## First Direct Evidence of Main Ion Flow Triggering the L-H Transition<sup>\*</sup>

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Simultaneous measurements of main ion flow (via main ion CER),  $E \times B$  flow, and turbulence level  $\tilde{n}/n$  (via Doppler backscattering) during transitions characterized by extended limit cycle oscillations (LCO [1]), show for the first time that the initial (transient) turbulence collapse [Fig. 1(a)] preceding the L-H transition is caused by turbulence-generated main ion flow and  $E \times B$  opposing the equilibrium (L-mode) edge plasma  $E \times B$  flow related to the edge ion pressure gradient. The formation dynamics of edge transport barriers is crucial for understanding the physics basis of the empirical L-H transition power threshold scaling, and for confidently extrapolating auxiliary heating requirements to burning plasmas. Figure 1(b) shows that the  $v_i \times B/B$  contribution to the  $E \times B$  velocity peaks as fluctuations are first suppressed. Fig. 1(c) shows that the  $E \times B$  shearing rate  $\omega_{E \times B}$  reverses at this time. The correlations between turbulence envelope, main ion flow, and pressure-gradient driven flow, and their detailed spatio-temporal evolution have been measured. The main ion poloidal velocity lags ñ early in the LCO, consistent with turbulence-driven poloidal ion flow [Fig. 1(d)]. As the LCO evolves, the periodic reduction in edge turbulence and edge transport enables a gradual increase (and periodic modulation) of the edge pressure gradient and ion diamagnetic flow. During the final phase of the LCO the pressure gradient diamagnetic flow) dominates the mean flow  $E \times B$  shearing rate, which becomes sufficiently large to sustain fluctuation suppression and secure the LCO-H-mode transition. A twopredator, one-prey model, similar to a previously developed model [2] but retaining opposite polarity of the turbulencedriven and pressure-gradient-driven  $E \times B$  flow, captures essential aspects of the transition dynamics, and is consistent with the direction of the  $(\tilde{n}, E_r)$  limit cycle observed in DIII-D and recently in JFT-2M. The scaling of the L-LCO transition threshold power and LCO frequency with edge plasma density, collisionality, and  $q_{95}$  will be presented.

[1] L. Schmitz, L. Zeng, et al., Phys. Rev. Lett. 108, 155002 (2012). [2] K. Miki and P.H. Diamond, Phys. Plasmas 19, 092306 (2012).

t٥ (a) 0.8 (au) پي ۵.4 153624 R=Rs-1.4 cm 0 (b) 4 (v<sub>i</sub>xB)/B (km/s) 2 0 8 R=Rs-1.1 cm (c) rad/s) R=R<sub>e</sub>-1.8 cm ω<sub>ExB</sub> (10<sup>5</sup> r, 2433 2434 2435 2436 Time (ms) 0.4  $C(\tilde{n}, v_{\theta}), C(v_{E\times B}, v_{\theta})$ 154287 ′<sub>ExB</sub>,,v⊕ ⊦1.5 ms  $C(\tilde{n},v_{\theta})$  to (2569 ms  $C(\tilde{n},v_{\theta}) t_0+1.5 ms$ (d) -0.4 0.4 -0.6 -0.4 -0.2 Ò 0.2 0.6 ∆t (ms)

Fig. 1. Time evolution of (a) relative density fluctuation level; (b) main ion **v**<sub>i</sub>×**B E**×**B** contribution to the velocity; (c) shearing rate  $\omega_{ExB}$ at two radii inside the LCFS. Time of initial turbulence suppression indicated: (d)Cross-correlation of  $\tilde{n}$ ,  $v_{ExB}$  with  $v_{i\theta}$ . R<sub>s</sub> is the LCFS radius.

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