HIGHLY LOCALIZED ELECTRON CYCLOTRON HEATING AND CURRENT DRIVE IN DIII-D

by

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- ECCD (electron cyclotron current drive) needed for
 - Current sustainment and profile control in tokamaks but high performance discharges have ELMs, tearing modes...
 - Stabilization of MHD activity in discharges which have MHD
- Need to validate the theory so codes can be predictive under realistic conditions and to take advantage of the unique localization properties of ECCD
 - MHD suppression
 - Transport barrier generation



ECH SYSTEM

- System with six 1 MW-class gyrotrons under construction
- Launchers for two gyrotrons have control of poloidal and toroidal angles (PPPL)



PPPL Launcher





CPI Diamond Window Gyrotron

- Independent scans of
 n_{||} and ρ (magnetic well
 depth) are possible
- Independent control over toroidal and poloidal launch angles facilitates science (independent $n_{||}$ and ρ scans) and applications (high $n_{||}$ gives high I_{CD} , while low $n_{||}$ gives high j_{CD})



OUTLINE OF RESULTS

- Analysis of ECCD experiments shows that driven current is highly localized, in agreement with theory
 - Behavior of j_{ECCD} with magnetic trapping consistent with theory
 - Successful experimental measurement of off-axis ECCD in ELMing H-mode
- Highly localized nature of ECCD is validated in experiments showing full stabilization of neoclassical tearing modes, under conditions where direct measurement of ECCD is not possible
- Localized nature of ECH may be involved in the observed generation of a strong electron transport barrier in discharges with reversed magnetic shear



ECCD CAN BE MEASURED DIRECTLY FROM MSE SIGNALS

- MSE (motional Stark effect) diagnostic measures magnetic field pitch angles at different major radii, so B_z = B_t tan⁻¹ (pitch angle)
- From Ampere's law $j_{\phi} \cong -\frac{1}{\mu_0} \frac{\partial B_z}{\partial R}$



so the local change in j_{φ} due to ECCD is proportional to the change in $\Delta B_z/\Delta R$, where ΔB_z is the difference in B_z between adjacent MSE channels and ΔR is the spatial separation

- The measured $\partial B_z / \partial R$ are compared to simulations to include the effects of small changes in bootstrap, NBCD, and Ohmic currents
- Total driven current is determined from a best statistical fit to the data, varying the location, width, and magnitude of the driven current in the simulation



MSE MEASUREMENTS SHOW THAT THE INCREASE IN CURRENT DENSITY FROM ECCD IS AS LOCALIZED AS RAY TRACING CALCULATIONS PREDICT



ELECTRON CYCLOTRON CURRENT DRIVE PROVIDES LOCALIZED CURRENT WITH GOOD CONTROL







ECCD EFFICIENCY DECREASES WITH RADIUS (FOR POLOIDAL ANGLE \approx 90 deg) AS EXPECTED FROM THEORY DUE TO TRAPPING EFFECTS



- Excellent agreement with theory except at largest radius
 - Need to test large ρ at higher ECH power with smaller error bars





POLOIDAL SCANS SHOW SYSTEMATIC INCREASE IN ECCD EFFICIENCY TO HIGH FIELD SIDE, IN GOOD AGREEMENT WITH THEORY



 Theoretically the increase in ECCD efficiency with poloidal angle (i.e., magnetic well depth), is due to (a) reduced trapping effects and (b) wave absorption on higher energy electrons from n_{II} upshift



LOCALIZED CHANGE IN CURRENT PROFILE DURING ECCD IS CLEARLY OBSERVED IN ELMING H-MODE PLASMAS



CO-ECCD RADIALLY LOCALIZED AT ISLAND CAN REPLACE THE "MISSING" BOOTSTRAP CURRENT AND COMPLETELY STABILIZE THE NEOCLASSICAL TEARING MODE

$$\frac{\tau_{\text{R}}}{r} \frac{dw}{dt} = \Delta \dot{r} + \epsilon^{1/2} \left(\frac{L_{\text{q}}}{L_{\text{p}}} \right) \beta_{\theta} \left[\frac{r}{w} - \frac{rw_{\text{pol}}^2}{w^3} - \frac{8qr\delta_{\text{ec}}}{\pi^2 w^2} \left(\frac{\eta j_{\text{ec}}}{j_{\text{bs}}} \right) \right]$$





HIGH CURRENT DENSITY OBTAINED THROUGH OPTIMIZATION OF ECCD



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FULL STABILIZATION OF NTM OBTAINED WITH MODEST ECH POWER



THE LOCATION OF ECCD IS CRITICAL TO FULL STABILIZATION



- Toroidal field was ramped down to scan ECCD past the island
- Alignment within 2 cm is required
- j_{ECCD} > j_{BS} is satisfied
- Sensitivity of effect to location implies that the width of the ECCD is less than the island size, in agreement with ray tracing calculation
- These results show that modeling is accurate even in ELMing H–mode with sawteeth and a tearing mode, at large ρ



Results similar to those obtained on ASDEX-U and JT-60U

ECH IN DISCHARGES WITH NEGATIVE MAGNETIC SHEAR LEADS TO FORMATION OF AN ELECTRON TRANSPORT BARRIER



- Electron transport barrier forms immediately upon application of ECH power of 0.5 MW
- Barrier location lies just outside heating location
- χ_e more than an order of magnitude smaller than χ_i in the barrier
- Stabilization of ETG mode by Shafranov shift believed responsible for decrease in χ_e
- Barrier found with co-ECCD, counter-ECCD, and radial ECH; also with no NBI

ECH EARLY IN DISCHARGE PRODUCES HIGH T_e AND LARGE PRESSURE GRADIENT WITHOUT NEUTRAL INJECTION

- 0.8 MW ECH applied at ρ ~0.4; no NBI
- Co–ECCD in this case; radial ECH also works
- Measurement of T_e ~15 keV by ECE roughly supported by Thomson scattering and pulse height analysis





CONCLUSIONS

- Modeling and experiment have substantially come together for ECH/ECCD
 - Best tested in quiescent L-mode
 - Also tested under realistic conditions of ELMs, sawteeth, and other turbulence
- The narrow current drive profile of ECCD is useful for stabilizing neoclassical tearing modes
 - Full suppression obtained, with increase in pressure and neutron rate
 - Effect very sensitive to location of current drive
 - Success indicates indirectly that the narrowness of the ECCD profile and its magnitude are close to those calculated by ray tracing
- Localized ECH generates an electron transport barrier in the vicinity of the power deposition
- These results strongly support ECCD for detailed control of the current profile needed to realize advanced tokamak discharges

