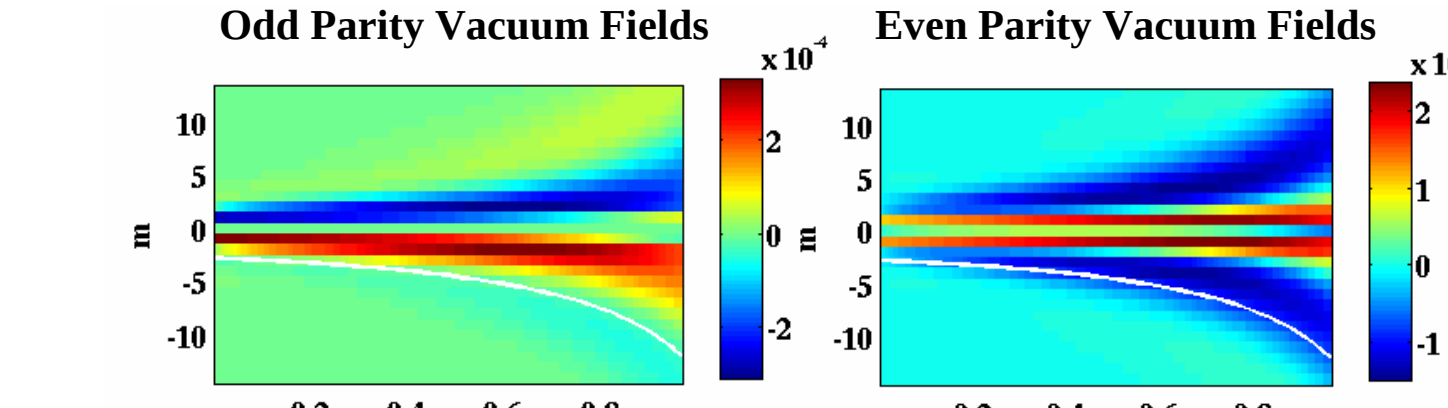
- Reduction by half for high edge rotation

- Pedestal density is reduced for no, low rotation, unchanged at high rotation

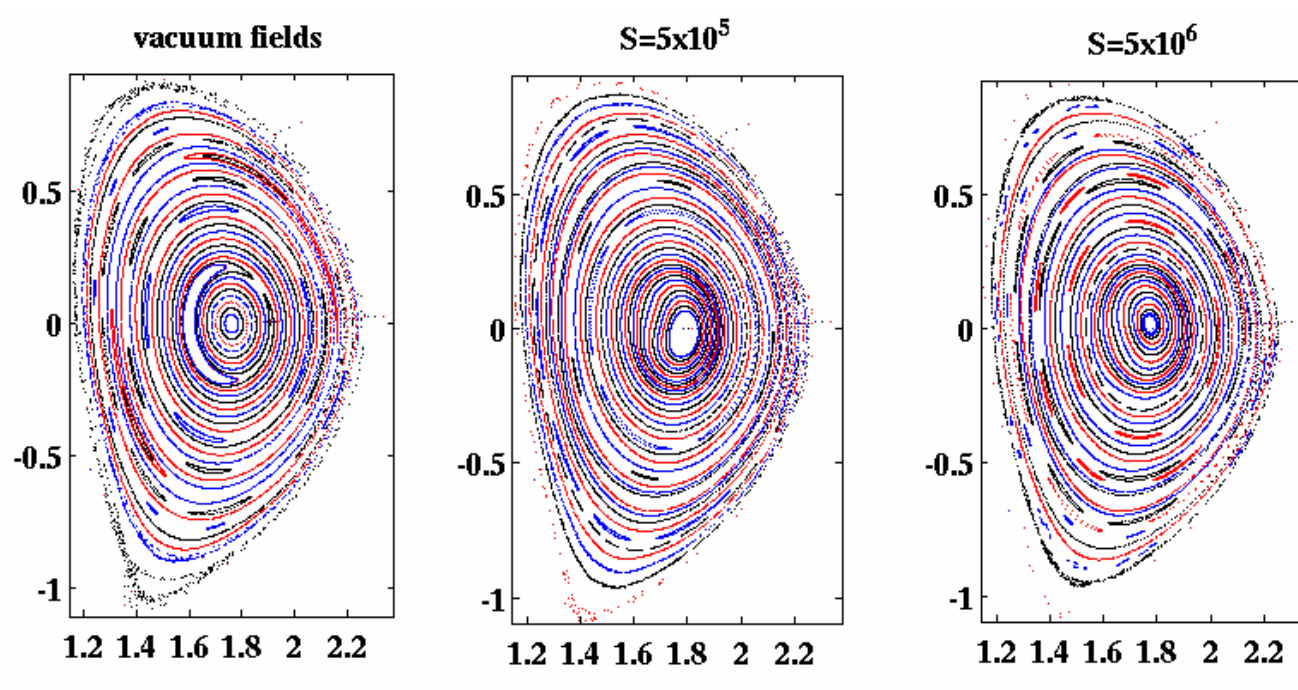


- Low collisionality odd parity I-coil experiment • Measured n_e and T_e profiles
- Measured rotation profile • 4.4 kA I-coil current (odd parity and even parity cases)
- Measured error fields, no I-coil current • $S=5 \times 10^5$ and 5×10^6

- Four simulations: even and odd parity I-coils; $S=5 \times 10^5$ and $S=5 \times 10^6$



- Even up-down I-coil parity has stronger resonant components than odd parity

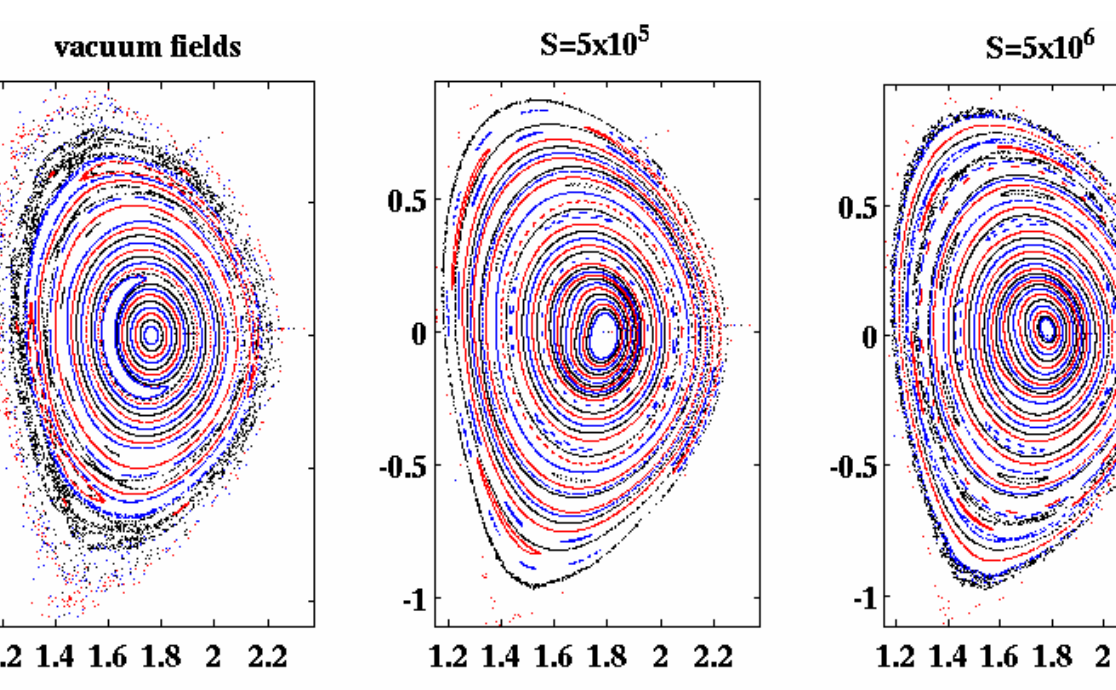
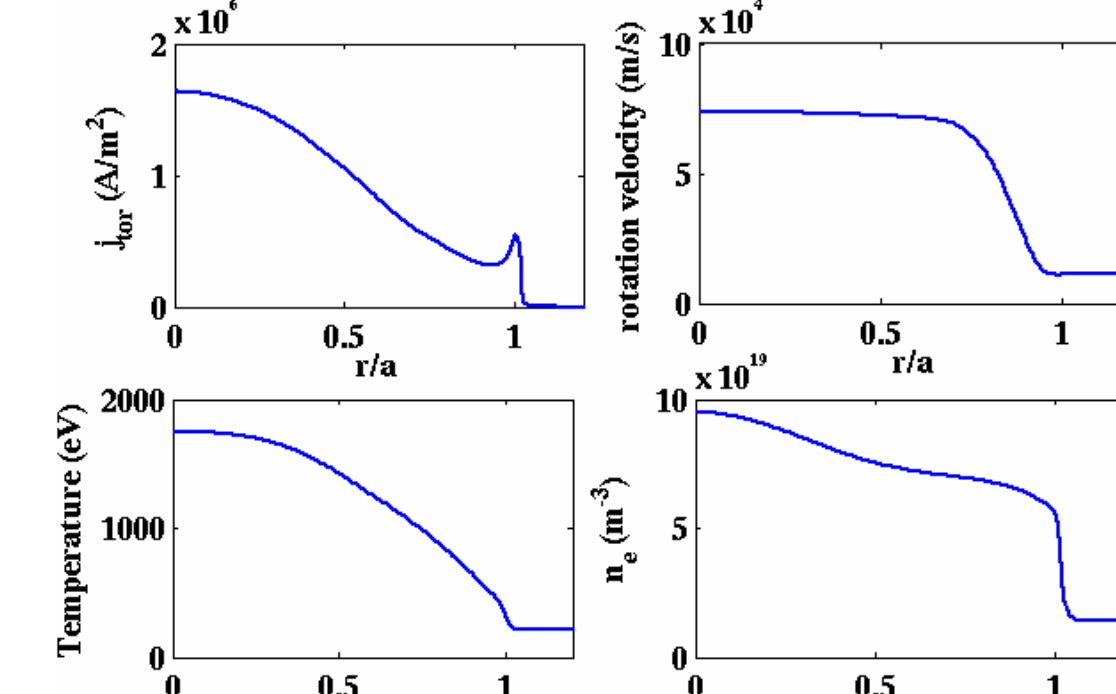


↑ Odd parity I-coil fields

- Mostly non-resonant fields produce small islands, little stochasticity

- Higher S case most closely resembles the 119690 experiment (still one order of magnitude too resistive).

Experimental Profiles

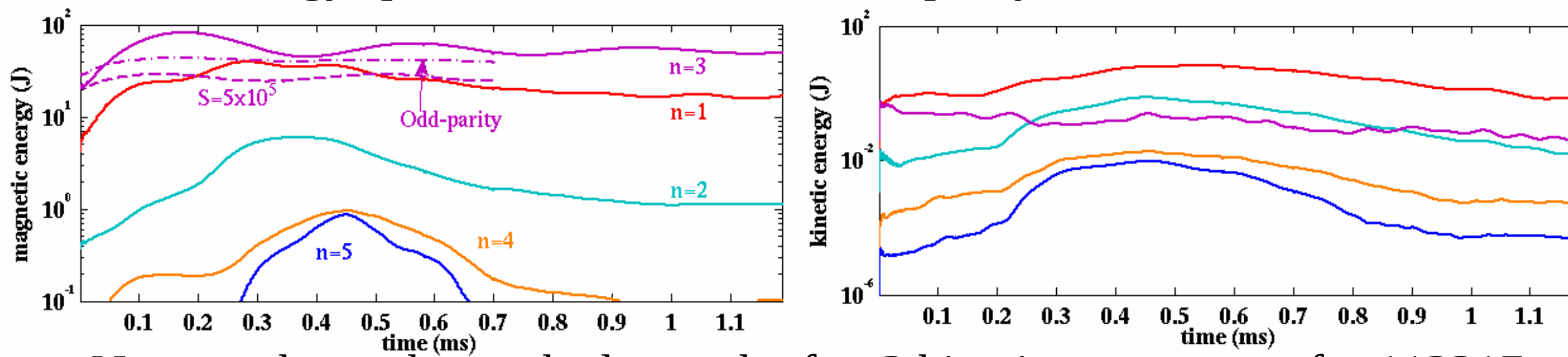


↑ Even parity I-coil fields

- Considerable vacuum field edge stochasticity is well screened by the plasma even at low S

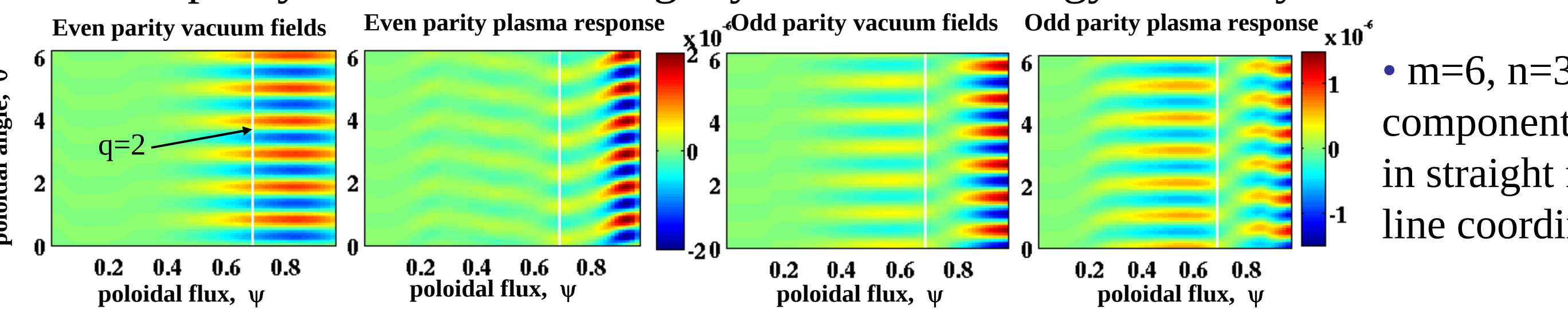
- In all cases, $m=1/n=1$ island is well screened by the plasma rotation

Energy spectrum evolution for even parity, $S=5 \times 10^6$ simulation



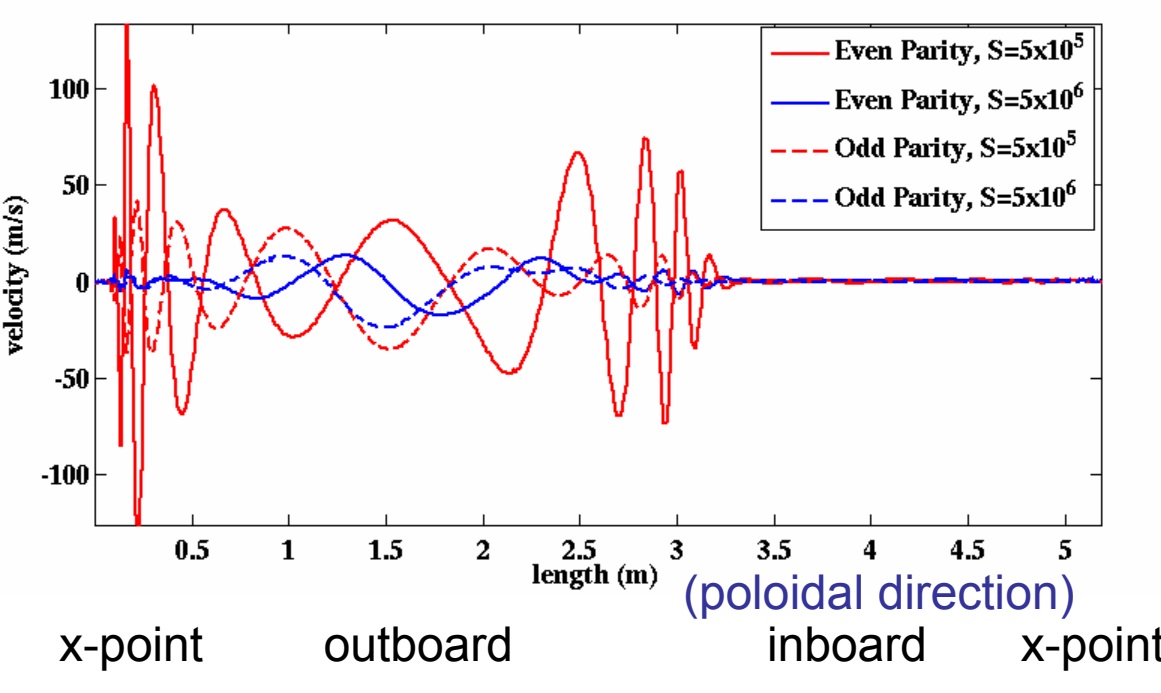
- No case shows the marked growth of n=3 kinetic energy seen for 113317

- Odd parity or lower S has slightly lower n=3 energy in steady state



- Not only amplitude but structure of resonant fields differs greatly between odd and even parity I-coil fields

- In odd parity case, opposite sign vacuum fields inside the resonant surface are amplified by the plasma response, while the fields outside the resonant surface are screened

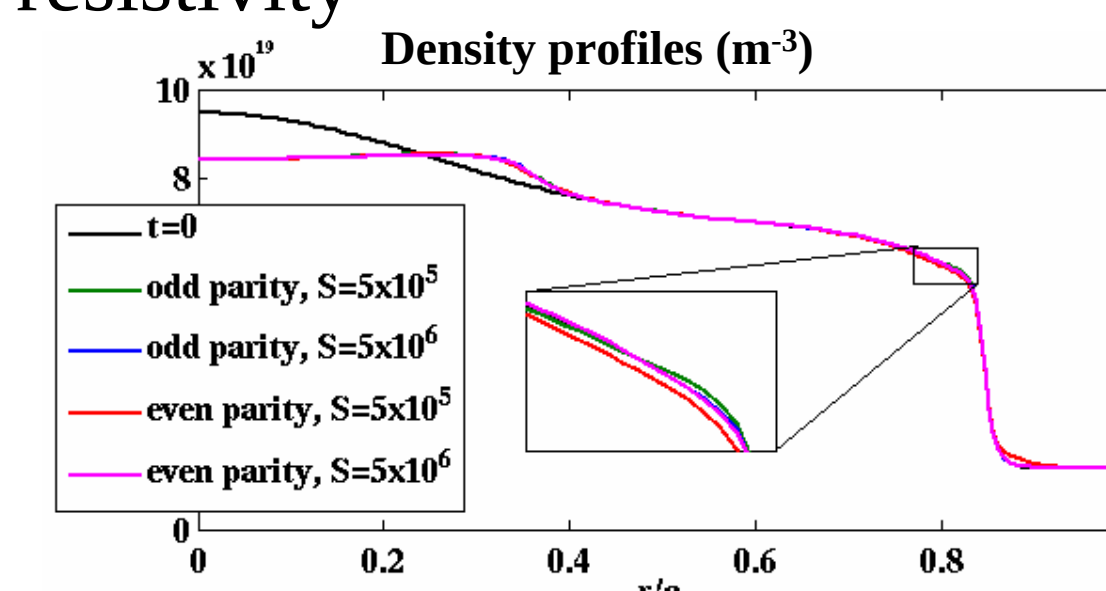


- Slight change in pedestal density is observed only in lowest S, even I-coil simulation

- Transport time scales are longer than simulation times— profiles are evolving

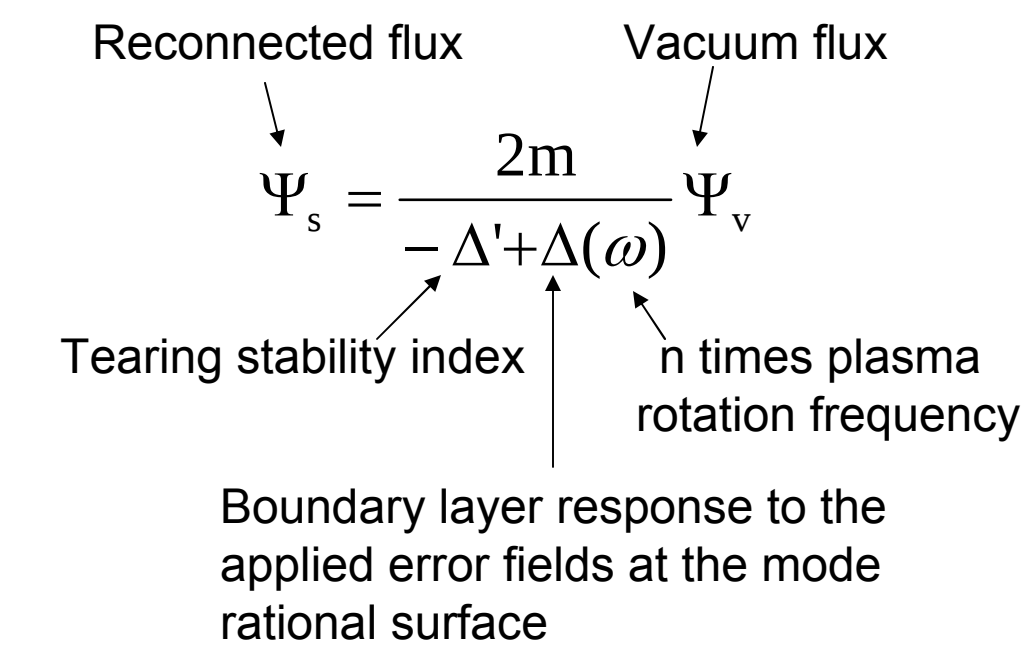
- Lower resistivity significantly reduces velocity at high and low-z locations, slightly reduces it at the outboard side

- Less resonant odd parity fields have lower convection velocity, but identical trend with decreasing resistivity



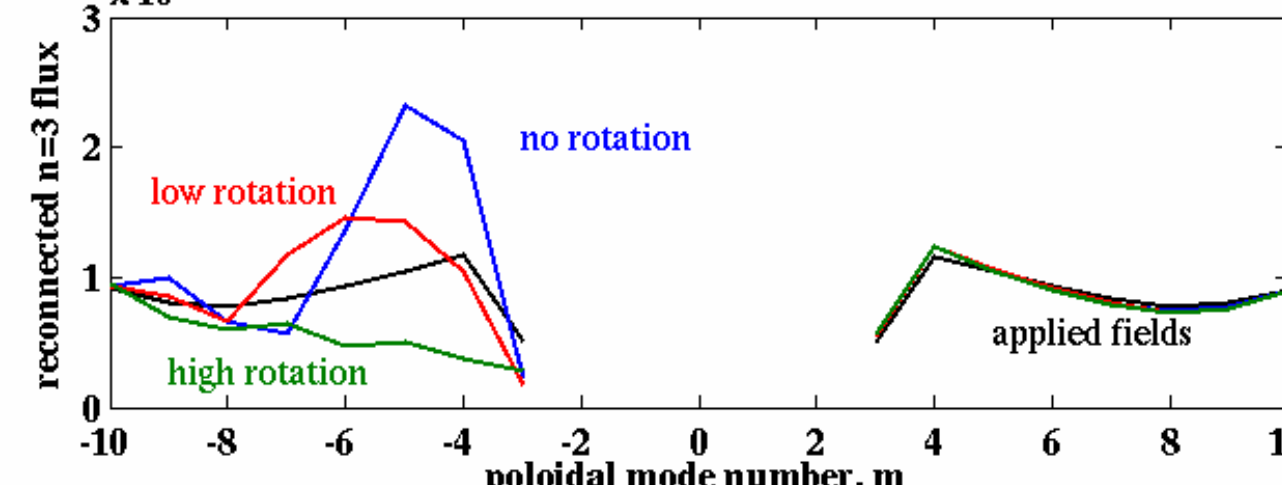
Comparison of simulations with cylindrical error field theory

- Theory does not include any toroidal effects such as mode coupling, hence the comparison can give some insight into the importance of realistic geometry



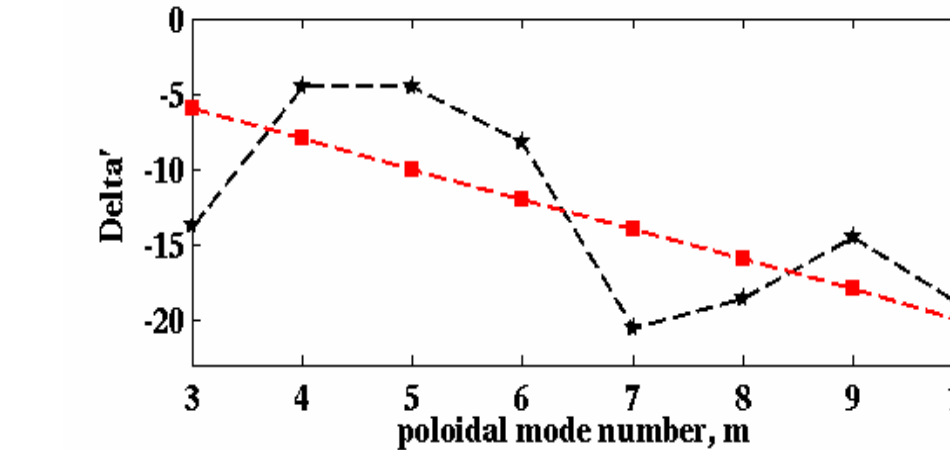
→ Without rotation, error fields are amplified in tearing stable plasmas with $-\Delta' < 2m$

Ψ_s and Ψ_v for 113317



- Phases of the resonant components evaluated at the resonant surfaces. A strongly screened mode in the visco-resistive regime has a phase shift of $\pi/2$

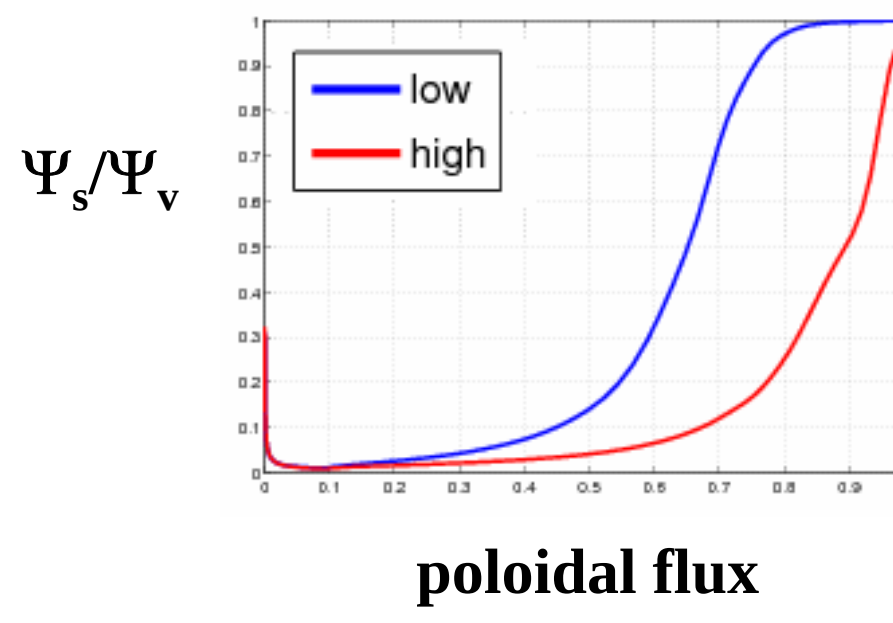
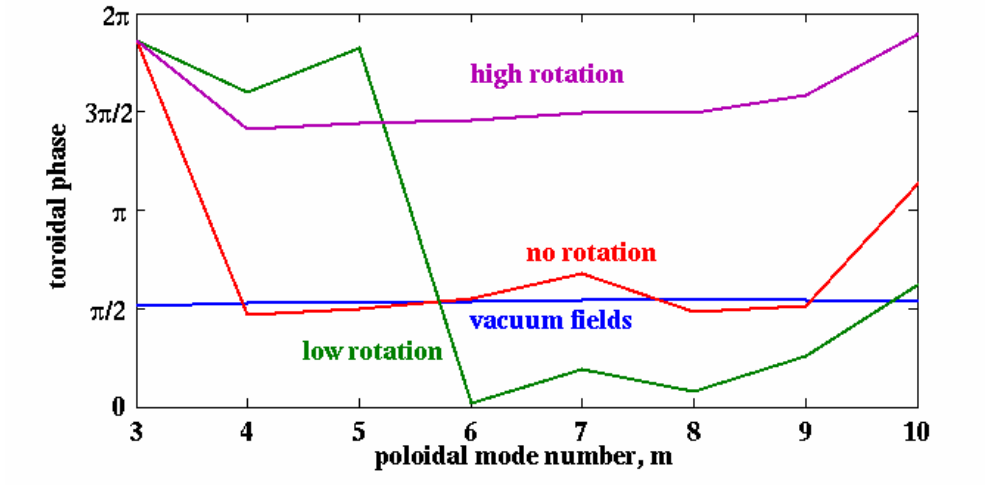
- Theoretically calculated screening factors calculated for the low and high rotation simulations assuming $-\Delta'=2m$ (no amplification)



- Mode amplification is an ideal response dependent only on the tearing stability parameter, Δ'

- Rotational screening is a resistive effect dependent on rotation velocity, resistivity, viscosity and other factors

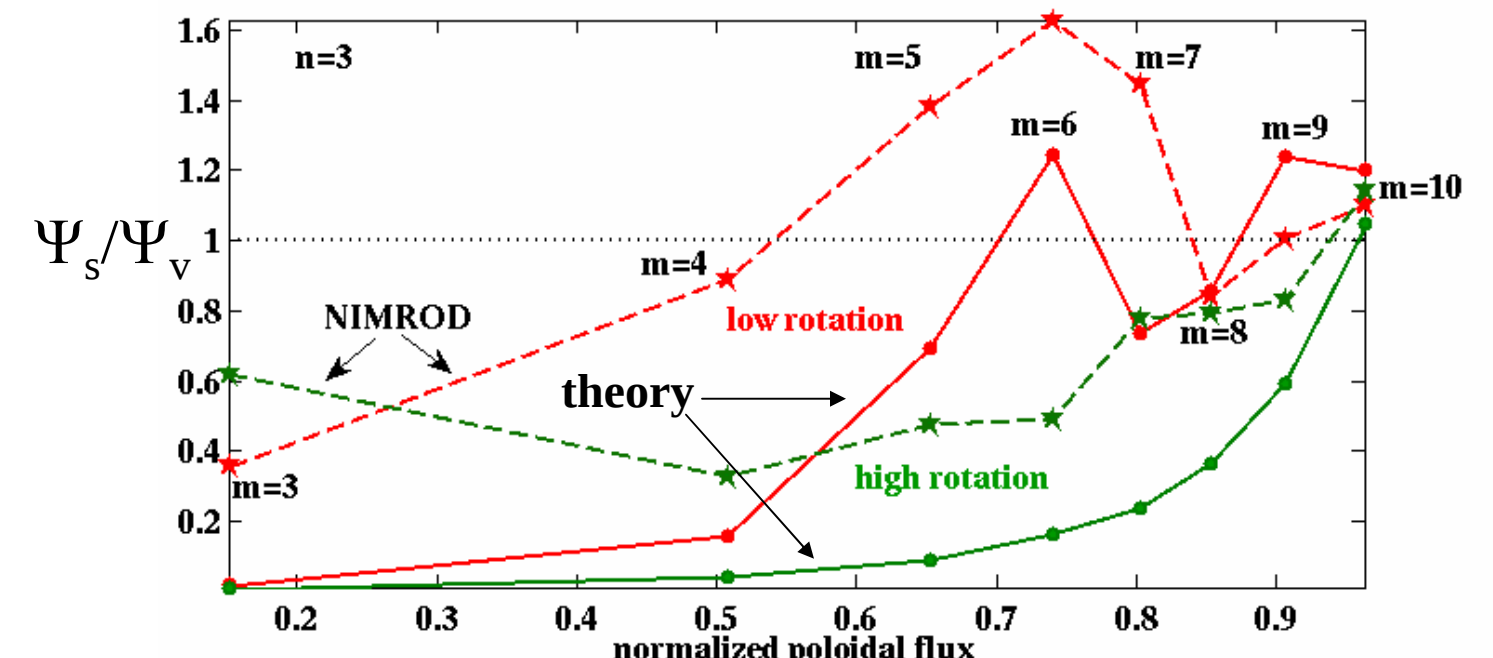
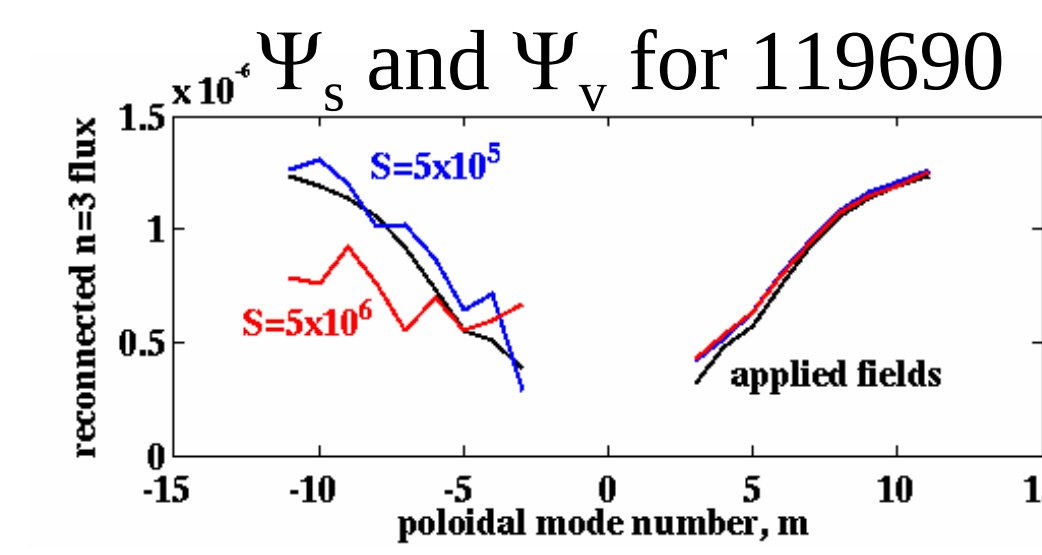
- Positive m are non-resonant modes (unaffected by plasma response), resonant negative m components are amplified, screened



- The non-rotating simulation is used to determine the amplification alone, which depends only on Δ'

- The estimated values of Δ' and the theoretically calculated screening factors are combined to predict the mode amplitudes in the rotating simulations

- Predicted amplitudes are consistently lower than NIMROD results

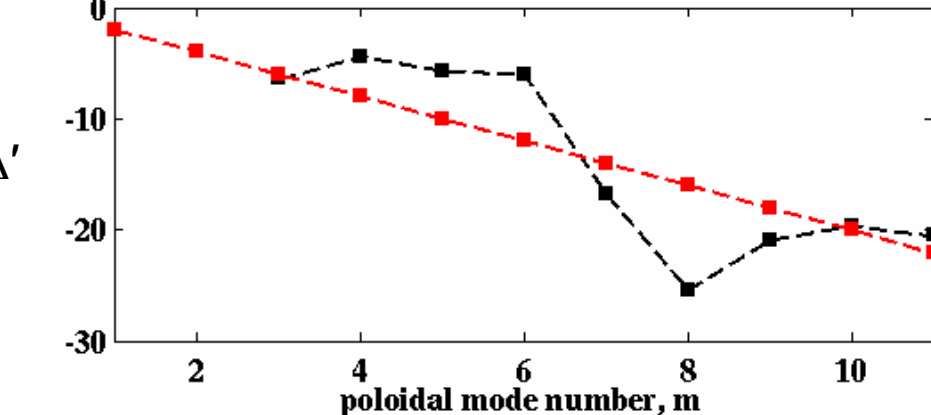


- Two odd-parity simulations

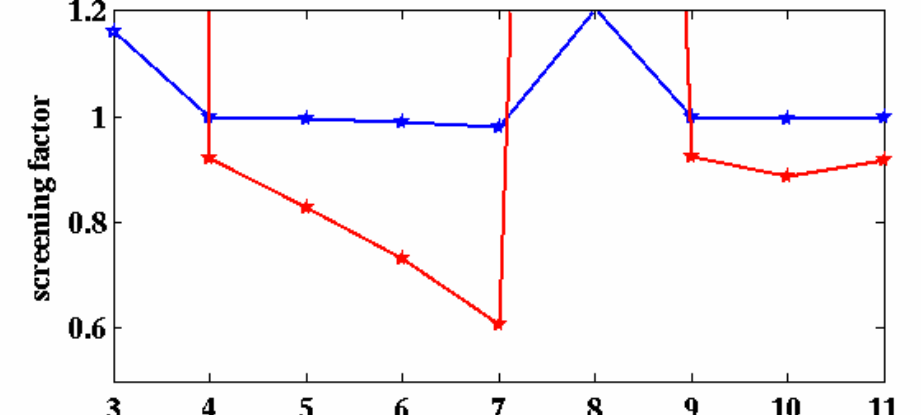
- Amplification is independent of S

- Screening factor $\Delta(\omega)$ is $\propto \eta^{-5/6}$ in the visco-resistive regime

- Estimates for Δ' and $\Delta(\omega)$ can be backed out from results of two simulations



- Screening results are inconsistent with theoretical dependence on resistivity and rotation



Conclusions- NIMROD simulations of DIII-D plasmas with RMP fields show:

- Increased screening of vacuum fields with increased edge rotation

- ExB convection near the separatrix can pump particles into the vacuum region

- A reduction of ExB convection cell velocity with: increased edge rotation, decreased resistivity, decreased resonant mode amplitude (odd parity I-coils)

- Compared with cylindrical error field theory, more realistic model tends to predict larger resonant mode amplitudes and shows qualitative anomalies likely linked to toroidal mode coupling effects