

# Effect of the electron temperature fluctuations on the turbulent flux measurements in the edge of DIII-D

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**D.L. Rudakov, J. A. Boedo,**

**R. A. Moyer, S. Krasheninnikov**

*University of California, San Diego*

**D.M. Thomas, W.P. West**

*General Atomics*

**J. G. Watkins**

*Sandia National Laboratories*

# Outline of the talk

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## I. Motivation

- Fluctuation-driven particle and heat fluxes in edge/SOL plasma
- Problems associated with the flux measurements
- Assessment of the possible effect of the temperature fluctuations on the measurements

## II. Experimental method on DIII-D

- Fast  $T_e$  measurements: Harmonic technique
- Probe head layout
- Different methods of calculating fluxes from probe data

## III. Experimental results

- Comparison of the spectral characteristics of  $n_e$ ,  $T_e$  and  $V_f$  fluctuations
- Comparison of the poloidal wave number spectra of  $T_e$  and  $V_f$  fluctuations
- Phase angles between  $n_e$ ,  $T_e$  and  $V_f$  fluctuations
- Results of applying temperature corrections to the measured fluctuation-driven fluxes

## IV. Discussion and summary

# Fluctuation-Driven Particle and Heat Fluxes

- Like many other aspects of SOL physics, mechanisms behind cross-field particle and heat transport are still poorly understood, but it is widely accepted that “anomalous” transport in the edge and SOL is largely driven by electrostatic fluctuations
- Fluctuating poloidal electric field  $E_\theta$  causes fluctuating radial  $E_\theta \times B_j$  drift. Providing the plasma density and temperature fluctuate, the following time-averaged fluxes result:

$$\text{Particle flux: } \Gamma_r^{ES} = \frac{1}{B_j} \langle \tilde{n}_e \tilde{E}_q \rangle \quad \text{Heat flux: } Q_r^{ES} = \frac{3}{2} k \frac{\langle \tilde{n}_e \tilde{T}_e \tilde{E}_q \rangle}{B_j}$$

- All plasma parameters needed to calculate those fluxes in the edge and SOL plasmas can be derived from probe data
- So far probes provide **the only available means** to directly measure local fluctuation-driven fluxes and are therefore widely used for that purpose
- However, interpretation is not straightforward and problems exist:
  - \* Measured fluxes sometimes exhibit unphysical behavior: heat flux reversal (DIII-D), particle flux continues to increase inside the separatrix into the pedestal (C-Mod)
  - \* Fluxes estimated from the probe data near LCFS are typically higher than surface-average values obtained from modeling and particle/energy balance considerations

# Possible source of errors: Neglecting $T_e$ fluctuations

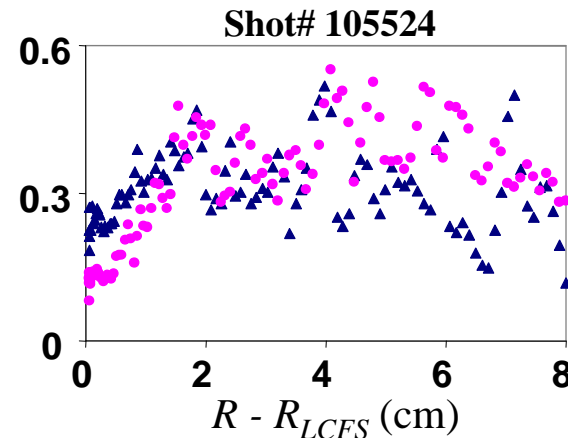
- The poloidal electric field ( $E_q$ ) is commonly estimated from the floating potential ( $V_f$ ) measurements at two poloidally separated locations.
- A proper estimate of  $E_q$  would require using two poloidally separated measurements of the plasma potential  $V_p = V_f + C(kT_e/e)$ , where  $C \sim 3$  for deuterium plasmas. Therefore, a large error can be introduced!
- Effect of  $T_e$  fluctuations can be neglected if:
  - \* Relative  $T_e$  fluctuation level is much smaller than potential fluctuation levels

This is not true in DIII-D

$$\tilde{T}_e / \langle T_e \rangle \approx \tilde{V}_f / \langle T_e \rangle$$

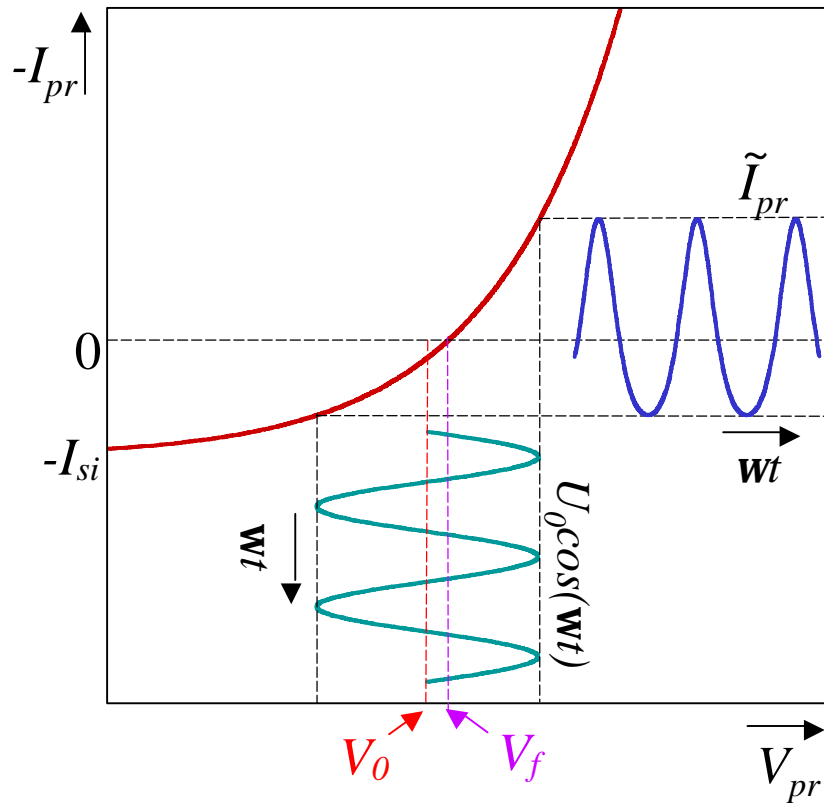
▲  $\tilde{T}_e / \langle T_e \rangle$   
●  $\tilde{V}_f / \langle T_e \rangle$

\*  $k_q(\tilde{T}_e) \ll k_q(\tilde{V}_f)$  - needs to be checked



- **Therefore:**
  - \* Fast measurements of the electron temperature are required
  - \* Poloidal structure of the temperature fluctuations needs to be studied

# Fast $T_e$ measurements: Harmonic technique



- A DC-floating probe is driven by high-frequency sinusoidal voltage  
 $U = U_0 \cos(\omega t)$
- Due to non-linearity of the probe I-V characteristic harmonics are generated in the current spectrum
- For  $eU_0/kT_e < 1$   $T_e$  can be determined from the ratio of the amplitudes of 1<sup>st</sup> and 2<sup>nd</sup> harmonics<sup>[1,2]</sup>

$$kT_e \approx \frac{eU_0}{4} \frac{I_w}{I_{2w}}$$

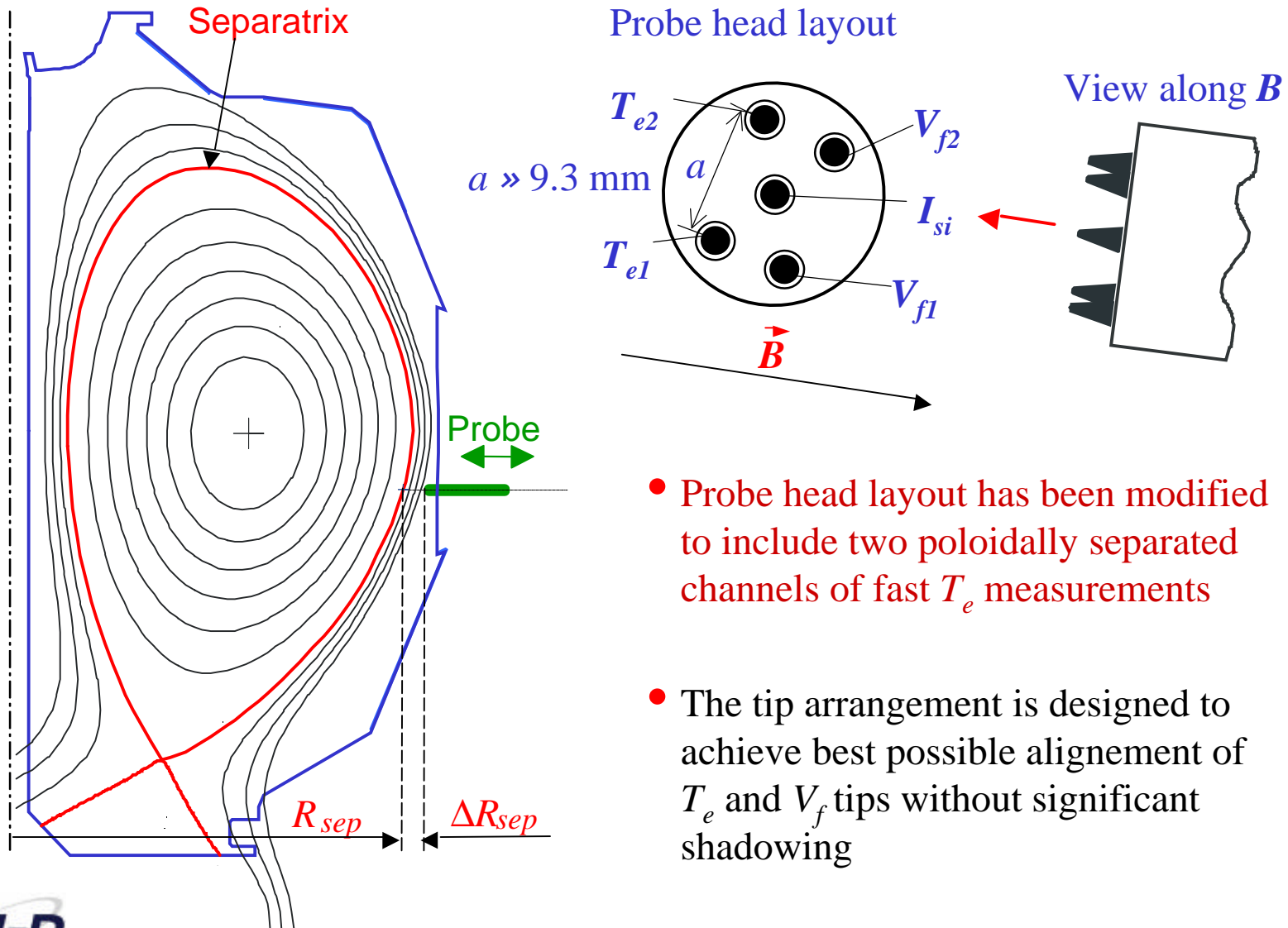
[1] Boedo *et al*, Rev. Sci. Instrum. **70** (1999)

[2] Rudakov *et al*, Rev. Sci. Instrum. **72** (2001)

On DIII-D:

- Drive frequency: 400 kHz
- Bandwidth: typical 100 kHz, maximum 200 kHz
- Novel features: fully digital analysis, active voltage feedback

# Experimental Arrangement on DIII-D



- Probe head layout has been modified to include two poloidally separated channels of fast  $T_e$  measurements
- The tip arrangement is designed to achieve best possible alignment of  $T_e$  and  $V_f$  tips without significant shadowing

# Calculating turbulent fluxes from probe data

Turbulent fluxes can be calculated in time domain:

$$\Gamma_r^{ES} = \frac{1}{B_j} \langle \tilde{n}_e \tilde{E}_q \rangle$$

$$Q_r^{ES} = \frac{3}{2} k \frac{\langle \tilde{n}_e \tilde{T}_e \tilde{E}_q \rangle}{B_j} = \frac{3}{2} k T_e \Gamma_r^{ES} + \frac{3}{2} \frac{n_e}{B_j} \langle k \tilde{T}_e \tilde{E}_q \rangle = Q_{conv} + Q_{cond}$$

**convective**      **conductive**

or in frequency domain:

$$\Gamma_r^{ES}(\omega) = \frac{1}{B_j} \int_0^\infty n_e^{rms}(\omega) E_q^{rms}(\omega) g_{nE}(\omega) \cos(\mathbf{a}_{nE}(\omega)) d\omega$$

$$Q_r^{ES}(\omega) = \frac{3}{2} k T_e \Gamma_r^{ES}(\omega) + \frac{3}{2} \frac{n_e k}{B_j} \int_0^\infty T_e^{rms}(\omega) E_q^{rms}(\omega) g_{TE}(\omega) \cos(\mathbf{a}_{TE}(\omega)) d\omega$$

**convective**      **conductive**

where:

$X^{rms}(\omega)$  - RMS fluctuation amplitude

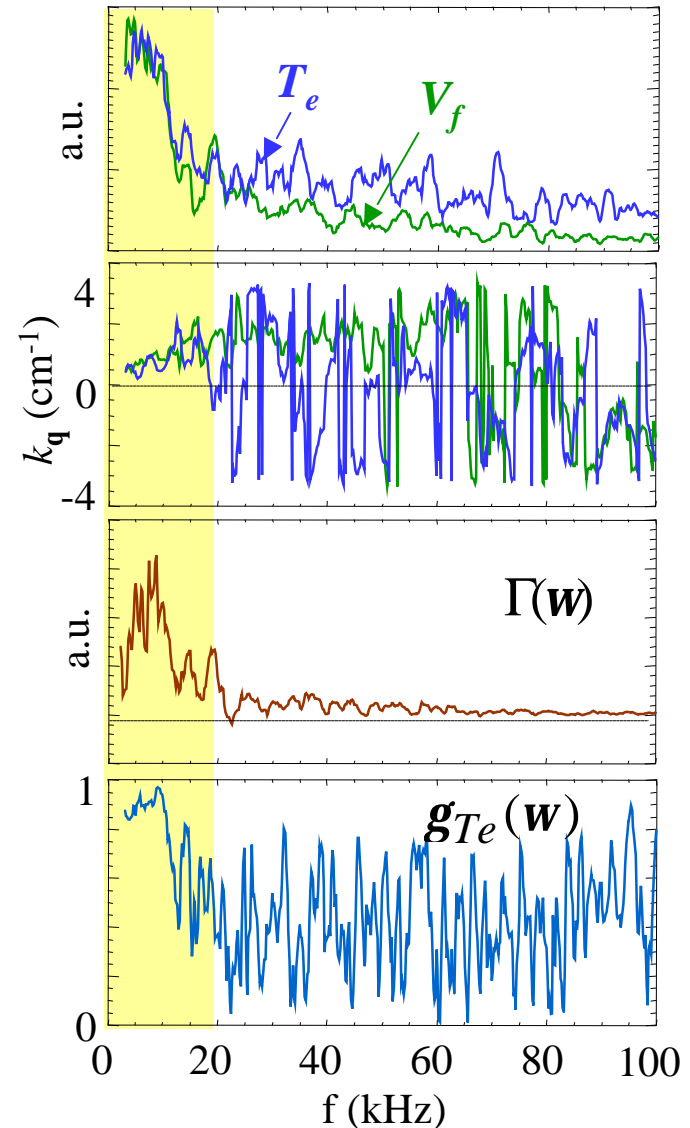
$\mathbf{a}_{XX}(\omega)$  - phase angle between  $n_e$  or  $T_e$  and  $E_q$

$\mathbf{g}_{XX}(\omega)$  - coherence

Can be obtained  
from raw signals  
using FFT

# Spectral features of fluctuations in L-mode (near SOL)

- The shapes of  $T_e$  and  $V_f$  spectra are generally close, but  $V_f$  spectra fall off more rapidly with frequency
- Poloidal wave numbers of the temperature fluctuations below 20 kHz are close but not equal to those of  $V_f$  fluctuations
- Most of the fluctuation power and of fluctuation-driven flux are concentrated at frequencies below 20 kHz. Fluctuations are most coherent in that range
- For the above conditions we can characterize the relative phasing of the fluctuations by an average phase angle calculated where the coherence and cross-power are maximum





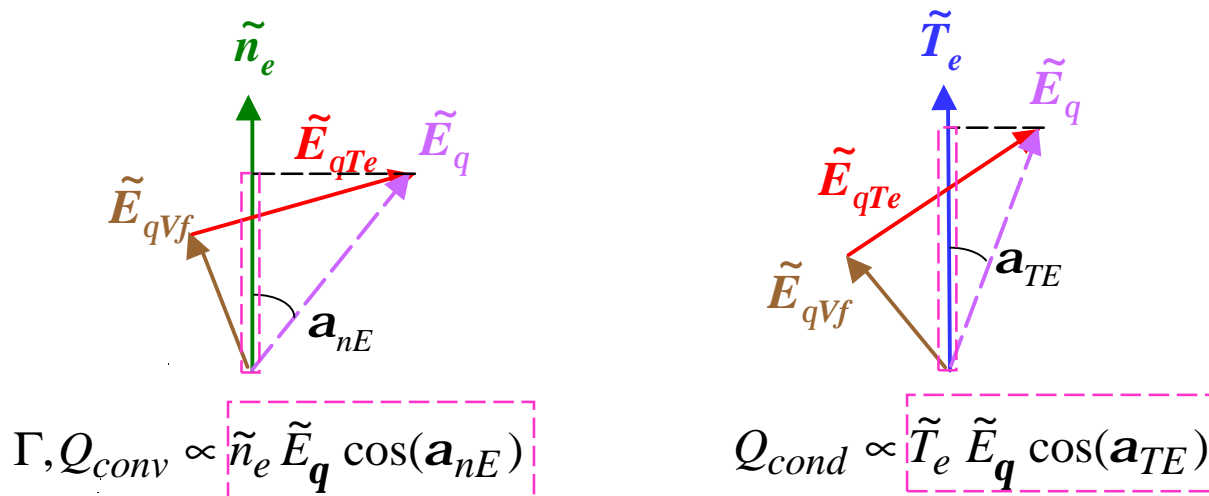
# Correcting the fluxes for $T_e$ fluctuations

Net poloidal electric field can be represented as a sum of two components:

$$\mathbf{E}_q = -\nabla_q(V_p) = -\nabla_q(V_f) - 3\nabla_q(kT_e/e) = \mathbf{E}_{qVf} + \mathbf{E}_{qTe}$$

Net field
 $V_f$  contribution  $\circ \mathbf{E}_{qVf}$ 
 $T_e$  correction  $\circ \mathbf{E}_{qTe}$

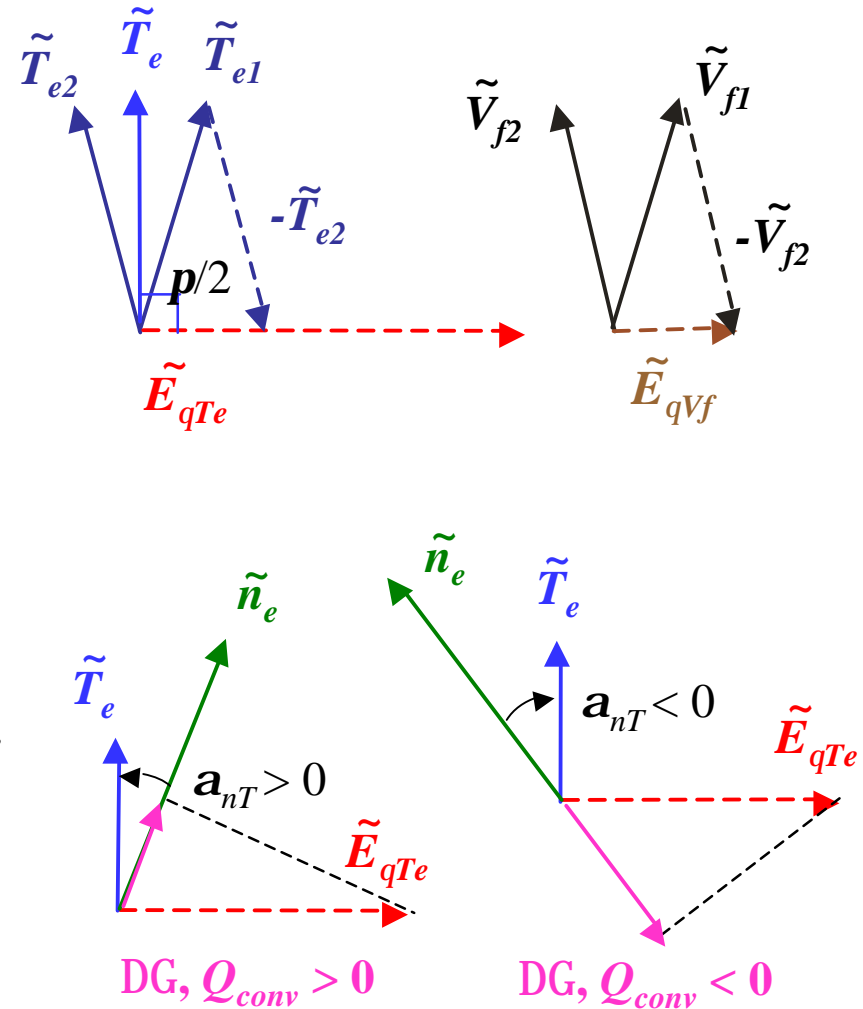
It is convenient to illustrate contribution of the two components of  $E_q$  to fluxes using phase diagrams (length = amplitude,  $\mathbf{a}$  = phase angle)



**Note:** if for any of  $E_q$  components  $\mathbf{a} = \pm 90^\circ$ , it will cause no flux

# $T_e$ correction to fluxes can be large and of either sign

- Since  $\tilde{T}_e / T_e \approx e\tilde{V}_f / kT_e$  and  $k_q(\tilde{T}_e) \approx k_q(\tilde{V}_f)$   $E_{qT_e}$  should be about 3 times larger than  $E_{qVf}$
- Phase angle between  $T_e$  and  $E_{qT_e}$  tends to be around  $p/2$ , so contribution of  $E_{qT_e}$  to  $Q_{cond}$  should be relatively small
- Contribution of  $E_{qT_e}$  to  $\Gamma$  and  $Q_{conv}$  depends on the phase angle between  $T_e$  and  $n_e$ . If  $T_e$  and  $n_e$  are in phase, this contribution is small, otherwise it may be quite large and of either sign



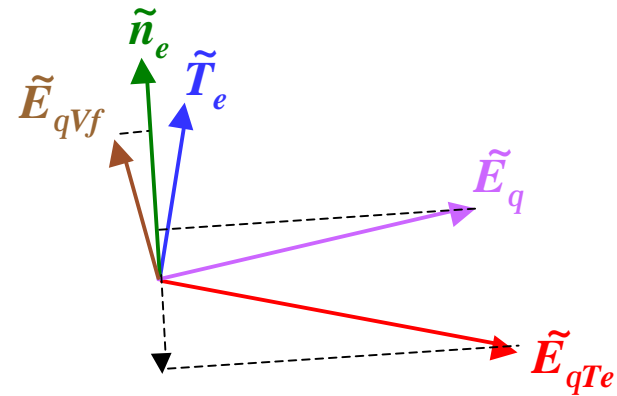
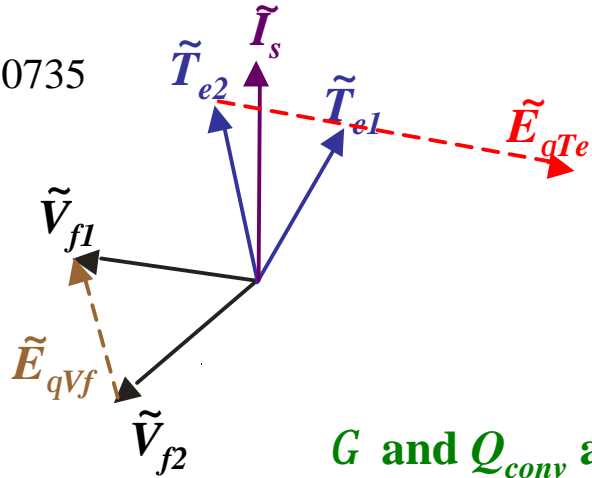
$$\Delta\Gamma, Q_{conv} \propto \sin(a_{nT})$$

# Experimental results: flux correction in L-mode

Measured parameters

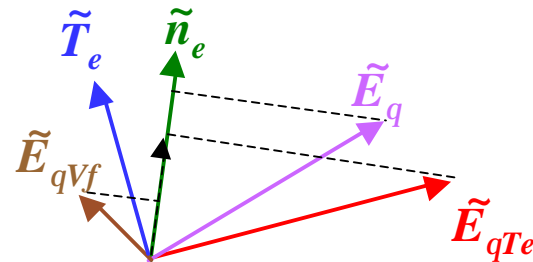
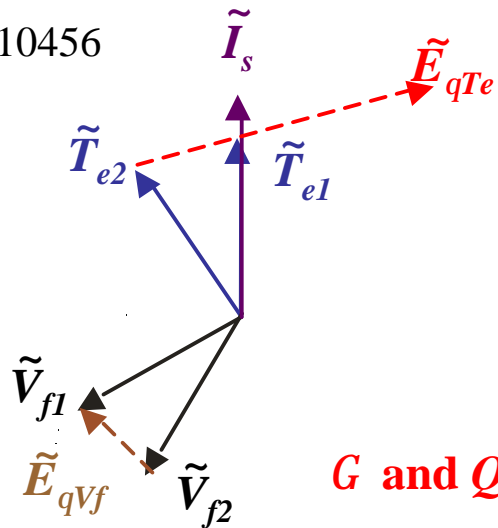
Derived parameters

Shot 110735



$G$  and  $Q_{conv}$  are reduced by a factor of  $\sim 3$

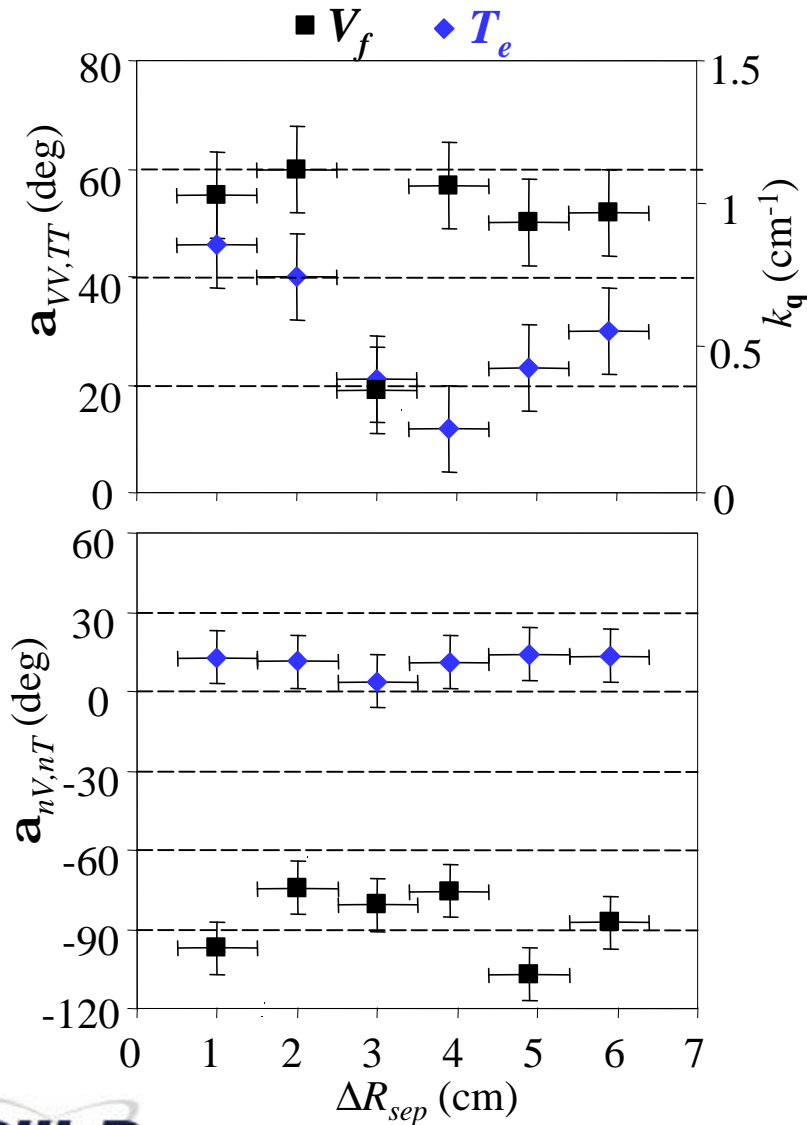
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$G$  and  $Q_{conv}$  are increased by a factor of  $\sim 3$

# Experimental results: radial dependence of phase angles

Shot 110735



- Poloidal wave numbers of the temperature fluctuations below in near SOL are close to those of  $V_f$  fluctuations. In far SOL they differ by a factor of 2 - 3
- Phase angle between  $T_e$  and  $n_e$  tends to be small, about  $15^\circ$ , while phase angle between  $V_f$  and  $n_e$  is around  $90^\circ$ .
- Since  $\Delta\Gamma$ ,  $Q_{conv} \propto \sin(\mathbf{a})$  and  $\sin(\mathbf{a}_{nV}) / \sin(\mathbf{a}_{nT}) \approx 5$ , relative effect of  $T_e$  fluctuations is offset in this case, but it can still be large

# Summary

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- SOL transport is still poorly understood but thought to be largely  $E \times B$  fluctuation- driven
- So far electric probes provide the only available means for measuring local fluctuation-driven fluxes in the edge and SOL
- Neglecting electron temperature fluctuations while deriving fluctuation-driven fluxes from the probe data can lead to large errors
- New experimental arrangement on DIII-D provides a unique opportunity to study poloidal structure of the electron temperature fluctuations
- Poloidal wave numbers of the electron temperature fluctuations below 20 kHz in L-mode are close to those of the floating potential fluctuations in near SOL, differ considerably in far SOL
- Depending on the phase angle between density and temperature fluctuations, contribution of the temperature fluctuations to the measured poloidal electric field can be large and lead to significant increase/decrease or even direction reversal of the measured particle and convective heat fluxes
- Conductive heat flux measurement is less affected by  $T_e$  contribution to  $E_q$  since phase angle between  $T_e$  and  $E_{qT_e}$  tends to be around  $\pi/2$

More experiments in L and H mode are needed!