

# Calculation of Linear Two-Fluid Plasma Response to Applied Non-Axisymmetric Fields

by

N.M. Ferraro

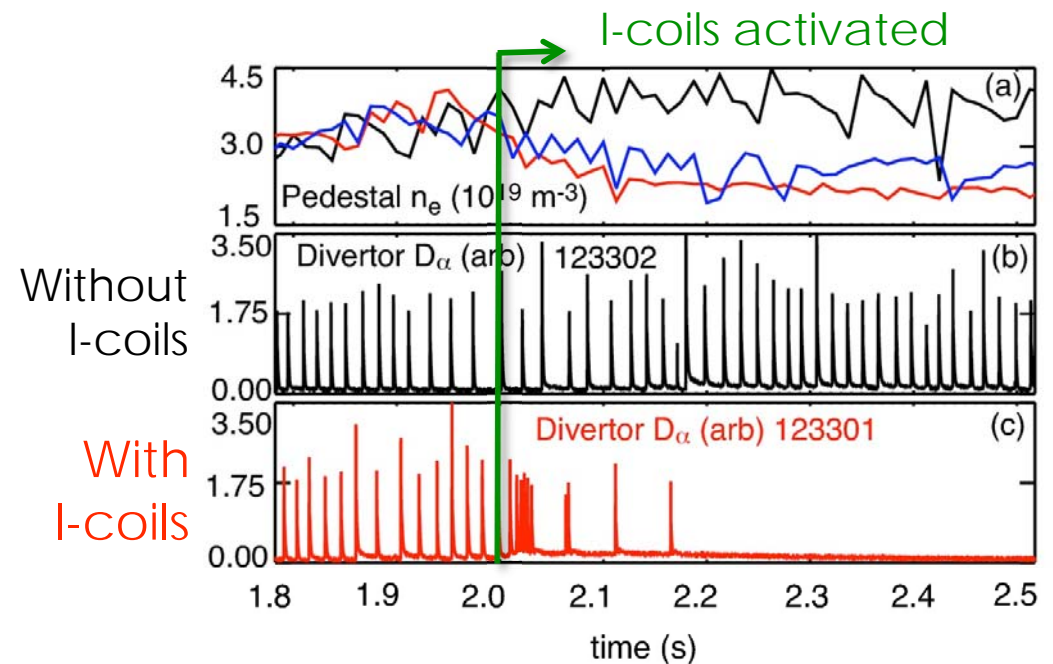
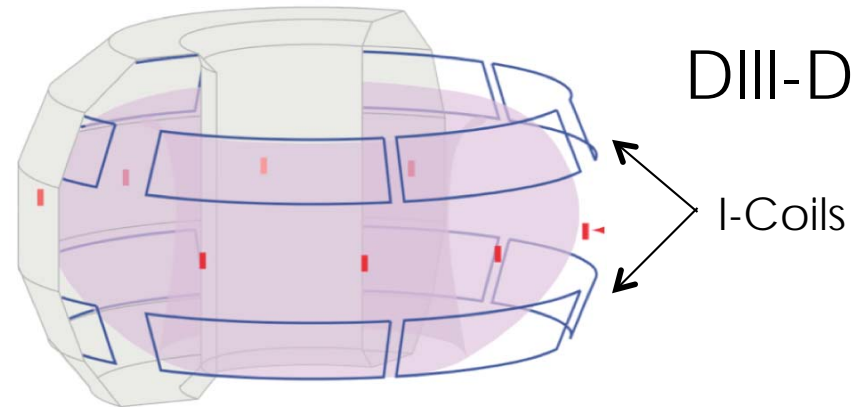
General Atomics

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# 3D Fields Significantly Affect Tokamak Performance

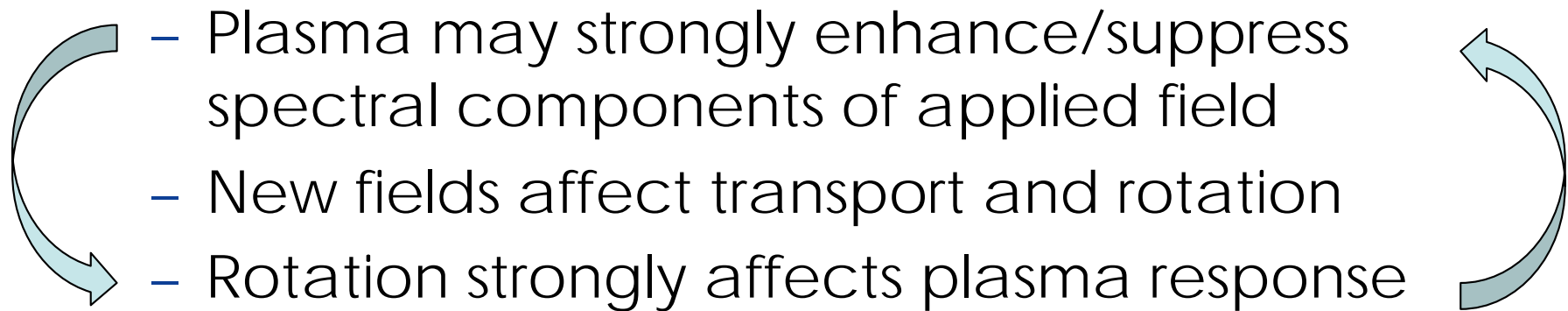
- **Edge-localized modes**
  - Mitigate/suppress ELMs in H-mode
- **Transport**
  - Density pump-out
- **Drive/brake rotation**
  - Affects RWM stability
  - Affects tearing mode stability (Buttery talk this session)
  - Allow access QH-mode without NBI (Burrell talk Friday)



Evans, *et al. Phys. Plasmas* **13** (2006)

# A Predictive Capability Requires Understanding Plasma Response

- **Predictive capability is challenging because plasma response is complicated**



- **New tools are being developed and applied to gain predictive understanding (M3D-C1)**

# Our Models of Plasma Response Are Evolving

- **“Vacuum” Fields (TRIP3D)**
  - Plasma does not respond to applied fields
  - Tells us degree to which applied fields are “resonant”
  - Doesn’t tell us dependence on plasma parameters
- **Ideal (IPEC, MARS-F, VMEC)**
  - Plasma responds such that magnetic surfaces remain intact → no islands
  - Tells us how strongly ideal modes respond to applied fields
  - Doesn’t explain dependence of plasma response on rotation; doesn’t directly determine island size

# Our Models of Plasma Response Are Evolving

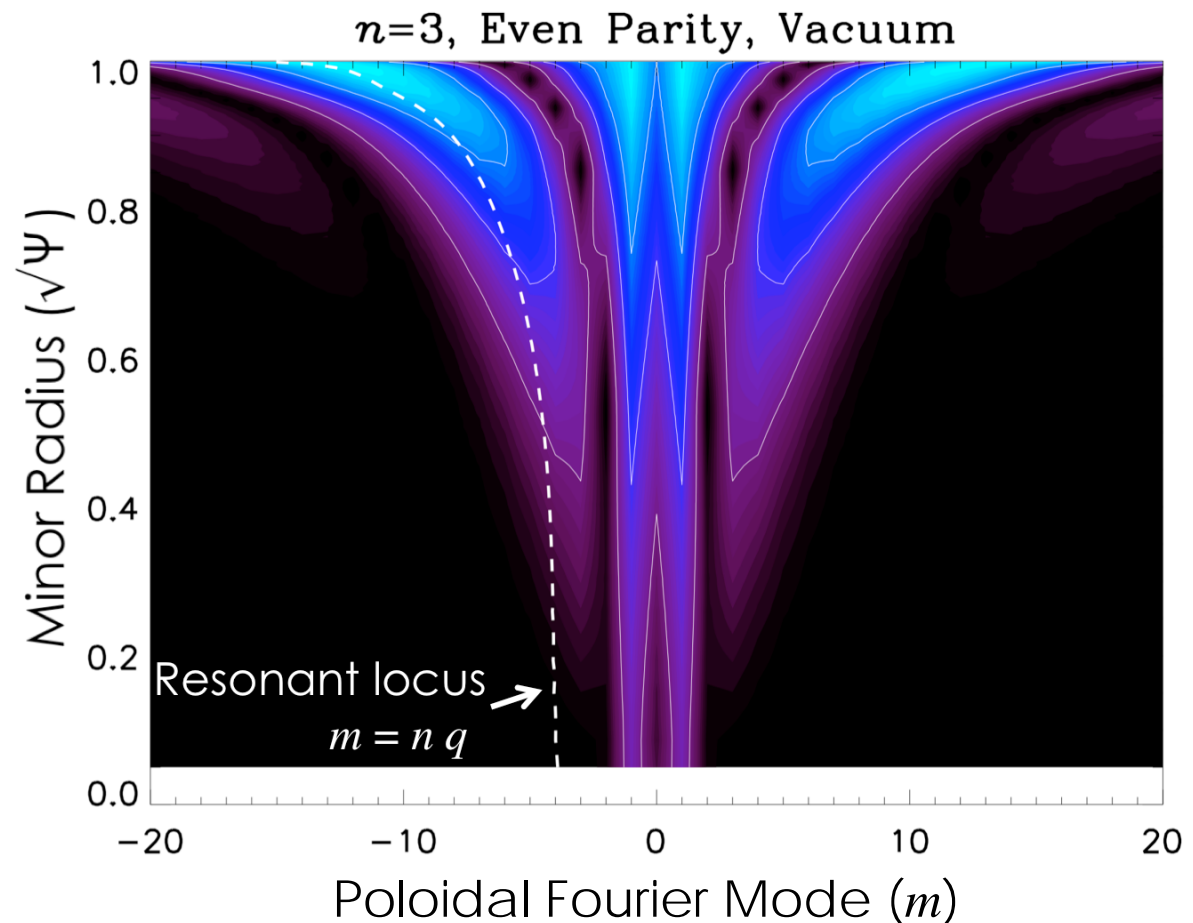
- **Resistive, Single-Fluid (MARS-F, JOREK)**
  - Describes tearing response → islands
  - Tells us how response depends on plasma parameters, especially **rotation**
  - When is response more like ideal? When is it more like vacuum? Does it smoothly transition between the two?
- **Two-fluid & “Extended” MHD (M3D-C1, NIMROD)**
  - Ion rotation ( $\Omega$ ) and electron rotation ( $\Omega^e$ ) are distinct
  - FLR effects, NTV, etc.

# Outline

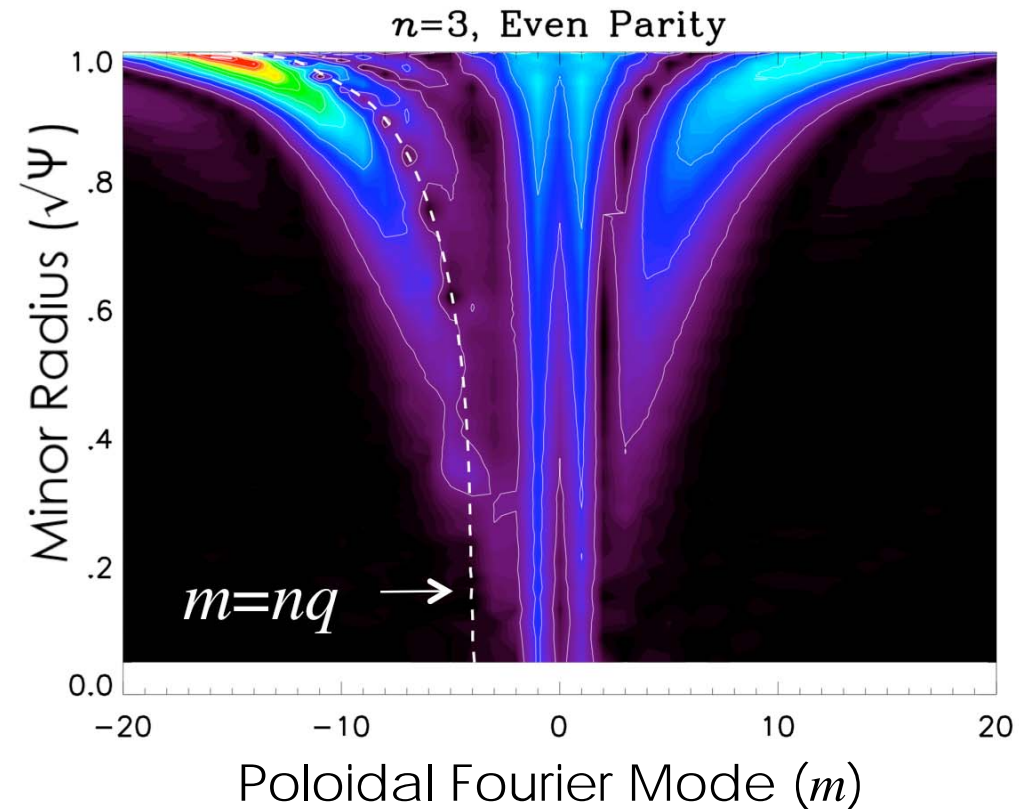
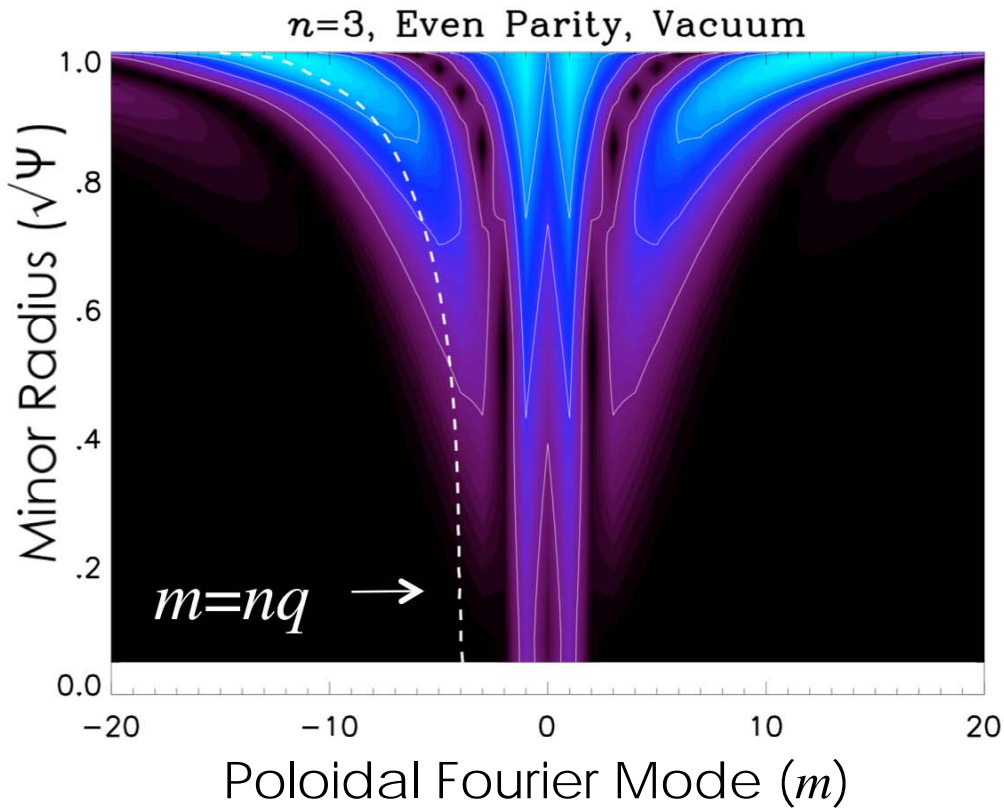
- **Basics of 3D response**
- **Introduction to M3D-C1**
- **Linear single-fluid results**
  - Rotation has strong effect on plasma response
- **Linear two-fluid results**
  - Electron rotation screens core islands
  - Sheared ion rotation affects edge
- **Linear vs. nonlinear**
  - Nonlinear calculations are required for some phenomena

# “Vacuum” Model Tells Us a Lot

- Plot shows Fourier spectrum of  $B_n$
- $B_n$  = component of applied field normal to equilibrium magnetic surfaces
- Resonant components (along dashed line) cause islands
- Non-resonant components cause bending of surfaces
- Poloidal spectrum of  $B_n$  depends on  $\Psi$



# Plasma Response Modifies Spectrum



- Ideal response  $\rightarrow$  no islands  $\rightarrow$  reduction in resonant components
- Excited ideal modes  $\rightarrow$  enhancement of non-resonant components



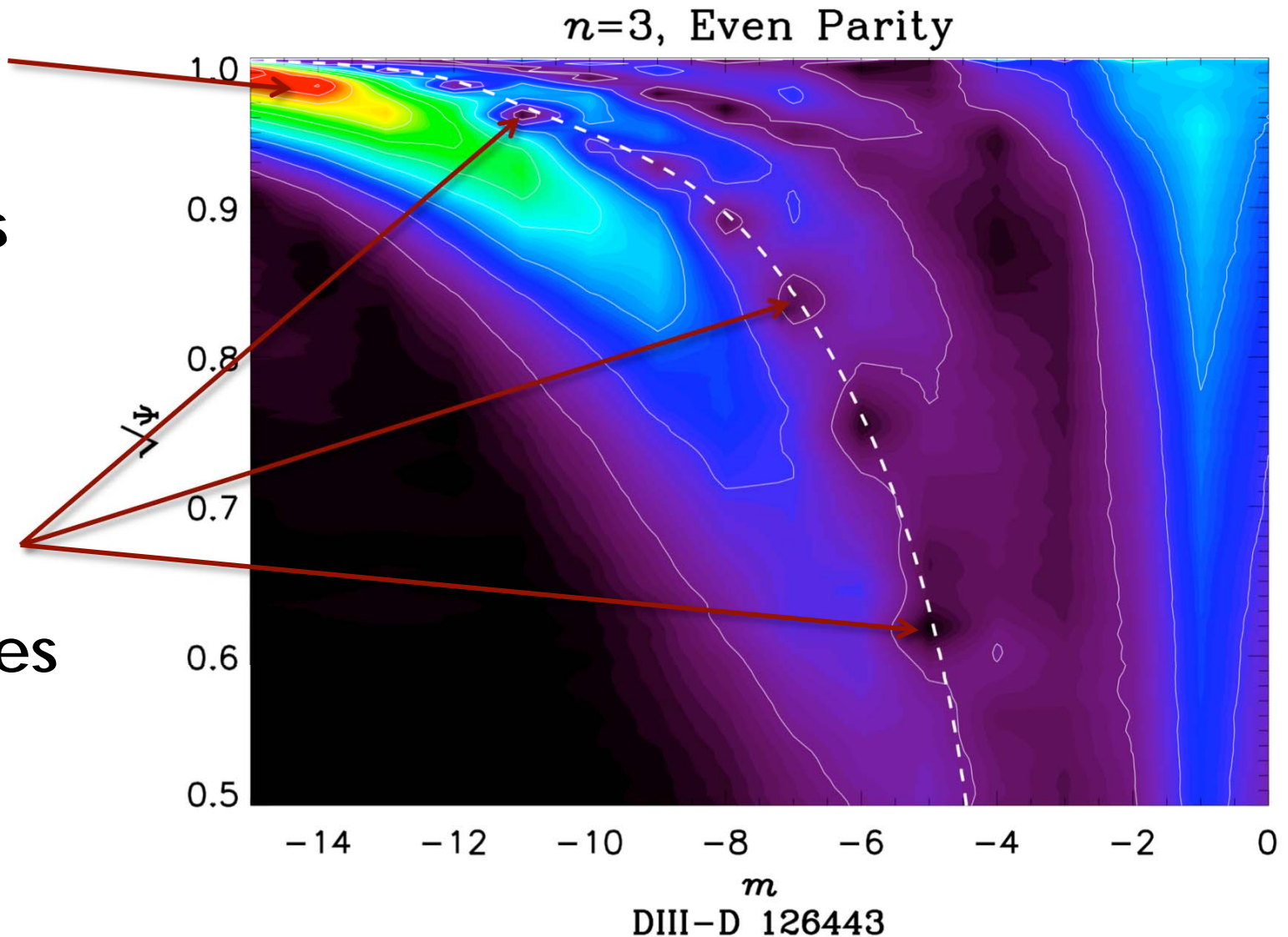
# Plasma Can Kink and Screen

“Kinking”

- Distorts surfaces

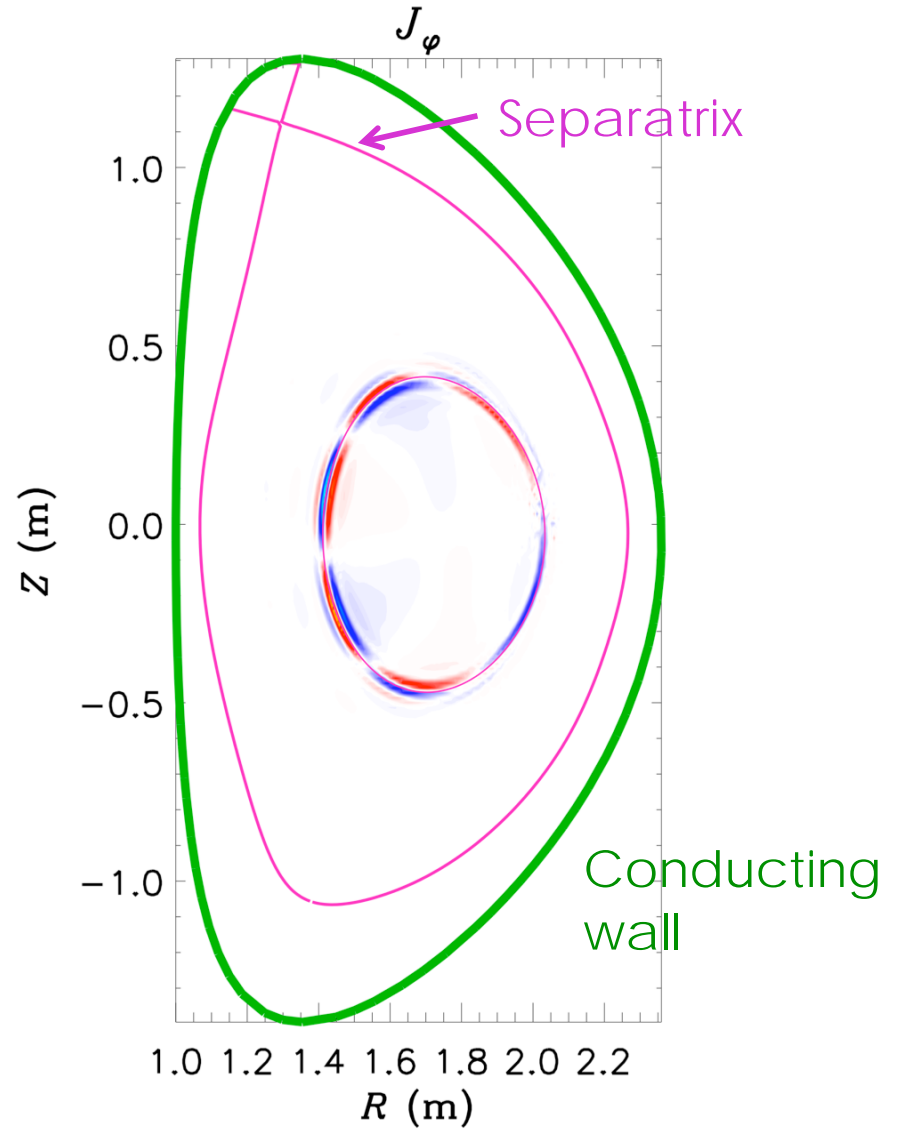
Screening

- Eliminates islands



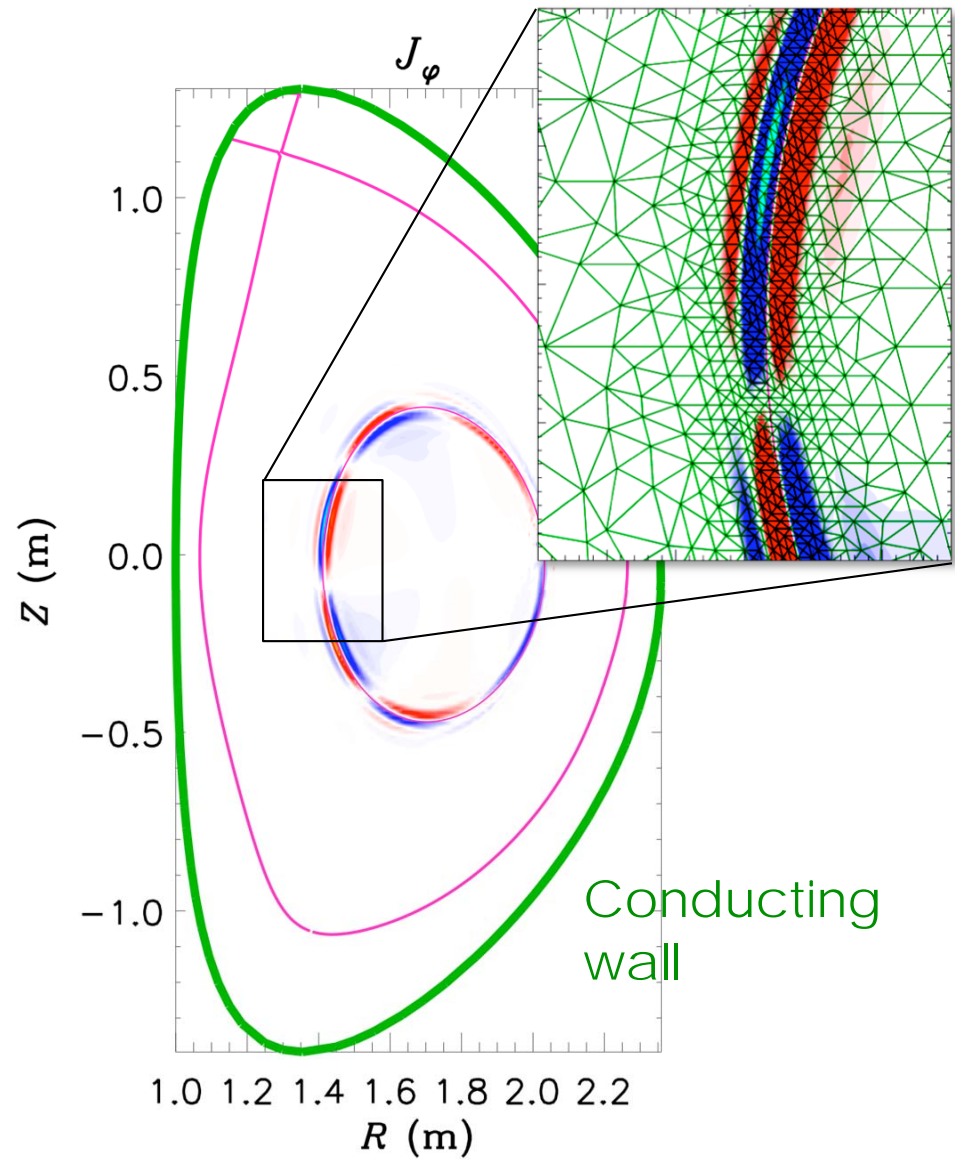
# M3D-C1 Can Calculate Two-Fluid Response

- M3D-C1 is a two-fluid resistive finite element code
  - Shares some design principles with M3D
  - (R,Z) coordinates (not spectral in poloidal angle)
- Computational domain includes plasma, separatrix, and open field-line region



# M3D-C1 Can Calculate Two-Fluid Response

- **M3D-C1 is a two-fluid resistive finite element code**
  - Shares some design principles with M3D
  - (R,Z) coordinates (not spectral in poloidal angle)
- **Computational domain includes plasma, separatrix, and open field-line region**
- **Unstructured mesh allows resolution packing at rational surfaces**
- **Both linear and nonlinear models are implemented**



# Two-Fluid Model Implemented in M3D-C1

$$\frac{\partial n}{\partial t} + \nabla \cdot (n\mathbf{u}) = 0$$

$$n \left( \frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} \right) = \mathbf{J} \times \mathbf{B} - \nabla p - \nabla \cdot \Pi$$

$$\frac{\partial p}{\partial t} + \mathbf{u} \cdot \nabla p = -\Gamma p \nabla \cdot \mathbf{u} - \frac{d_i}{n} \mathbf{J} \cdot \left( \Gamma p_e \frac{\nabla n}{n} - \nabla p_e \right) - (\Gamma - 1) \nabla \cdot \mathbf{q}$$

$$\frac{\partial \mathbf{B}}{\partial t} = -\nabla \times \mathbf{E}$$

$$\mathbf{E} = -\mathbf{u} \times \mathbf{B} + \eta \mathbf{J} + \frac{d_i}{n} (\mathbf{J} \times \mathbf{B} - \nabla p_e)$$

$$\Pi = -\mu \left[ \nabla \mathbf{u} + (\nabla \mathbf{u})^T \right]$$

$$\mathbf{q} = -\kappa \nabla \left( \frac{p}{n} \right) - \kappa_{\parallel} \mathbf{b} \mathbf{b} \cdot \nabla \left( \frac{p_e}{n} \right)$$

$$\mathbf{J} = \nabla \times \mathbf{B}$$

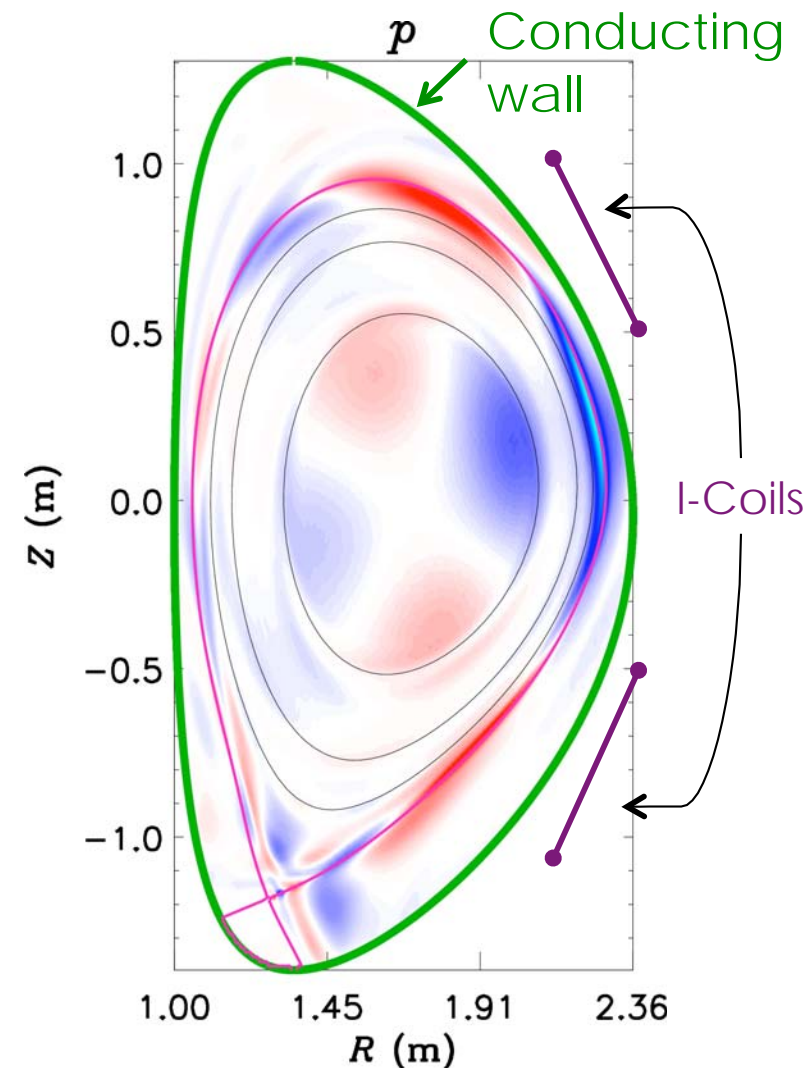
$$\Gamma = 5/3$$

$$p_e = p/2$$

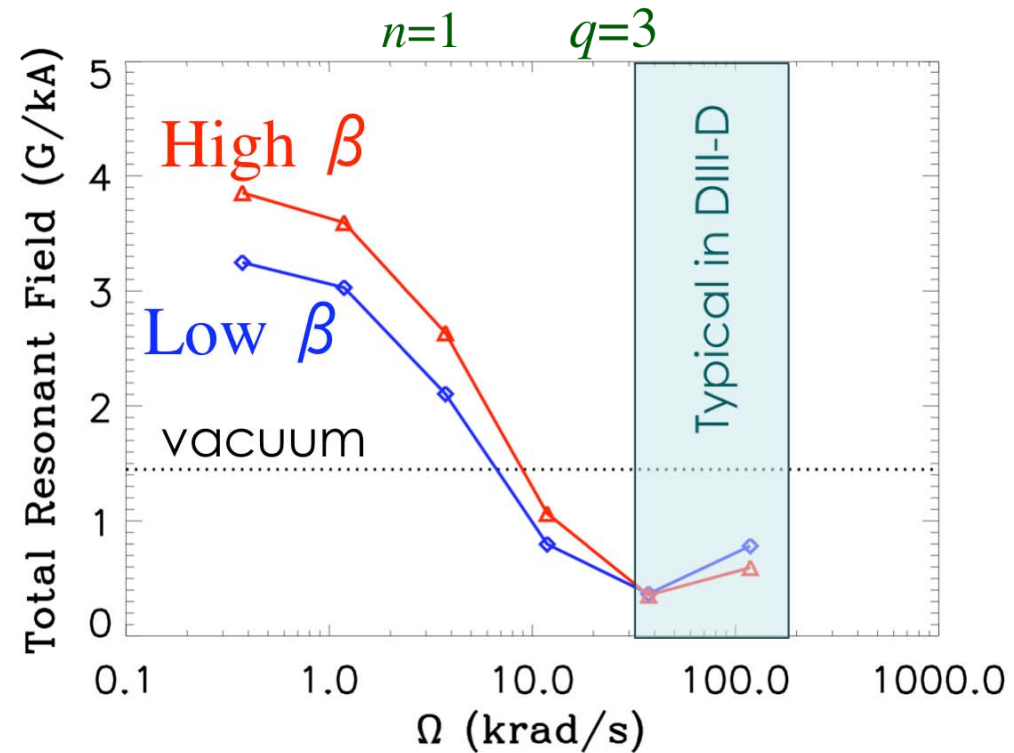
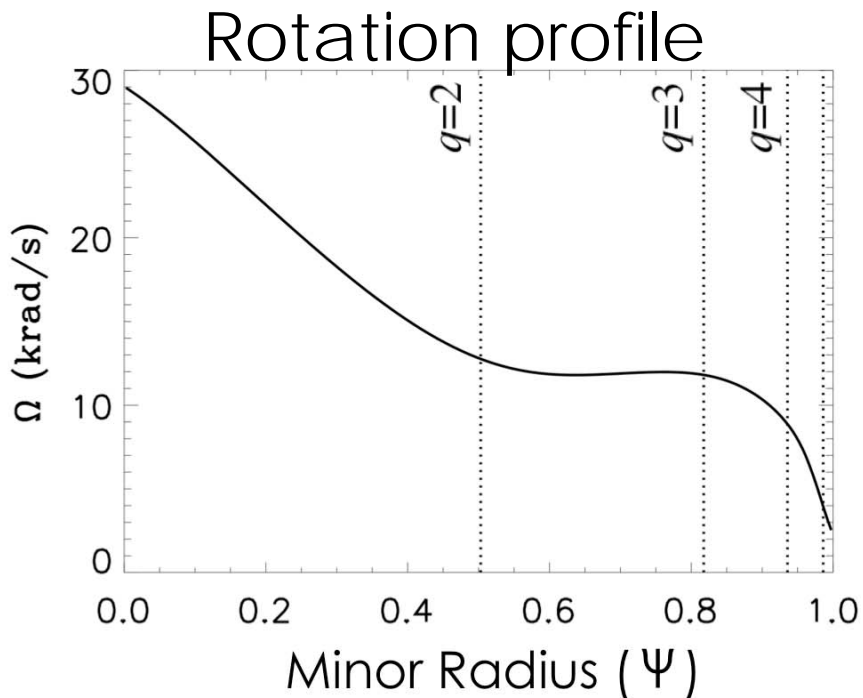
- Complete (not reduced) **two-fluid** model is implemented
- **Time-independent** equations may be solved directly for linear response

# Analysis Considers Reconstructed DIII-D Equilibria

- Vacuum fields generated by DIII-D I-coils
- Boundary conditions
  - Vacuum  $B_n$  is held constant at the boundary
  - No-slip ( $\mathbf{v}=0$ )
- Realistic transport parameters
  - Lundquist number  $\sim 10^9$
- Toroidal rotation
  - Rotation is added self-consistently:  $p \neq p(\psi)$



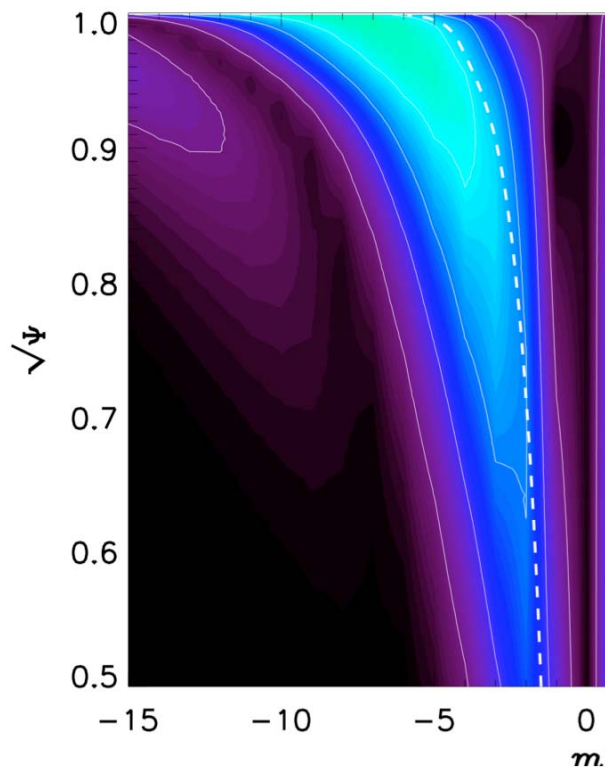
# Single-Fluid Result — Rotation (Usually) Improves Screening



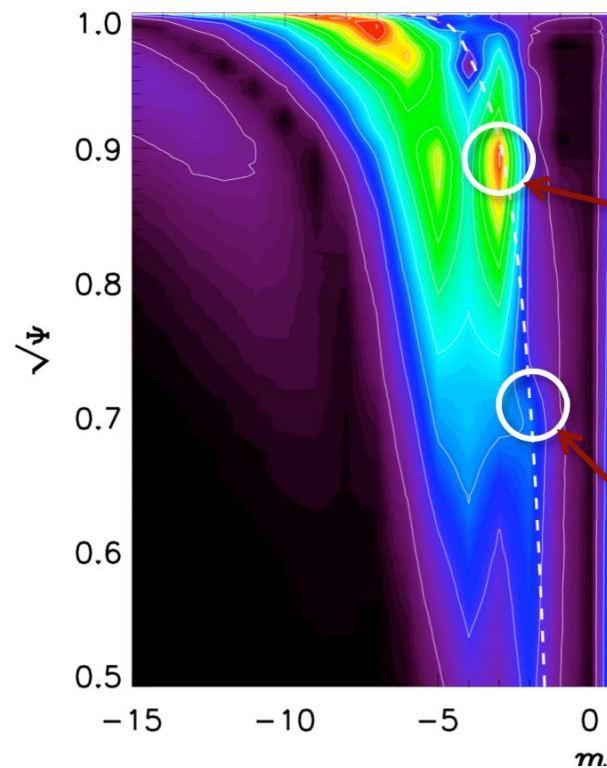
- Plasma may enhance resonant fields at low rotation
- Large rotation screens resonant fields
- Response depends on beta

# Tearing of 3/1 and 2/1 Surfaces is Driven by External Fields, But Suppressed by Rotation

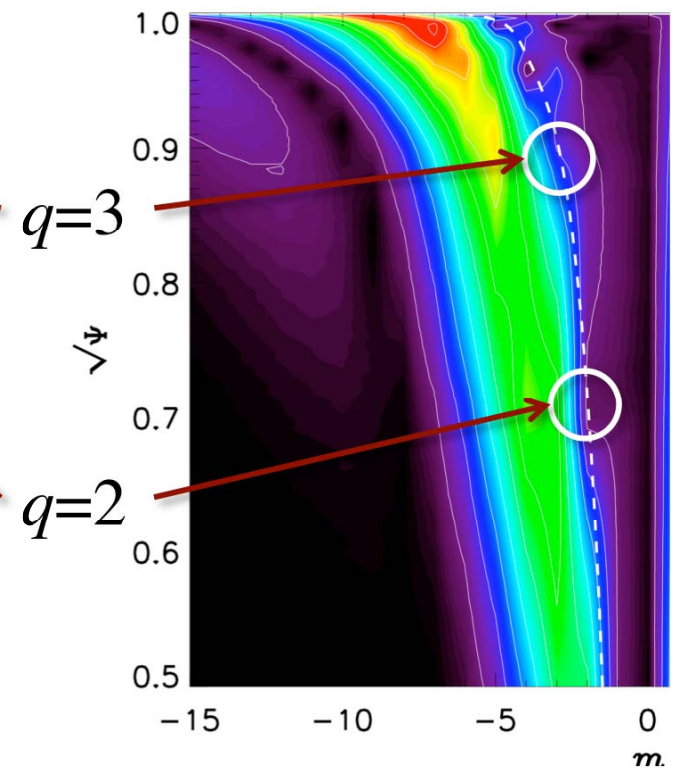
Vacuum



$\Omega = 0$



$\Omega_0 \approx 30$  krad/s



$n=1$

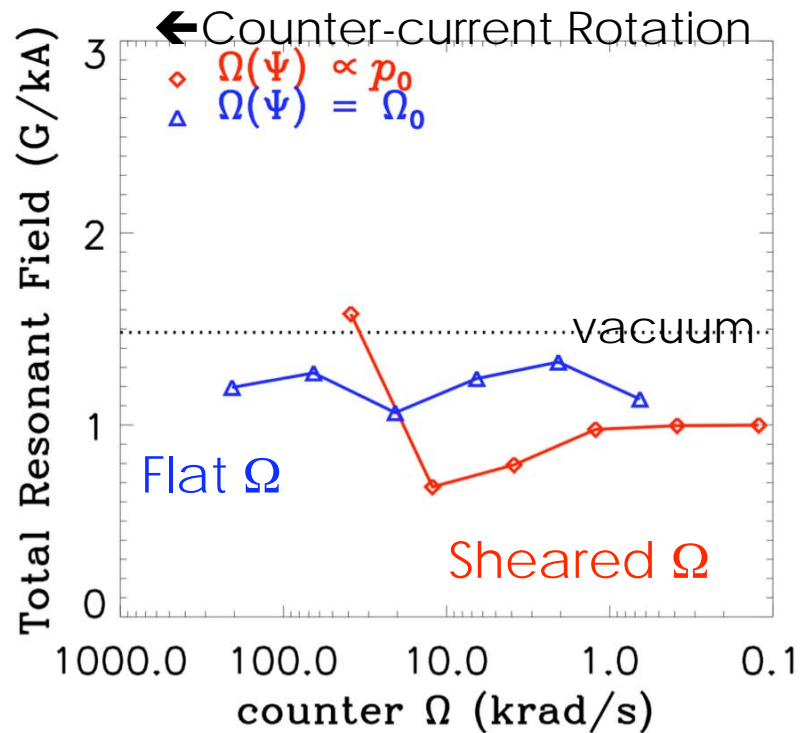
# Why is Plasma Response Sensitive to Rotation? Why is it Sensitive to Beta?

- **From a (rotating) plasma's perspective, the static external fields are oscillating**
  - If field is oscillating faster than tearing response, plasma won't tear
- **Rotation drives static tearing modes away from marginal stability**
- **Higher beta moves modes closer to marginal stability**
  - At marginal stability, an infinitesimal perturbation yields an infinite response

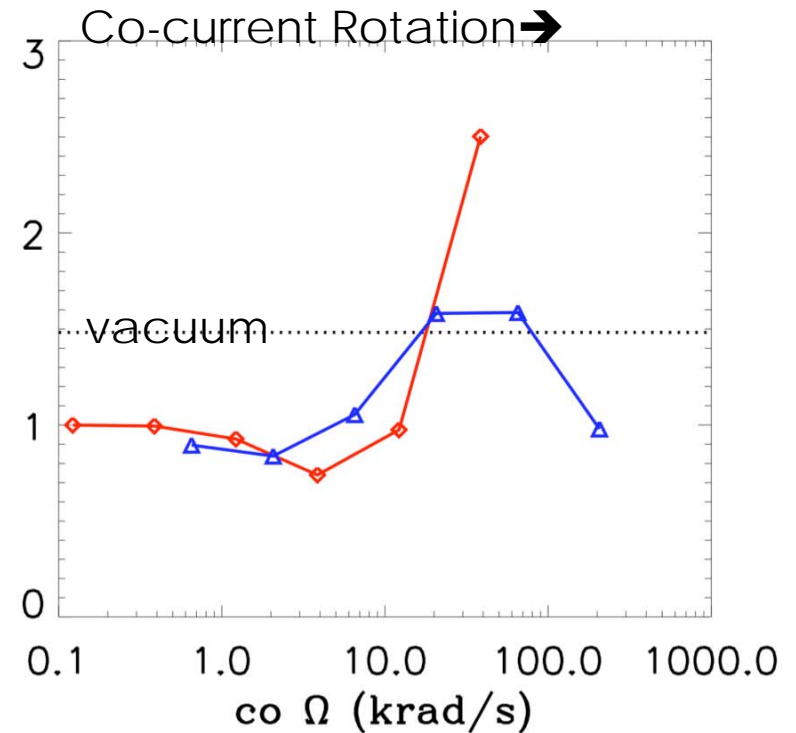


# Single-Fluid Result — Rotation Shear Increases Edge Response

DIII-D 135762



$q=5$

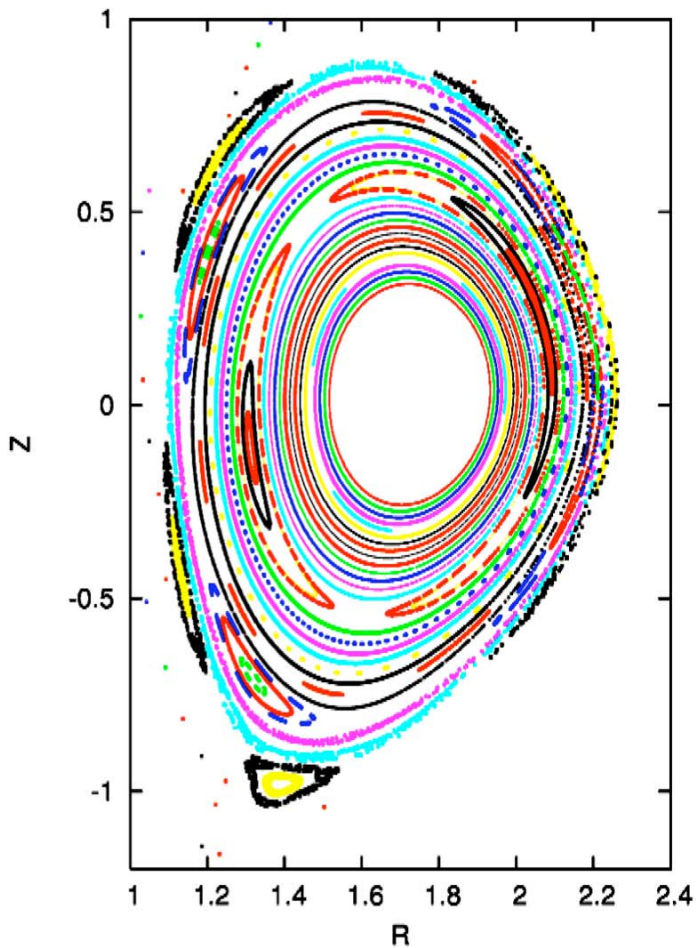


- Large rotation shear seems to increase edge response
- **Why?** Theory predicts  $\Omega'$  is destabilizing to low- $n$  peeling-ballooning modes\*

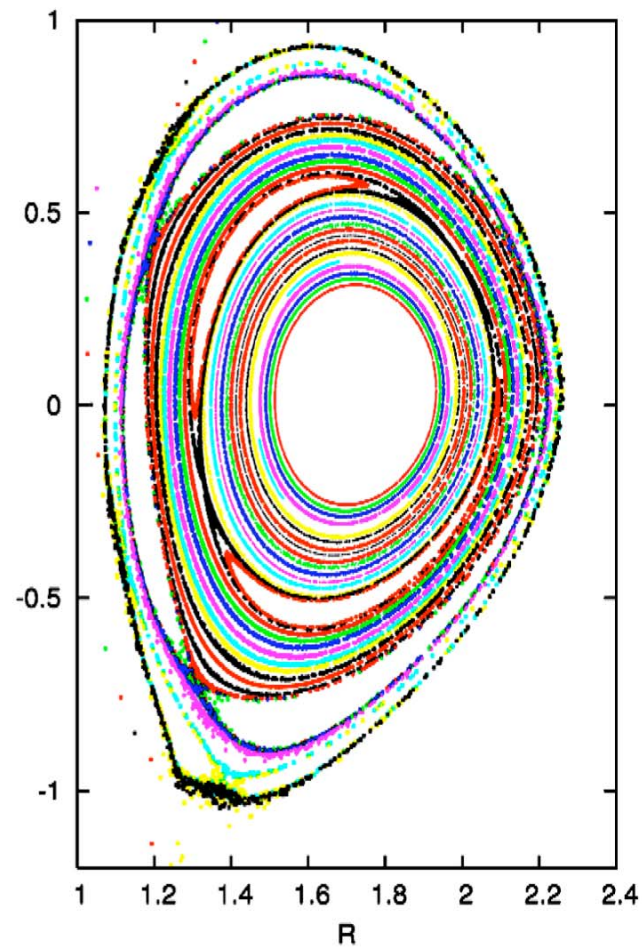
\*Snyder, *et al.*, *Nucl. Fusion* **47** (2007); Aiba, *et al.*, *Nucl. Fusion* **50** (2010); Ferraro, *et al.*, *Phys. Plasmas* **17** (2010)

# Rotation Improves Core Screening; But Sheared Rotation Stochasticizes Edge

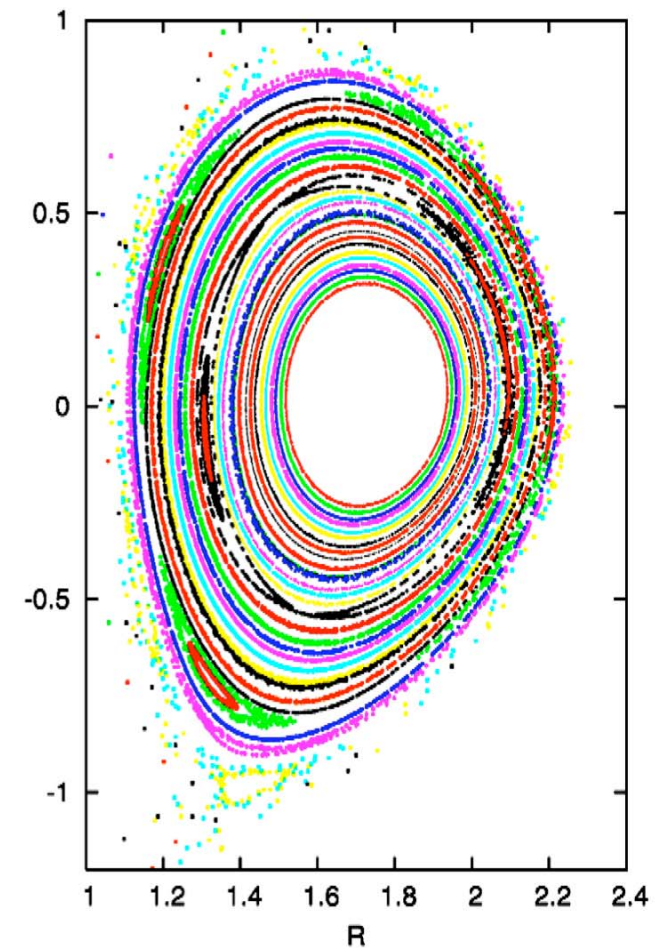
Vacuum



Plasma, Static



Plasma, Rotating



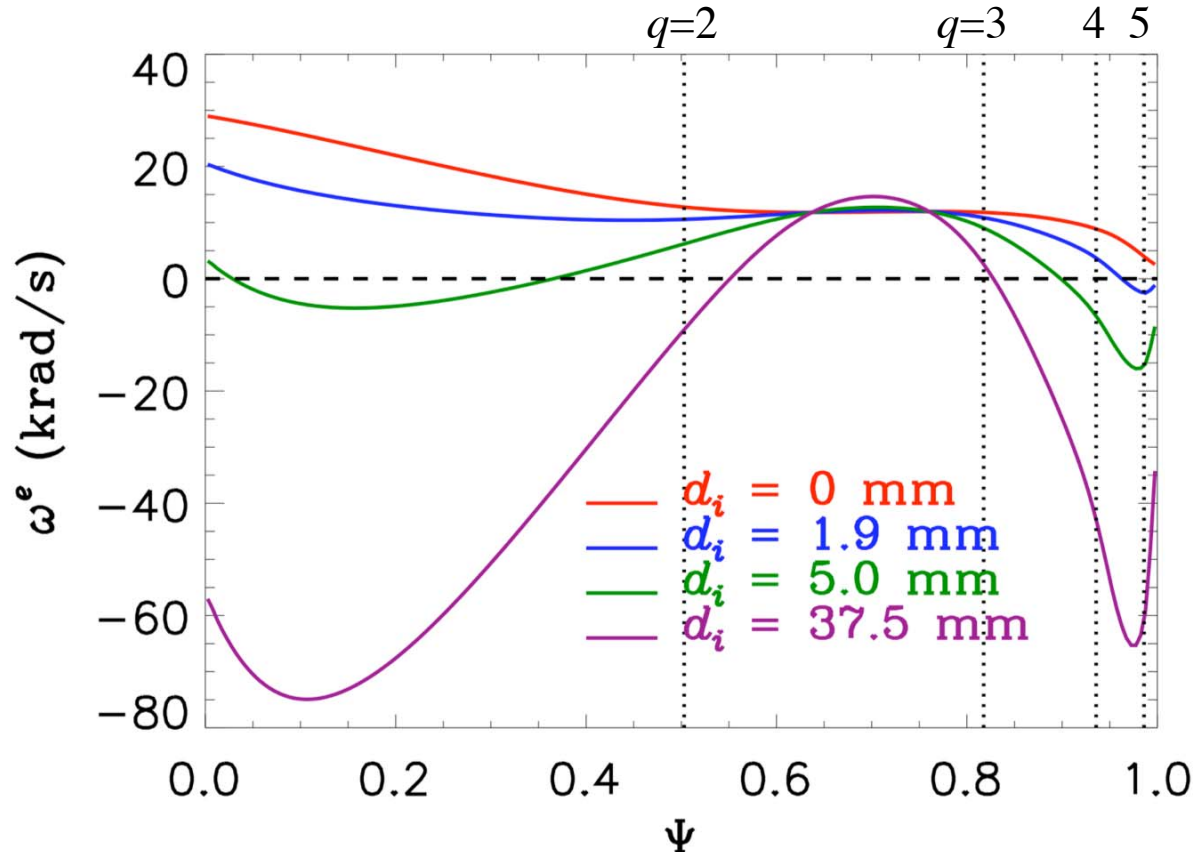
$$\Omega_0=0$$

$$\Omega_0=300 \text{ krad/s}$$

# Two-Fluid Results

# Two-Fluid Results — Ion and Electron Rotations are Distinct

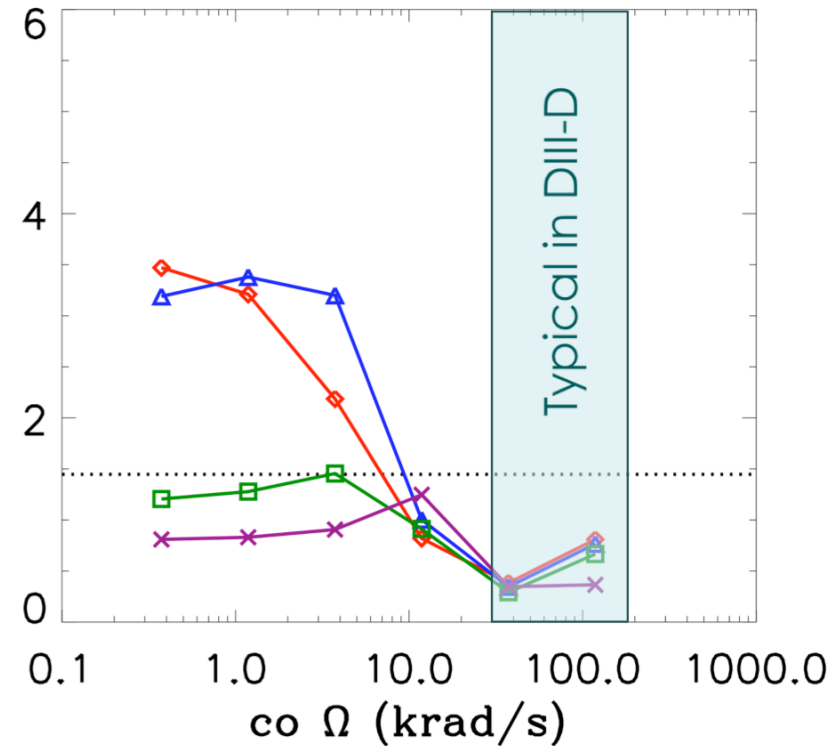
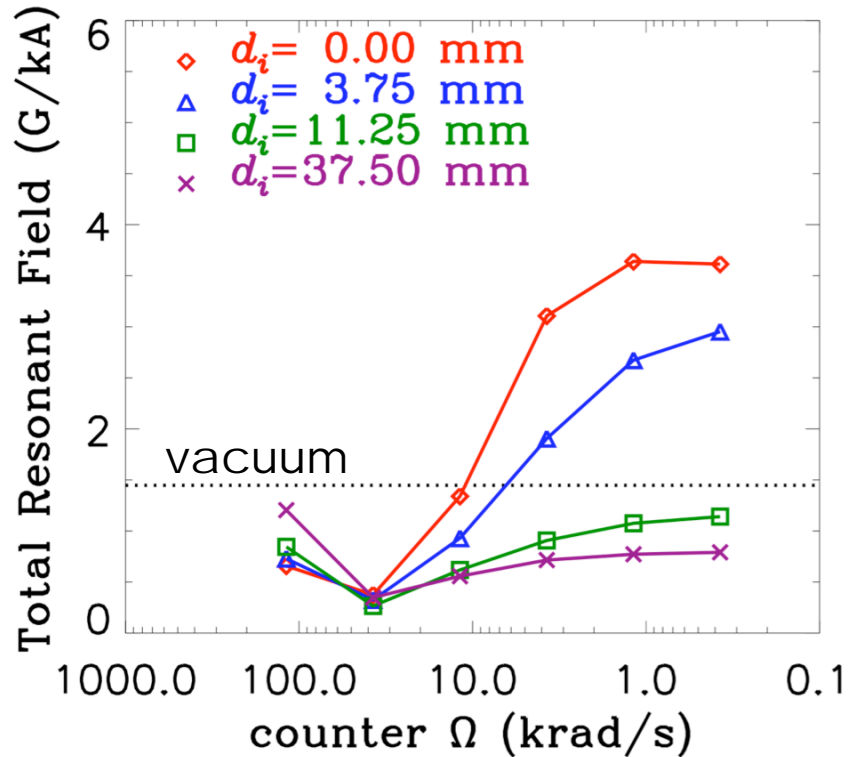
- Given  $\Omega$ , we can change  $\Omega^e = \Omega + \omega_*$  by adjusting  $\omega_* = d_i p' / n$



For this equilibrium,  $d_i = 37.5$  mm is the physical value

# Two-Fluid Effects Shift Resonance

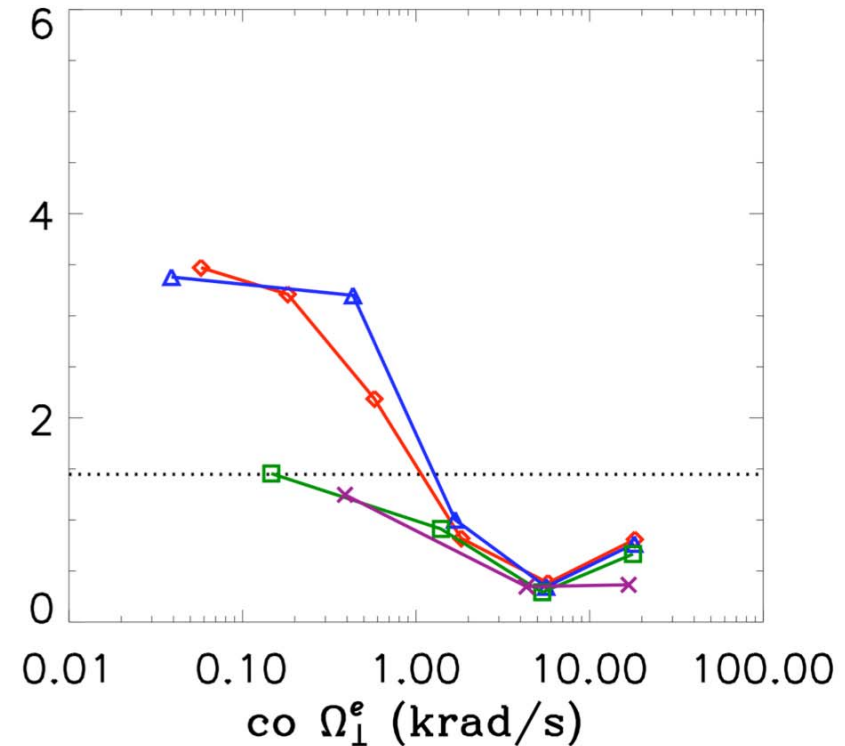
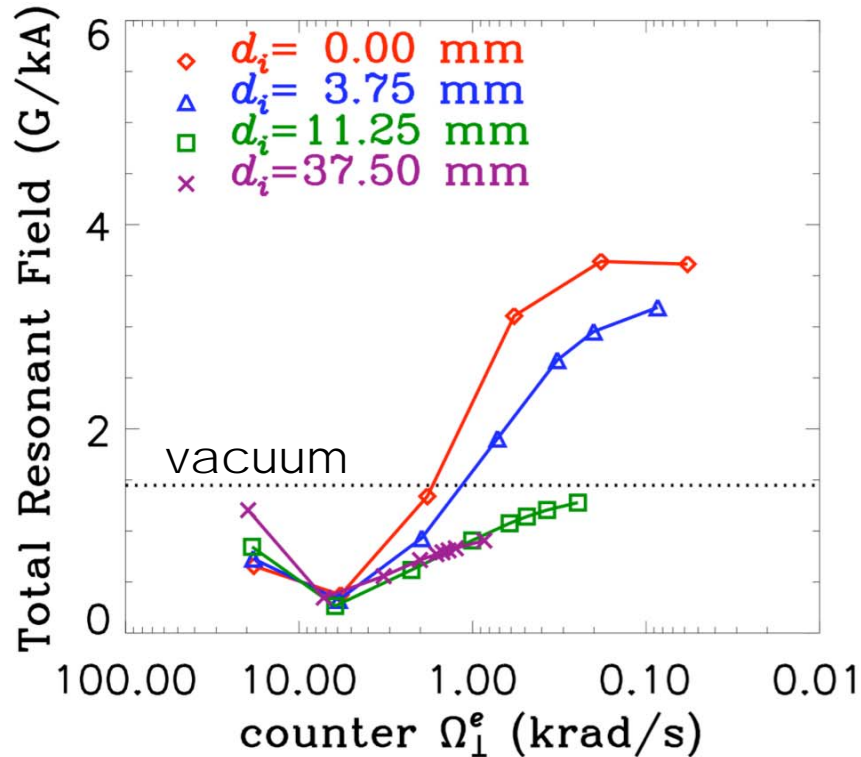
(Mass) rotation at  $q=3$



- Strongest tearing no longer occurs at  $\Omega = 0$

# Penetration In Core Depends on Electron Rotation

Perpendicular electron rotation at  $q = 3$



- Screening of  $q=3$  island clearly depends more on  $\Omega^e$  than  $\Omega$

# What is “Perpendicular” Electron Velocity?

- The perpendicular angular velocity is defined as

$$\Omega_{\perp}^{e,i} = \frac{\mathbf{v}^{e,i}}{R} \cdot \frac{\mathbf{B} \times \nabla \psi}{|\mathbf{B} \times \nabla \psi|}$$

- To lowest order,  $\mathbf{v}^{e,i} = R^2 \omega^{e,i}(\psi) \nabla \varphi + \lambda^{e,i} \mathbf{B}$ . Thus:

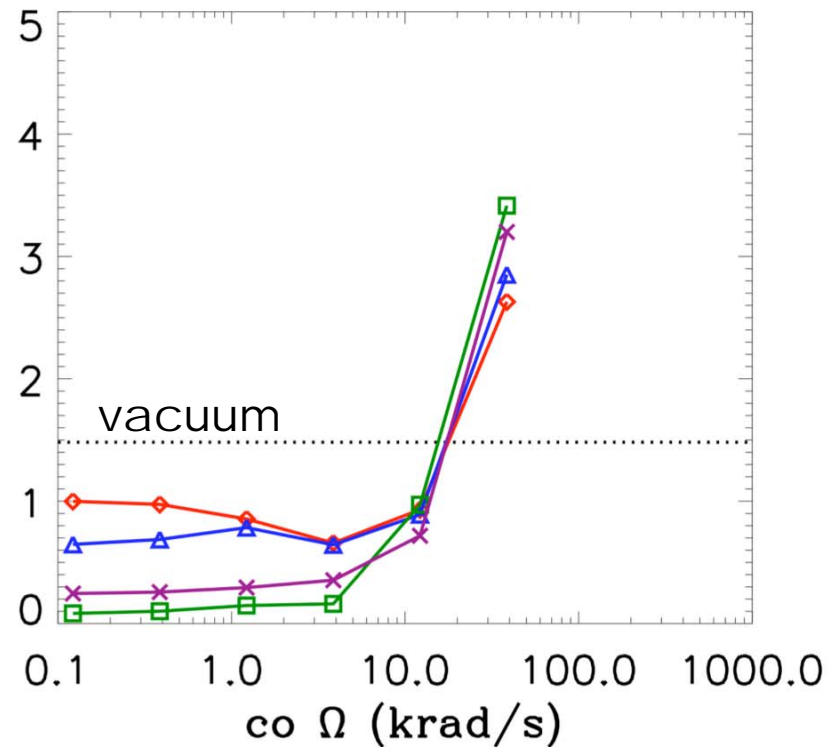
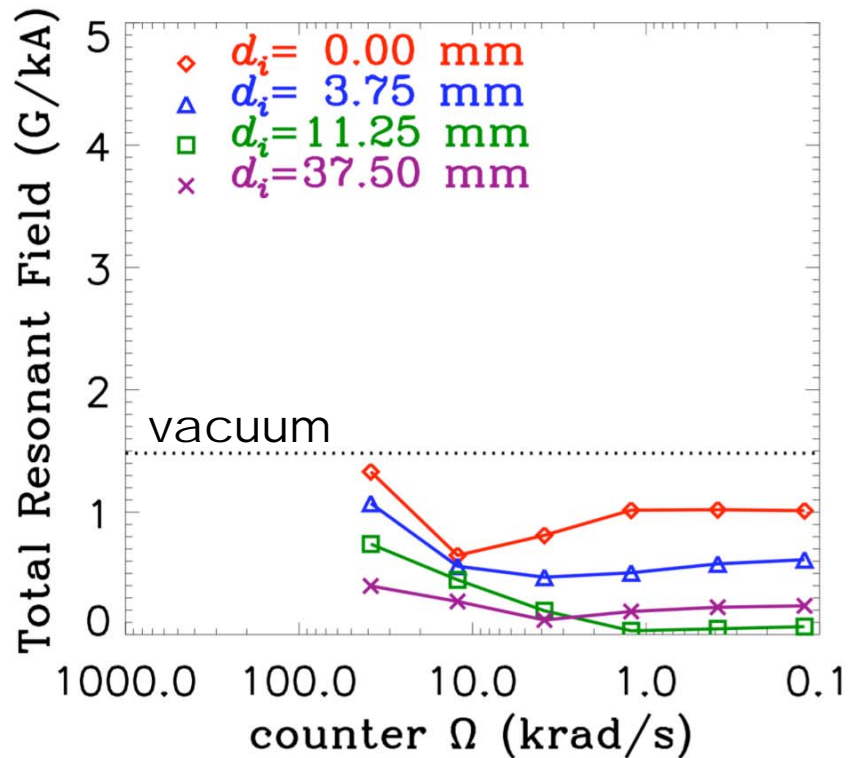
$$\Omega_{\perp}^{e,i} = \frac{|\nabla \psi \times \nabla \varphi|}{|B|} \omega^{e,i}(\psi)$$

- From radial force balance:  $\omega^{e,i}(\psi) = \phi'(\psi) + \frac{p_{e,i}'(\psi)}{n_{e,i} q_{e,i}}$
- $\Omega_{\perp}^e$  vanishes wherever  $\omega^e$  vanishes, but also at  $\mathbf{B}_{pol}$  nulls

# Edge Response Depends on Mass Rotation Shear

- Tearing of edge modes is dependent on ion, not electron, rotation shear

(Mass) rotation at  $q=5$

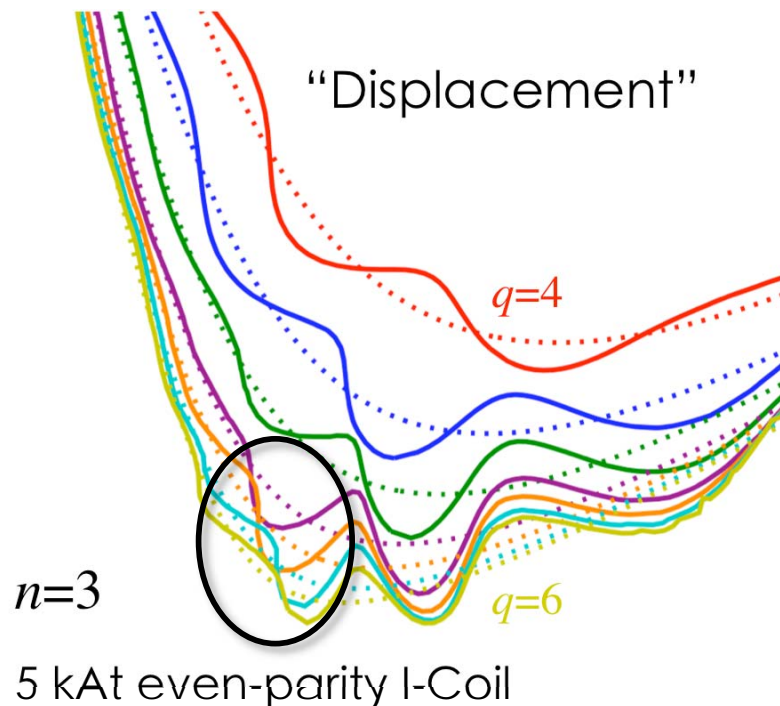




# Linear vs. Nonlinear

# Is Linear Response Appropriate?

- For typical experimental parameters, linear response may not be strictly valid in some regions
  - Large current density near rational surfaces
  - Back-reaction on rotation is important

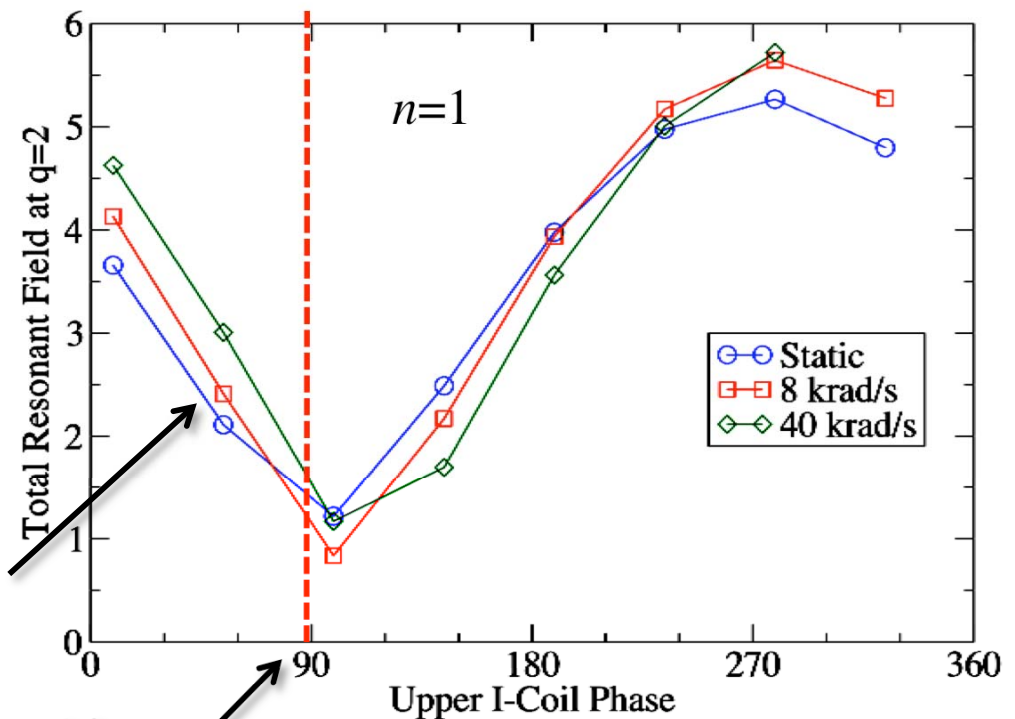


- “Displacement” shows overlapping surfaces near separatrix!

- Quantitative predictions of island size, stochasticity from linear calculations are suspect

# Linear Response Gets Some Things Right

- Which modes are most sensitive
- How parameters (rotation, viscosity, etc.) affect sensitivity
- How to optimize coil design



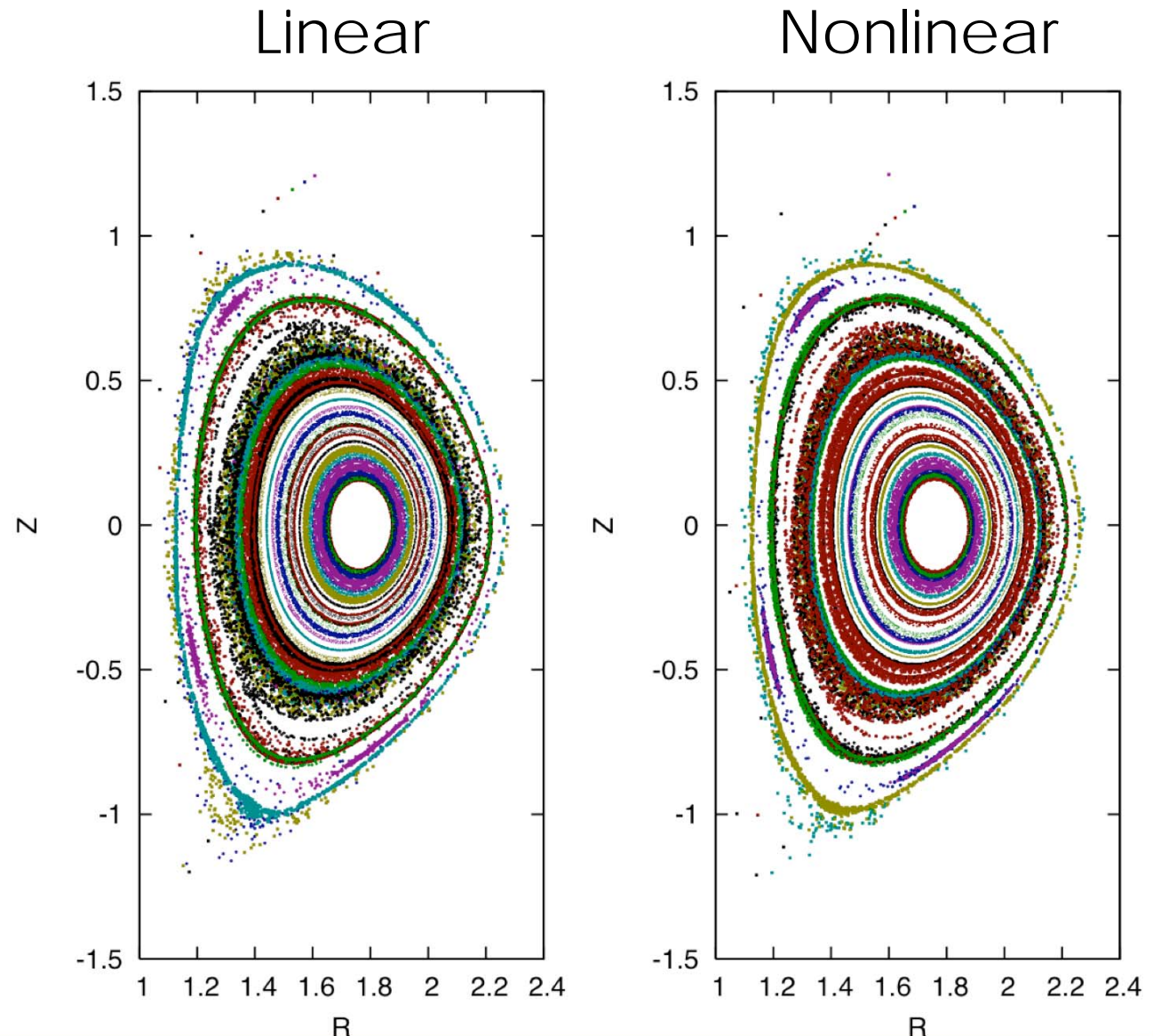
Calculated resonant field  
(proxy for resonant torque)

Empirical phase least prone to locking

DIII-D 144182

# Nonlinear Calculations are Underway

- **Nonlinear calculations are necessary for some things**
  - Rotation/locking
  - Transport
  - Large islands
- **Preliminary non-linear results agree with linear results for non-rotating plasma**



# Summary

- We can calculate resistive two-fluid plasma response for diverted equilibria with realistic transport parameters
- Rotation (usually) improves screening
- Perpendicular electron velocity is the most relevant rotation for core islands
  - ELM suppression may correlate with intersection of  $\omega^e \approx 0$  and rational surface ( $q_{95}$  windows?)
- (Mass) rotation seems to enhance edge response
  - Edge rotation may be crucial to ELM suppression (depending on mechanism)

# Challenges to Understanding Plasma Response Remain

- Mode locking depends sensitively on scaling of viscous torques
- ELM suppression requires interplay between field response and transport (probably)
- Need better understanding of transport in 3D fields
- **3D equilibrium properly requires nonlinear calculation**
  - $n = 0$  rotation and  $n \neq 0$  response are strongly coupled
  - Island saturation is nonlinear
  - There is healthy debate how to do this efficiently!
- **We are moving quickly to overcome these challenges**