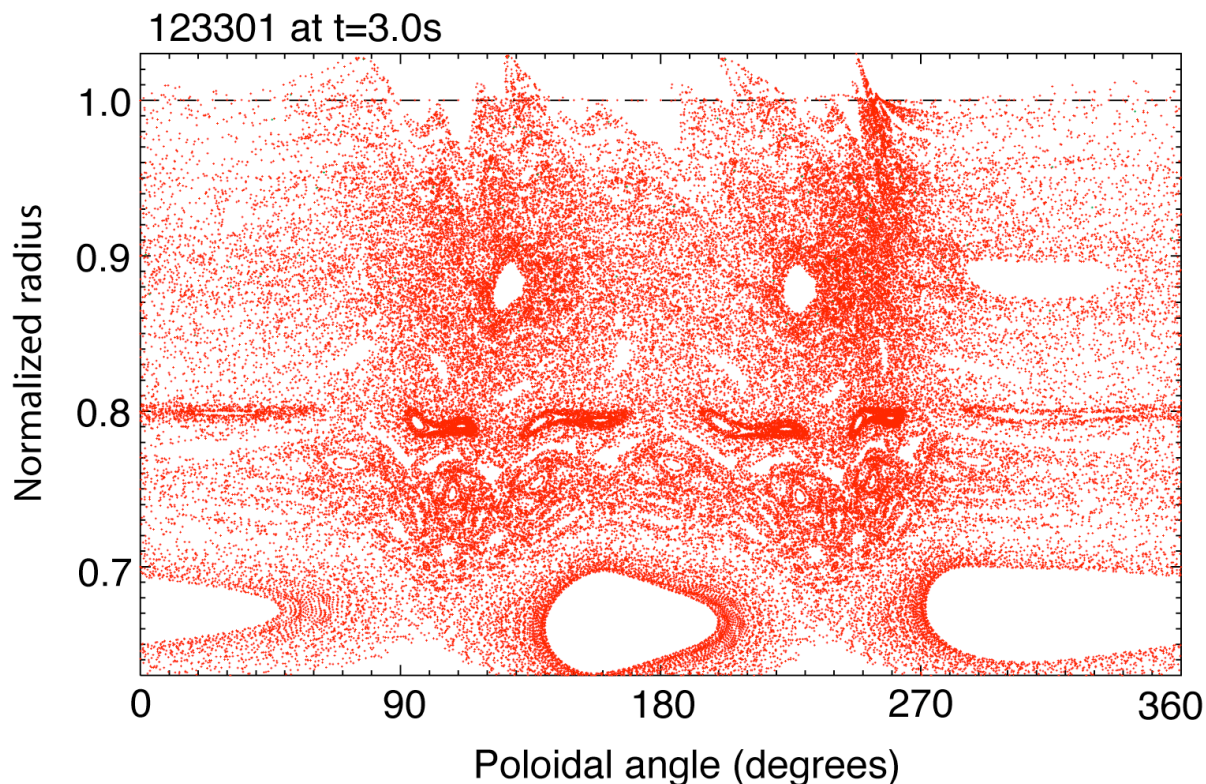


The Physics of Edge Resonant Magnetic Perturbation in Hot Tokamak Plasmas

Presented by
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General Atomics

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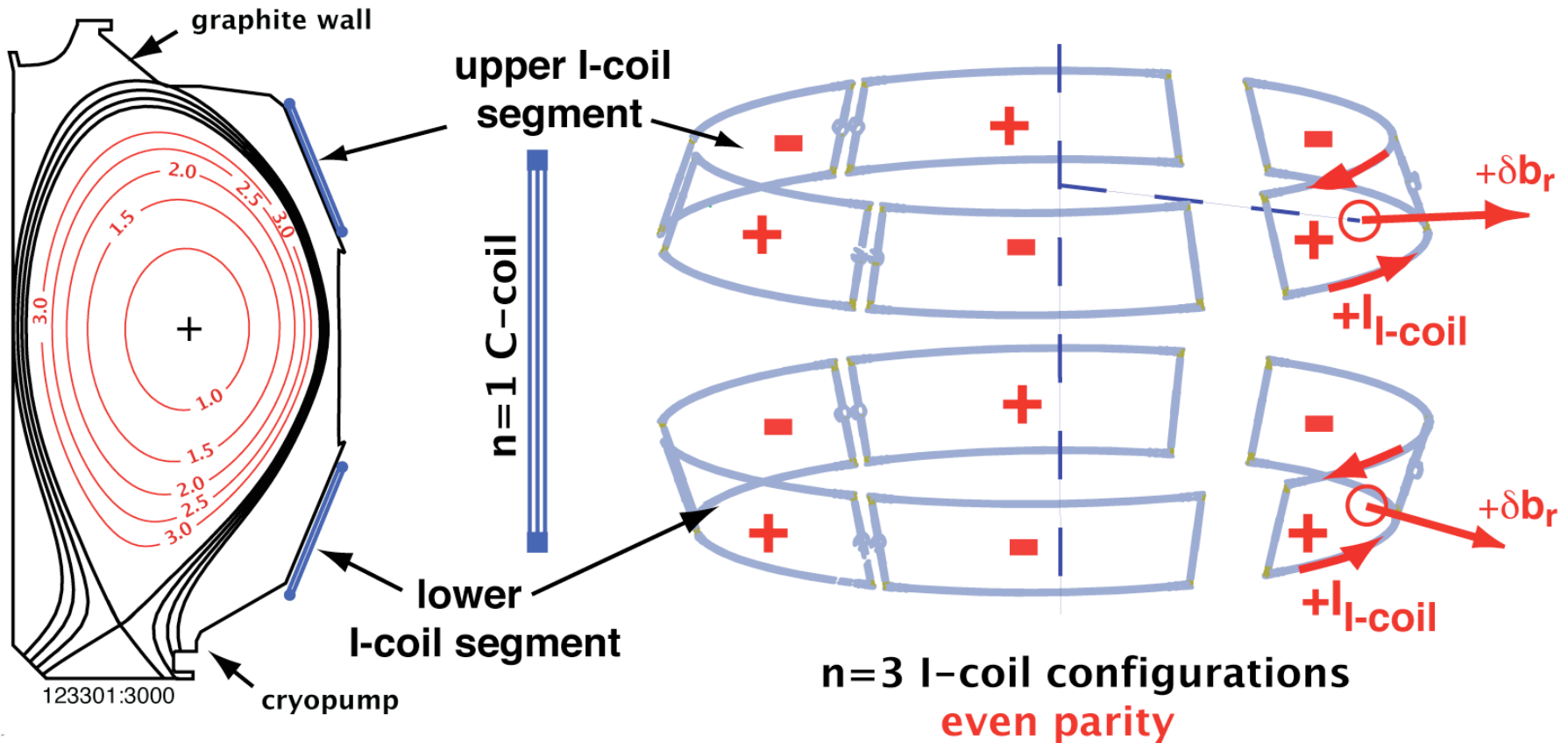
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Edge resonant magnetic perturbations: a promising approach for pedestal control in burning plasmas

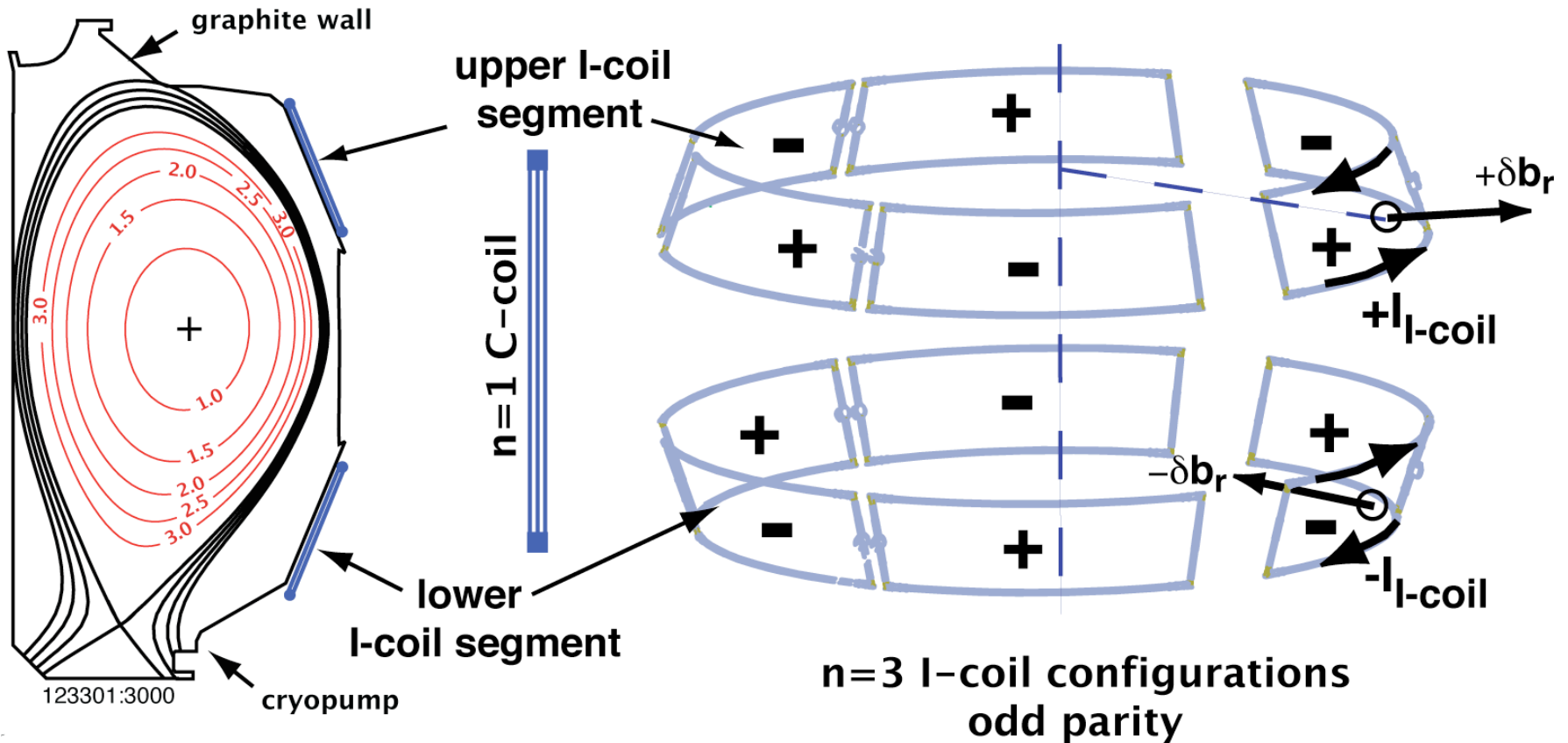
- Small edge Resonant Magnetic Perturbations (*RMPs*) used to:
 - > Control pedestal profiles in high confinement plasmas
 - ∇p control \rightarrow edge bootstrap current (j_{edge}) control?
- Type-I ELMs completely eliminated in low collisionality (ν_e^*) burning plasma relevant conditions:
 - > Consistent with peeling-ballooning ∇p , j_{edge} stabilization for all cases tested to date
 - > ∇p (j_{edge} ?) operating point controlled with *RMP* coil current
- ∇p changes primarily due to increased particle rather than energy transport:
 - > Challenges stochastic transport theory
 - > Suggests particle convection dominates stochastic open field thermal conduction in low ν_e^* pedestal plasmas

The DIII-D I-coil produces a variety of edge localized resonant magnetic perturbations (RMPs)



- Flexible control of poloidal (m) spectrum with $n=3$, **even** parity

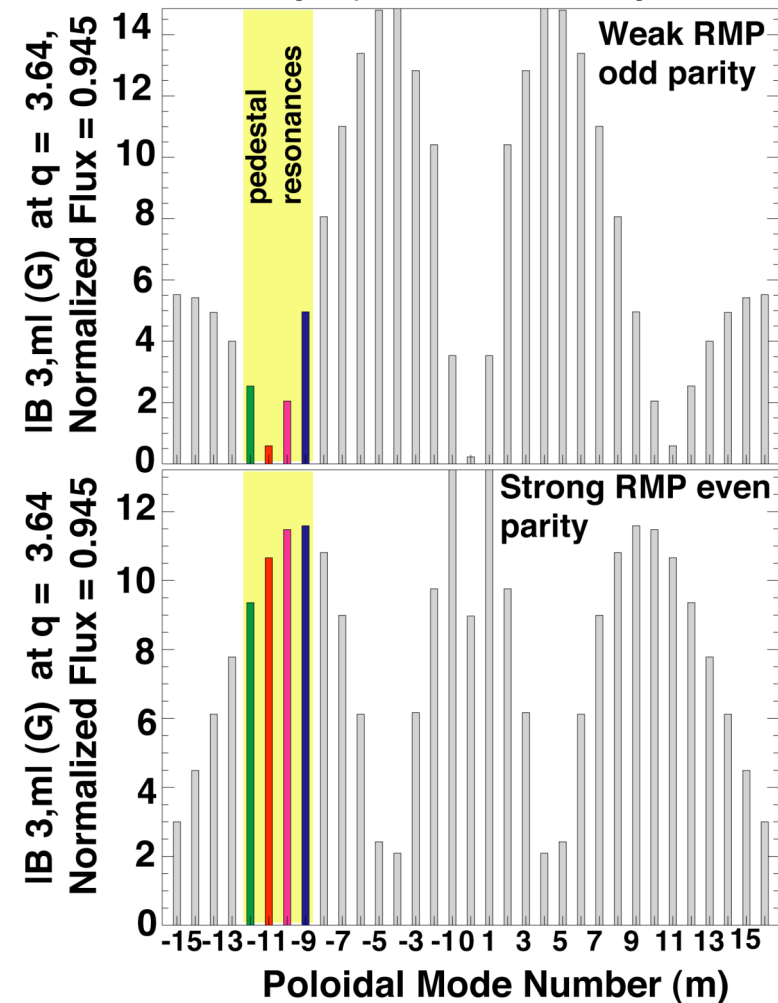
The DIII-D I-coil produces a variety of edge localized resonant magnetic perturbations (RMPs)



- Flexible control of poloidal (m) spectrum with $n=3$, **even** and **odd** parity

I-coil parity sets maximum pedestal *RMP* amplitude

ITER Shape (119690@2600, $q_{95}=3.72$)



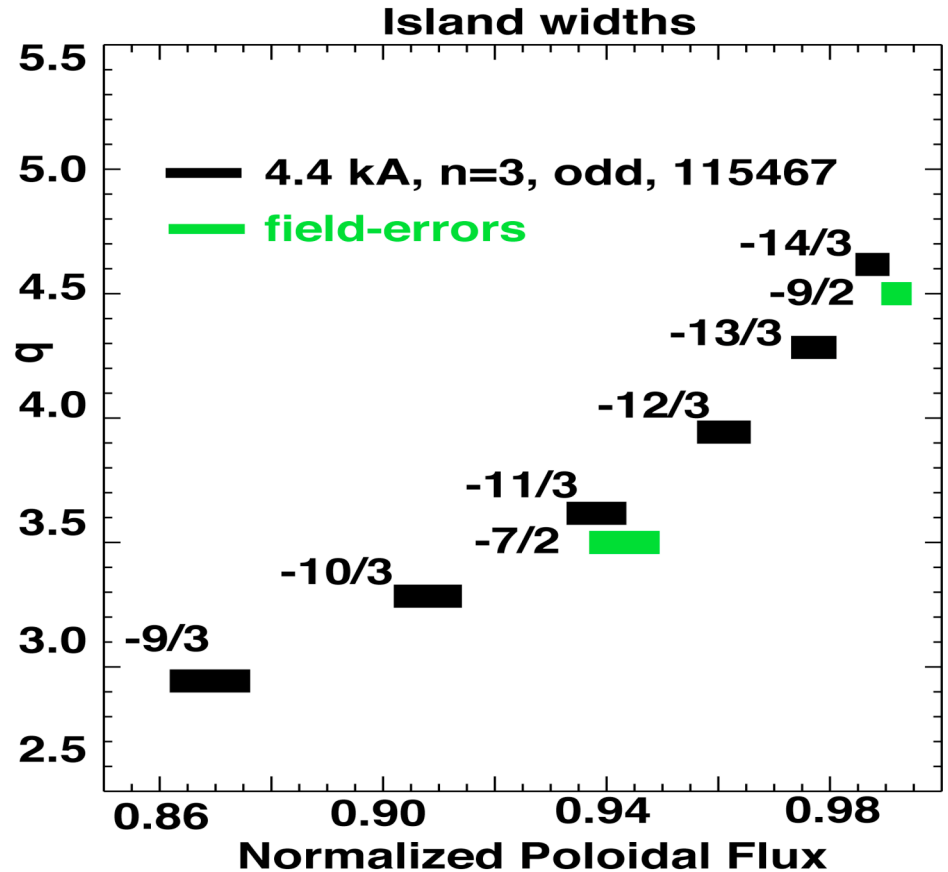
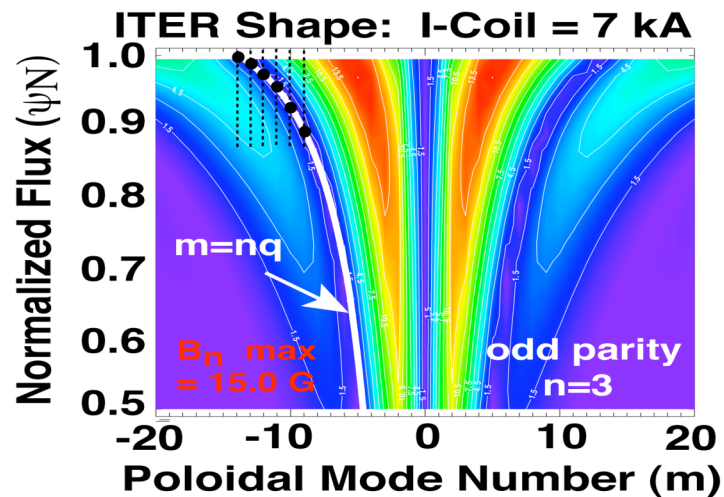
2004 Results:

odd parity → weak edge *RMP*
 $\delta b_r/B_T \sim 2 \times 10^{-5} \sim$ size of field-errors
ELMs suppressed by increasing fluctuations
little or no profile changes
high collisionality (ν_e^*)

New 2005 Results:

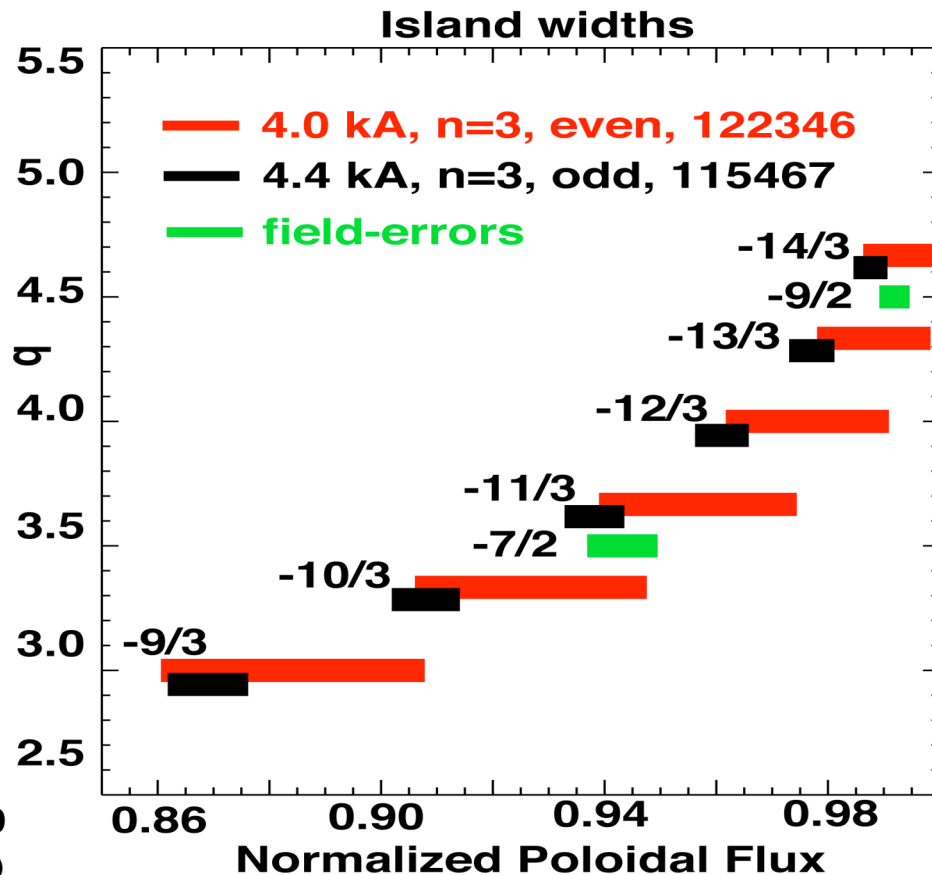
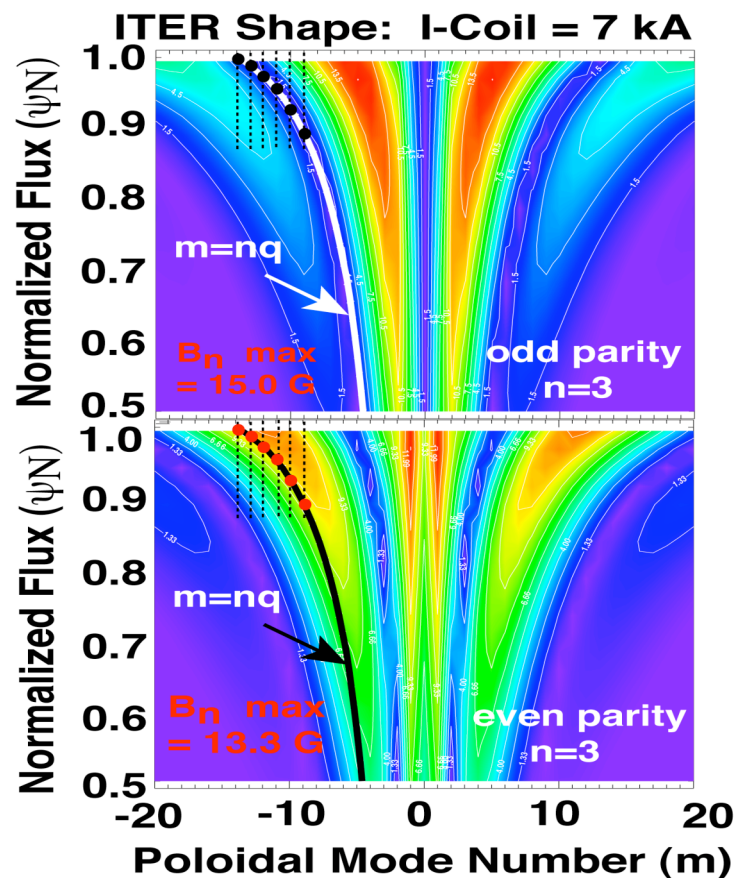
even parity → strong edge *RMP*
 $\delta b_r/B_T \sim 3 \times 10^{-4} \sim 10\text{-}20\times$ size of field-errors
→ ELMs eliminated by controlling ∇p
low collisionality (ν_e^*)
first burning plasma relevant experiments

I-coil parity controls pedestal island overlap



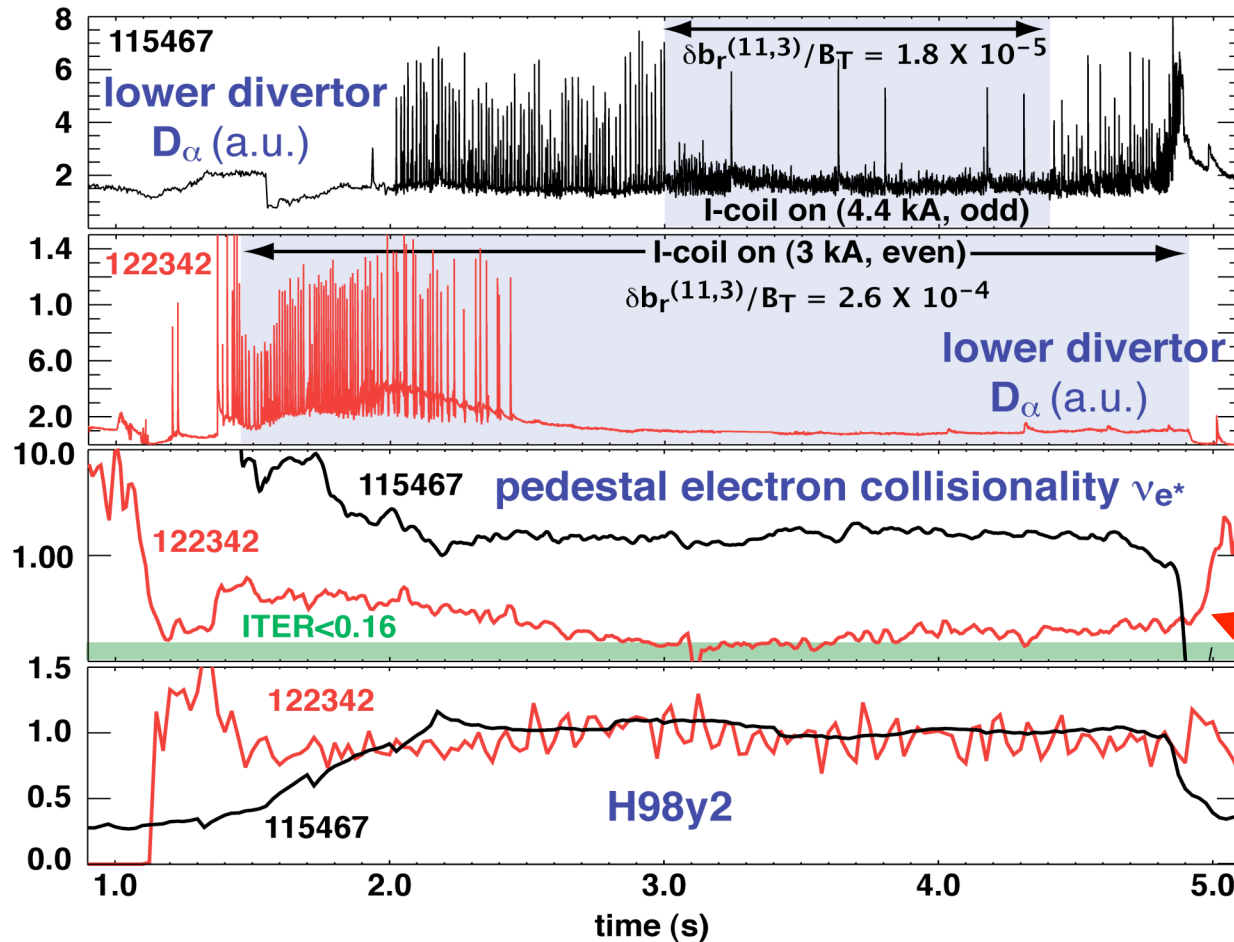
- Both parities suppress ELMs
 - > Odd (weak RMP) \rightarrow small islands \rightarrow little or no change in pedestal

I-coil parity controls pedestal island overlap



- Both parities suppress ELMs
 - > Odd (weak RMP) \rightarrow small islands \rightarrow little or no change in pedestal
 - > Even (strong RMP) \rightarrow stochastic \rightarrow transport / pedestal control

Strong RMP configuration results in long quiescent ELM-free H-modes at ITER ν_e^* operating point



2004: weak edge RMP (odd parity) results -
 → Some intermittent ELMs remain

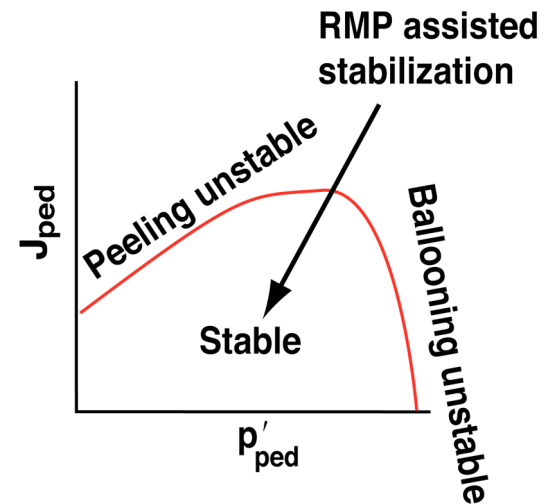
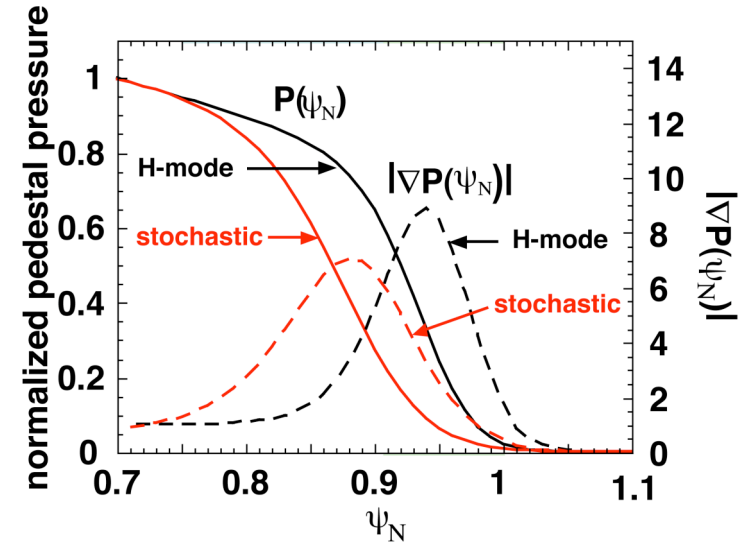
New 2005: strong edge RMP (even parity) results
 → ELMs completely eliminated at low ν_e^*

ν_e^* at burning plasma target in strong (even parity) edge RMP case

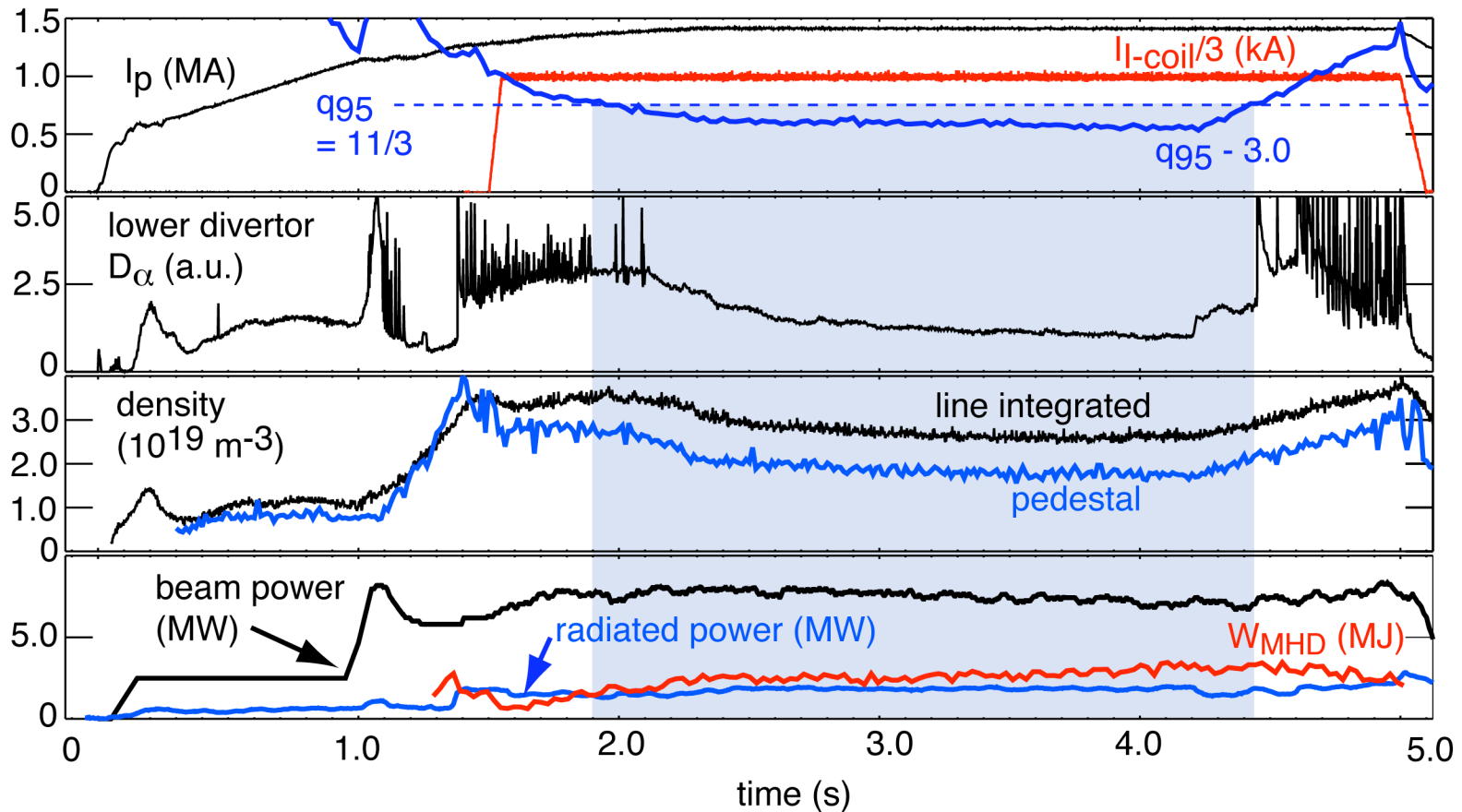
- With $\nu_e^* < 0.2$ ELMs are eliminated for 2.6 s ($\sim 17\tau_E$) - limited only by hardware constraints

Original ELM control concept: $RMP \rightarrow$ stochasticity \rightarrow increased energy transport \rightarrow ELM stabilization

- Edge $RMP \rightarrow$ stochastic magnetic field across pedestal:
 - > Increased electron energy transport \rightarrow reduced pedestal $T_e \rightarrow$ reduced pedestal pressure gradient
- Reduced pedestal pressure gradient \rightarrow stable peeling-ballooning operating point
 - > Operating point controlled with RMP amplitude
 - > Maintain good H-mode confinement (wider T_e pedestal: comparable height)
- ELM impulses eliminated

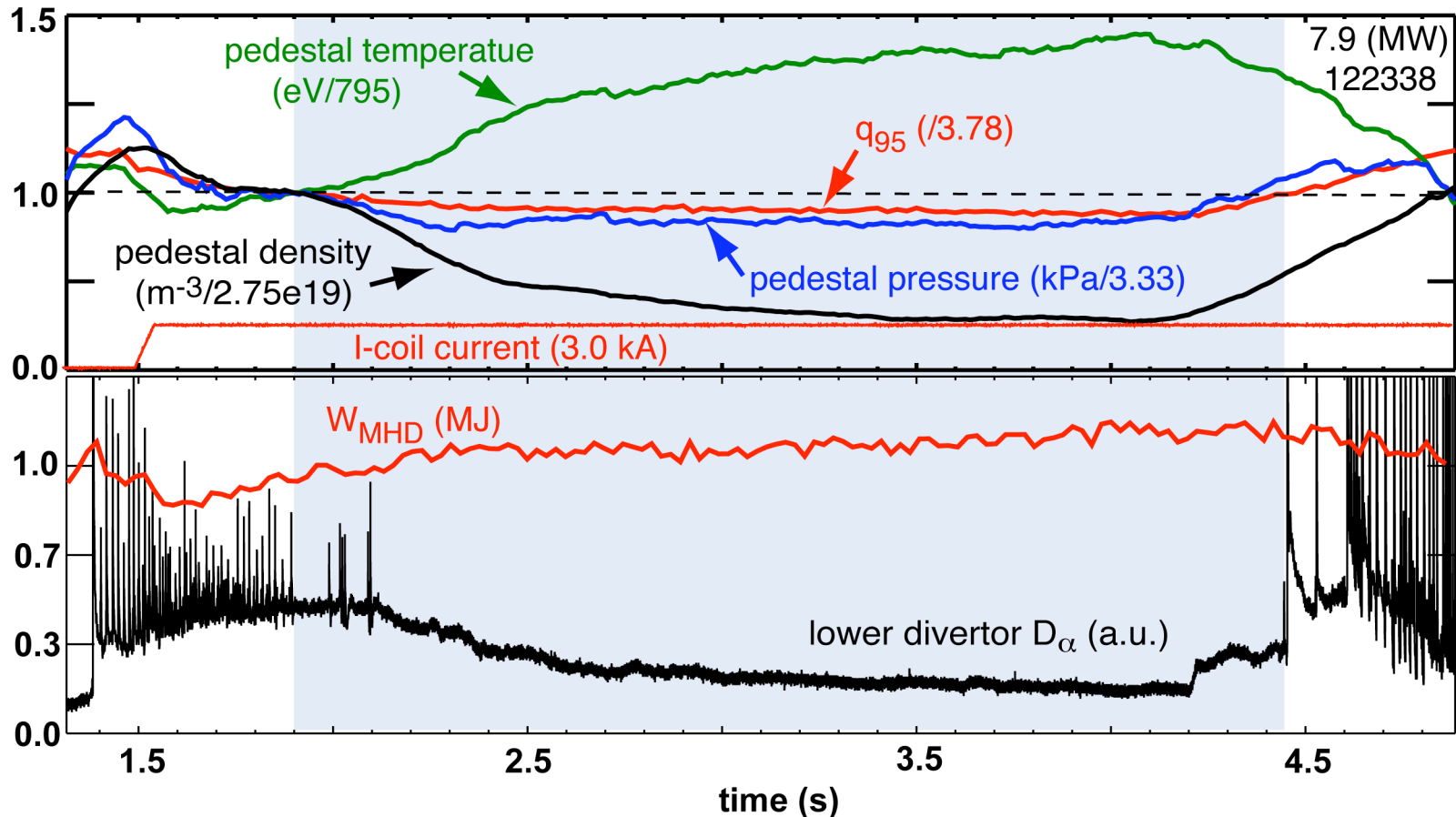


I-coil RMPs have the largest effect on pedestal and ELMs inside resonant window $11/3 \leq q_{95} \leq 7/2$



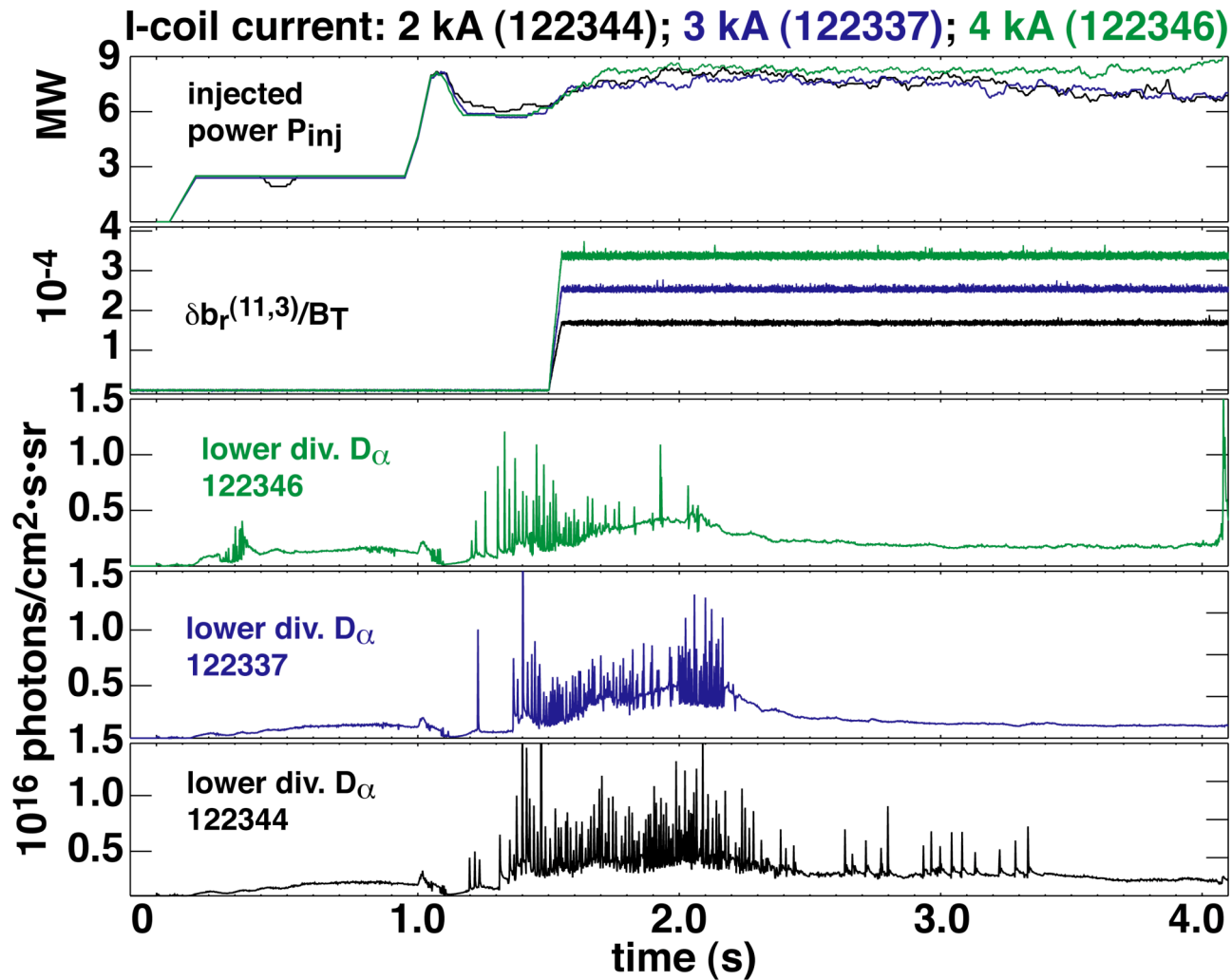
- NBI heating power, total radiated power and stored energy (W_{MHD}) remain relatively constant inside q_{95} resonant window

$T_{e,ped}$ increase with $n_{e,ped}$ decrease contradicts original expectation but ELMs are eliminated

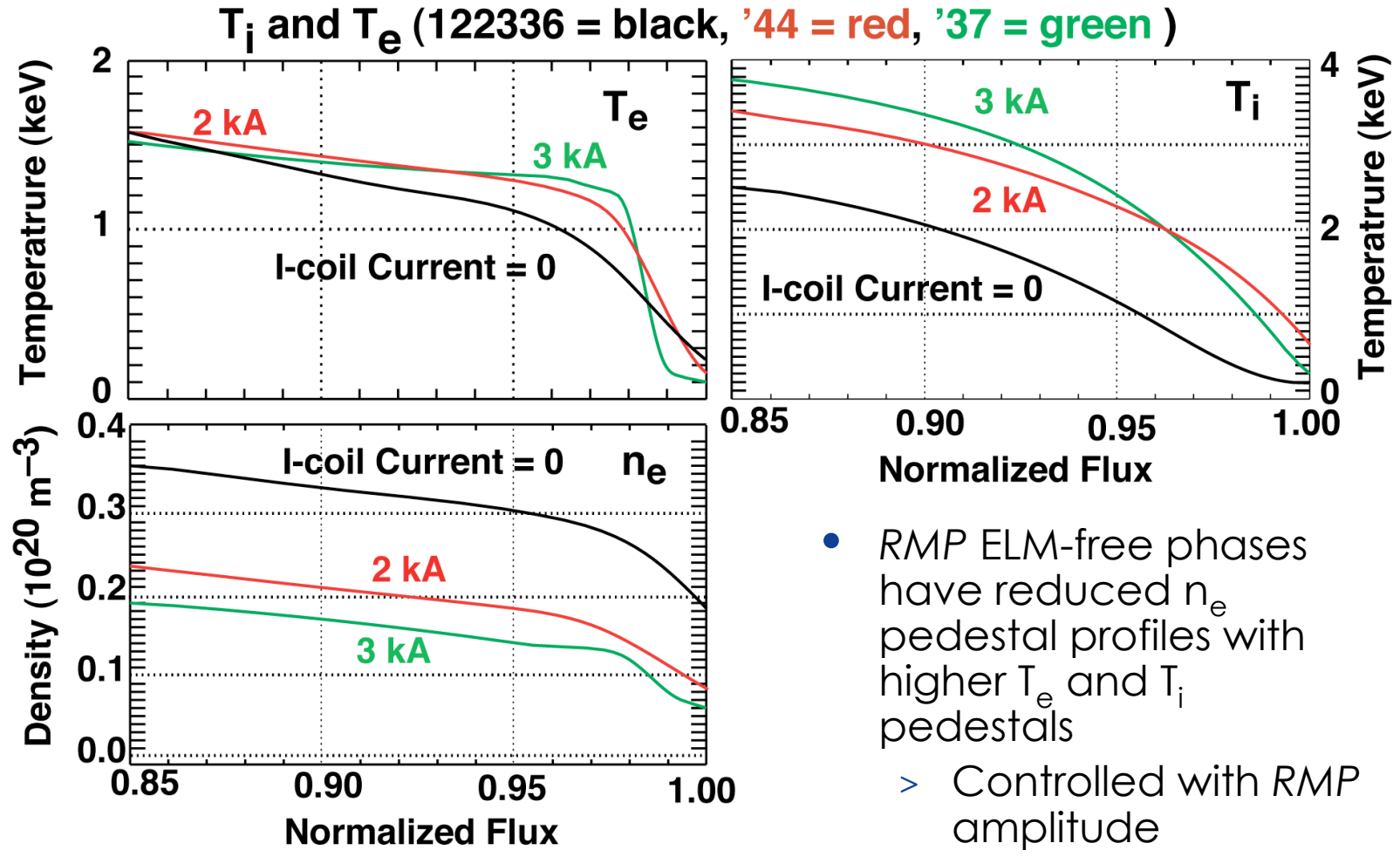


- Energy confinement, stored energy (W_{MHD}), relatively unaffected by RMP
 - > ELMs suppressed and particle confinement reduced when $q_{95} \leq 3.78$

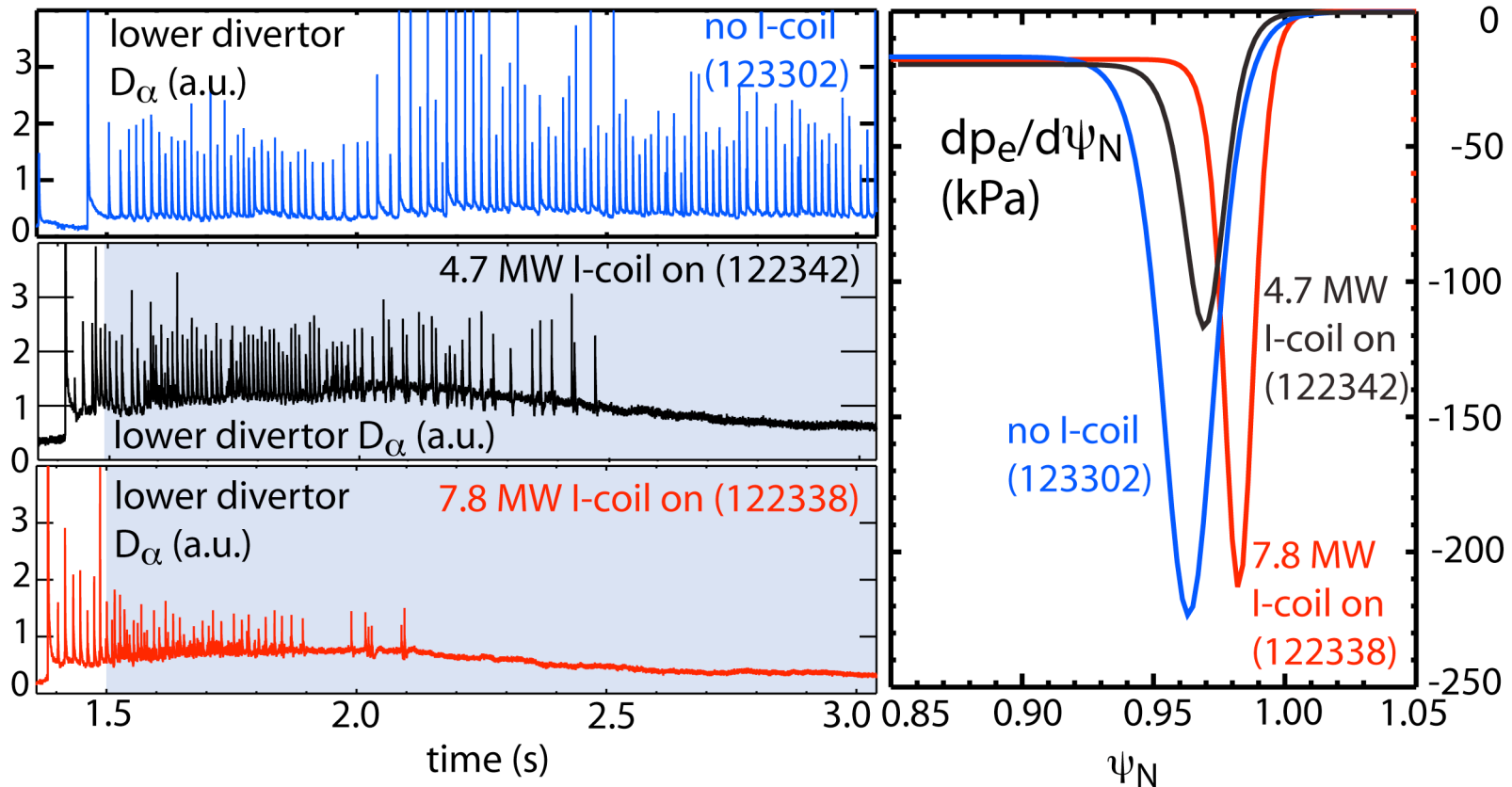
ELMs completely eliminated above RMP amplitude threshold



RMP controls pedestal without destroying H-mode

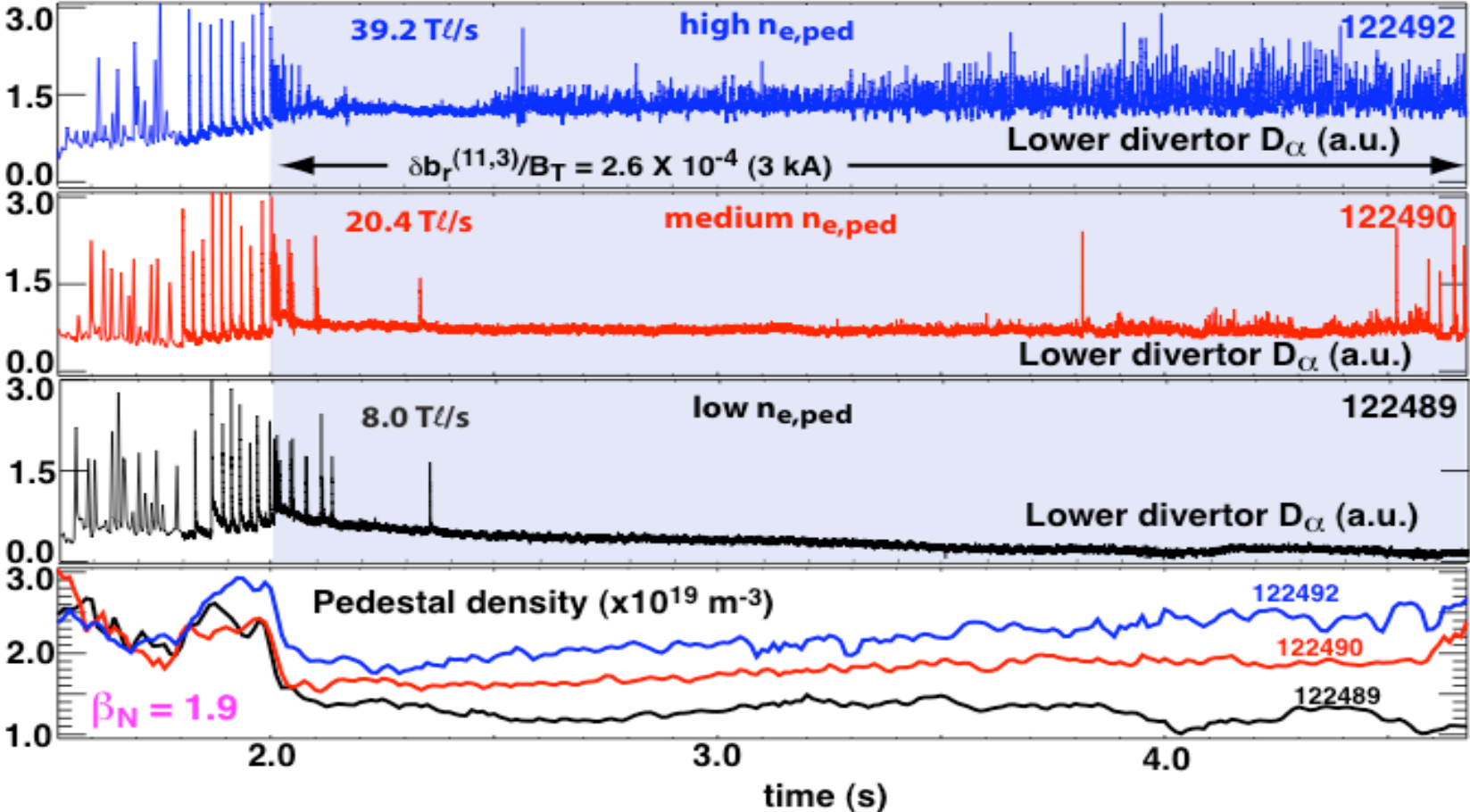


The amplitude, width and radial position of the pedestal ∇p_e are controlled with the RMP and P_{NBI}



- Complete ELM suppression is obtained with ∇p_e comparable to that in ELMing plasmas but narrower and shifted outward

ELMs eliminated below a critical pedestal density

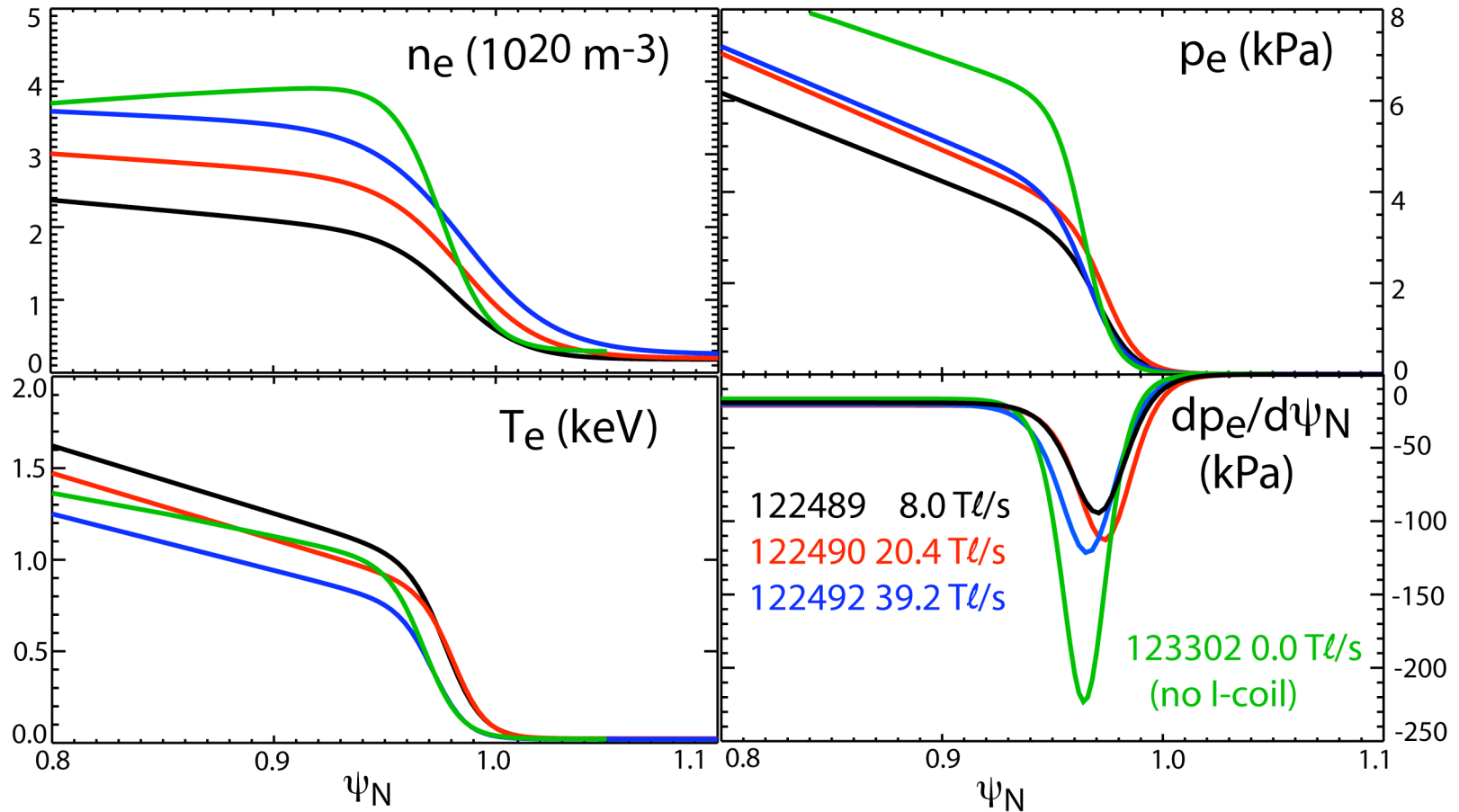


- Small, high frequency ELMs return as pedestal n_e increases
 - > Similar to previous weak RMP (odd parity) with high n_e and v_e^*



Even parity, low v_e^* , high, medium and low $n_{e,ped}$
 See: CP1.00008 - J. Watkins, et al.

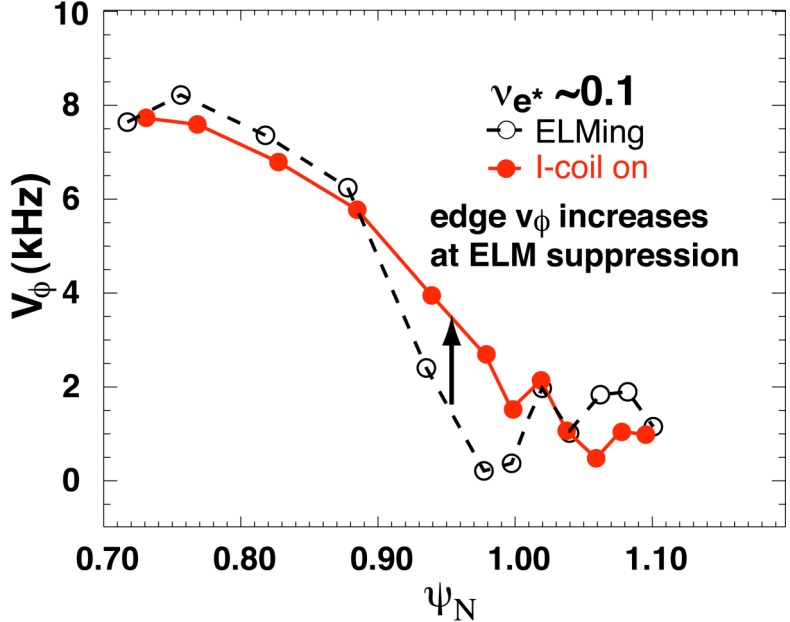
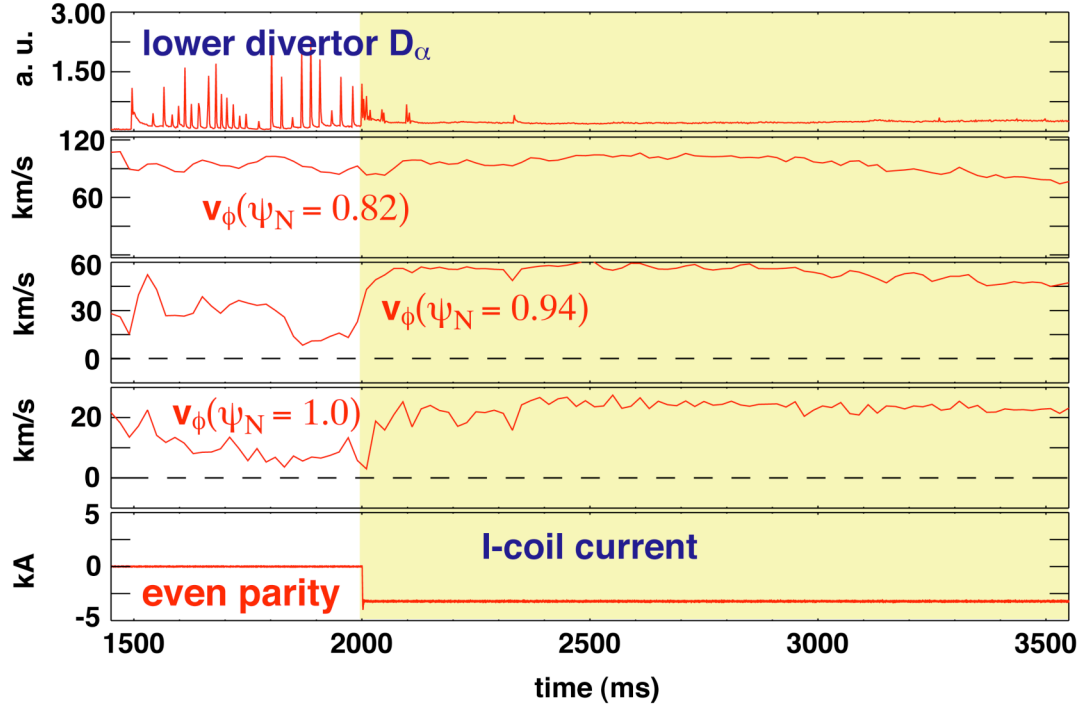
$n_{e,ped}$ increases with increased fueling rate while $\nabla n_e, \nabla T_e, \nabla p_e$ remain approximately constant



- $n_{e,ped}$ increases while $T_{e,ped}$ decreases $\rightarrow v_e^* \sim \text{constant}$

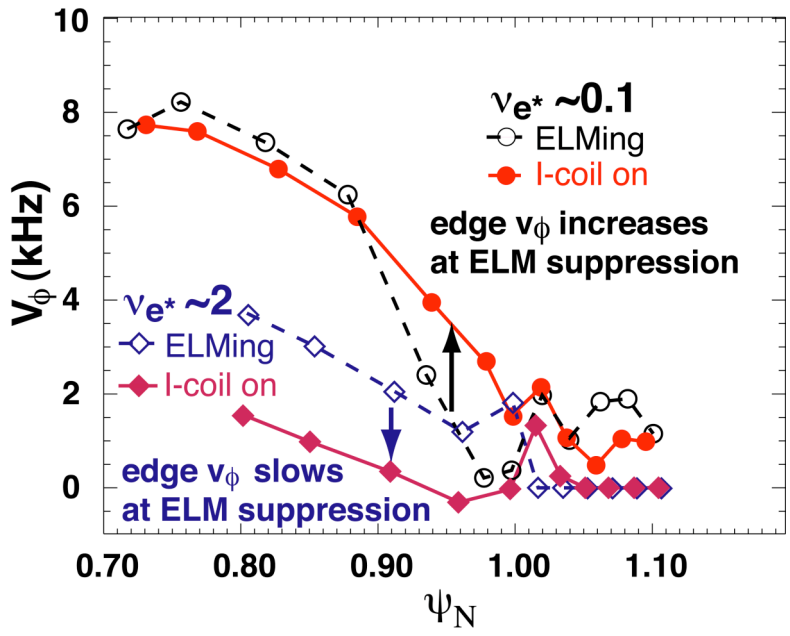
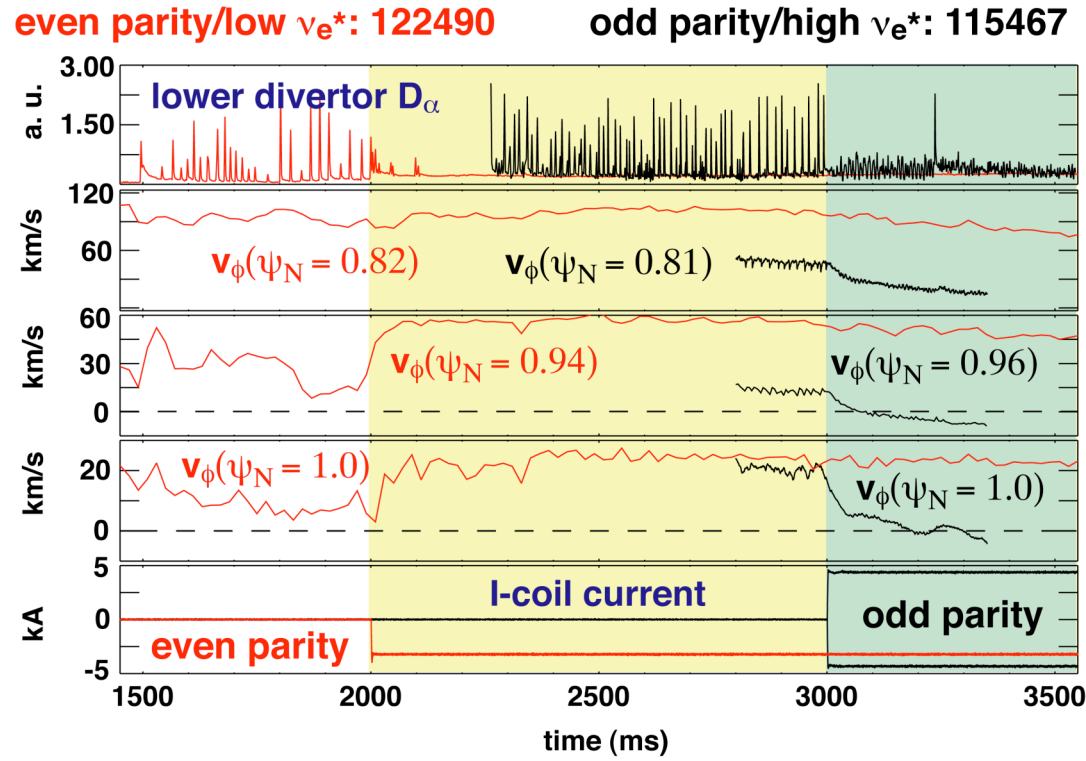
Edge rotation increases in low v_{e^*} RMP ELM-free phase

even parity/low v_{e^*} : 122490



- Performance preserved in low v_{e^*} (high RMP) cases

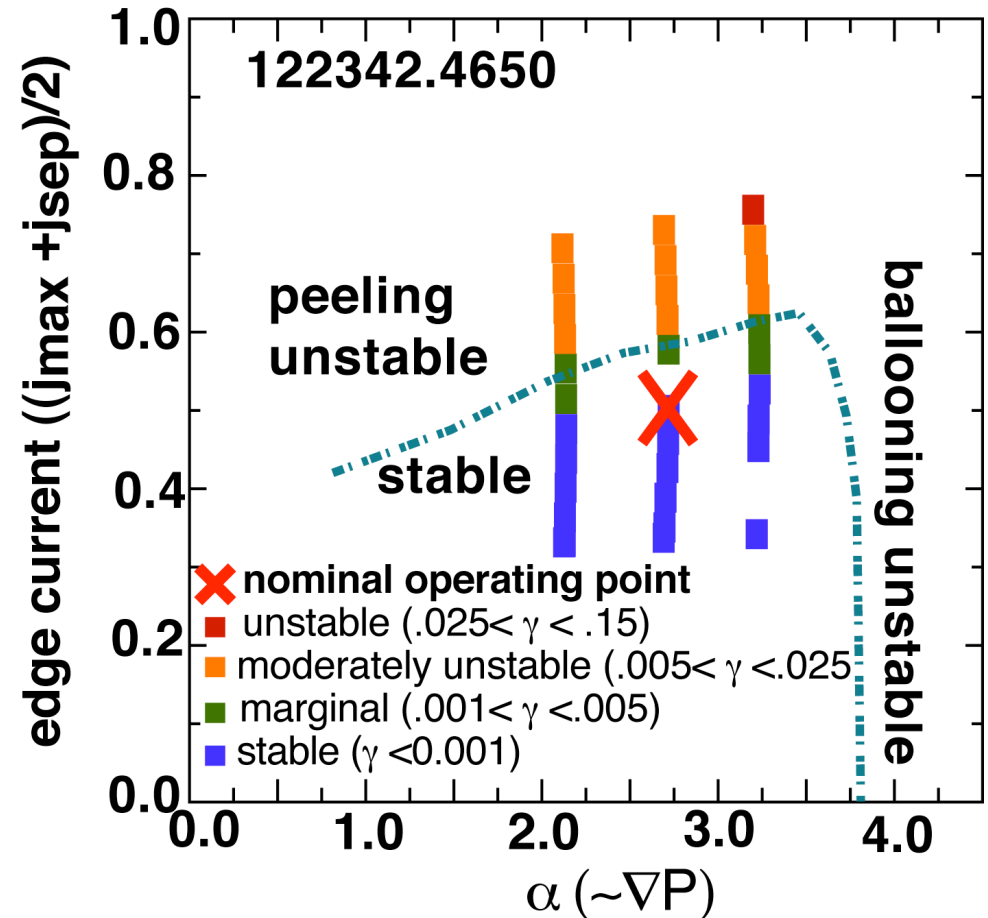
Edge rotation increases in low ν_{e^*} RMP ELM-free phase



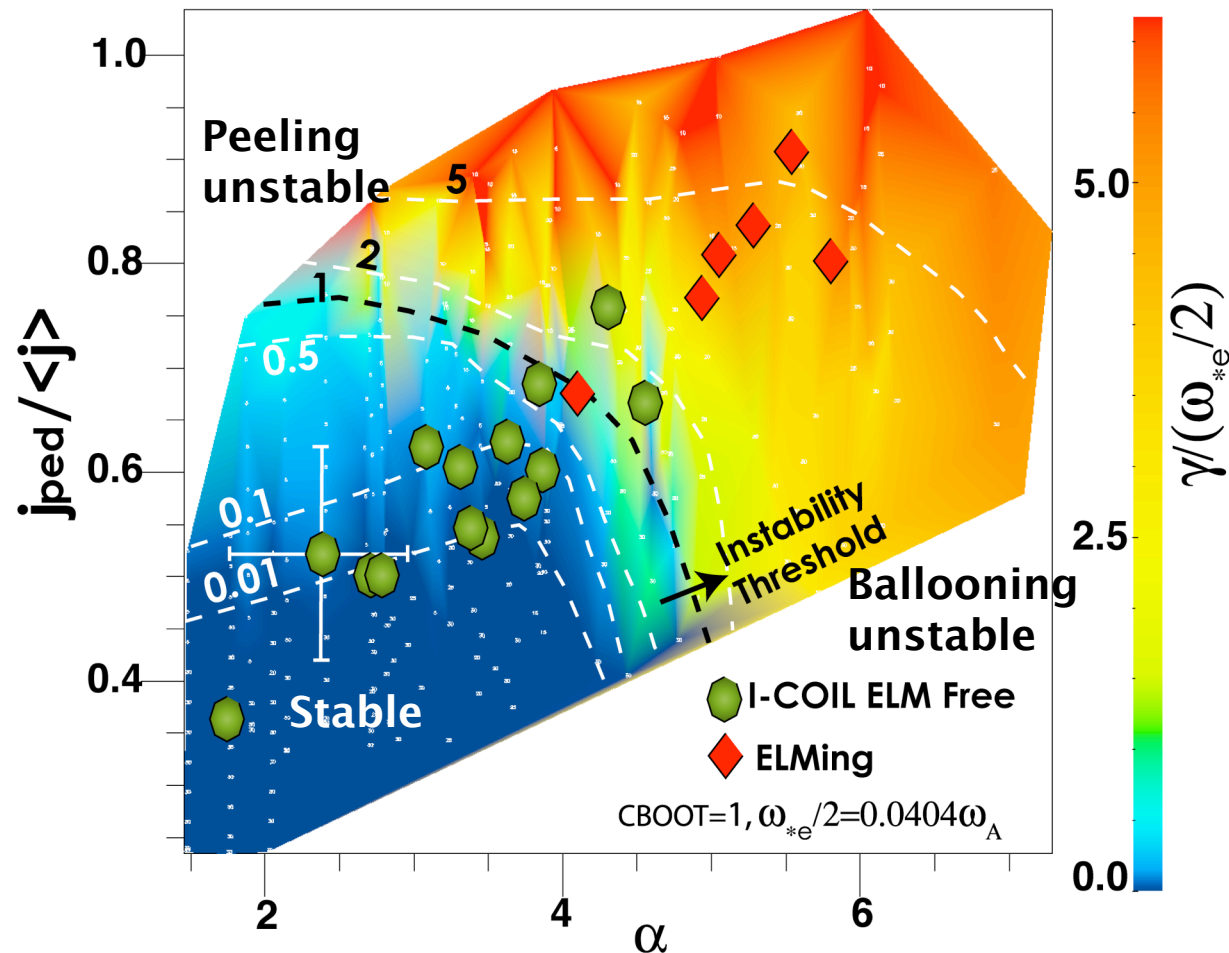
- Performance preserved in low ν_{e^*} (high RMP) cases
- Previous high ν_{e^*} (weak RMP, odd parity) cases had a large decrease in rotation

RMP ELM-free H-modes can be operated well away from the ELM instability boundary

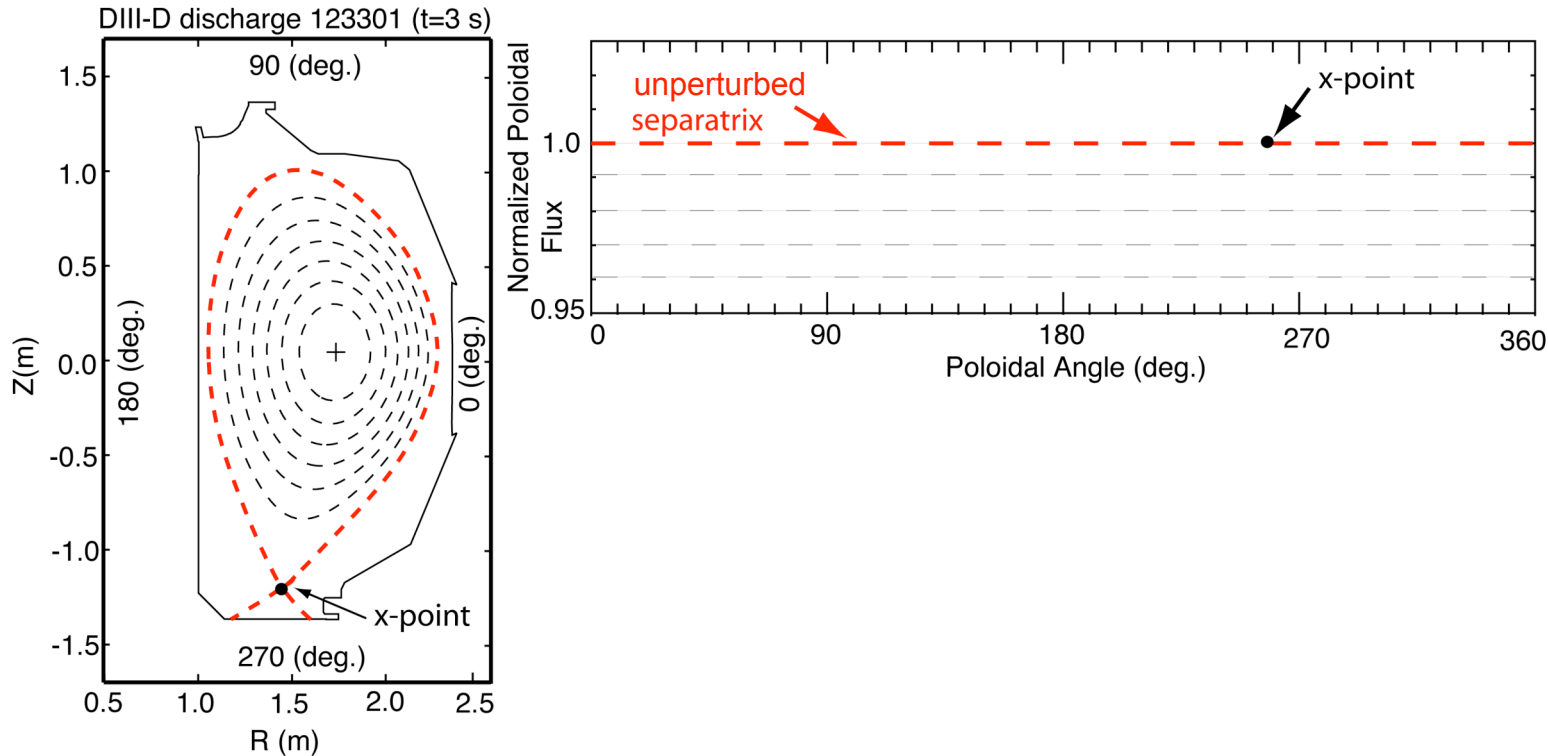
- P-B Model: ELMs triggered by pedestal pressure gradient and current driven MHD modes
 - > ELMing discharges go unstable across the ballooning or coupled P-B regions
- Stable RMP ELM-free point **X** well inside both the peeling and ballooning instability boundaries



RMP ELM-free H-modes can be pushed deep into stable region by increasing the RMP amplitude



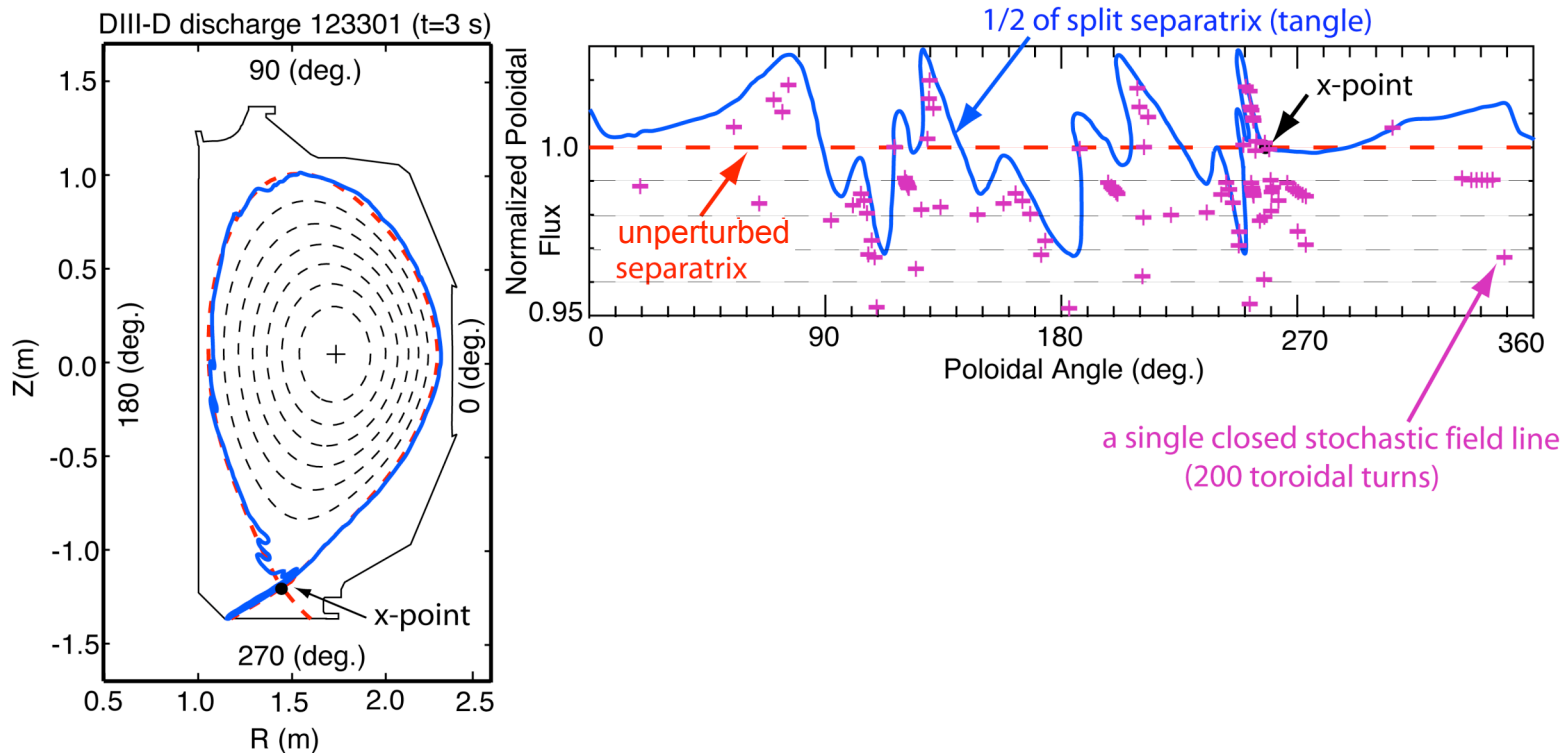
Small edge RMPs significantly alter the magnetic topology of the pedestal



TRIP3D code results:
no plasma response

- Ideal axisymmetric separatrix and flux surfaces are smooth

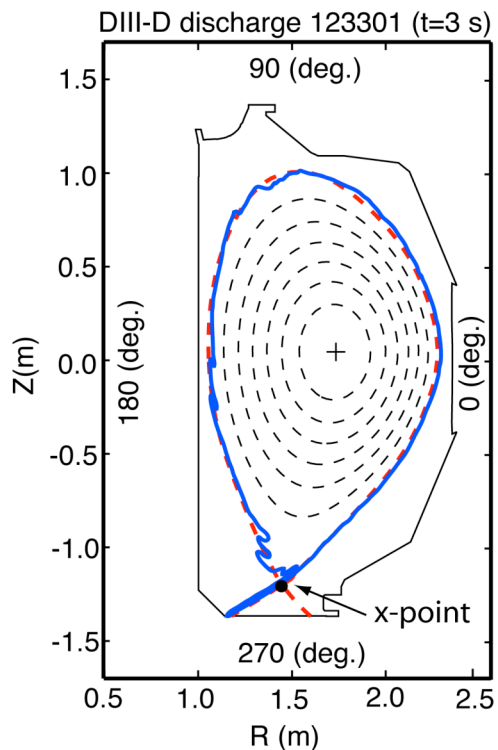
Small edge RMPs significantly alter the magnetic topology of the pedestal



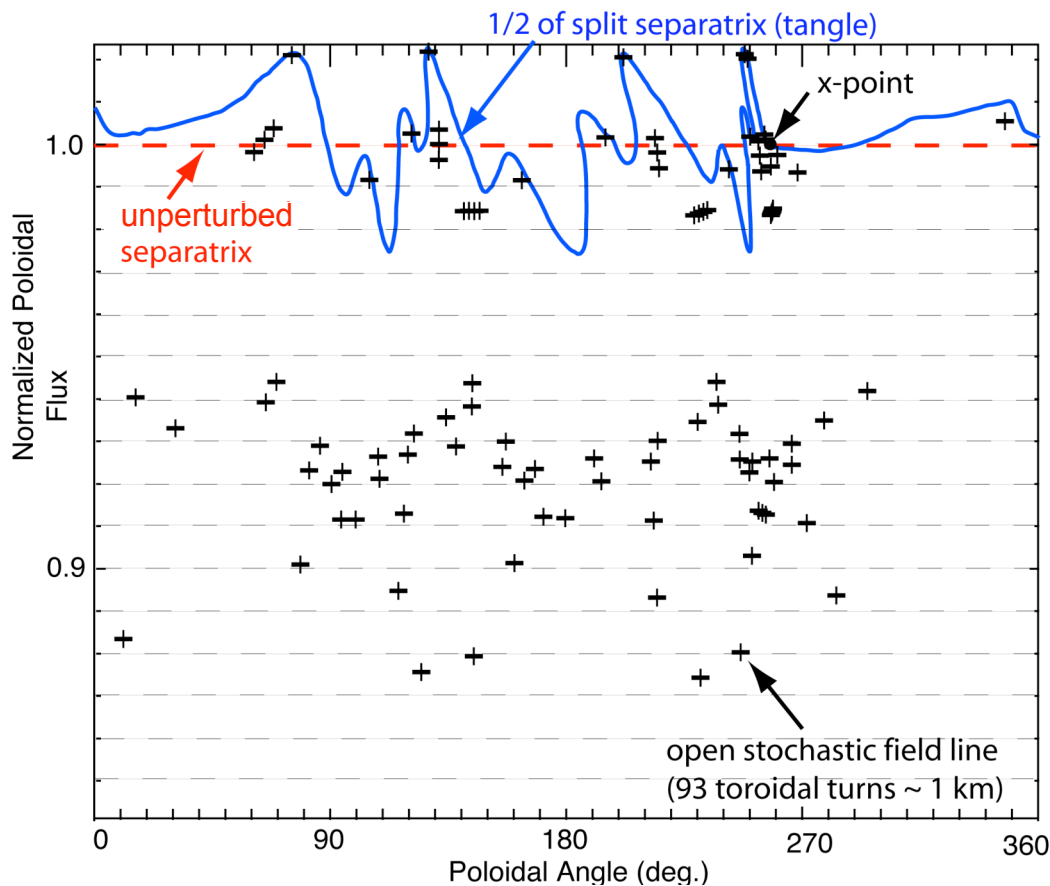
TRIP3D code results:
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- Ideal axisymmetric separatrix and flux surfaces are smooth
- Non-axisymmetric perturbation split the separatrix → homoclinic tangle

Small edge RMPs significantly alter the magnetic topology of the pedestal



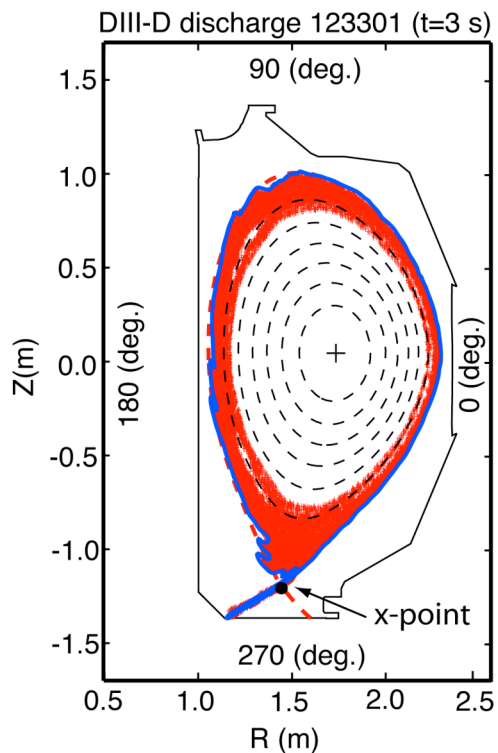
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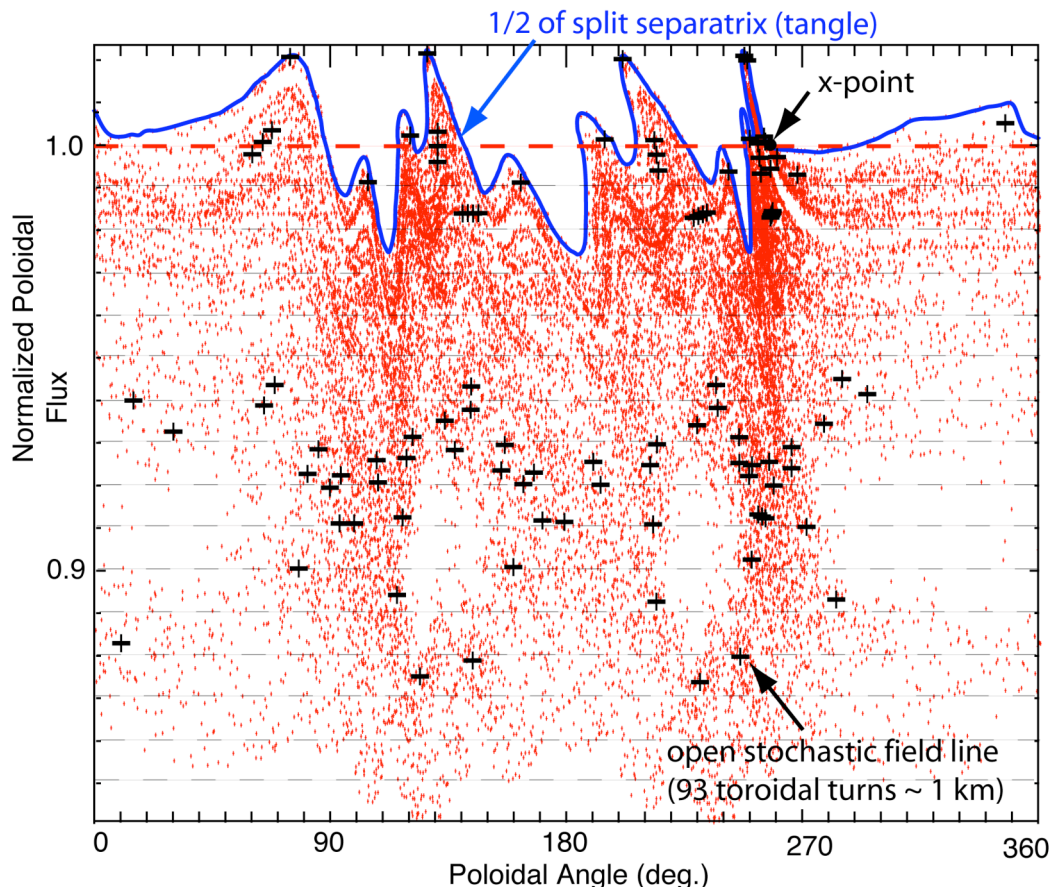
Open stochastic field lines connect to divertor through tangle at 240 deg.

- Ideal axisymmetric separatrix and flux surfaces are smooth
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- Stochastic field lines cross the pedestal and hit solid surfaces

Small edge RMPs significantly alter the magnetic topology of the pedestal



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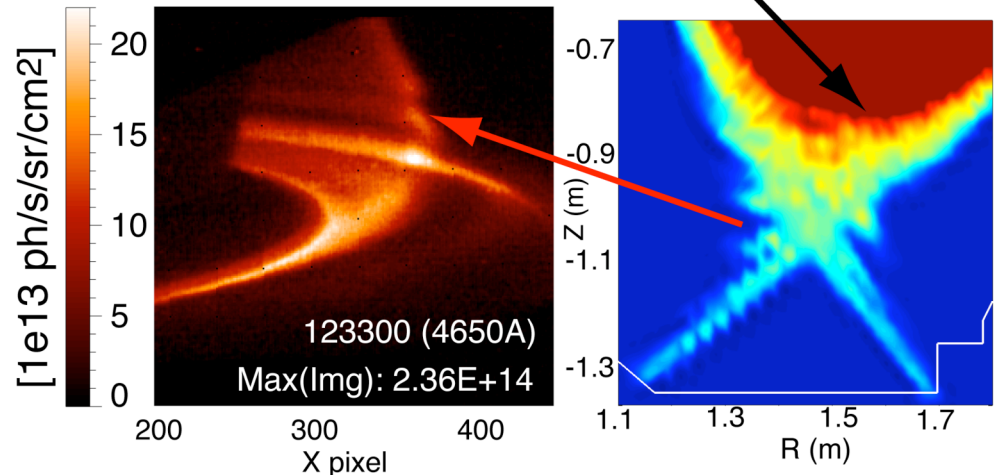
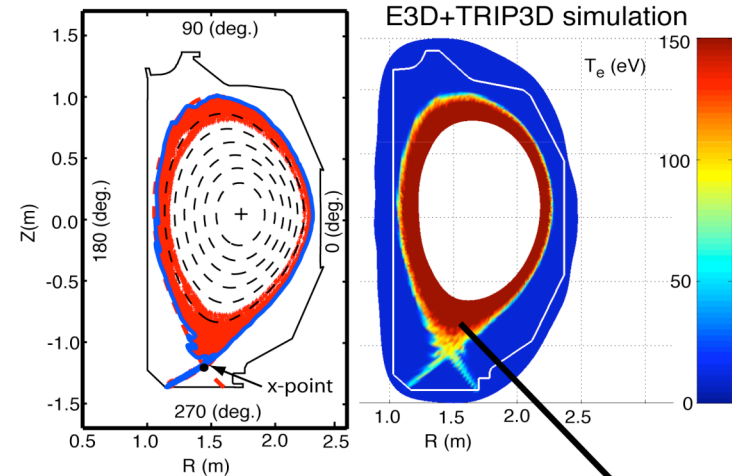
Open stochastic field lines connect to divertor through tangle at 240 deg.

Open stochastic field lines circle islands: changes island potential and $E_r E_\theta \rightarrow$ increase convective particle transport?

- Ideal axisymmetric separatrix and flux surfaces are smooth
- Non-axisymmetric perturbation split the separatrix \rightarrow homoclinic tangle
- Stochastic field lines cross the pedestal and hit solid surfaces - strong mixing

3D energy transport modeling with E3D+TRIP3D codes shows heating of x-point tangle structure

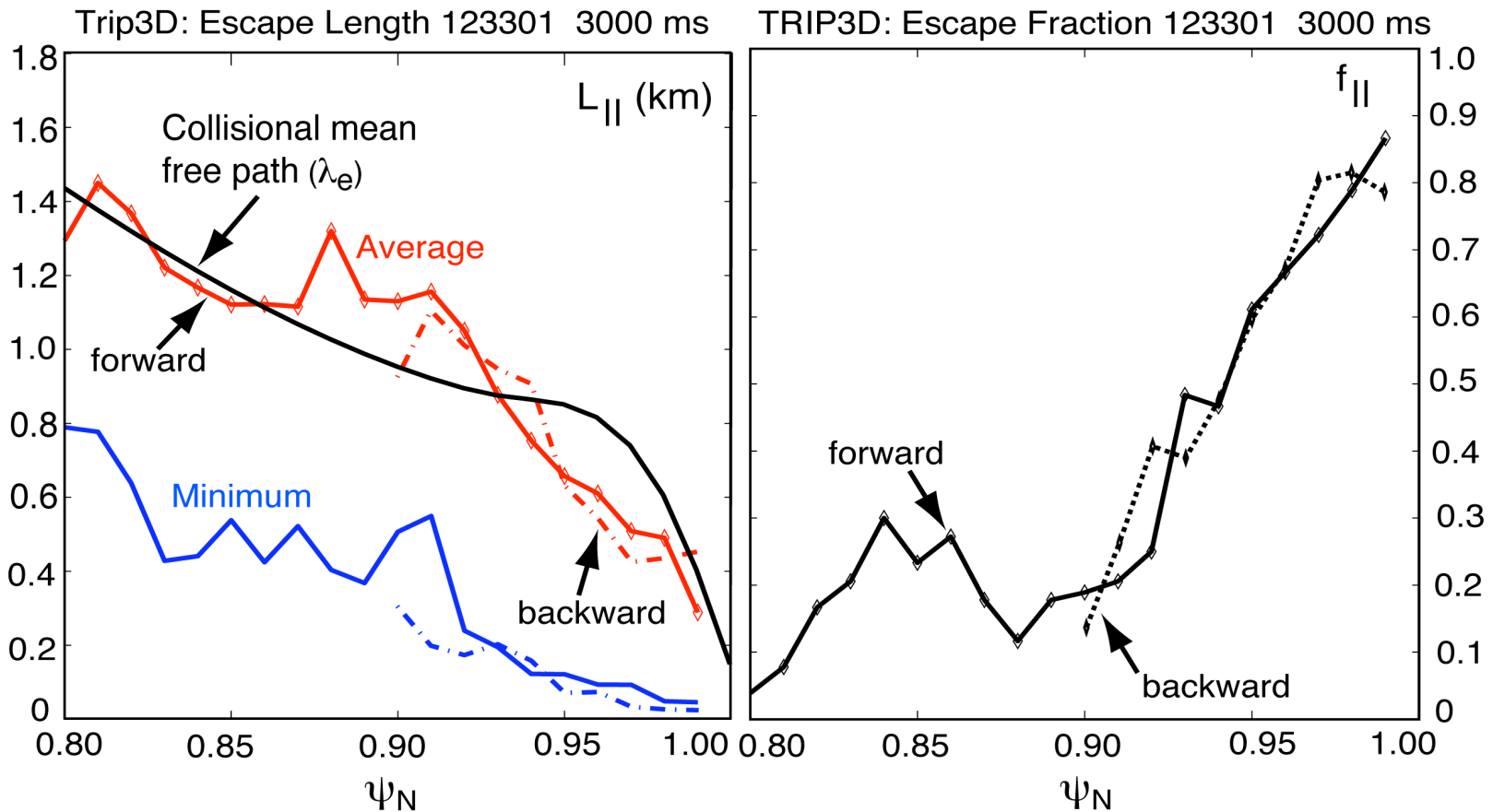
- Non-axisymmetric (homoclinic) tangle appears as a filament-like structure in 2D image
- E3D+TRIP3D heat transport simulations reproduce temperature distribution consistent with observed X-point carbon emission



E3D+TRIP3D heat transport results:
 A. Runov, R. Schneider (MPI Greifswald), S. Kasilov (Kharkov IPT) and I. Joseph (UCSD)
 - see CP1.00006 I. Joseph, et al.,

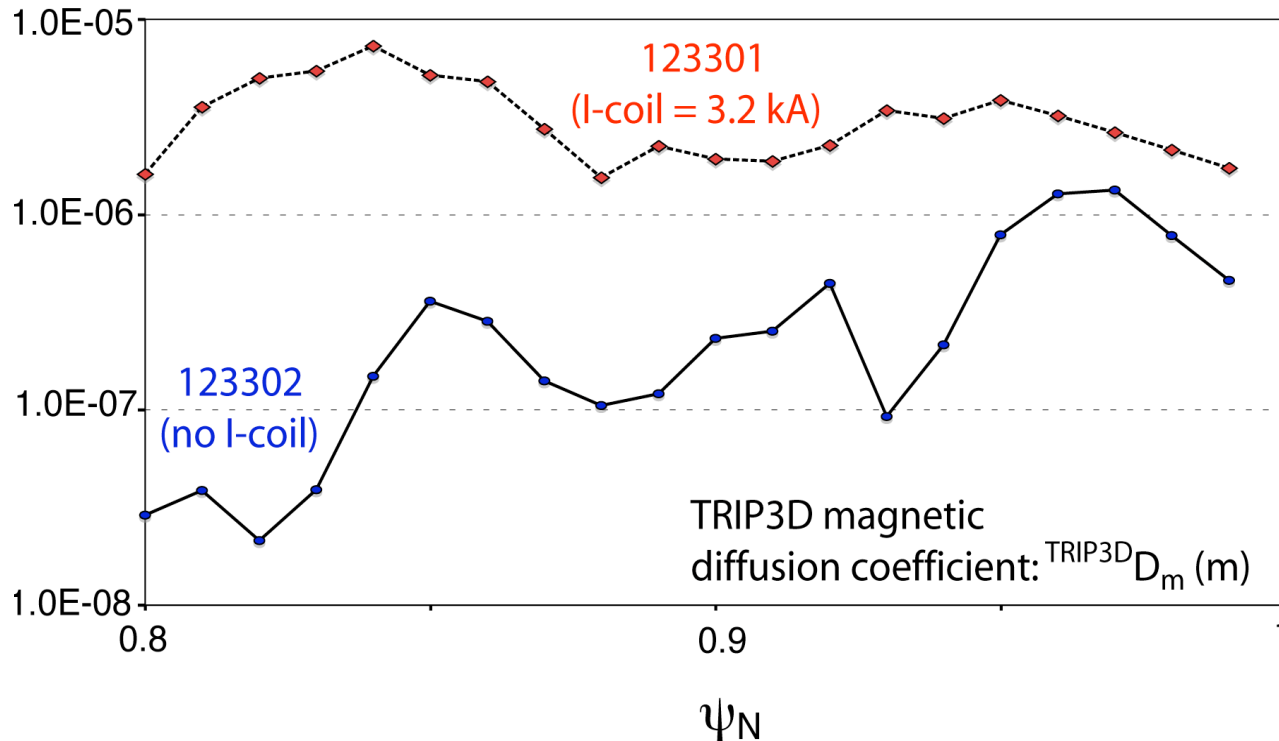
X-point carbon images:
 - see CP1.00003 - M. Fenstermacher, et al.,

The field line escape fraction increases rapidly with ψ_N outside 0.9



- A significant number of escaping field lines have lengths exceeding the electron collisional mean free path length λ_e

Calculated field line diffusion implies thermal diffusivity that is two orders of magnitude too large

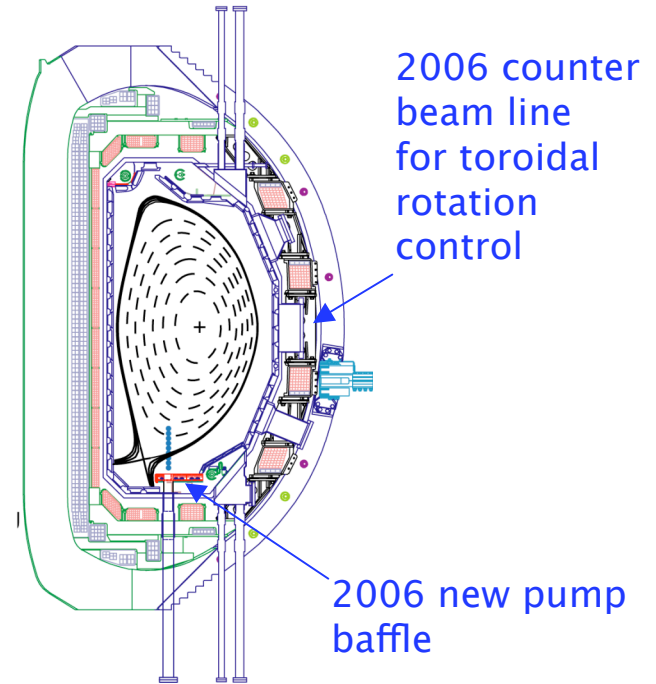
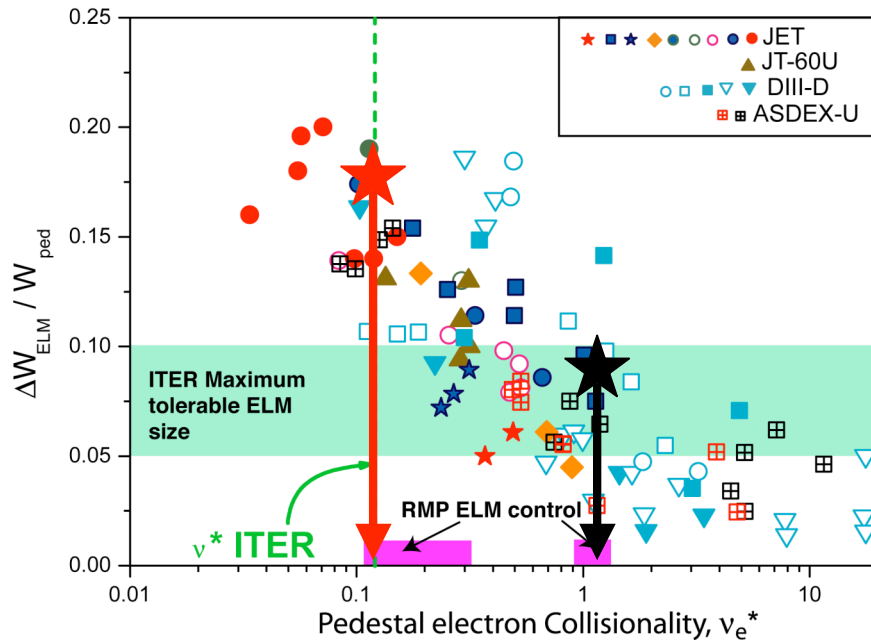


C-coil and field-errors included in each case:
I-coil = 3.2 kA (upper)
 and
no I-coil (lower).

- At $\psi_N = 0.95$, ${}^{\text{TRIP3D}}D_m = 3.9\text{E-}6$ m and $\text{quasi-linear } D_m = 3.5\text{E-}6$ m:
 - > $\text{quasi-linear } \chi_e = v_{Te} D_m \sim 49$ m²/s but to match experimental $T_{e,\text{ped}}^{\text{simulation}} \chi_e \sim 0.2$ m²/s
- Need more comprehensive edge RMP transport theory
 - > Is RMP screening due to plasma rotation or pressure a significant factor?

Significant progress made toward burning plasma pedestal and ELM control in DIII-D using edge RMPs

- **2004 results** : $n=3$ RMPs suppress ELMs at high collisionality in ITER shape



- **New results in 2005** : $n=3$ ELM suppression at ITER relevant low collisionality using DIII-D pumping in low triangularity (δ) plasmas
- **Next year**: new DIII-D hardware allows pumping in higher δ , ITER-like, shapes \rightarrow $n=3$ RMPs in low collisionality ITER-like shapes with low rotation

Summary and Conclusions

- Small edge Resonant Magnetic Perturbations (*RMPs*) used to:
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- Type-I ELMs completely eliminated in low collisionality (ν_e^*) burning plasma relevant conditions:
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