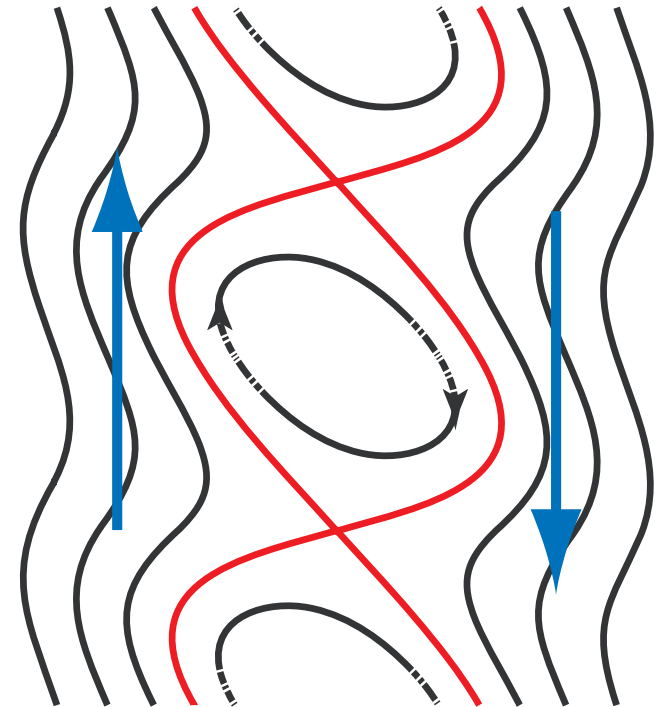


Islands in the Stream: The Effect of Plasma Flow on Tearing Stability

by
R.J. La Haye, General Atomics,
with many acknowledgments given at end of talk

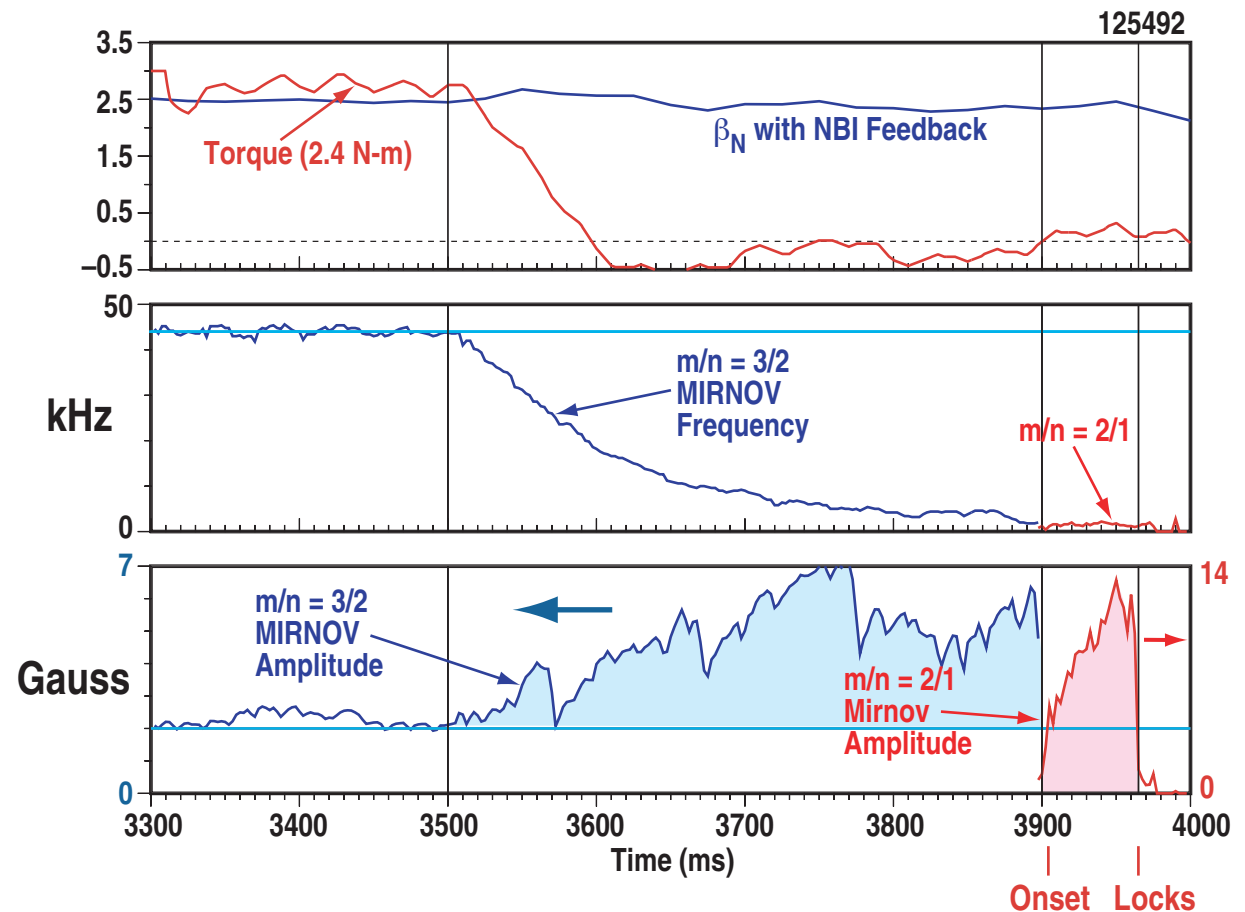
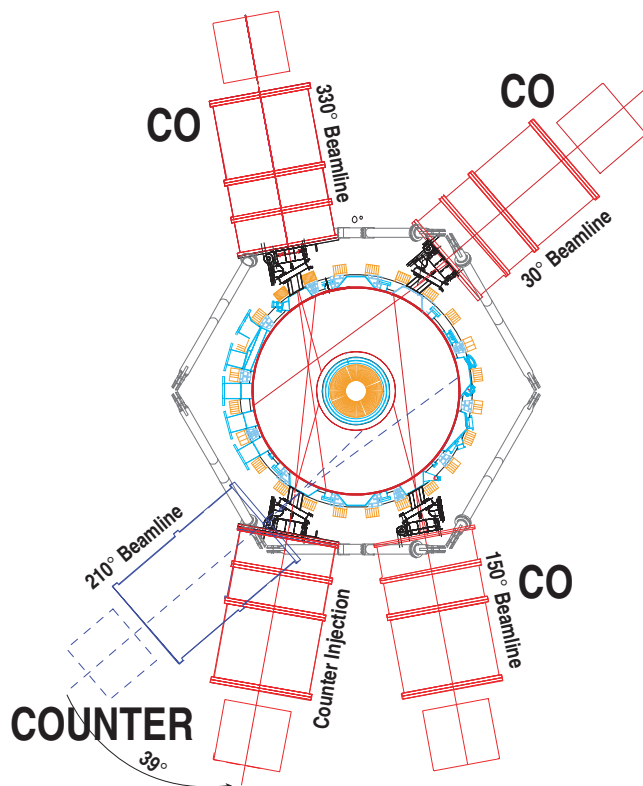
Presented at the
51st APS Annual Meeting of
the Division of Plasma Physics
Atlanta, Georgia

November 2–6, 2009



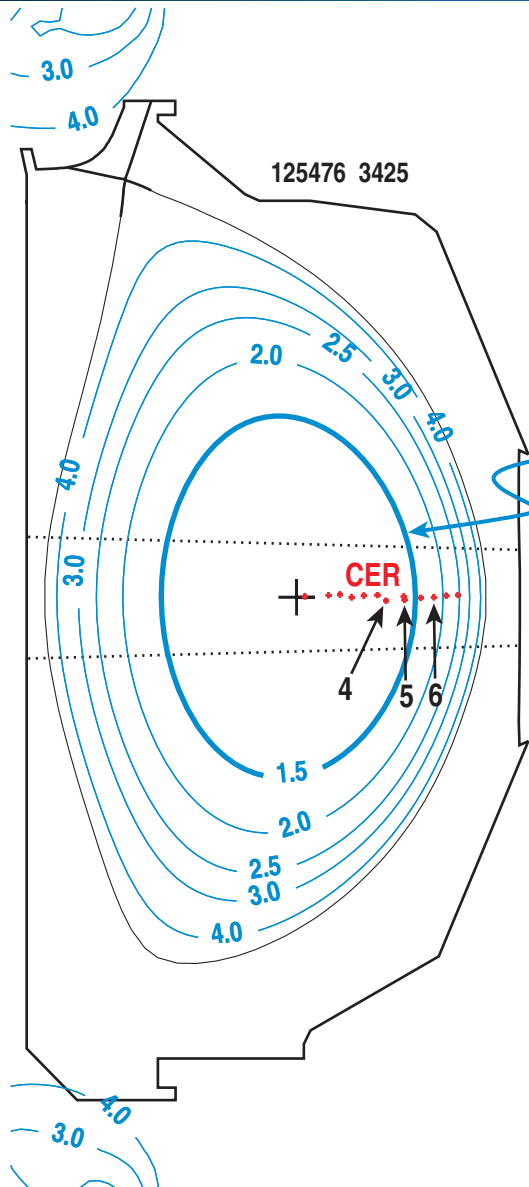
Prologue: Decreasing Plasma Flow Impairs Tearing Stability

- Unanticipated consequences that impact performance for future tokamaks with low rotation



- Existing tearing modes ($m/n = 3/2$ shown) get bigger \Rightarrow degrades confinement
- Otherwise stable $m/n = 2/1$ mode can become unstable \Rightarrow lock, disrupt

Plasma Rotation Measured by Charge Exchange Recombination (CER) of CVI Line

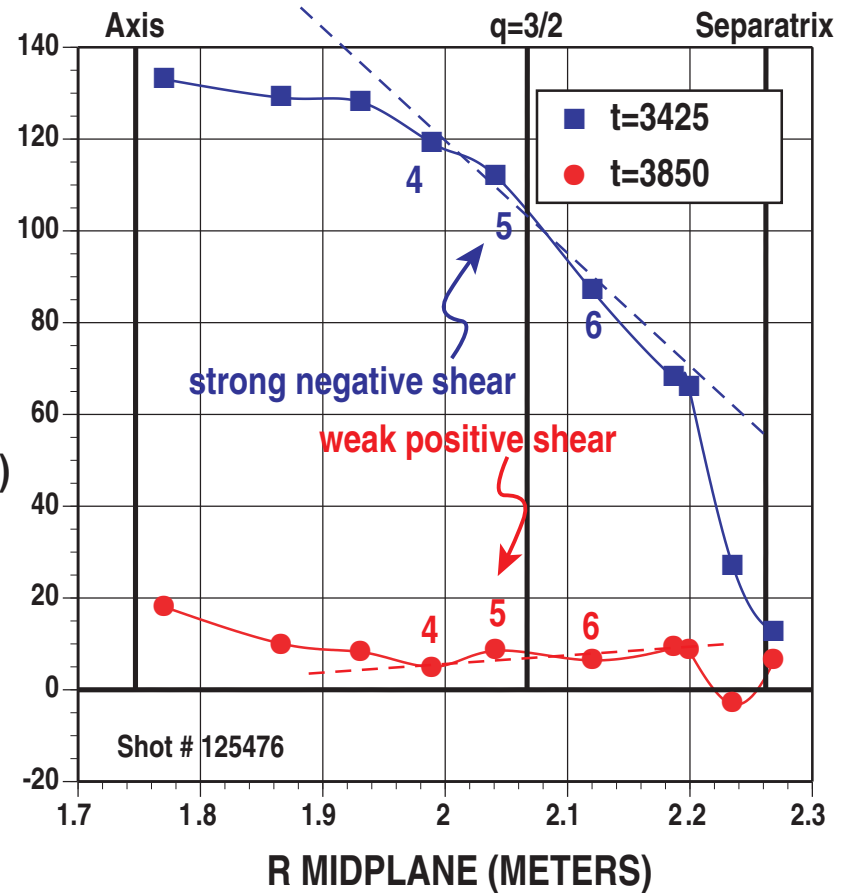


Locate $q = m/n$ with MHD reconstruction code EFIT with MSE

poloidal mode m , toroidal mode n , surface ($q=3/2$)

Ω_ϕ (krad/s)

Toroidal Angular Rotation



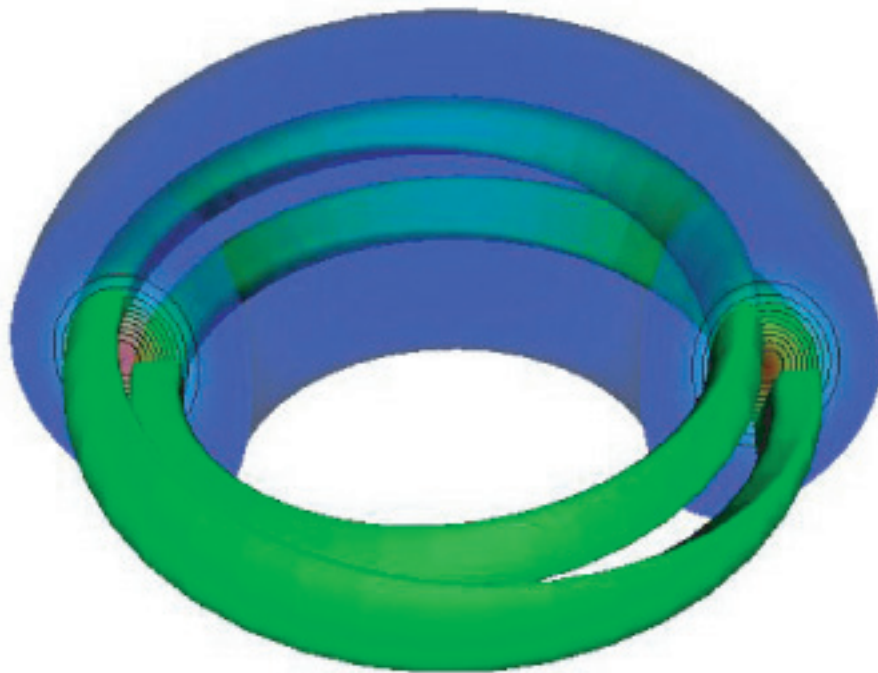
Classical Tearing Mode Island Stable ($\Delta' < 0$) with a **Stable** Total Plasma Current Density Profile

- Rutherford Equation (P.H. Rutherford 1973) $\frac{\tau_R}{r^2} \frac{dw}{dt} = \Delta'$, the effective classical tearing index

★ for $\Delta' < 0$, the helically perturbed current has higher energy

– linearly stable as the perturbed energy $\delta W_{\text{mag}} \approx -\frac{1}{4\mu_0} \left(\frac{r\tilde{B}_r}{m} \right)^2 \Delta' r > 0$

$$\dots \Delta' = \frac{d\tilde{B}_r/dr}{\tilde{B}_r} \Big|_{r_s - \epsilon}^{r_s + \epsilon} \quad \text{calculable from magnetic equilibrium}$$



- Island in DIII-D simulated (NIMROD) for an $m/n = 2/1$ tearing mode
 - ★ decays away if classically stable
 - ★ but can be sustained by helically perturbed bootstrap current (NEOCLASSICAL TEARING MODE)

Postulate: Flow Shear Effect On Classical Tearing Stability is a Key Feature of Neoclassical Tearing Mode Behavior

$$\frac{\tau_R}{r^2} \frac{dw}{dt} = \Delta' + \varepsilon^{1/2} \frac{L_q}{L_{pe}} \beta_{\theta e} \left[\frac{1}{w} - \frac{w_{\text{marg}}^2}{3w^3} \right] \quad \text{Modified Rutherford Equation (MRE)}$$

effective
classical
tearing
index

$$\varepsilon = r/R$$

$$L_q \equiv q / (dq/dR)$$

$$L_{pe} \equiv -p_e / (dp_e/dR) \frac{2\mu_0 p_e}{B_{\theta}^2}$$

lumped small island
threshold effects

- For large saturated islands ($\dot{w} = 0$)
- For onset of islands ($\dot{w} > 0$)

$$0 \approx \Delta' + \varepsilon^{1/2} \frac{L_q}{L_{pe}} \frac{\beta_{\theta e}}{w} \left[1 - \frac{w_{\text{marg}}^2}{3w^2} \right]$$

$$\beta_{\theta e} > \frac{3}{2} \frac{(w_{\text{marg}}/r) (-\Delta' r)}{\varepsilon^{1/2} (L_q/L_{pe})}$$

- Look for influence of flow shear on effective classical tearing Δ'

$$\Delta' r \approx -C_0 - C \times (-d\Omega_{\phi}/dr) L_s \tau_A$$

Experimental Data Examined for Trends with the Normalized Flow Shear

- existing islands ($\dot{w} \approx 0$)
 $w \propto (\tilde{B}_\theta)^{1/2} \propto \frac{1}{-\Delta' r}$
- onset of islands ($\dot{w} > 0$)
 $\beta \propto -\Delta' r$

with $\Delta' r$ a function of the flow shear
normalized to both the parallel magnetic shear length
and the Alfvén time

When you follow two separate chains of thought, Watson, you will find some point of intersection which should approximate the truth.

— Sherlock Holmes, *The Disappearance of Lady Francis Carfax*
Sir Arthur Conan Doyle

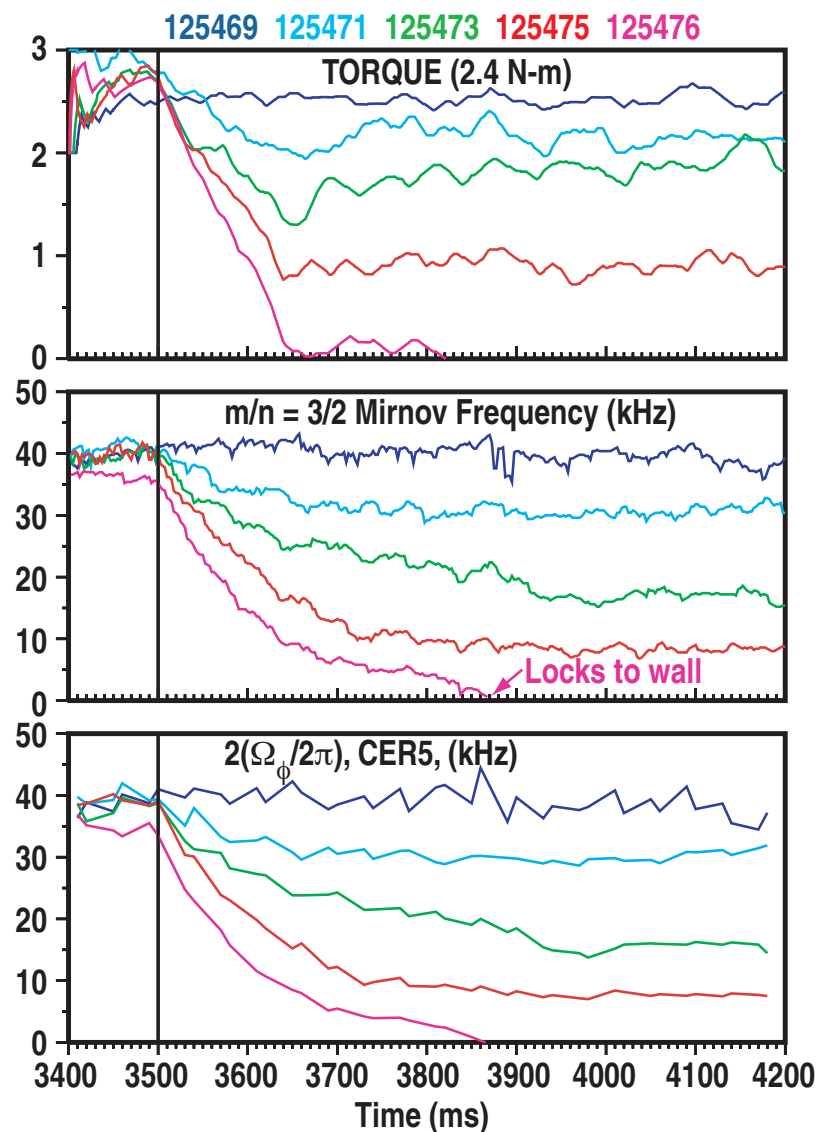
MODE AMPLITUDE VERSUS FLOW SHEAR

$m/n = 3/2$ Modes Get Larger with Lower Rotation (and Shear) as “Inferred” Δ' Less Negative

- All co-NBI has counter beams added ($\beta_N \approx 2.6$, $\bar{n}_{13} \approx 3.5$)

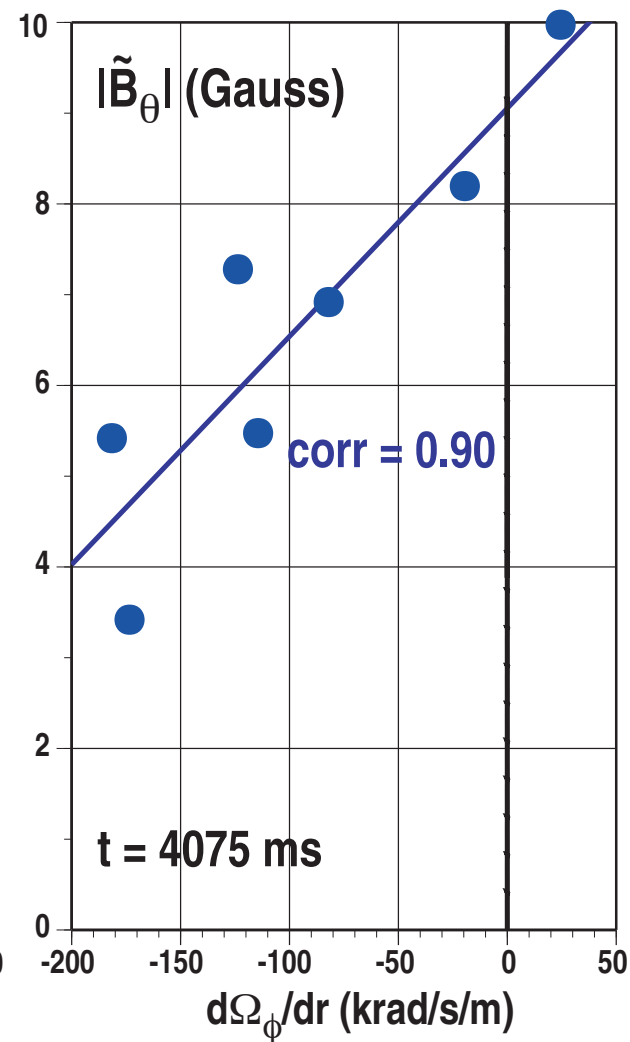
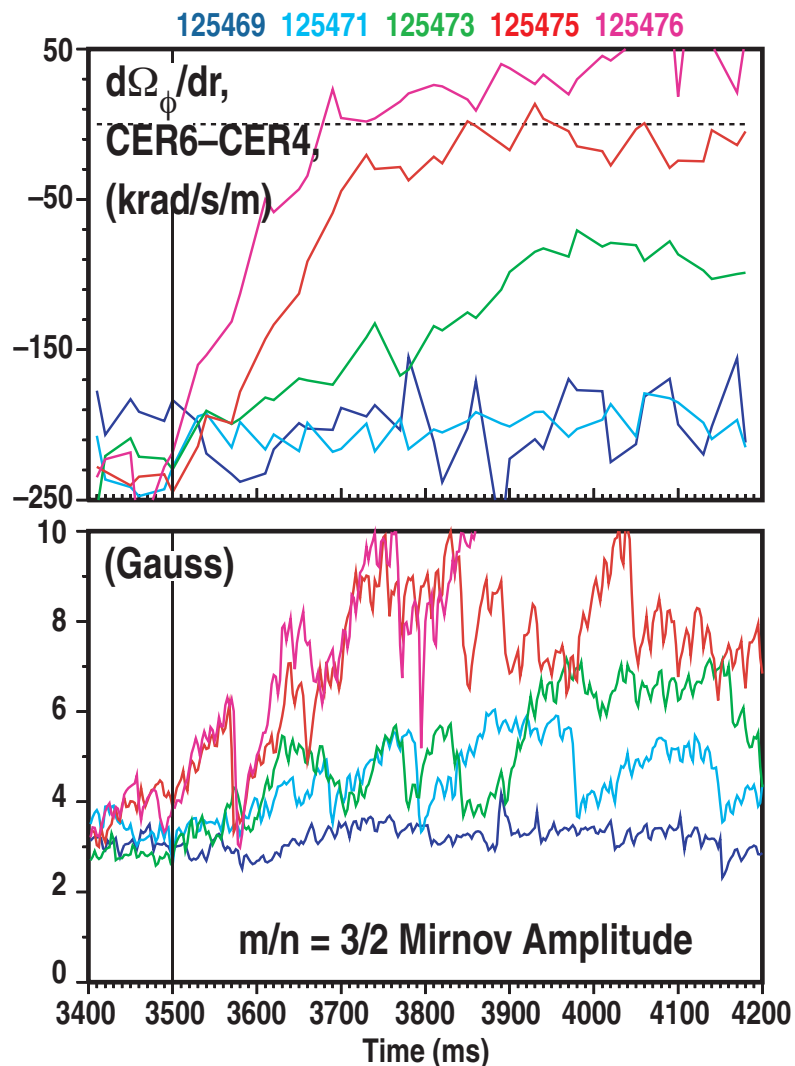
★ rotation reduced

...CER CVI chord 5 toroidal rotation is good measure of MIRNOV rotation



$m/n = 3/2$ Modes Get Larger with Lower Rotation (and Shear) as “Inferred” Δ' Less Negative

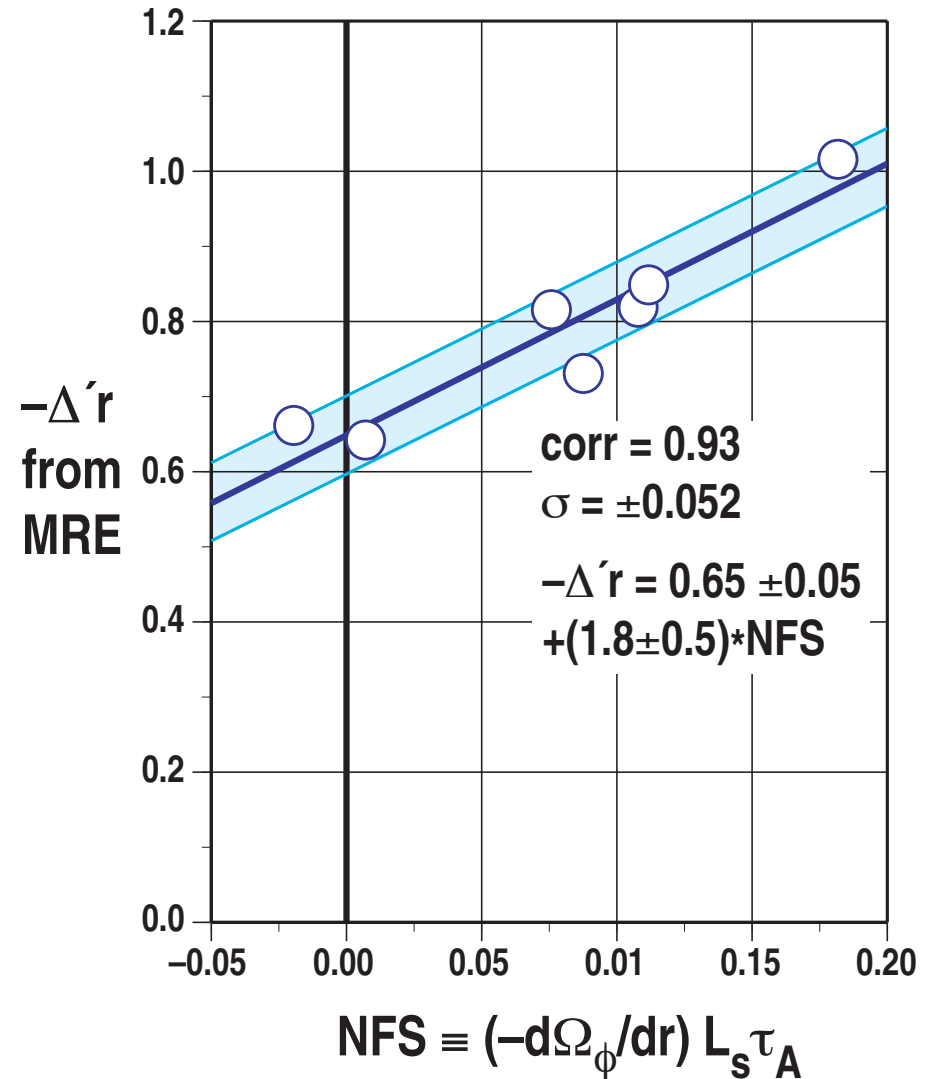
- As flow shear is reduced ...
 - ★ $3/2$ amplitude increases



$m/n = 3/2$ Modes Get Larger with Lower Rotation (and Shear) as “Inferred” Δ' Less Negative

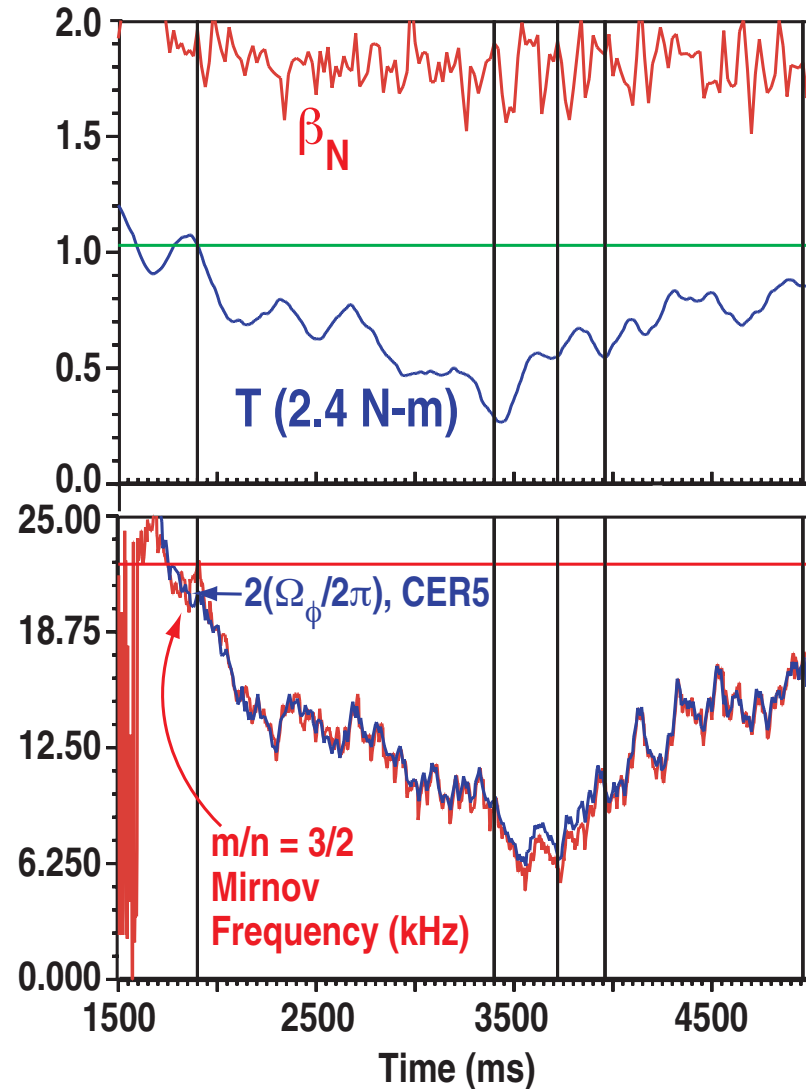
- “Inferred” Δ' **less negative** with reduced positive normalized flow shear

★ NFS ≈ 0.4 for a factor of two increase in stability



... and “Inferred” Δ' Less Negative with Inverted Flow Shear with $m/n = 3/2$ Modes

- Reduced β_N , less initial all co-torque (#135867) to avoid $m/n=2/1$ mode
 - ★ torque programmed down and back ($\sim 8 \tau_E$ each way)



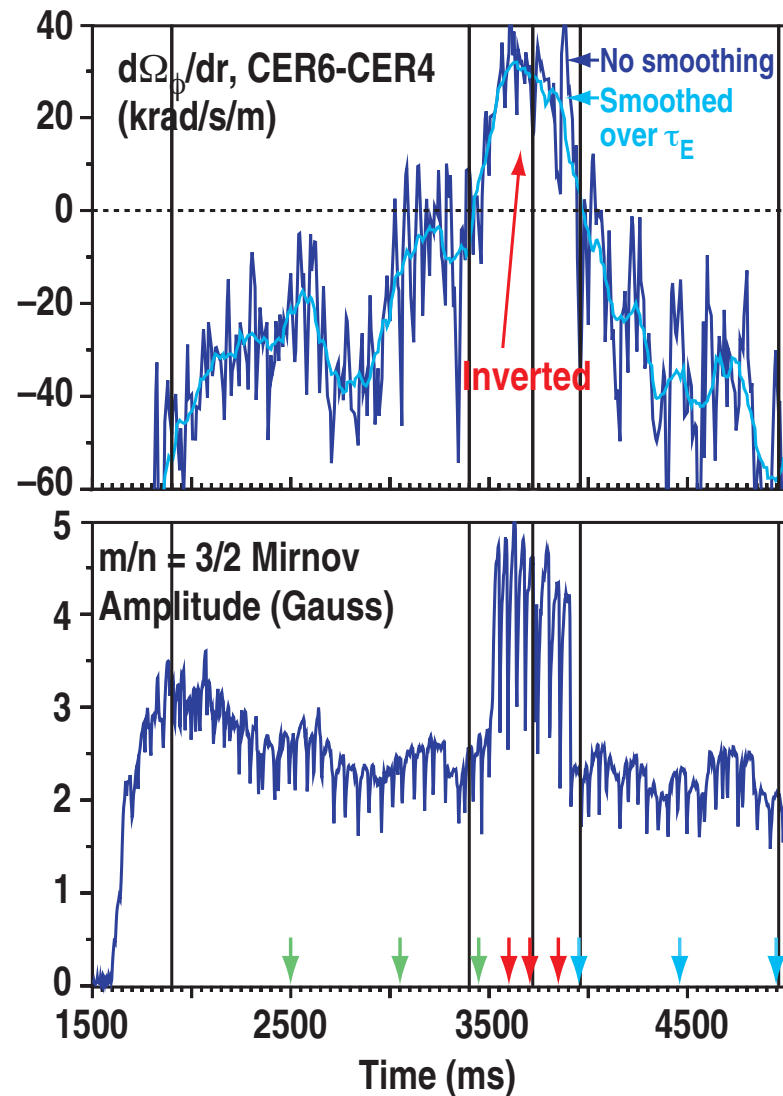
... and “Inferred” Δ' Less Negative with Inverted Flow Shear with $m/n = 3/2$ Modes

- Initial flow shear small

- ★ $d\Omega_\phi/dr \approx 0$
“forbidden” state
... inverts to hollow

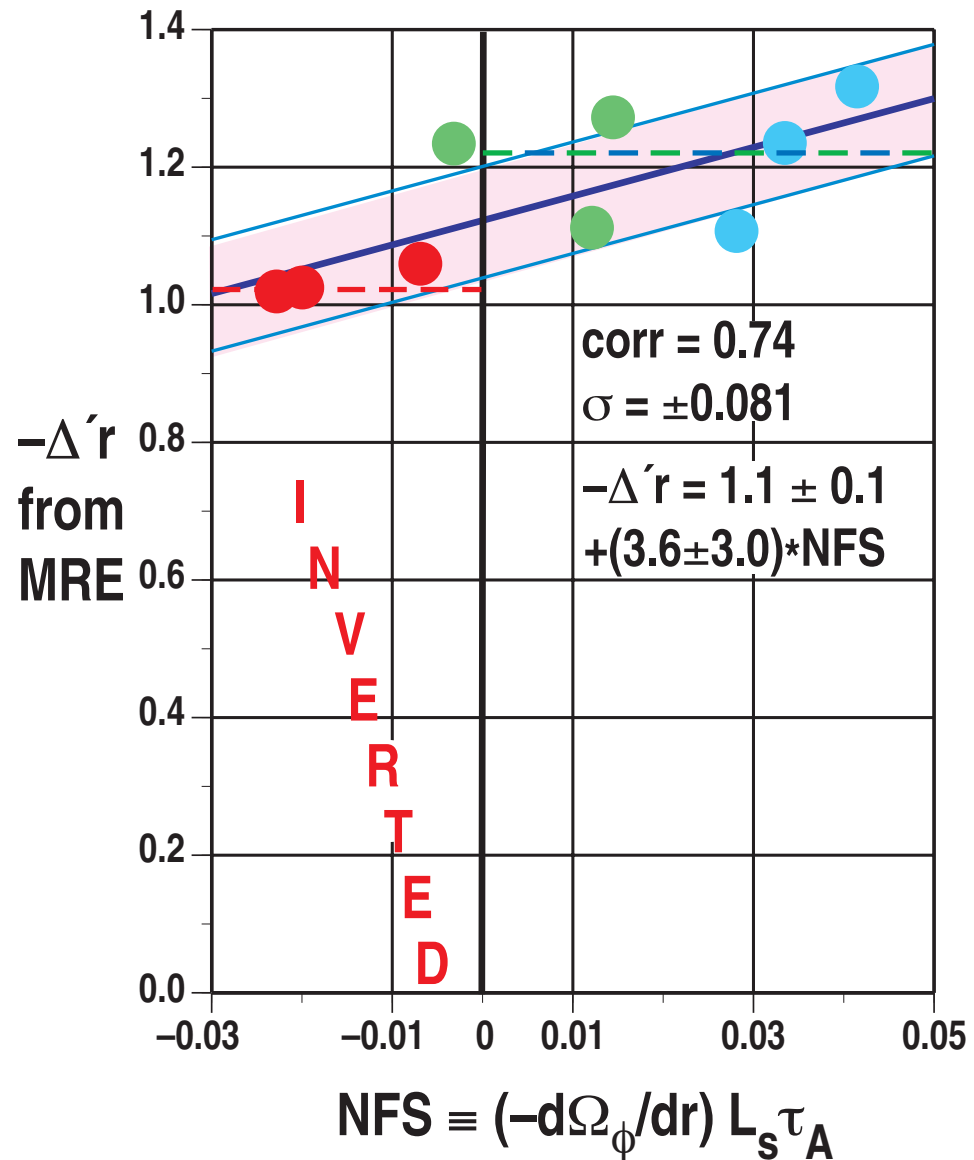
- Little change in Mirnov amplitude

- ★ until inversion



... and “Inferred” Δ' Less Negative with Inverted Flow Shear with $m/n = 3/2$ Modes

- “Inferred” Δ' depends on sign of $-d\Omega_\phi/dr$?
 - ★ or shift of zero to left?
 - ★ NFS ≈ 0.3 for a factor of two increase in stability



BETA AT ONSET VERSUS FLOW SHEAR

Reduced Plasma Rotation Experimentally Destabilizes Neoclassical Tearing Modes at Lower Beta in DIII-D, NSTX and JET

- Each device with slow ramp up in beta to instability
 - ★ note mixing of seeding and small island effects
 - ... makes analysis of Δ' by itself problematic

Also see R.J. Buttery PP8.00050 We PM

“Sensitivity of Tearing Mode Beta Limits to Rotation and Current Profile”

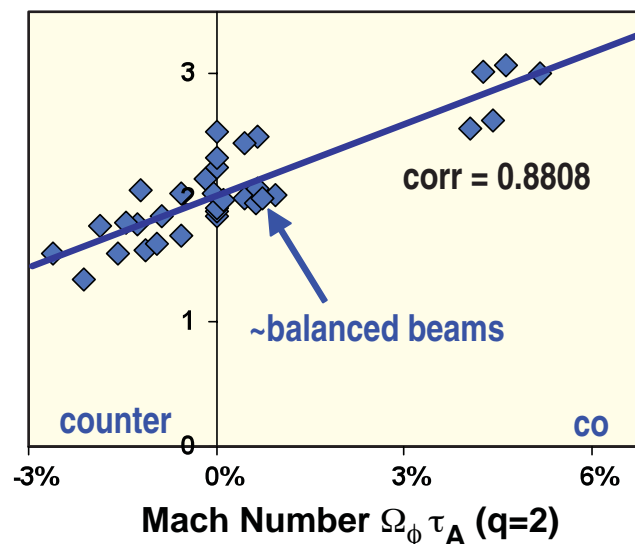


Onset β Continues to Fall with Counter Plasma Rotation in **DIII-D**

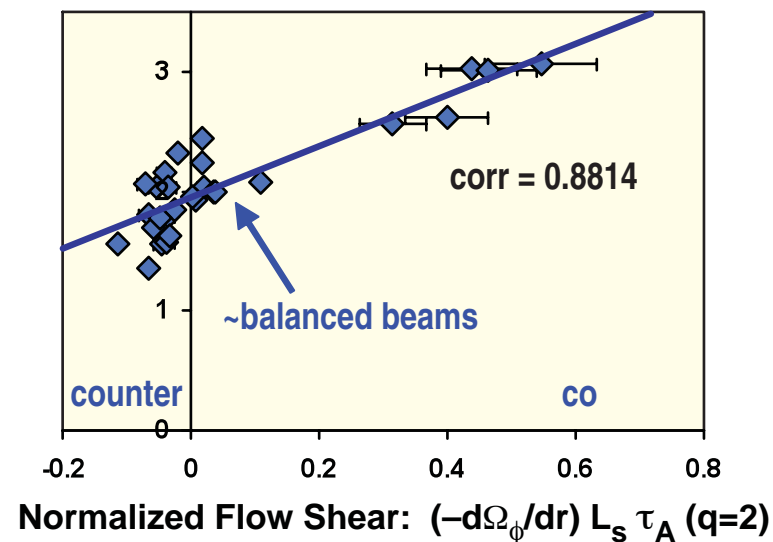
- No systematic variations or trends in key profiles
 - ★ rotation degenerate with flow shear
 - ... profile unchanged with co/ctr beam mix
- For counter rotation, flow shear reverses sign wrt magnetic shear
 - ★ relative sign important?
 - ... or “shift” in zero to left?

(R.J. Buttery et al, PoP 2008)

2/1
Onset
 β_N
 $\left(\frac{\beta}{I/aB}\right)$

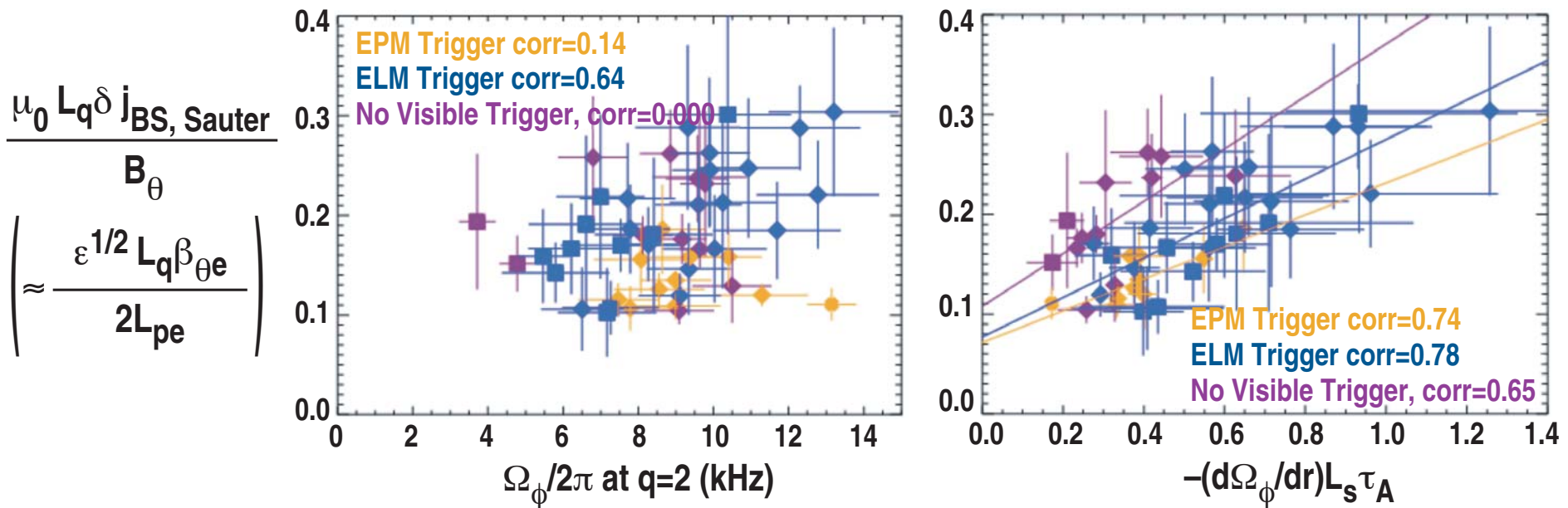


2/1
Onset
 β_N
 $\left(\frac{\beta}{I/aB}\right)$



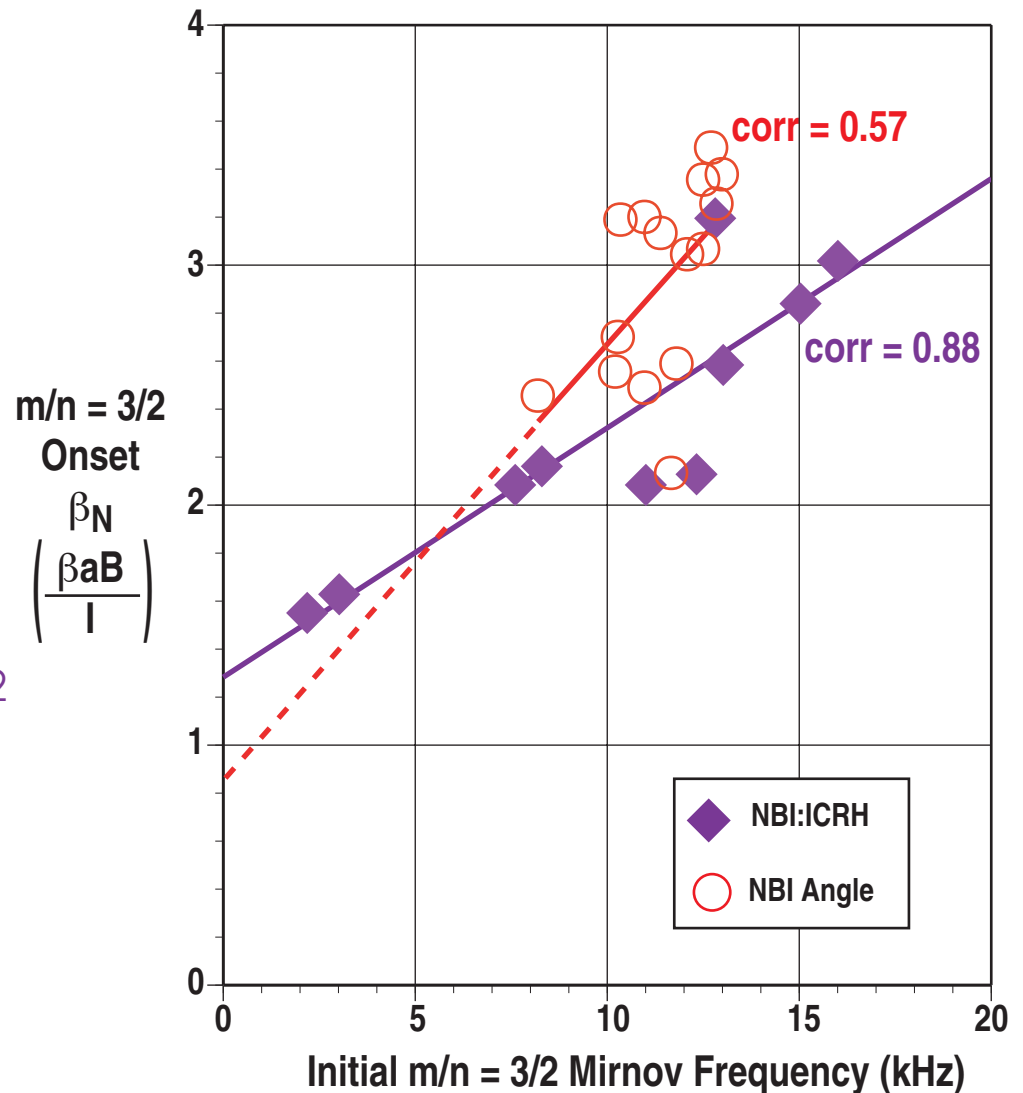
Strongest Correlation Found Between Local Mode Drive at Onset and Normalized Flow Shear in Low Aspect Ratio **NSTX**

- NTM “drive” at onset only poorly correlated with $q=2$ (carbon) rotation, “CHERS”
- Essentially no correlation for Energetic Particle Modes and “Triggerless” cases
 - ELM correlation better
- Degeneracy between Ω_ϕ and $d\Omega_\phi/dr$ broken in NSTX with use of $n = 3$ magnetic braking from external coil
- Strongest correlation between “noise-prone” parameters (i.e., gradients)
 - Provides confidence that this is the correct physics



Distinct Onset β Decline in JET with ICRH: NBI Scan or NB Only Injection Angle for Sawtooth Triggered 3/2 NTM

- Ramp up power for different mixes of ICRH:NBI power or NBI only angle varied
 - ★ Phase ICRH to avoid large sawteeth
- Clear result obtained:
 - ★ Removing NBI momentum drops NTM onset
 - ... normalized flow shear at $q=3/2$ $(-d\Omega_\phi/dr) L_S \tau_A \leq 0.27$



PHYSICS OF FLOW SHEAR EFFECT ON TEARING

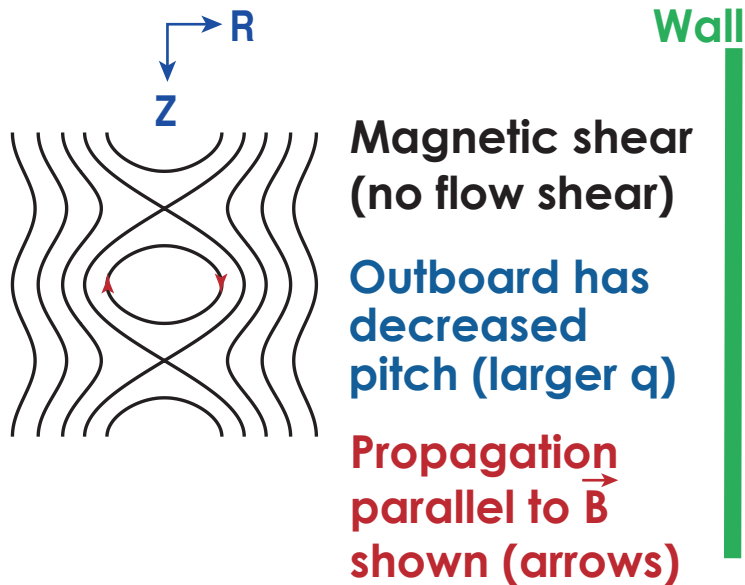
Flow Shear Has Been Invoked for Increasing Effective Classical Tearing Stability of Total Current Profile

“I am afraid that I rather give myself away when I explain. Results without causes are much more impressive.”

Sherlock Holmes, The Stock-Broker's Clerk
Sir Arthur Conan Doyle

- X.L. Chen & P.J. Morrison, PF 1990, “Resistive tearing instability with equilibrium shear flow” . . . **but zero viscosity**
 - ★ when the flow shear is larger than the magnetic shear..., the flow freezes the magnetic field and stabilizes the tearing mode.”
- L. Ofman, P.J. Morrison & R.S. Steinolfson, PF 1993, “Nonlinear evolution of resistive tearing mode instability with shear flow **and viscosity**”
 - ★ “shear flow decreases the saturated magnetic island width”
- R. Coelho & E. Lazzaro, PoP 2007, “Effect of sheared equilibrium plasma rotation on the classical tearing mode...”
 - ★ “above a given threshold in the rot. shear, a tearing mode, unstable in the absence of rotation, can be stabilized... the local derivative affects stability”
 - stabilizing for $P \equiv \tau_R/\tau_v \gg 1$ and $S \equiv \tau_R/\tau_A \gg 1$

How Magnetic Shear Influences Tearing



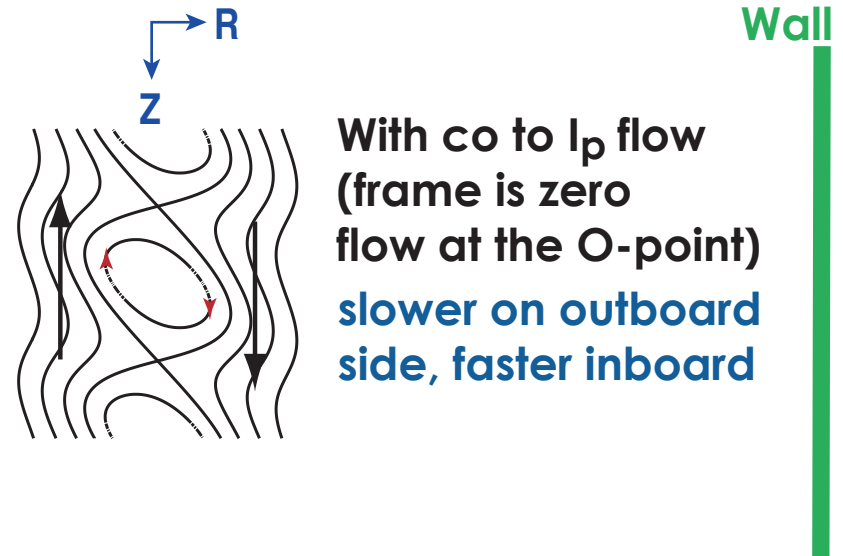
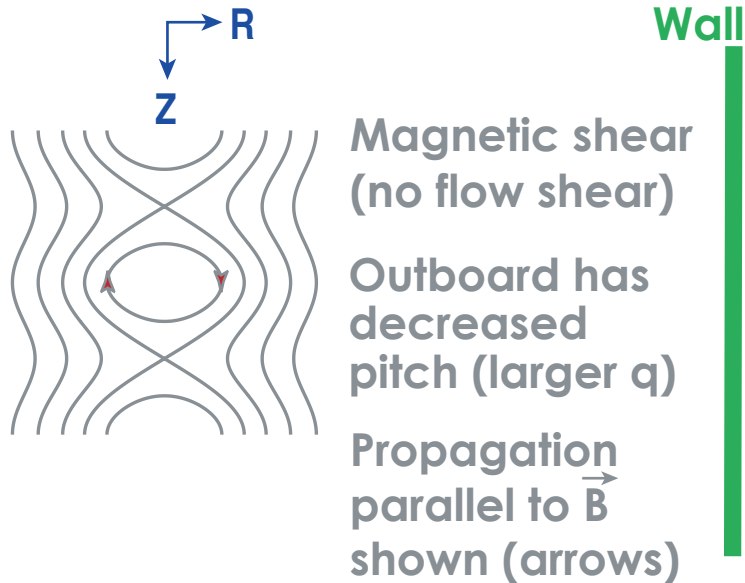
- **Magnetic shear varies field line pitch**

- ★ singularity at rational surface limited

$$\dots \left(1 - \frac{nq}{m}\right) = - \frac{(r - r_s)}{L_q}$$

- makes tearing “harder”
(more negative Δ')

How Flow Shear Influences Tearing



- Magnetic shear varies field line pitch
 - ★ singularity at rational surface limited

$$\dots \left(1 - \frac{nq}{m}\right) = - \frac{(r - r_s)}{L_q}$$

- makes tearing “harder”
(more negative Δ')

- Flow shear & viscosity make gradient of phase $\phi = n(Z + V_z t)/R - mx/r$

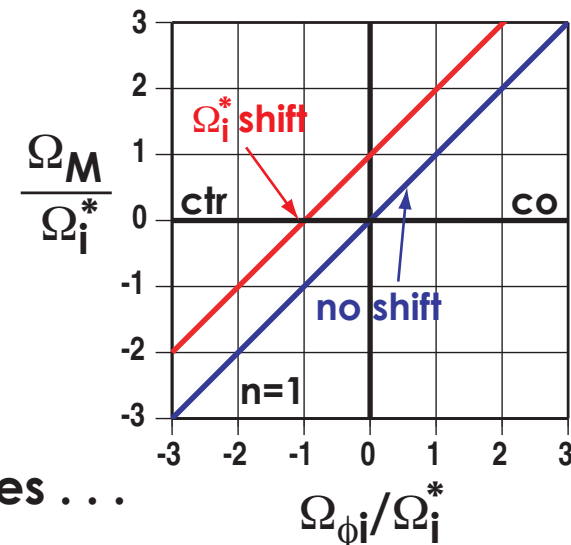
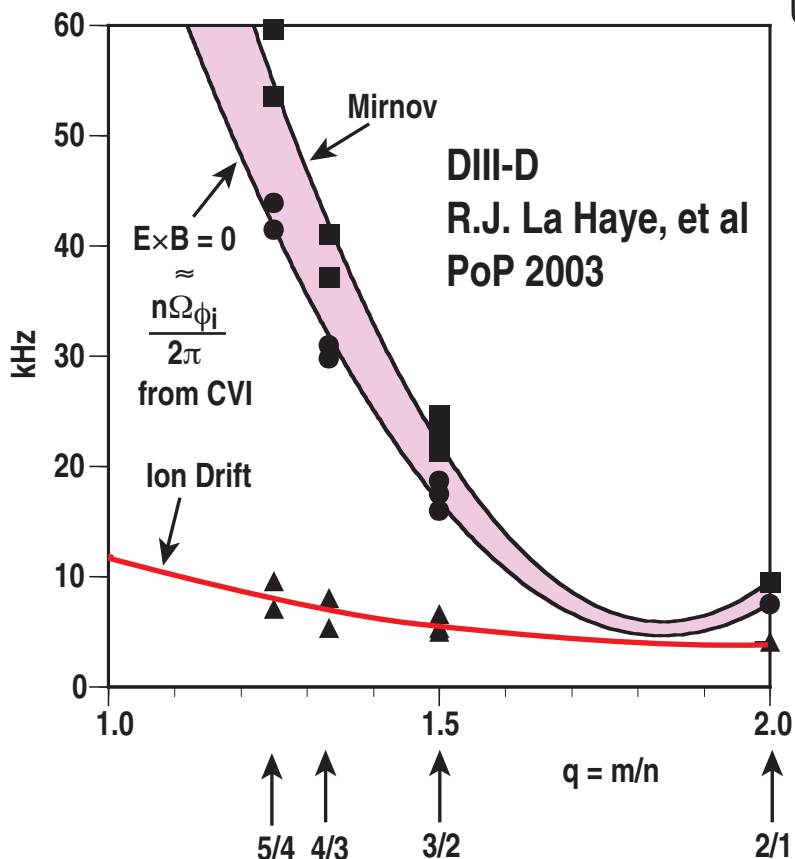
- ★ tearing resonance limited radially
 - $\dots \delta\phi \propto v_{\perp} d(V_z/R)/dr (r-r_s)$

- makes tearing “harder”
(inner layer effect dominant?)

As Islands Tend to Propagate Faster than the Plasma, Shear in This Rotation May be Mimicking a Sign Dependence

- viscosity forces island to propagate faster than n times the plasma toroidal rotation

★ $\Omega_{\text{Mirnov}} \approx n\Omega_{\phi i} + \Omega_i^*$ with $\Omega_i^* = n k_B T_i (dp_i/d\Psi)/e p_i$
 (ion diamag. drift freq.)
 usually in co-direction



- Flow shear now becomes ...

$$\approx \frac{d\Omega_{\phi i}}{dr} - \frac{\Omega_i^*}{nL_{Ti}}$$

\uparrow can change sign \uparrow "always" negative

★ null shifted to counter direction
 ... need deeper counter data to confirm

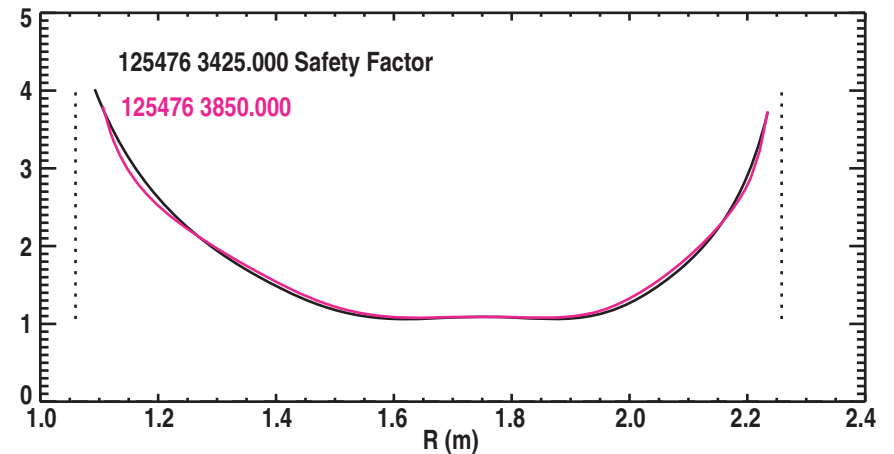
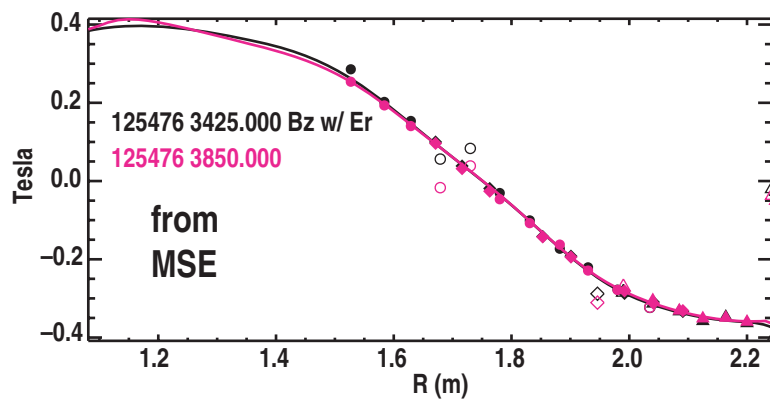
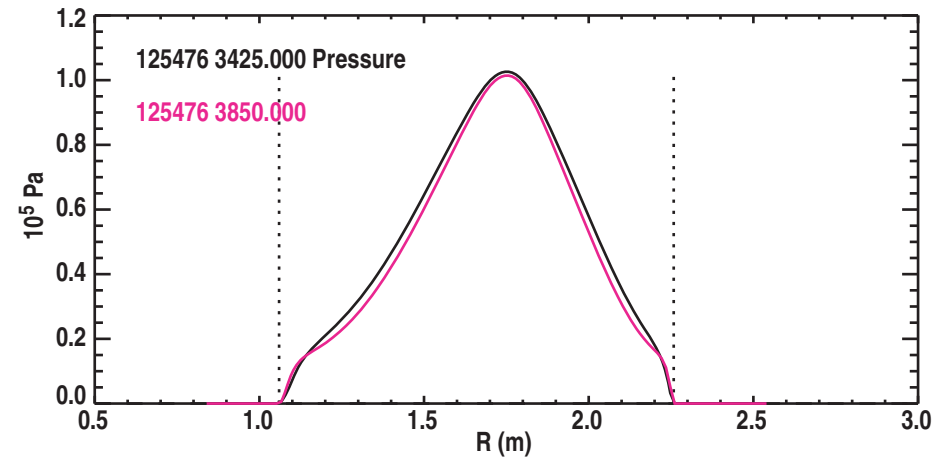
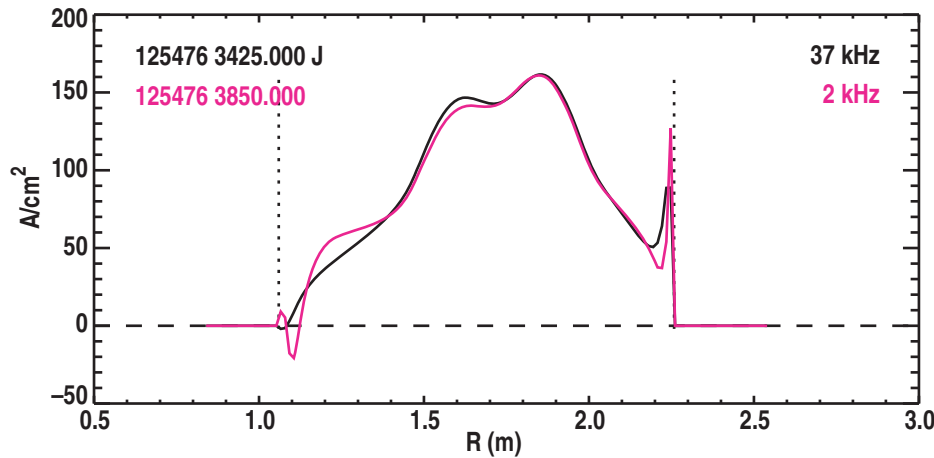
Conclusions and Acknowledgements

- A flow shear $-\frac{d\Omega_\phi}{dr} \approx \mathcal{O}\left(\frac{\tau_A^{-1}}{2L_s}\right)$ @ $q=m/n$
 - ★ has a significant stabilizing effect on tearing
 - thus an advantage for all co beams
 - . . . interpreted as an effect in classical Δ' of total plasma current profile
 - ★ but positive flow shear (counter peaked, co inverted) destabilizing?
 - may not be a sign effect but an offset?
- Acknowledgements to Dylan Brennan (University of Tulsa), Richard Buttery (formerly UKAEA Culham now General Atomics), Ming Chu (General Atomics), D. Raju and A. Sen (Institute for Plasma Research, Gujarat), Stefan Gerhardt (PPPL), Chris Hegna (University of Wisconsin, Madison), Scott Kruger (Tech-X Corp), Craig Petty (GA), Peter Politzer (GA), Ted Strait (GA), Francois Waelbroeck, (IFS Austin), and the DIII-D, NSTX and JET experimental teams

BACKUPS

Lower Rotation in Hybrid Scenario and NBI Feedback on Beta Produce Very Similar J and q Profiles

- **Difficult to discern significant change in Δ' from change in equilibrium**
 - ★ PEST-3 runs from “kinetic” EFITs with MSE \Rightarrow 3/2 stable at both times
 - ... rotation is not included in the code

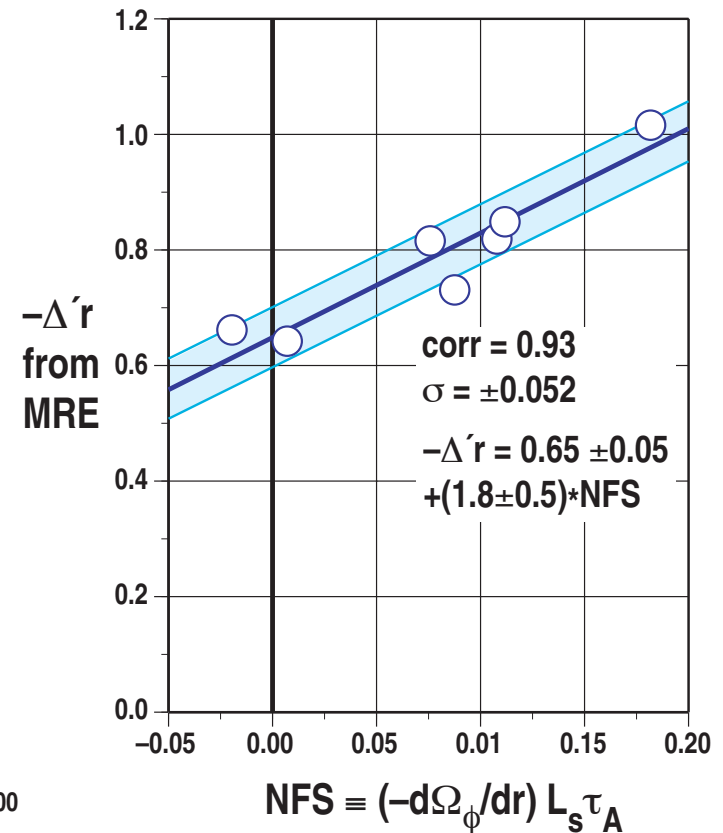
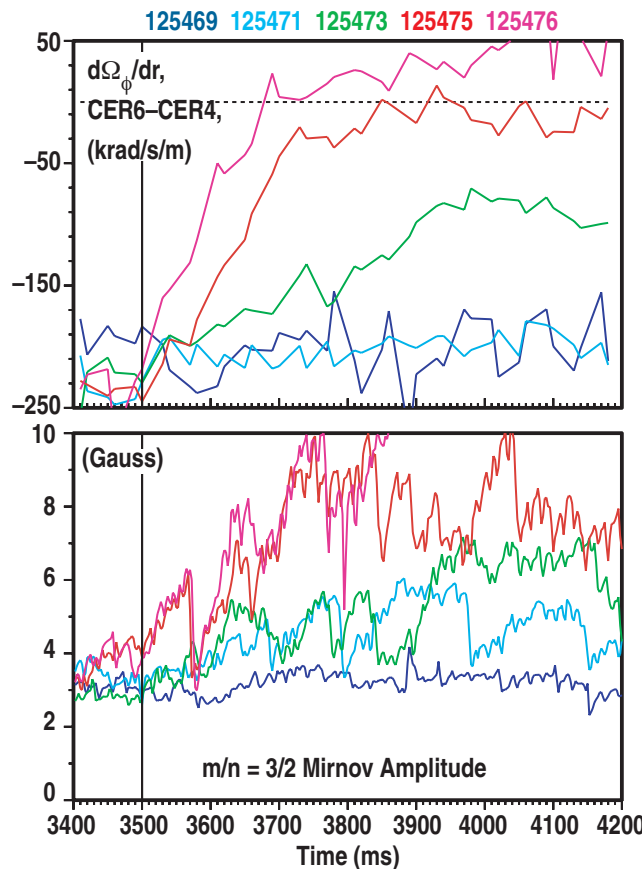
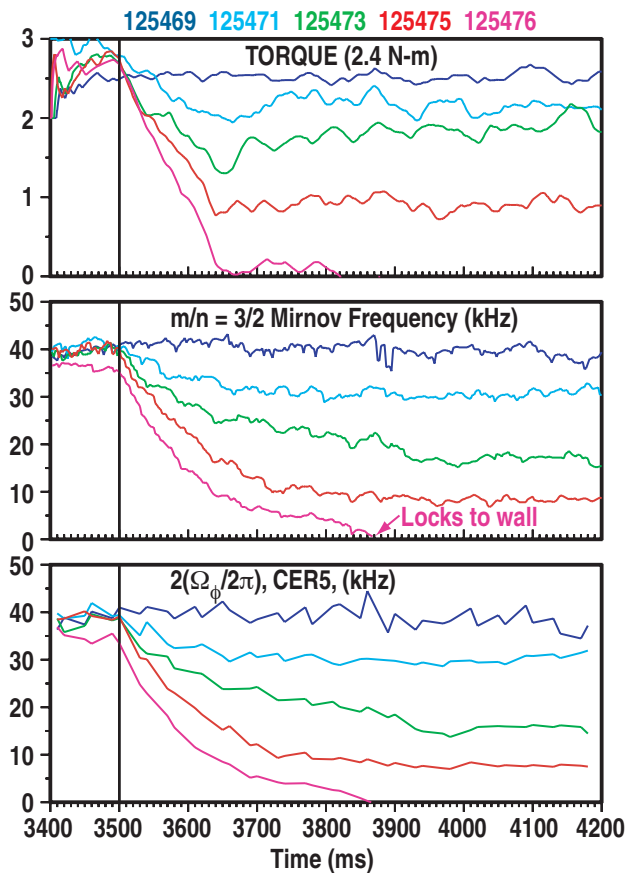


$m/n = 3/2$ Modes Get Larger with Lower Rotation (and Shear) as “Inferred” Δr Less Negative

- All co-NBI has counter beams added ($\beta_N \approx 2.6$, $\bar{n}_{13} \approx 3.5$)
- ★ rotation reduced

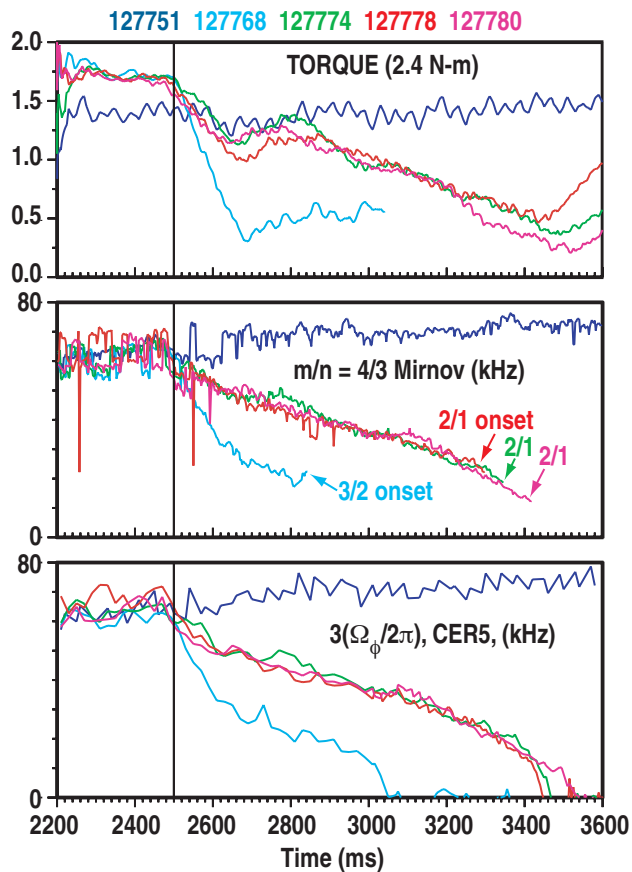
- As flow shear is reduced . . .
- ★ 3/2 amplitude increases

- “Inferred” Δr less negative with reduced positive normalized flow shear

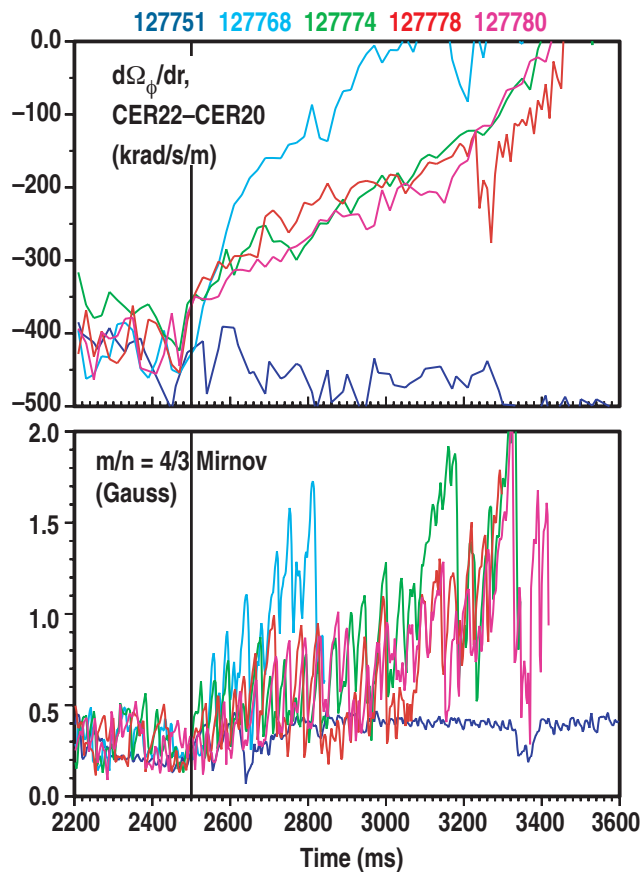


$m/n = 4/3$ Modes Also Larger with Lower Rotation (and Shear) as “Inferred” Δr Less Negative

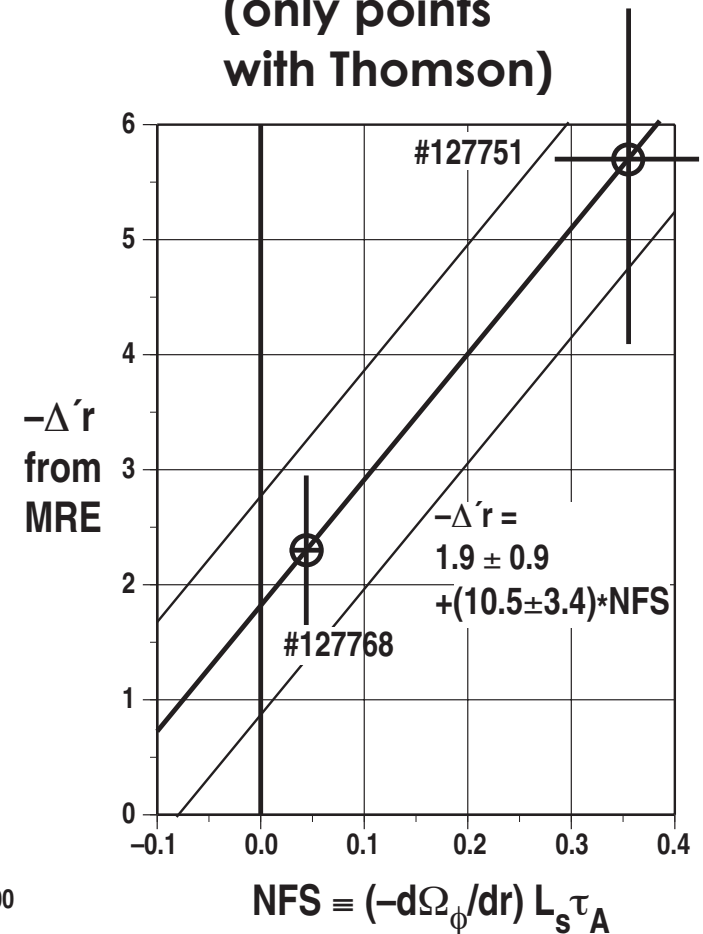
- All co-NBI has counter beams added ($\beta_N \approx 2.7$, $\bar{n}_{13} \approx 3.7$)
- ★ rotation reduced



- As flow shear is reduced . . .
- ★ $4/3$ amplitude increases



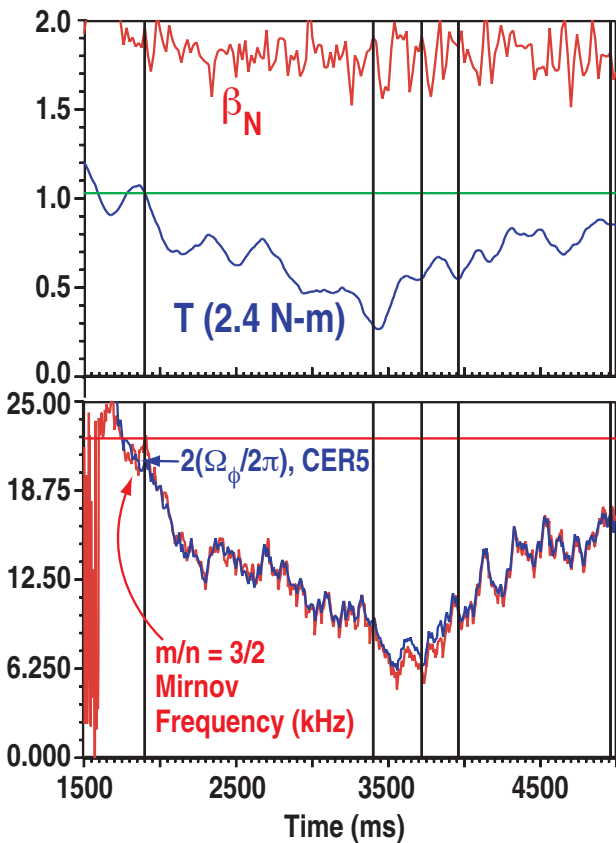
- “Inferred” Δr less negative with reduced flow shear (only points with Thomson)



... and "Inferred" Δ' Less Negative with Inverted Flow Shear with $m/n = 3/2$ Modes

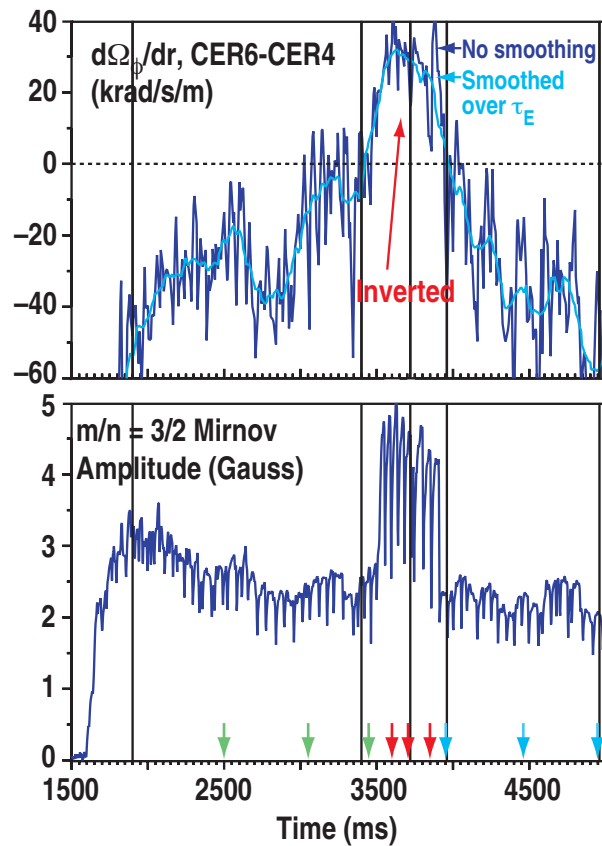
- Reduced β_N , less initial all co-torque (#135867)

★ torque programmed down and back ($\sim 8 \tau_E$ each way)



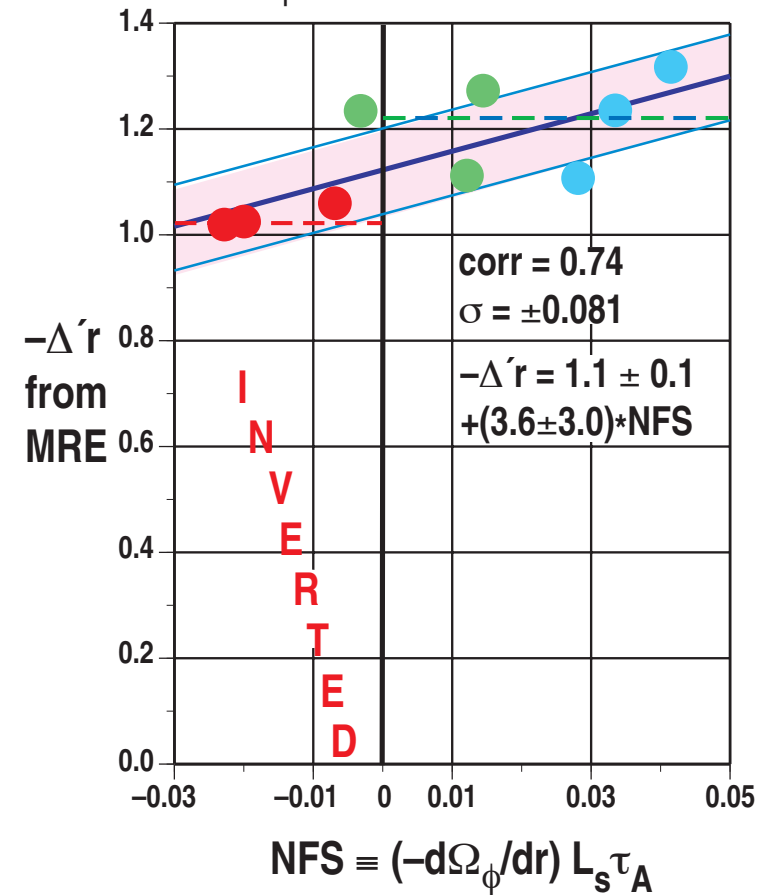
- Little change in $n=2$ Mirnov amplitude

★ until $d\Omega_\phi/dr \approx 0$ "forbidden" state



- "Inferred" Δ' depends on sign of $-d\Omega_\phi/dr$?

★ Ω_ϕ co



Flow Shear Can Be Put Into an Analytic Expression for Δ' for Guiding Comparison to Experiments: I

- C.C. Hegna & J.D. Callen (PoP 1994) simple analytic expression for Δ'

★ large aspect ratio circular cylinder $m \geq 2$ asymptotic solution

$$\dots \Delta' = -(2m/r) \lambda \pi \cot(\lambda \pi) \text{ with } \lambda = - \frac{r q}{m(dq/dr)} \frac{\mu_0}{B_\theta} \left. \frac{dj_{\parallel}}{dr} \right|_{r=r_s} \text{ with } L_q = \frac{q}{dq/dr}$$

– note dj_{\parallel}/dr (negative) destabilizing
and magnetic shear dq/dr (positive) stabilizing

... expanding in λ , $-\Delta' r/2m \approx 1 - (\lambda \pi)^2/2 \Rightarrow 1$ as $\lambda \Rightarrow 0$, $\Rightarrow 0$ as $\lambda \Rightarrow \frac{\sqrt{2}}{\pi} \approx 0.5$

- Add radial magnetic shear and “parallel” projection of toroidal flow shear

$$\star \frac{1}{L_q} \rightarrow \frac{1}{L_q} - \left(\frac{q}{r/R} \right) \frac{d\Omega_\phi}{dr} \tau_A \quad \star \frac{\Delta' r}{-2m} \approx 1 - \frac{\pi^2}{2} \frac{\lambda^2}{1 - 2(d\Omega_\phi/dr)L_s \tau_A + (d\Omega_\phi/dr)^2 L_s^2 \tau_A^2}$$

– L_s arises in parallel connection length along an island of full width w
[$\lambda_{//} = L_s/(k_\theta w/2)$ with $k_\theta = m/r$,
 $L_s = q L_q/(r/R)$ and $L_q = q/dq/dr$]

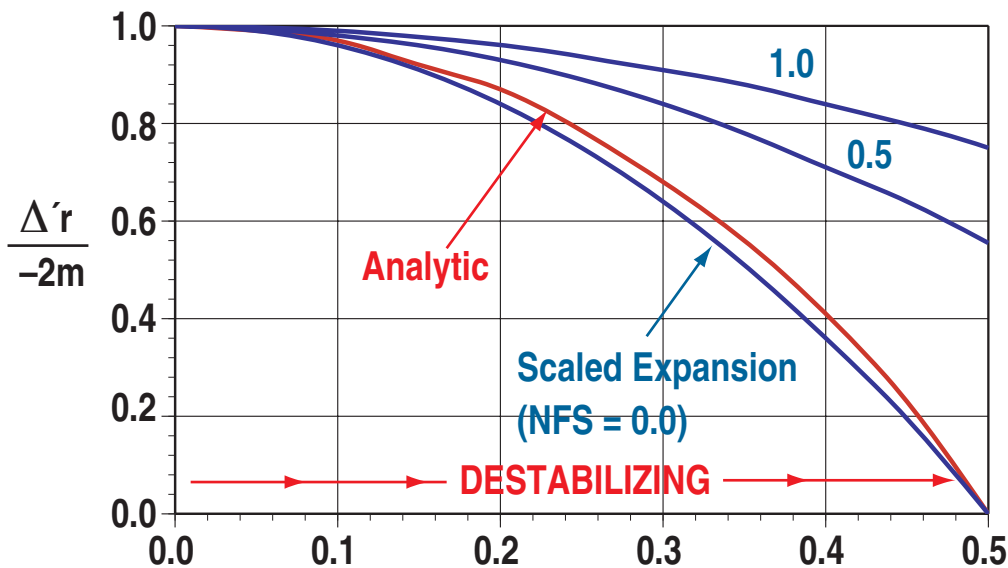
– $\tau_A = R/V_A$ with $V_A = B_z/\sqrt{\mu_0 n_e m_i}$
[resistive MHD is akin to ideal MHD field line bending]

Flow Shear Can Be Put Into an Analytic Expression for Δ' for Guiding Comparison to Experiments: II

- Expansion without flow shear is in good agreement with analytic expression for λ scaled by 0.901 times ($\lambda = 0.0$ to 0.5)

$$\star \frac{\Delta' r}{-2m} \approx 1 - \frac{\pi^2 \lambda^2}{2} + \pi^2 \lambda^2 \left(-\frac{d\Omega_\phi}{dr} L_s \tau_A \right) + \dots \text{ to leading order in NFS}$$

and normalized flow shear NFS $\equiv -\left(\frac{d\Omega_\phi}{dr}\right) L_s \tau_A$



$$\lambda = -\frac{rq}{m(dq/dr)} \frac{\mu_0}{B_\theta} \left. \frac{dj_{||}}{dr} \right|_{r=r_s}$$

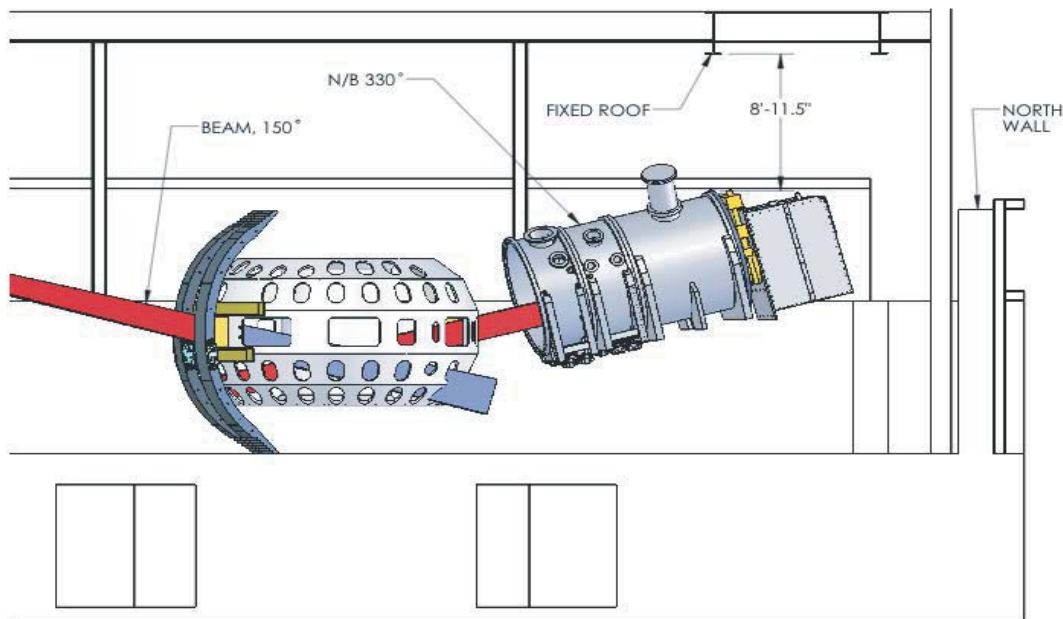
- Flow shear makes stable classical tearing yet more stable
- and insensitive to details in profiles ... i.e., λ

Numerical Modeling is in Process

- NEAR code (D. Chandra, D. Raju, and A. Sen, IPR, Gujarat, India)
 - ★ dominant effect comes from the outer layer modification to Δ'
 - ... scaling of Δ' to velocity shear is quite close to heuristic model
 - clearly shows dependence on sign of flow shear
 - but sign and magnitude in question
- MARS code (M. Chu, GA)
 - ★ MARS has sheared flow but is not time dependent
 - ... computes linear growth rates if unstable
- NIMROD code (D.Brennan, Univ. Tulsa and S. Kruger, Tech-X Corp)
 - ★ recently added single fluid sheared plasma rotation
 - ... computes linear and nonlinear states from initial values

DIII-D 5-Year Plan Has Off-Axis Neutral Beam Injection Allowing Independent Control of Plasma Rotation Profile

NEUTRAL BEAM @ 14.8°
Elevated 330° Beamline, Shield Roof, 14.8° Vertical Rotation Angle



Poloidal projection of fast ion birth points on a single toroidal plane

