

A Mechanism for Tearing Mode Onset Near Ideal Stability Boundaries

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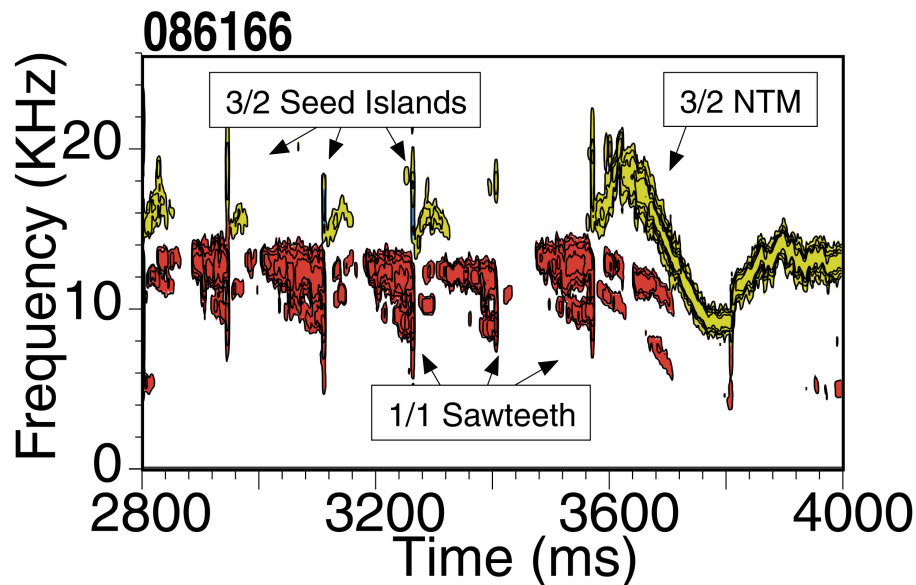


New Onset Mechanism Addresses Two Common NTM Onset Puzzles

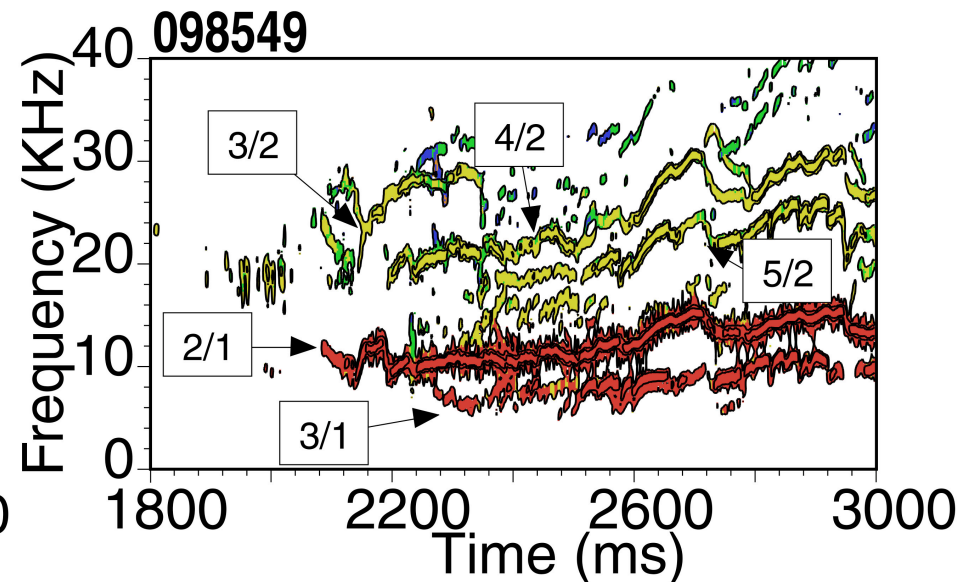
NTMs are linearly stable but nonlinearly unstable, ie. Metastable

- Need “seed” or triggering event

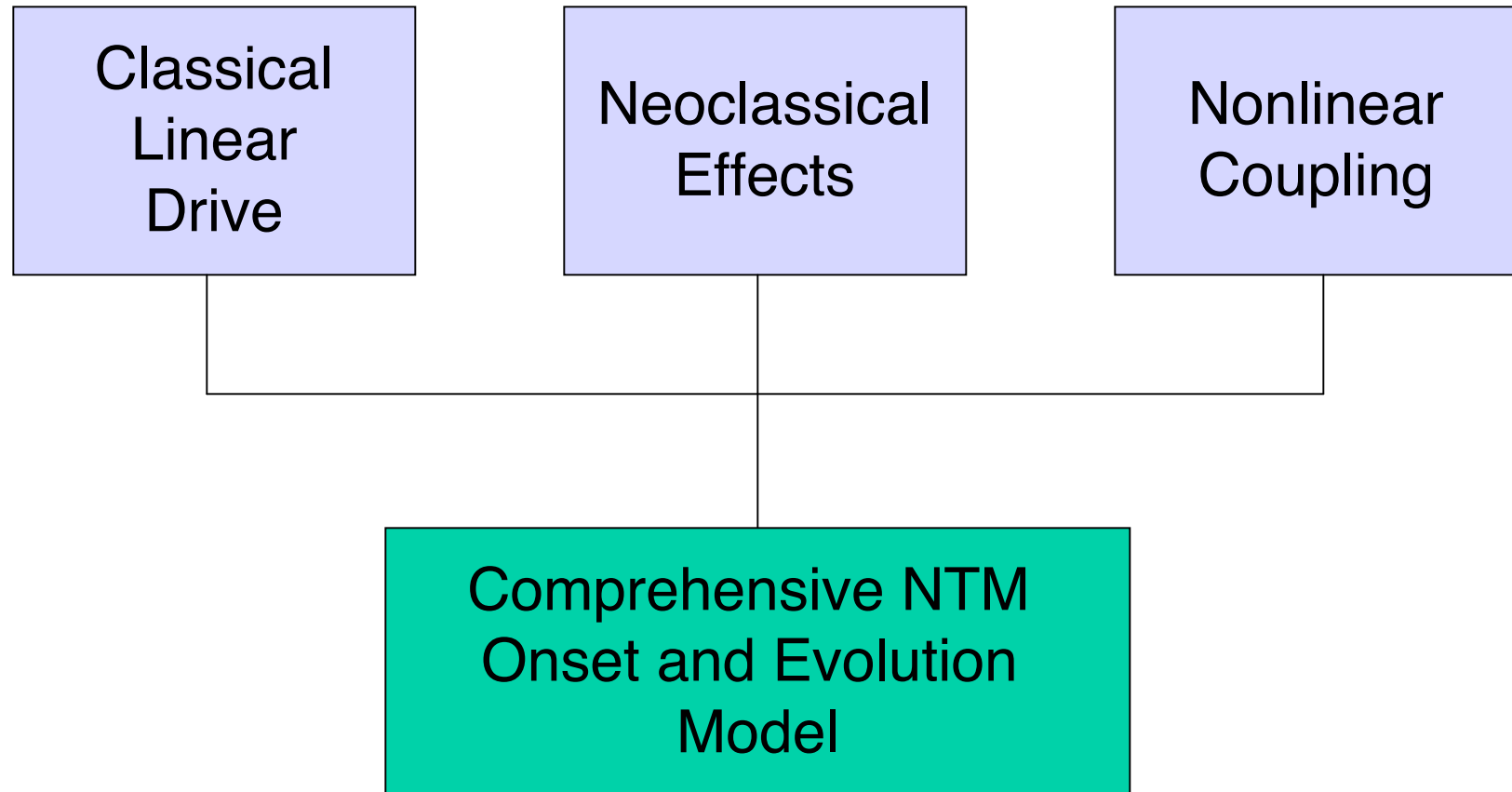
Why Does The n^{th} Sawtooth Seed an NTM?



What Causes Seedless NTMs?

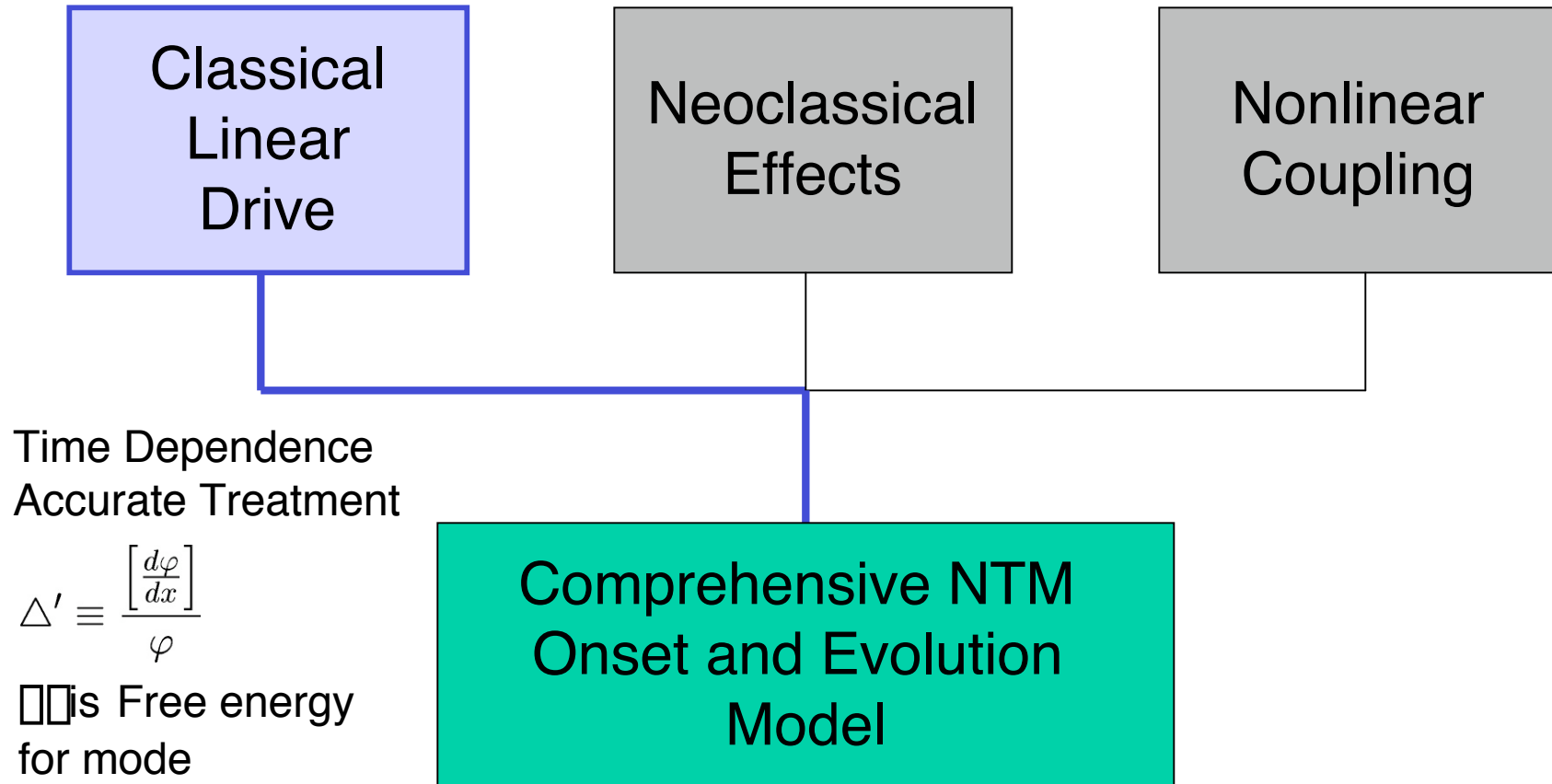


A Comprehensive Model of NTM Onset Emerges From a Combination of Theories



All three components are necessary to explain the onset and evolution

A Comprehensive Model of NTM Onset Emerges From a Combination of Theories



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Outline

Thesis: $\beta(t)$ depends strongly on the equilibrium parameters near ideal limit _ key physics for NTM onset in some cases.

Island Evolution Equation : essential physics of nonlinear evolution

Analytic model : the basic physics of the time dependence of $\beta(t)$

Sawtooth Seeding of 3/2 NTM : DIII-D discharge which encompasses **nonlinear**, **neoclassical** and **classical** effects is comprehensively analyzed showing the critical time dependence in $\beta(t)$.

Pole Experiment : theoretical predictions of the effects of time dependence in $\beta(t)$ are compared to results from a DIII-D experiment to isolate these effects.

Island Evolution Equation Captures Essential Physics Necessary to Describe Onset and Early Evolution

Coupling not included

□□ often assumed
constant negative

**With accurate time
dependence and
weak coupling,
onset and early
evolution can be
predicted**

$$\frac{dw}{dt} = \frac{\eta^*}{k_0} \left[\Delta^* + \frac{k_1}{w} \left(D_{nc} + \frac{D_R}{\alpha_s - H} \right) + \frac{D_{pol}}{w^3} \right]$$

where

$$\Delta^* = \Delta' \left(\frac{w}{2} \right)^{-2\alpha_l} (-4D_I)^{1/2}$$

$$\frac{\eta^*}{k_0} \sim \frac{r_s^2}{\tau_r}$$

Classical drive

$$k_1 = \frac{w^2}{w^2 + w_d^2}$$

Neoclassical drive

$$D_{nc} = k_2 \sqrt{\epsilon} \beta_\theta \frac{L_q}{L_p}$$

Polarization drive

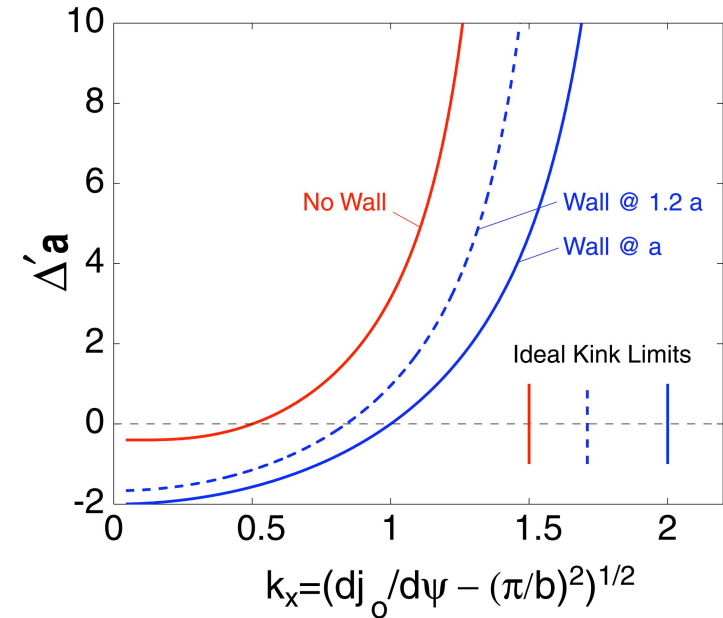
$$D_{pol} = D_{nc} \frac{\rho_{i\theta}}{a^2} \frac{L_q}{L_p} g$$

Analytic Models of Linear Tearing Stability Show Strong Variation of $\Delta'a$ Near Ideal Instability

Small change in equilibrium and/or conducting wall dramatically affects $\Delta'a$ near ideal limit due to

POLE DISCONTINUITY at ideal limit.

The Pole Form is Ubiquitous
Among Derivations, Numeric and Analytic



With a conducting wall at a:

j_o constant,
slab geometry:

$$\Delta'a = -2k_x a \cot(k_x a)$$

$$k_x = \sqrt{j_o' - k_y^2}$$

With a vacuum region from a to b:

$$\Delta'a = 2k_x a \frac{E \sin(k_x a) + \cos(k_x a)}{E \cos(k_x a) - \sin(k_x a)}$$

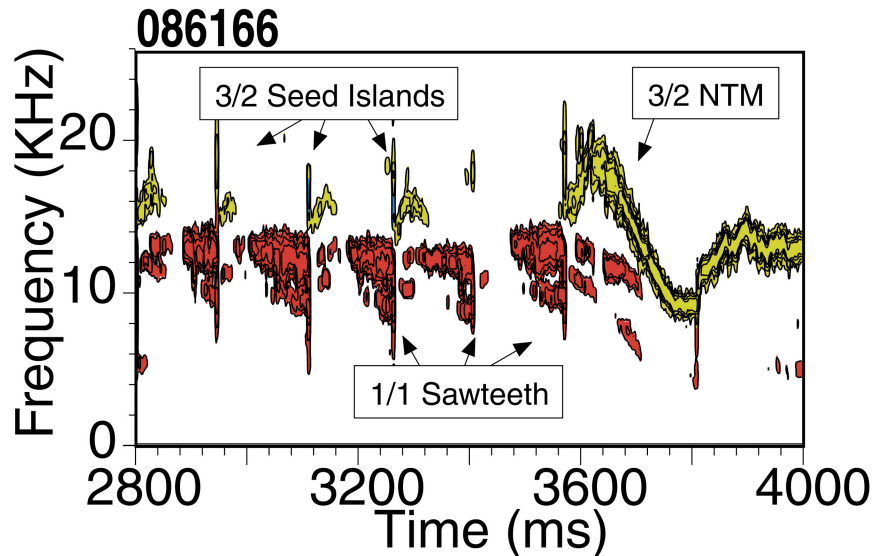
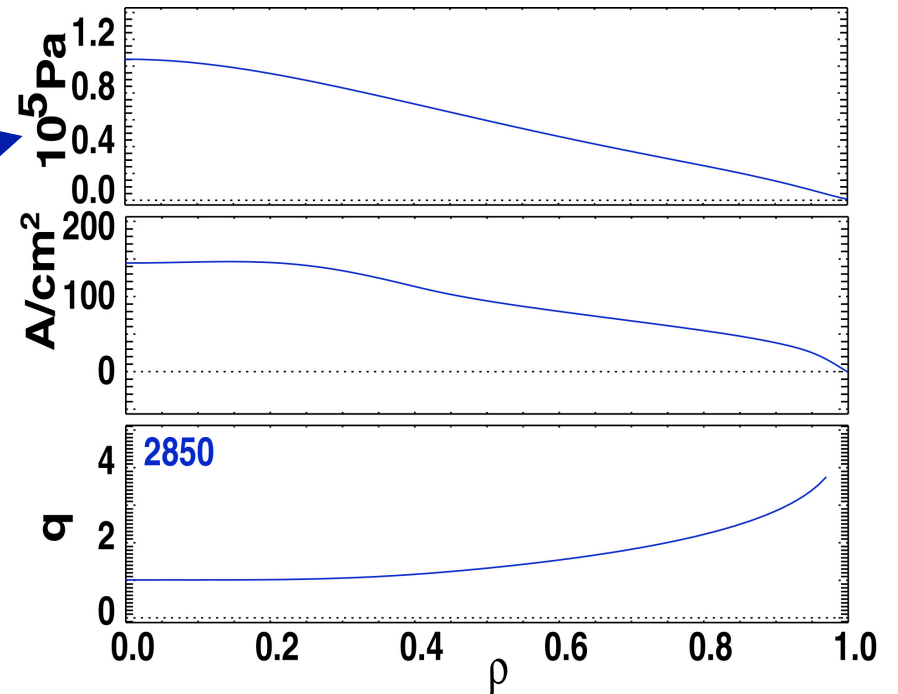
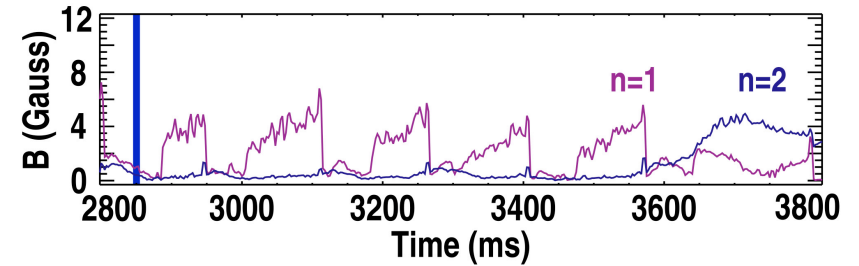
$$E = \frac{e^{k_x a} - e^{k_x(2b-a)}}{e^{k_x a} + e^{k_x(2b-a)}}$$

$$k_x = \sqrt{j_o' - k_y^2}$$

Slight Increase in Core Pressure Destabilizes a Seed Island from a Sawtooth and Causes the Onset of a 3/2 NTM

Equilibrium reconstructions between sawteeth

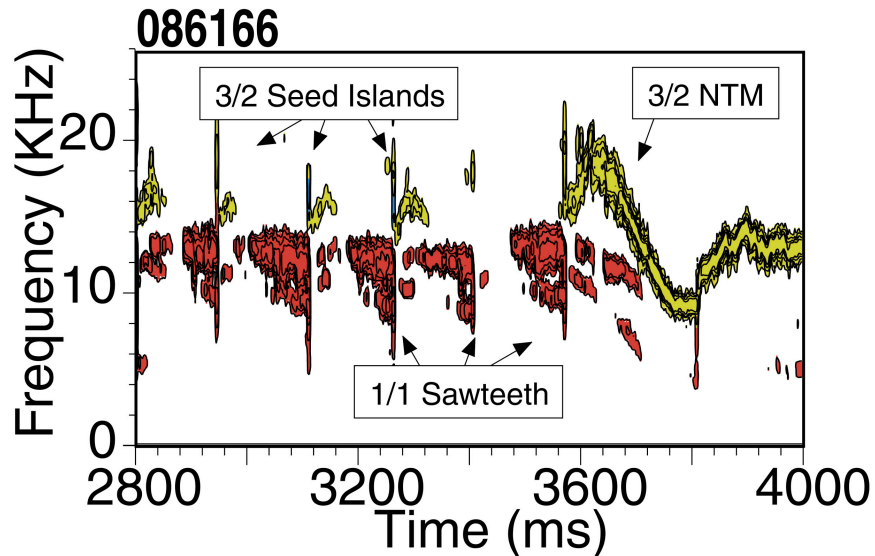
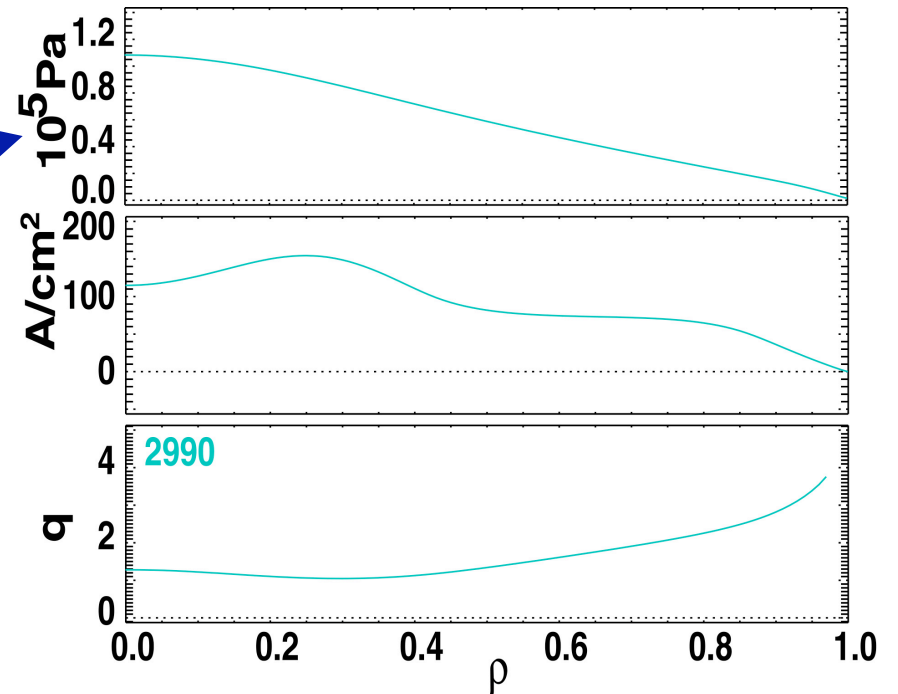
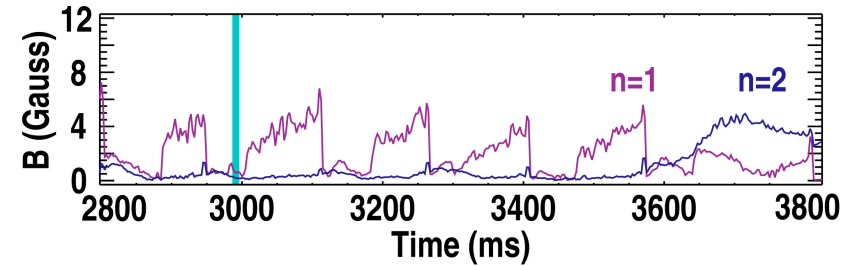
Most pronounced profile variation is core pressure just before onset.



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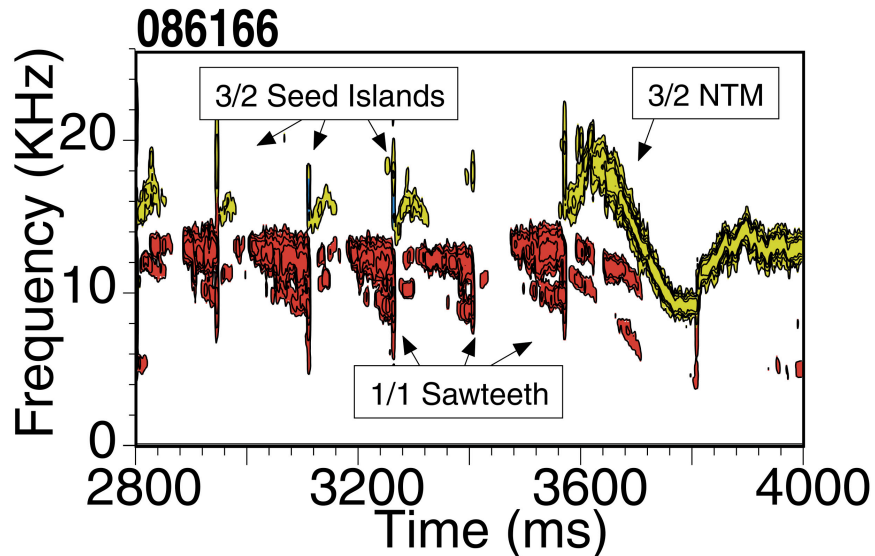
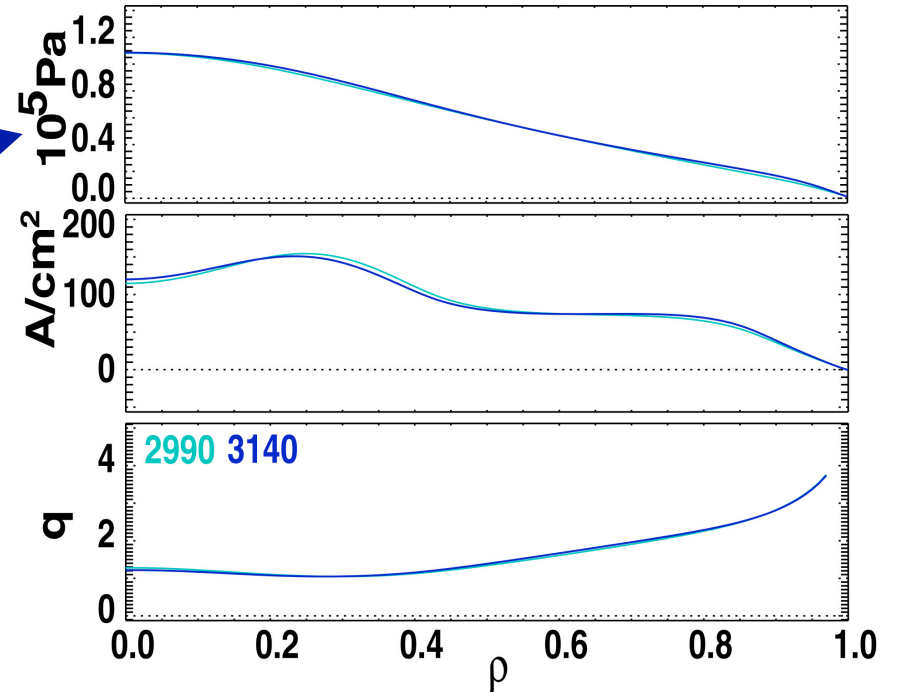
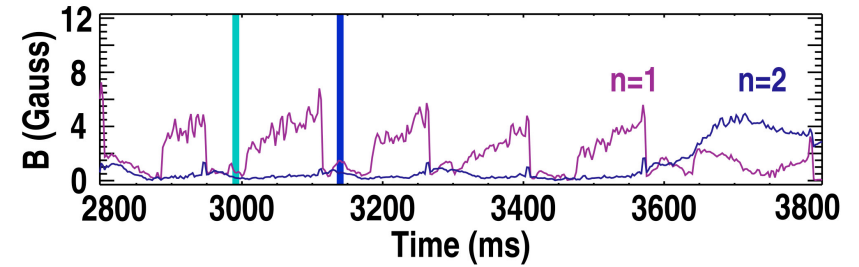
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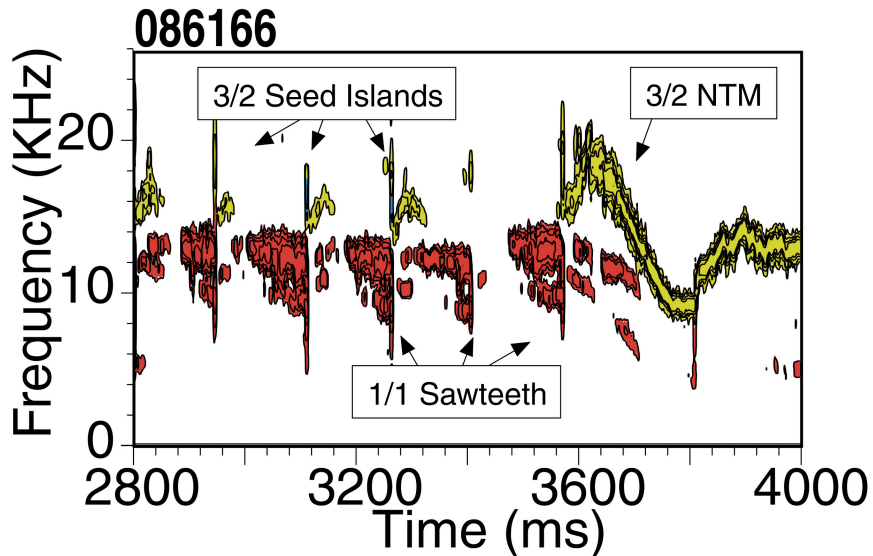
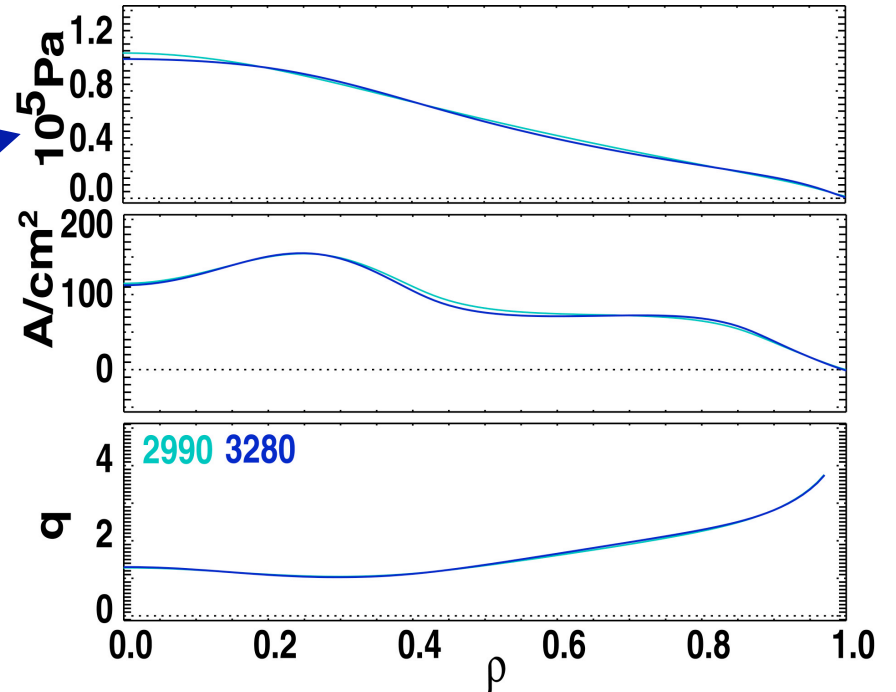
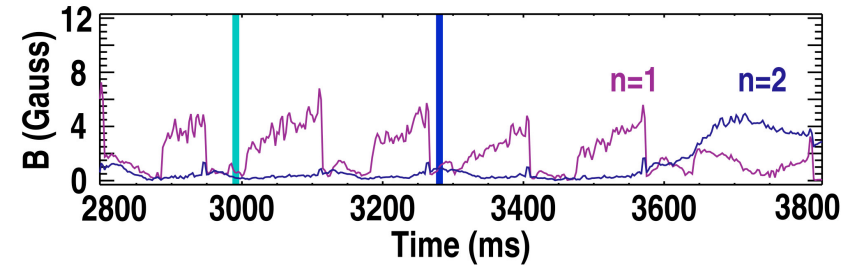
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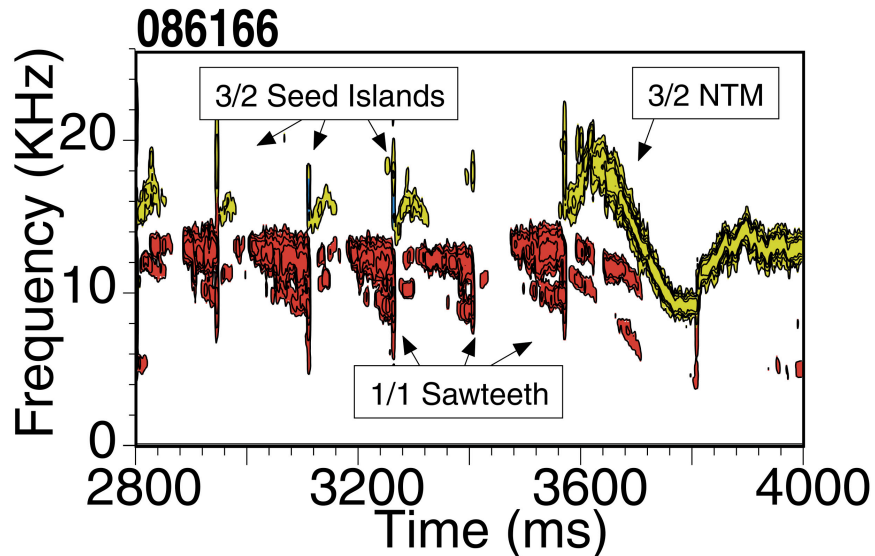
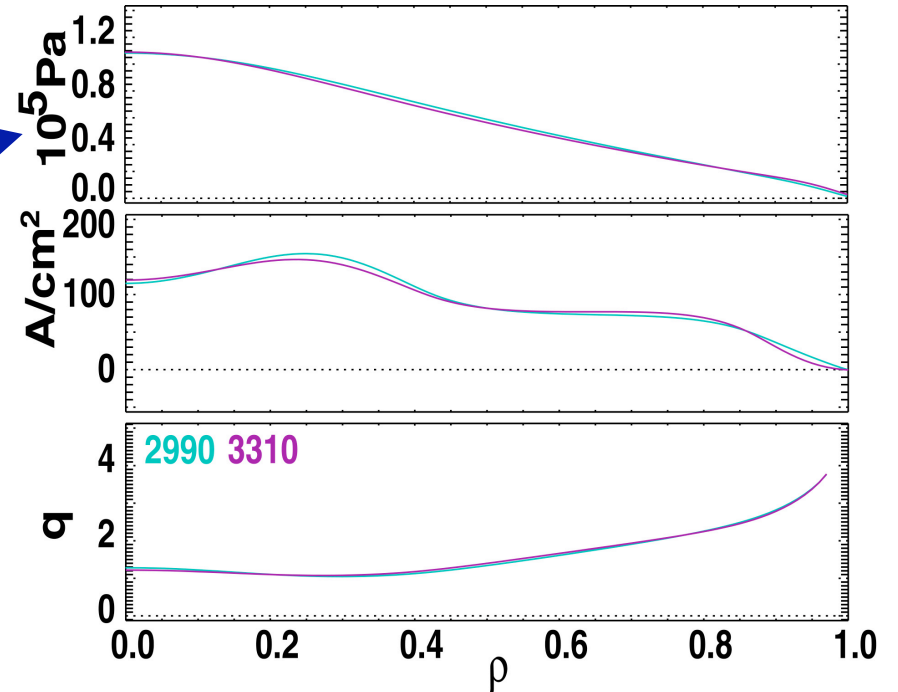
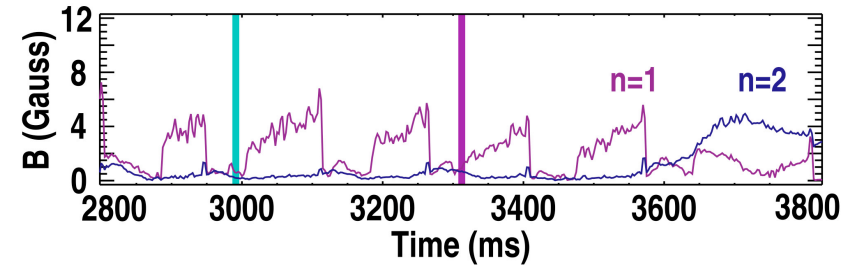
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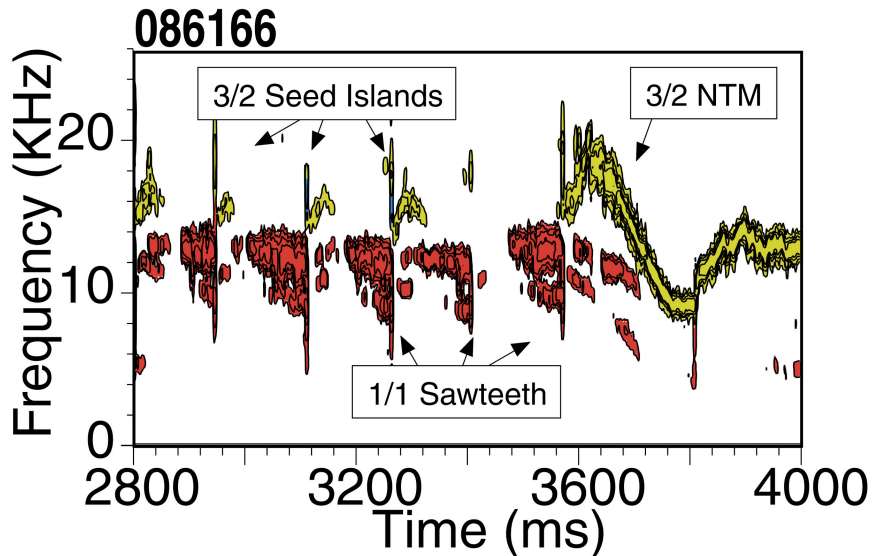
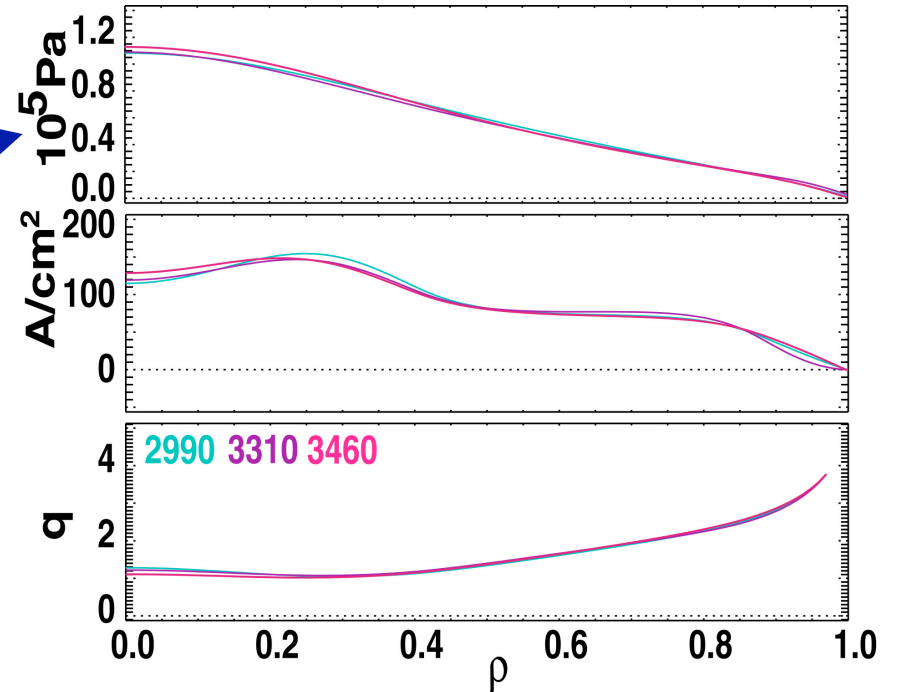
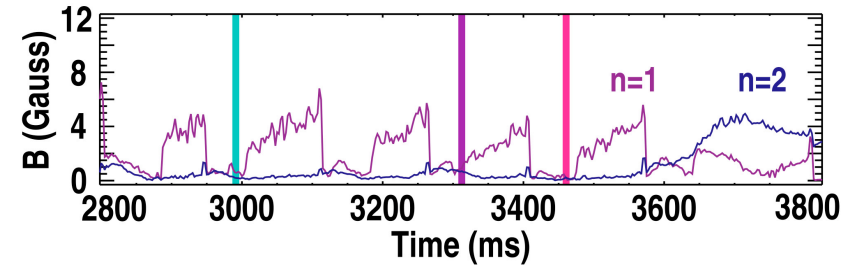
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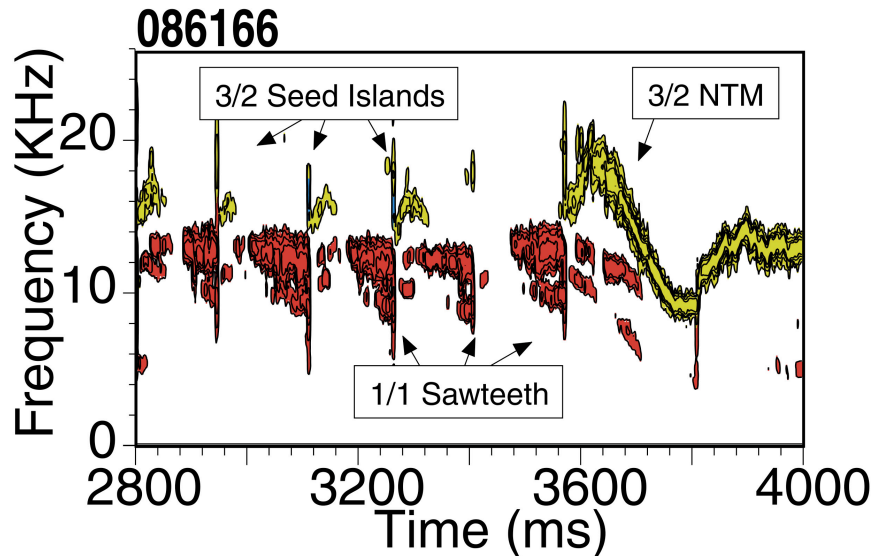
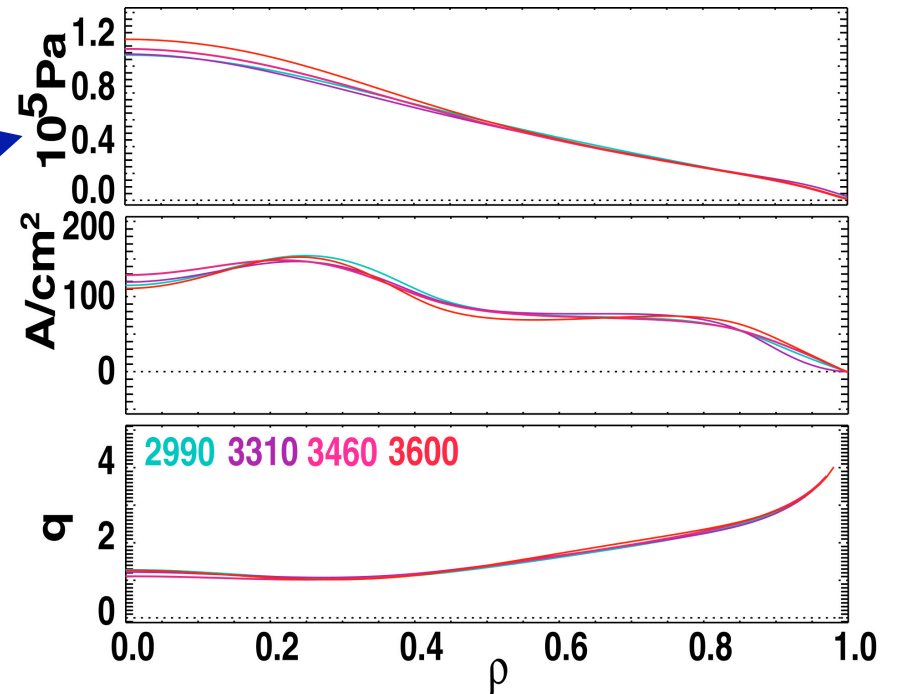
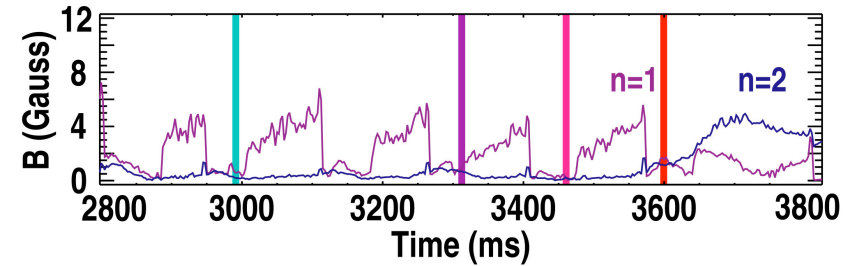
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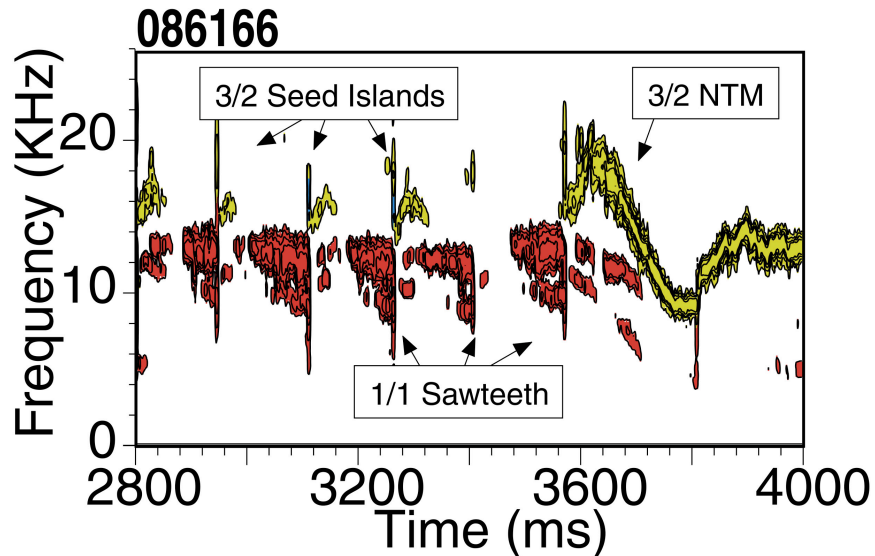
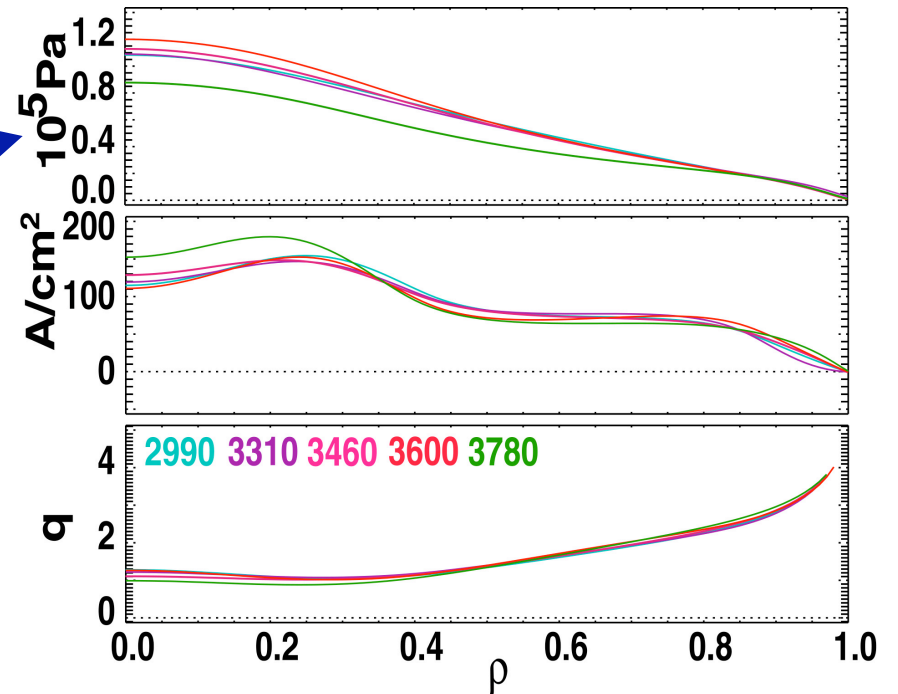
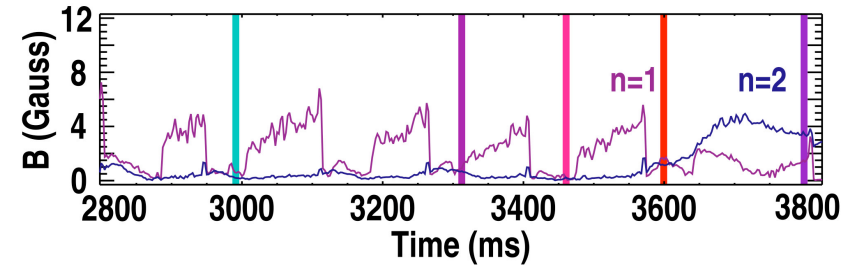
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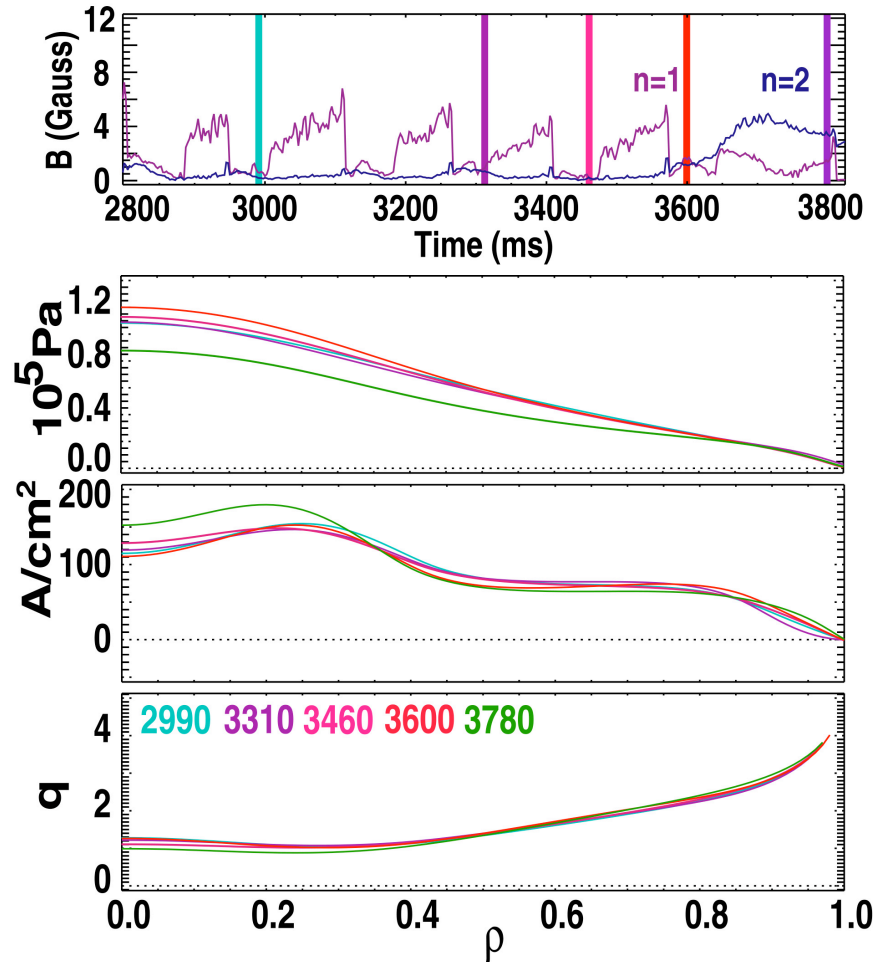
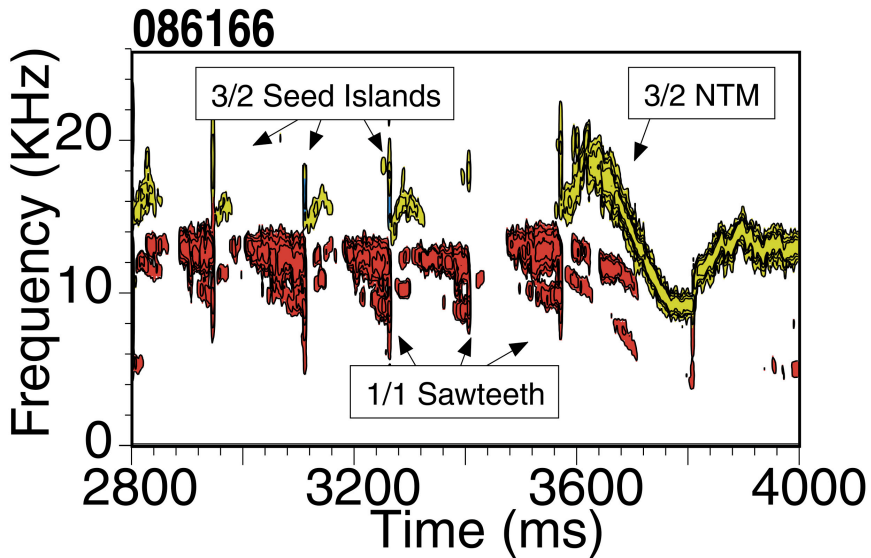
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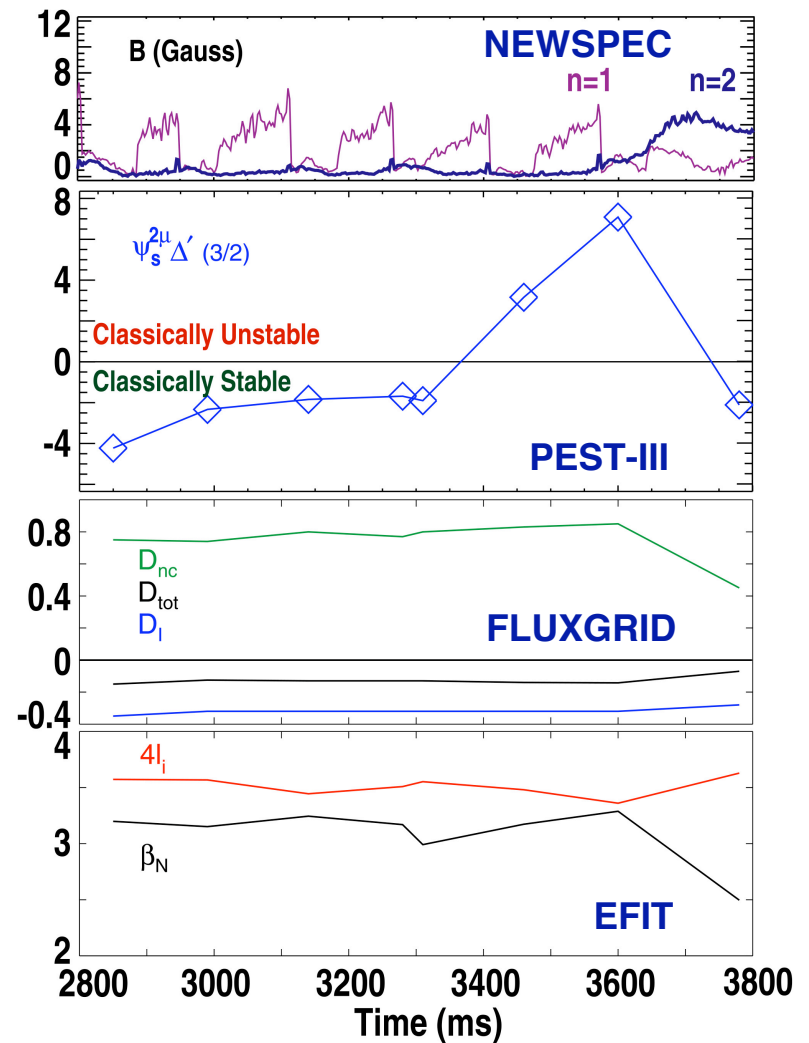
Point is to answer: why did the last sawtooth set off the 3/2 NTM while the several preceding similar sawteeth did not.



Neoclassical Terms Remain Constant Between Sawtooth Periods, While ΔI Increases Sharply

Destabilization of seed island and transition to NTM state is driven primarily by change in ΔI

β_N approaches no-wall limit at onset of 3/2 NTM _ ΔI pole is underlying mechanism for destabilization

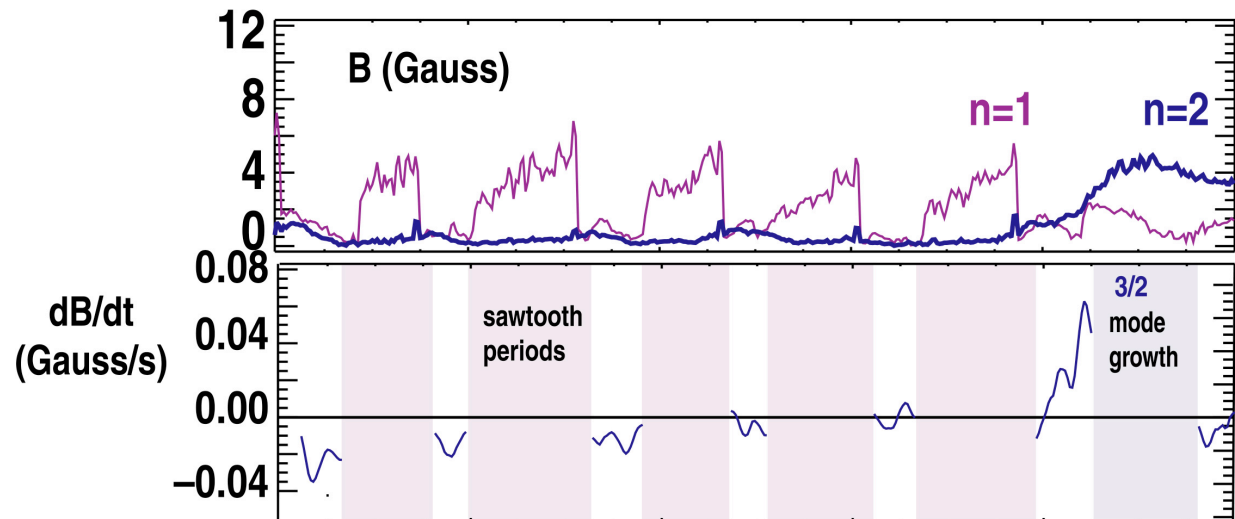


To Test This Proposition Island Evolution is Reconstructed From Experiment and Compared to Theory

$$w = \left(\frac{16rR_o|\tilde{B}_r|}{msB_{\phi o}} \right)^{1/2}$$

$$|\tilde{B}_r| = \frac{1}{2} \left(\frac{b}{r} \right)^{m+1} |\tilde{B}_\theta|$$

w from Magnetic Probe Signals used to compare dw/dt from theory and experiment

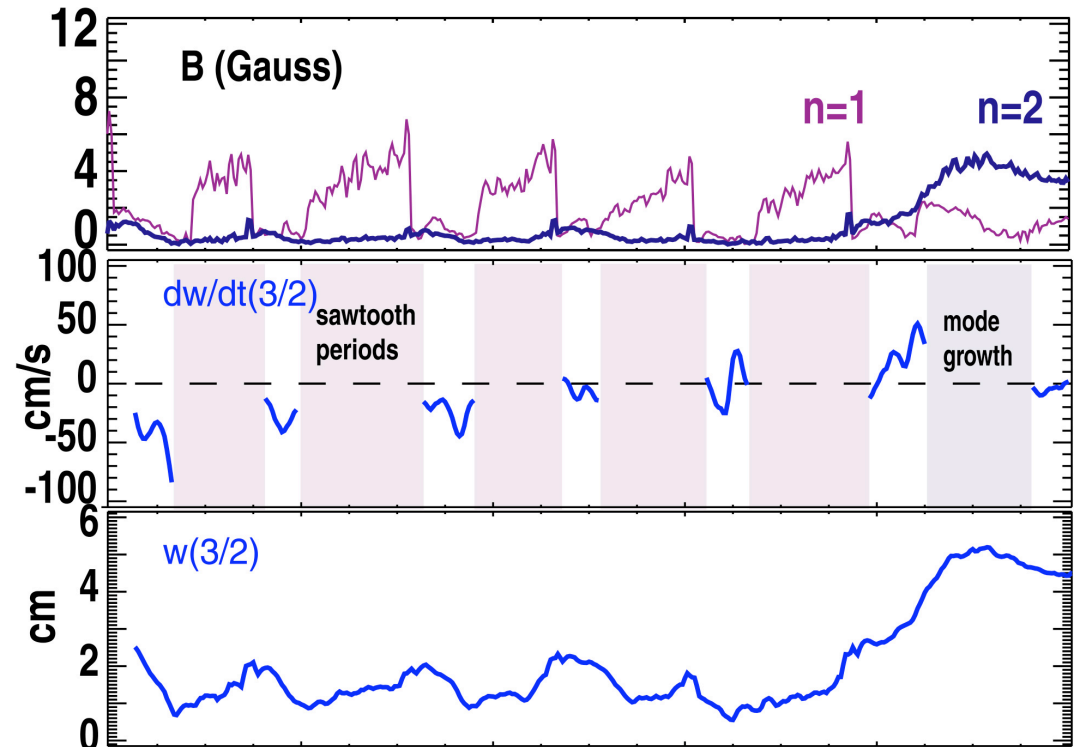


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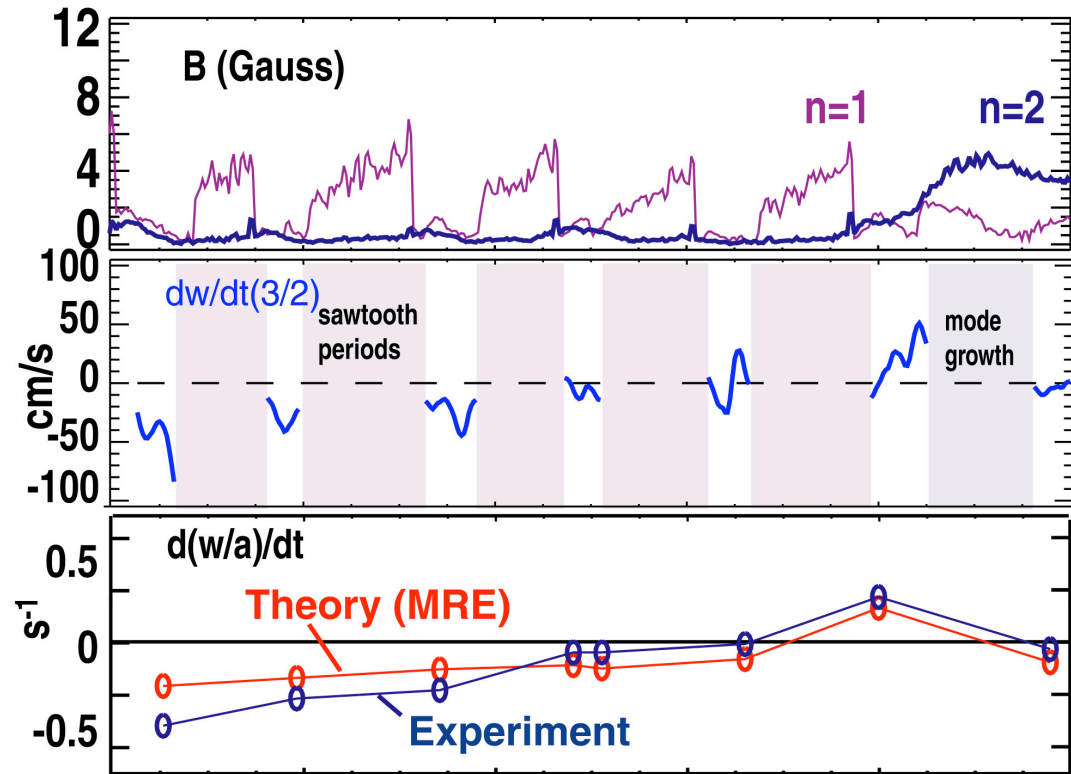


dw/dt from Island Evolution Equation Agrees with dw/dt from Experimental Data

Island Evolution Equation gives right answer even without nonlinear coupling and with axisymmetric \square .

D_{pol} determined from \square^2 minimization agrees in sign with analytic value, but is smaller.

$$D_{pol} = -2.3e-5$$

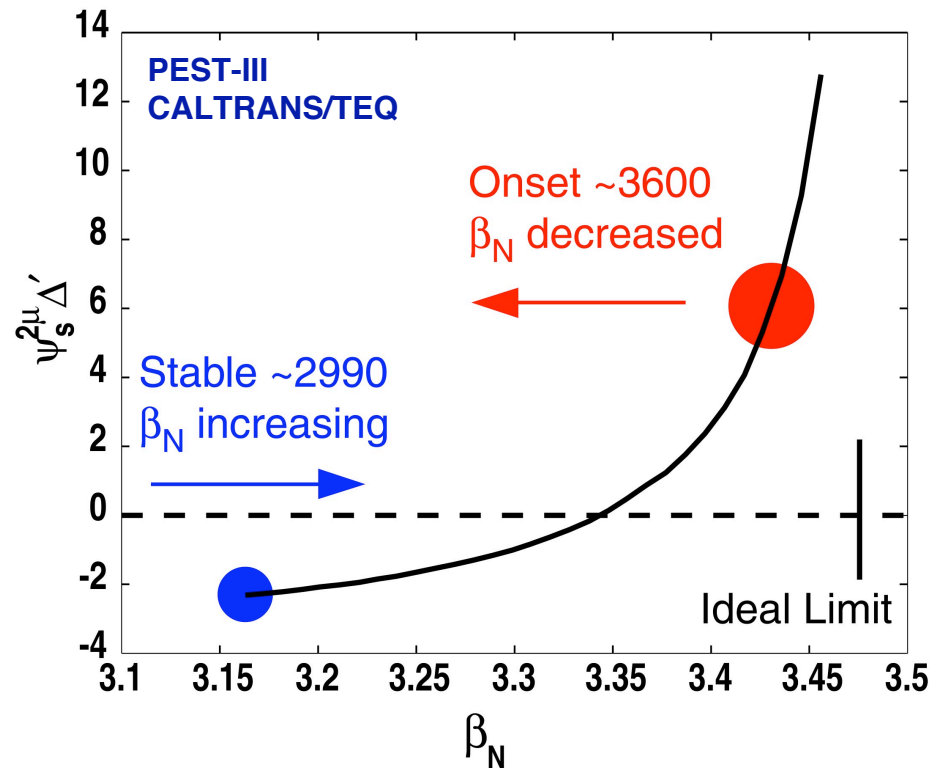


Increase in Core Pressure Causes β_N to Sharply Increase Due to Approach of Ideal Limit

Pressure profile at 2990 was gradually modified to approximate profile at 3600, showing transition.

Approach to $n=2$ ideal limit increases $3/2 \beta_N$

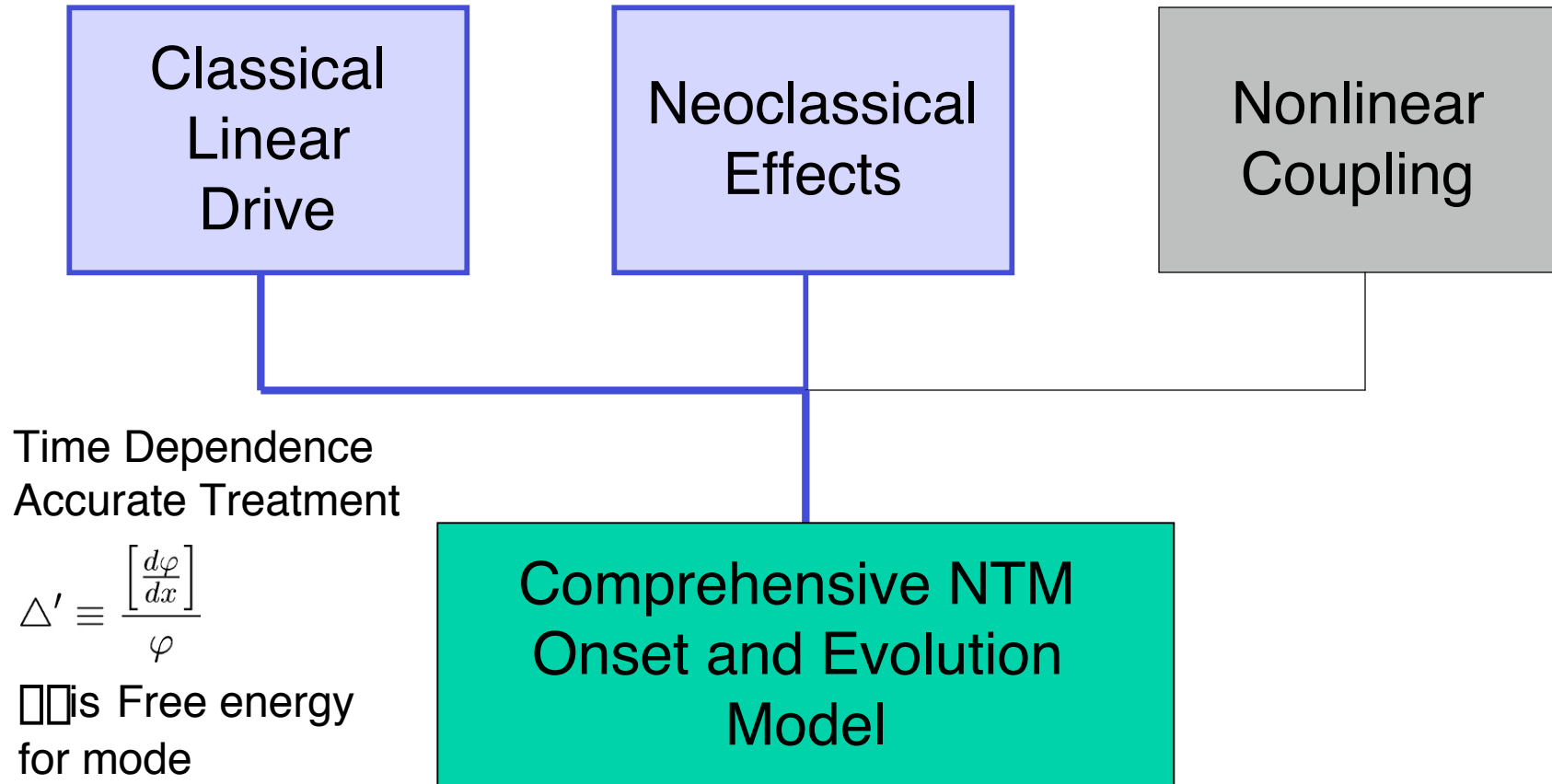
Pole causes rapid change in β_N as β_N changes slightly



Nonlinear coupling between $n=1$ and $n=2$ modes not addressed by this model

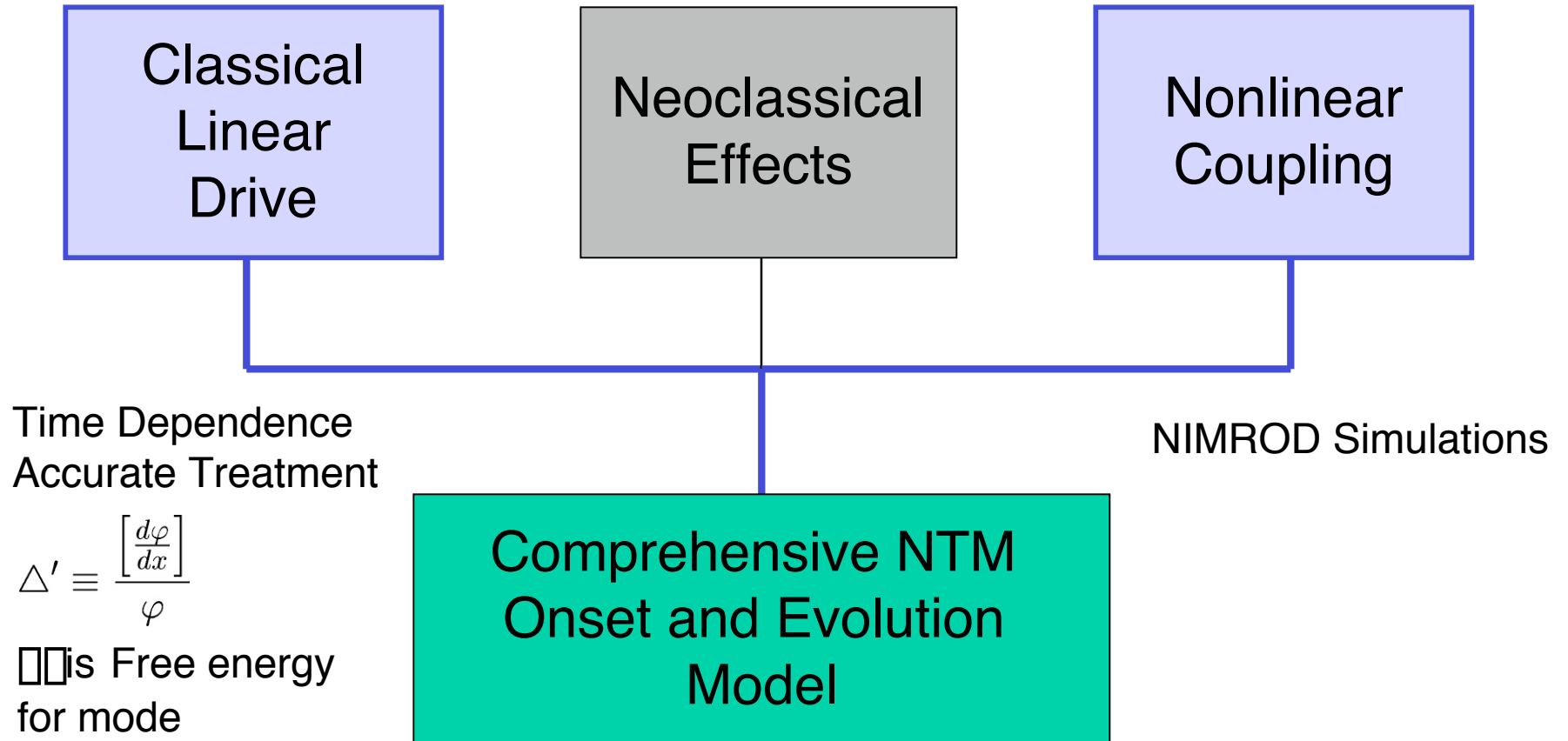
How will this affect the stability and evolution?

A Comprehensive Model of NTM Onset Emerges From a Combination of Theories



All three components are necessary to explain the onset and evolution

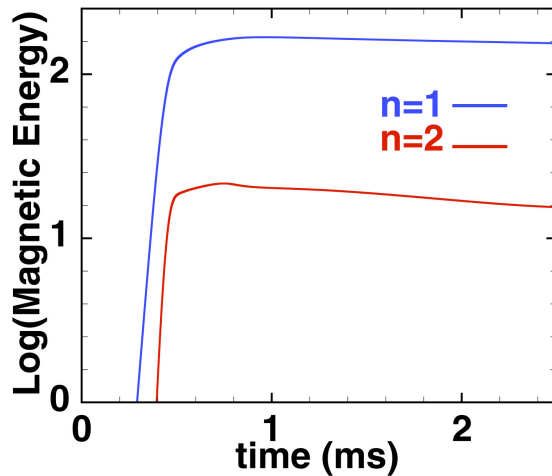
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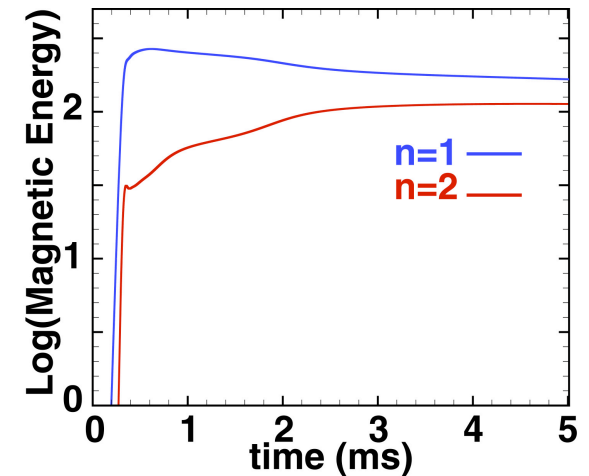
NIMROD Simulations of two discharge times show unstable n=1 mode and driven n=2 mode, in agreement with experiment

$t_{\text{exp}}=2990$

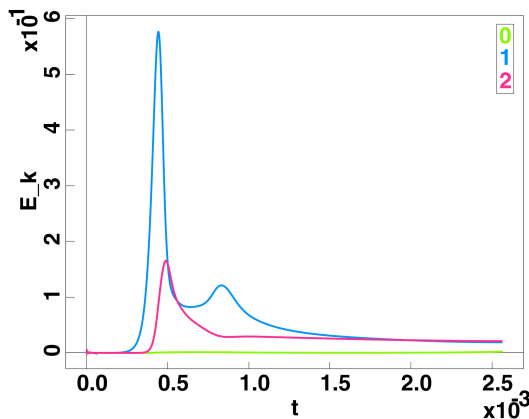


Stable
n=2 decays
after n=1
drive begins
to reduce

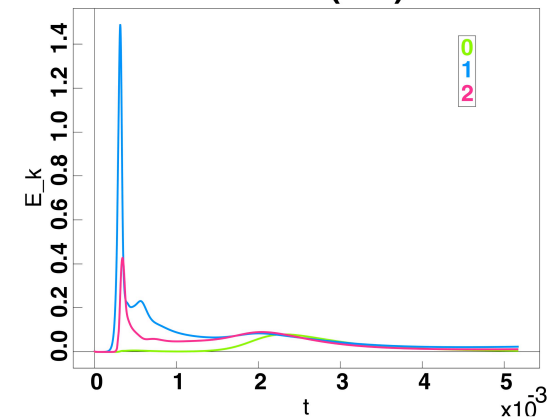
$t_{\text{exp}}=3600$



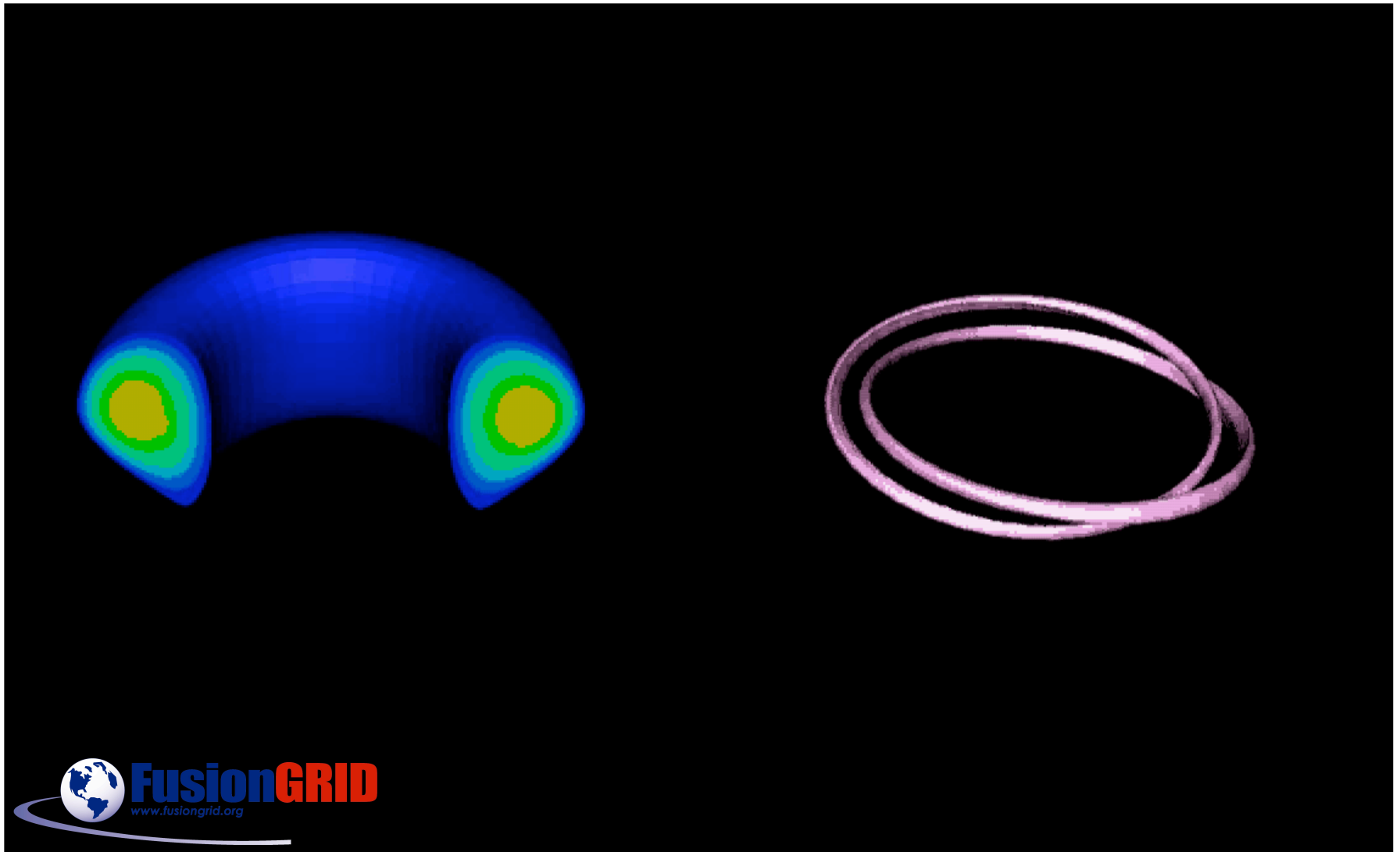
Unstable
n=2 grows
after n=1
drive begins
to reduce



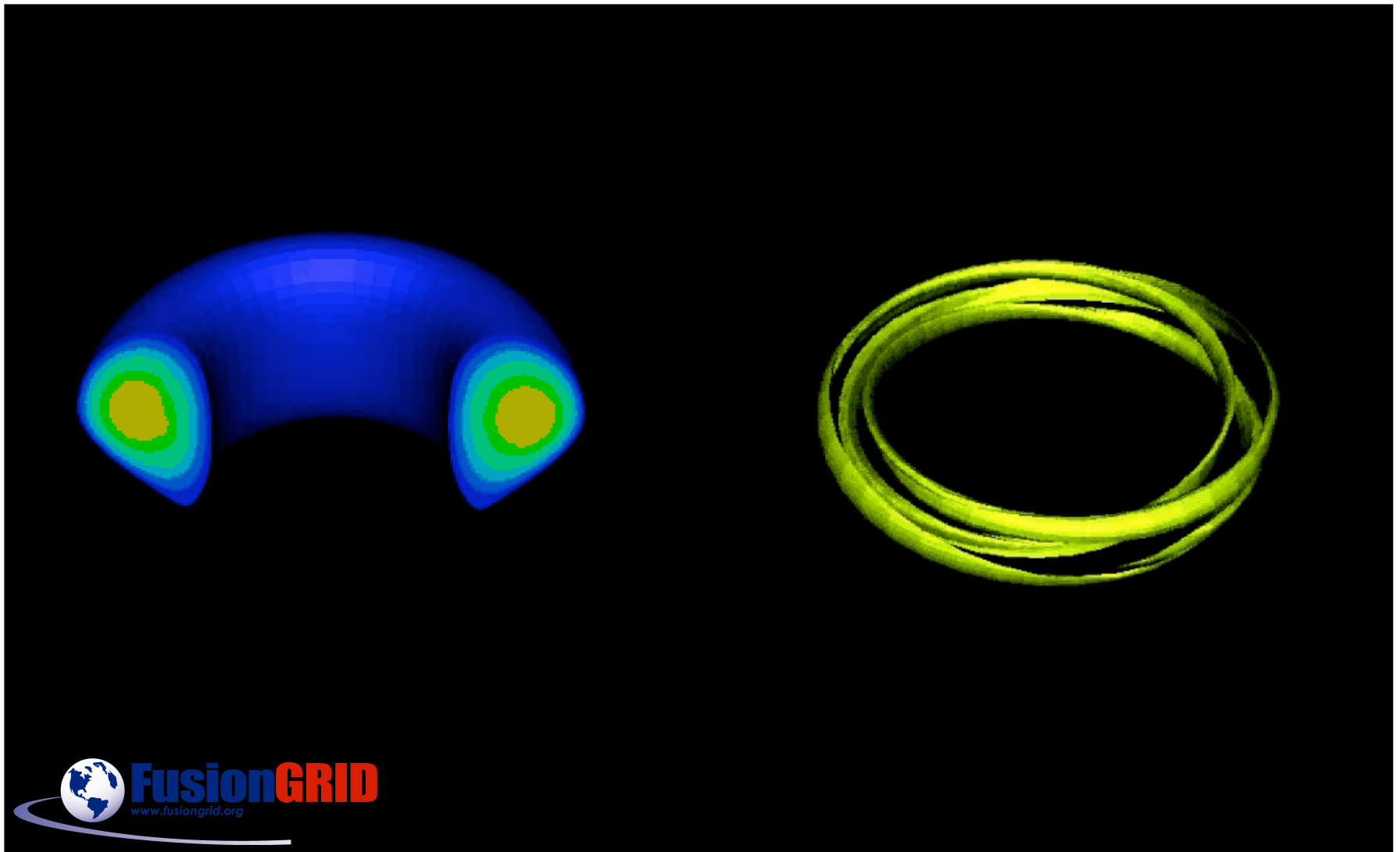
$S = \frac{\sigma_R}{\sigma_A} = 2.3e6$
approaching
realistic conditions



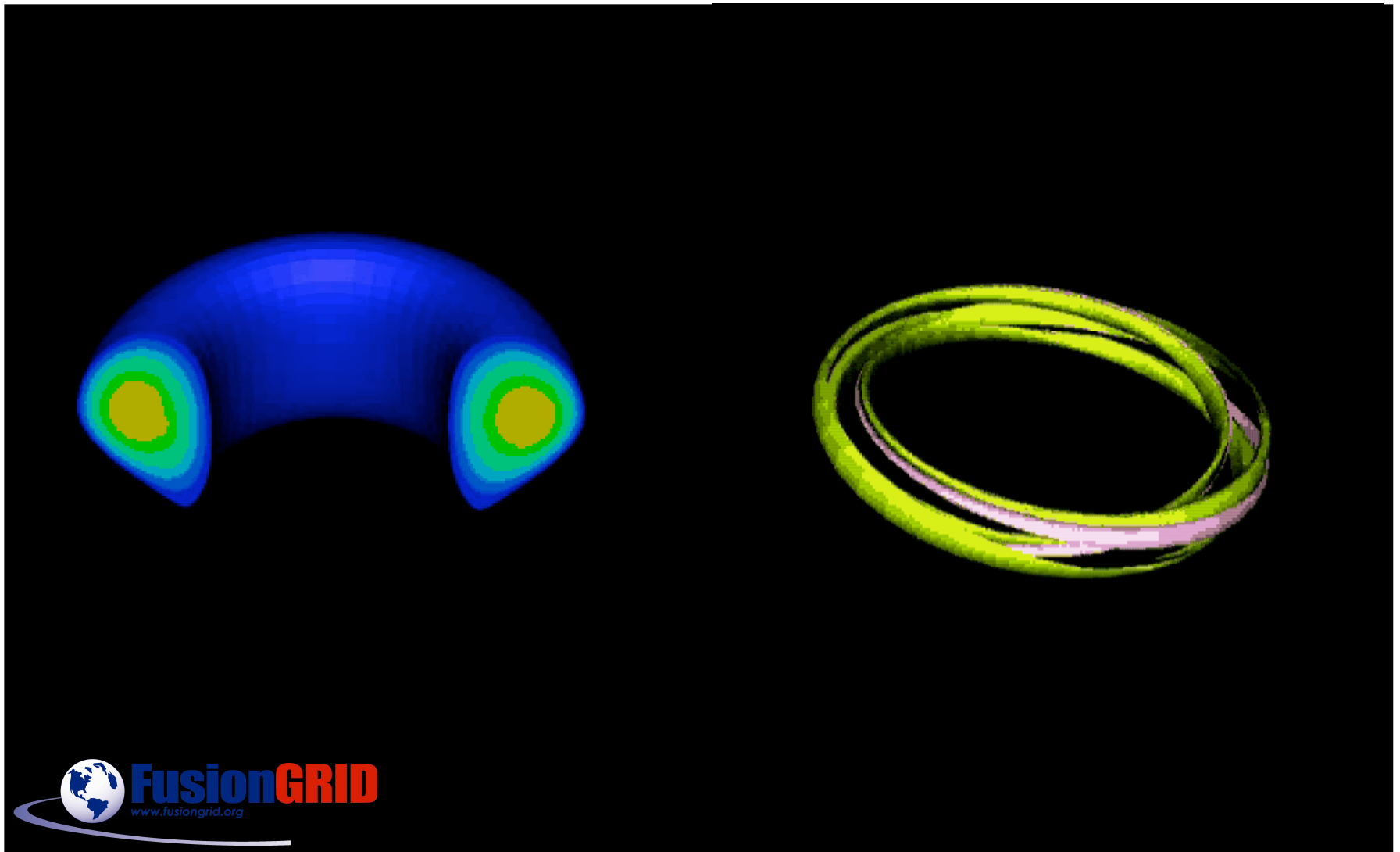
Studying the structure and nonlinear evolution of these modes with NIMROD can lead to new intuition and new physics



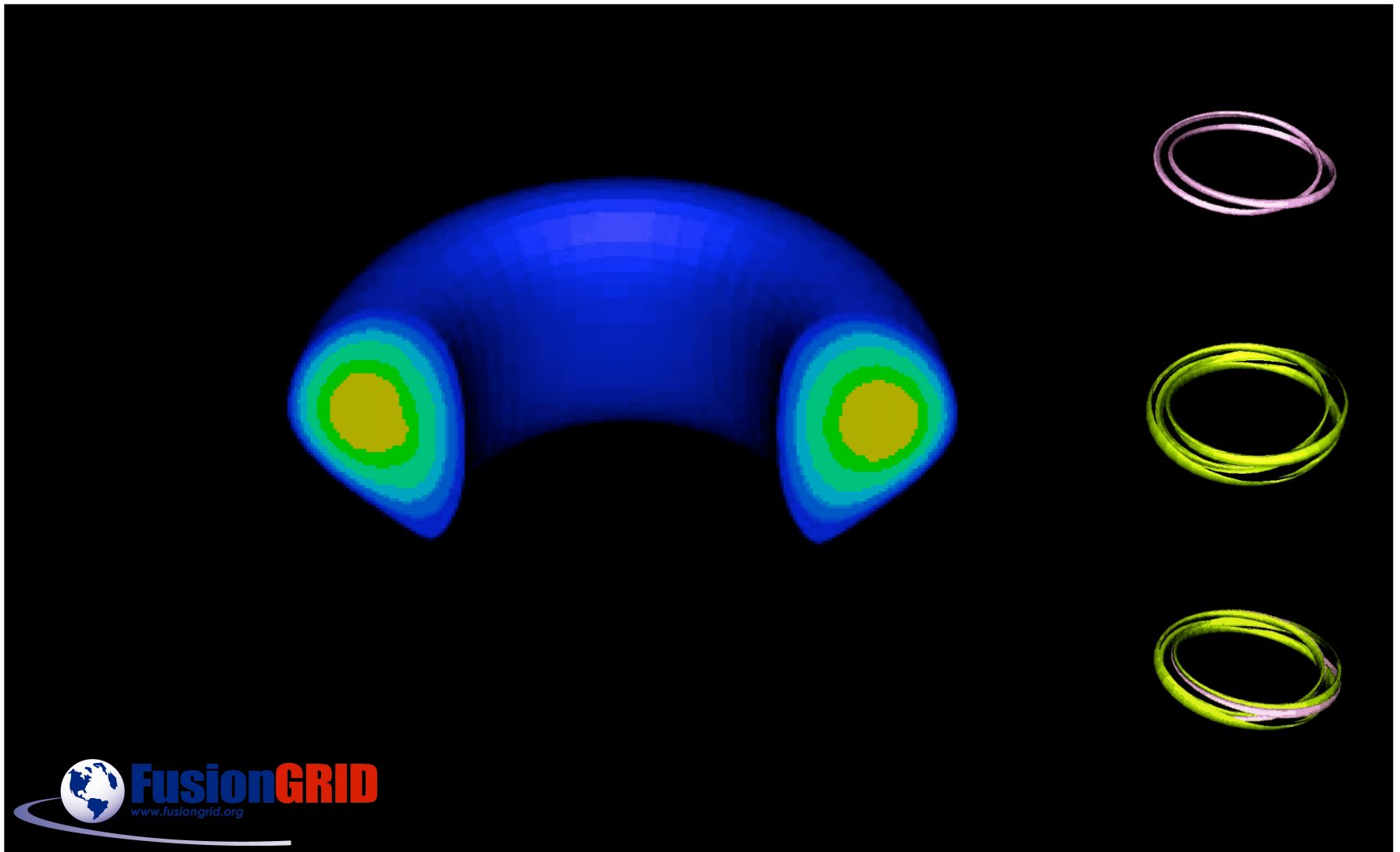
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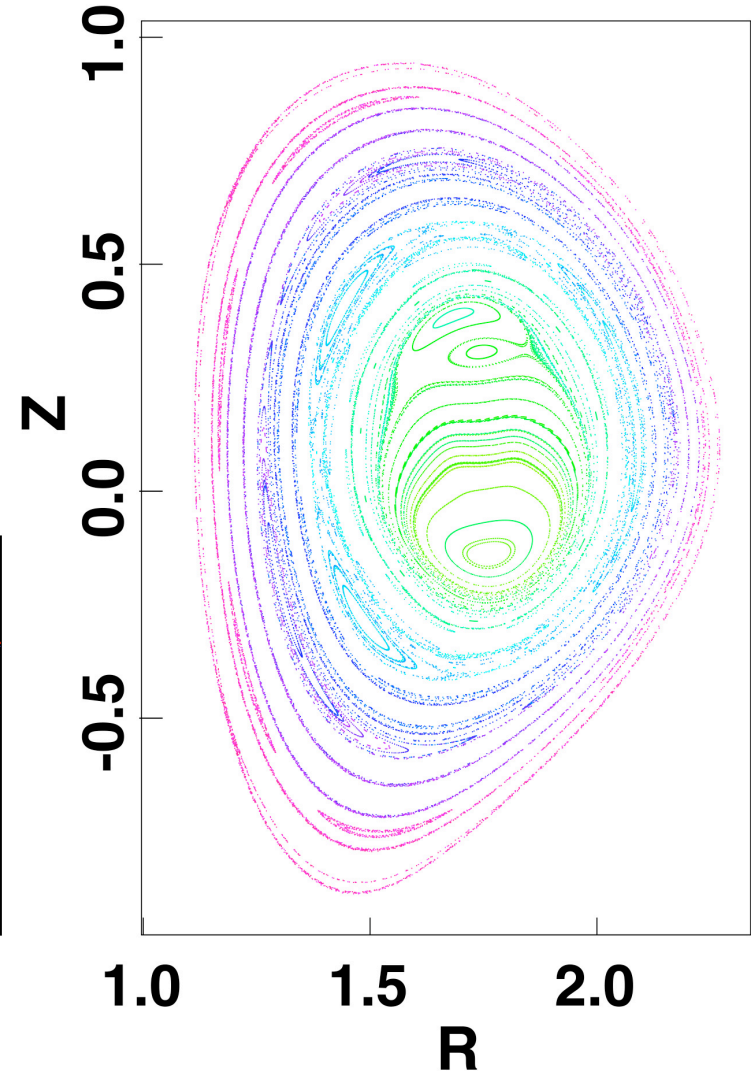
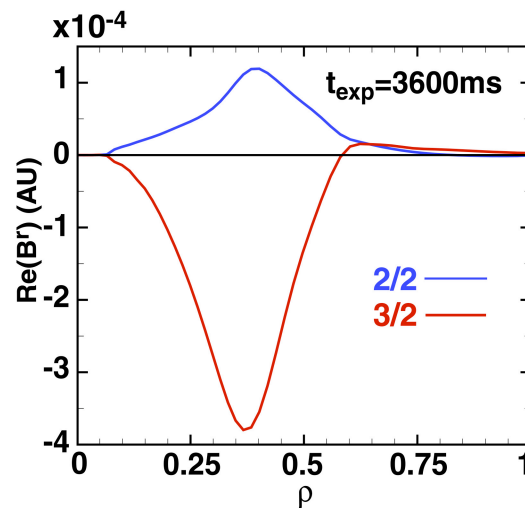
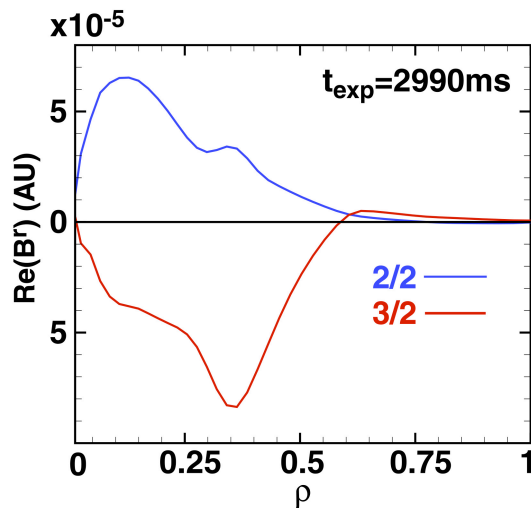
At later times, after saturation of $n=1$, the $m=3$ component of $n=2$ is dominant

$3/2$ decays in simulation of 2990ms

$3/2$ is larger and grows to saturation in simulation of 3600ms

Eigenfunctions show $m=3$ to be dominant $n=2$ component.

$m=1$ is dominant $n=1$.

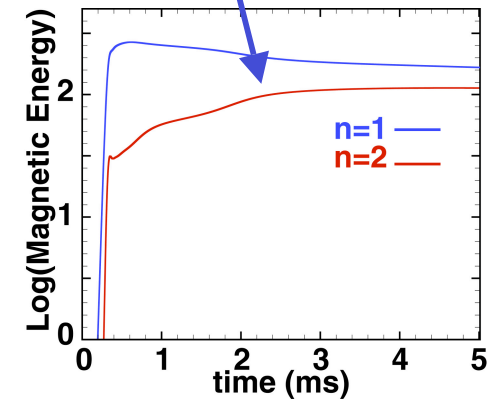
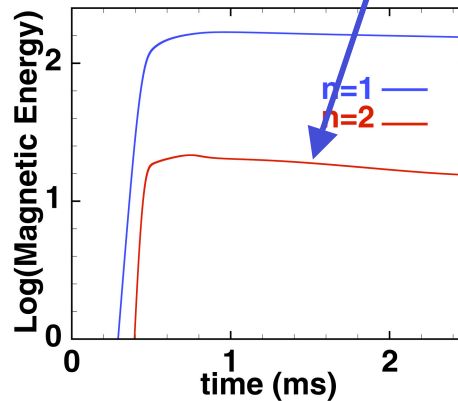
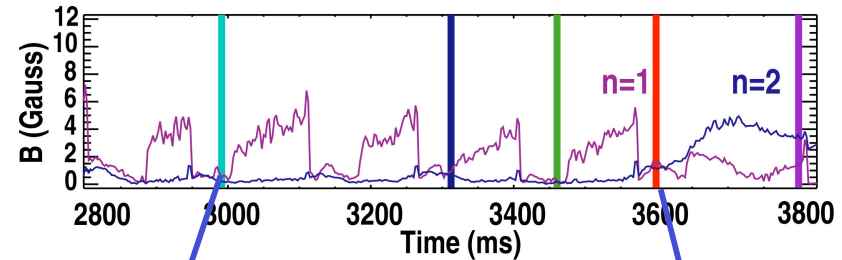


Dynamics of NIMROD Simulations Indicate That Linear Drive Affects Islands Of Finite Width

$m/n=3/2$ is unstable for 3600ms, when $n=1$ mode decays.

Nonlinear coupling effects do not prevent later seed island growth.

Equilibrium at 3600ms is closer to ideal limit and has larger linear drive.



The pole in ω affects the stability of NTM seed islands

DIII-D Experiment was Designed and Performed to Determine Effect of Poles in ω_N on Tearing Stability

Experiment designed to isolate ω_N pole mechanism

avoid other modes

vary $d\omega_N/dt$ on approach to onset of 2/1

Prediction was made that for spontaneous NTMs, evolution should depend on rate of approach to ideal boundary.

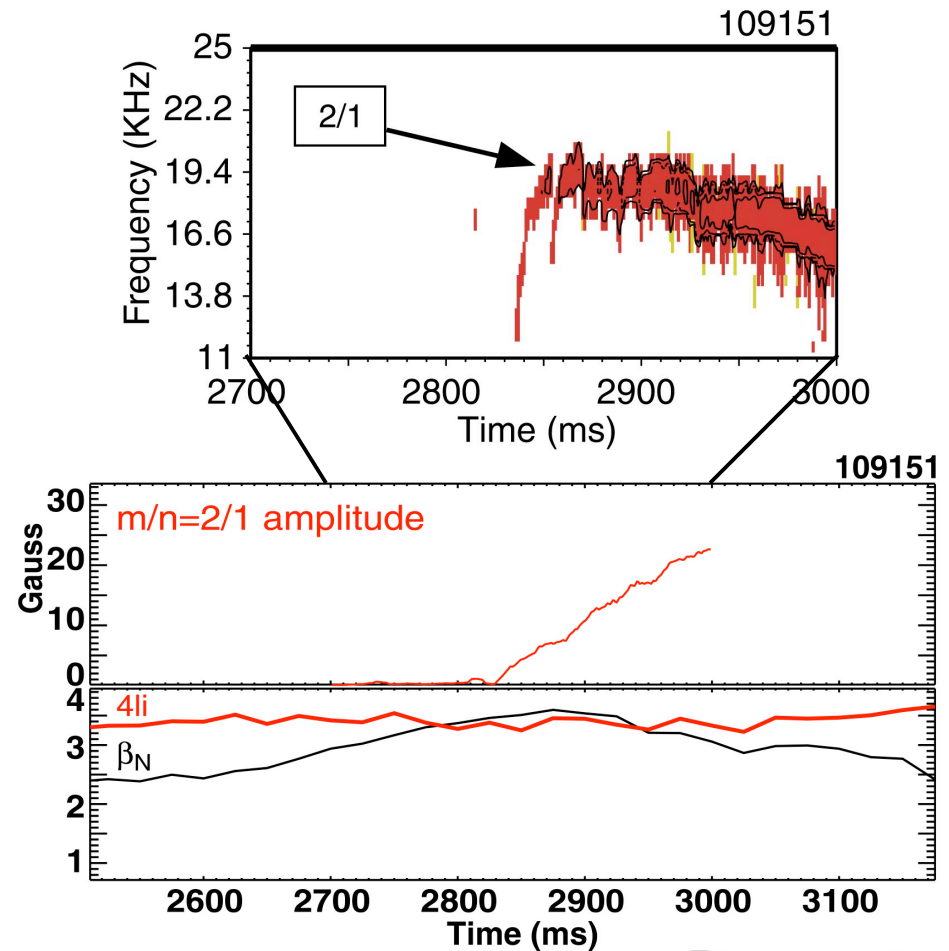
Measure ω_N at point where $w \sim w_d$ as a function of $d\omega_N/dt$

Forced reconnection gives random w at random t , and should not show a correlated function like $\omega(\omega)$

Spontaneous/Seedless NTMs were Generated to Isolate Effect of β_N Pole Onset Mechanism

Neutral beam injection applied at increasing rates up to β_N limit for 2/1 NTM onset.

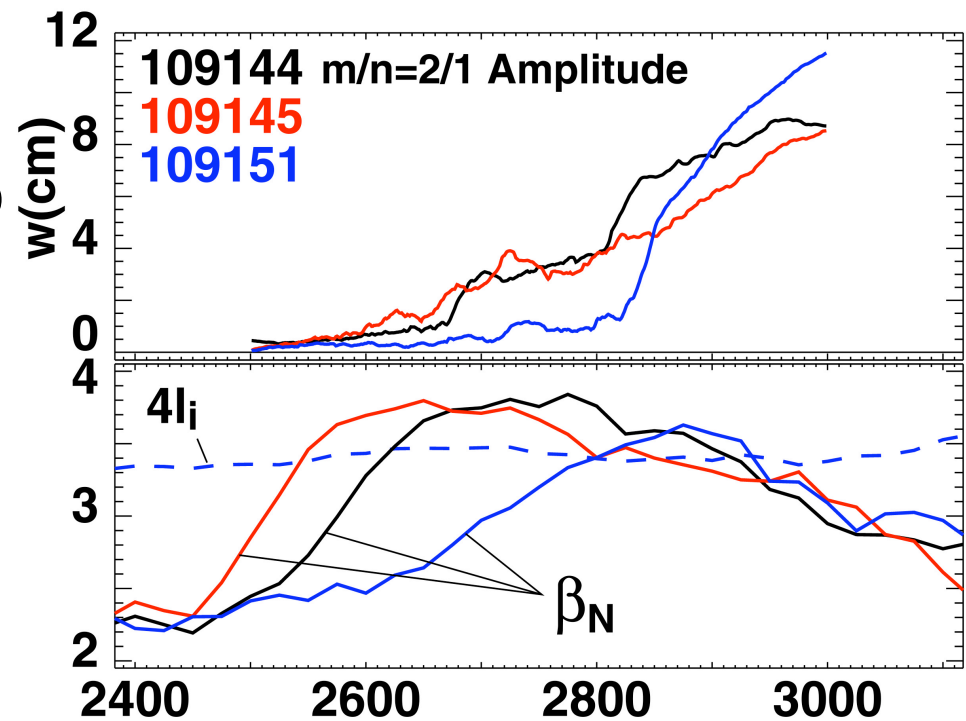
Experimental data indicates spontaneous onset in many cases.



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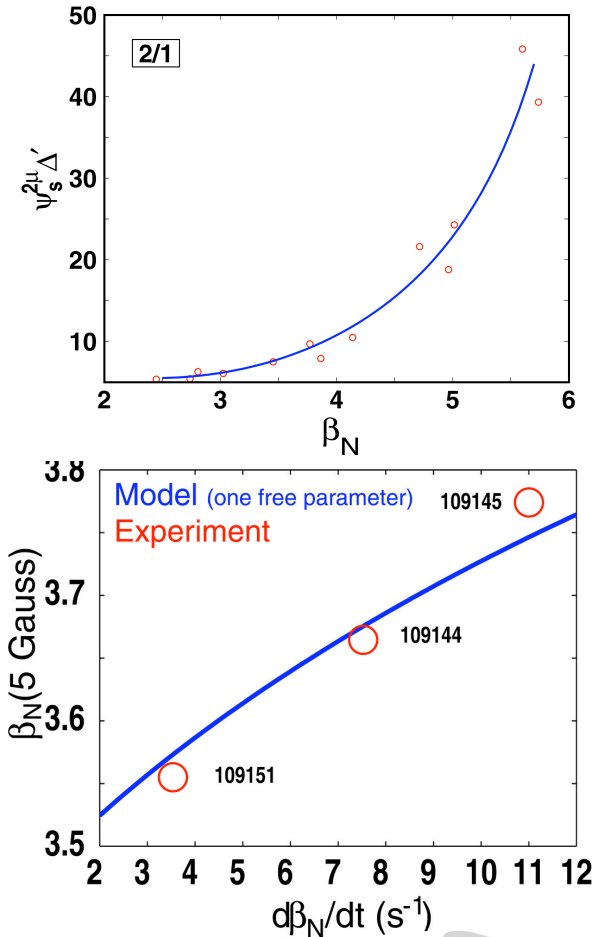
Experimental Data Confirms New Theoretical Prediction from $2/1$ Pole Model

Model uses $\beta_N(t)$ from experiment and $\psi_s(\beta_N)$ in island evolution equation.

Pressure profiles modified in model equilibria and ψ_s at $2/1$ surface calculated. Result is a function $\psi_s(\beta_N)$.

D_{pol} is single constant free parameter fit to find β_N vs. $d\beta_N/dt$ at mode onset.

Results support hypothesis that β_N is increasing rapidly in time, consistent with theoretical pole model.



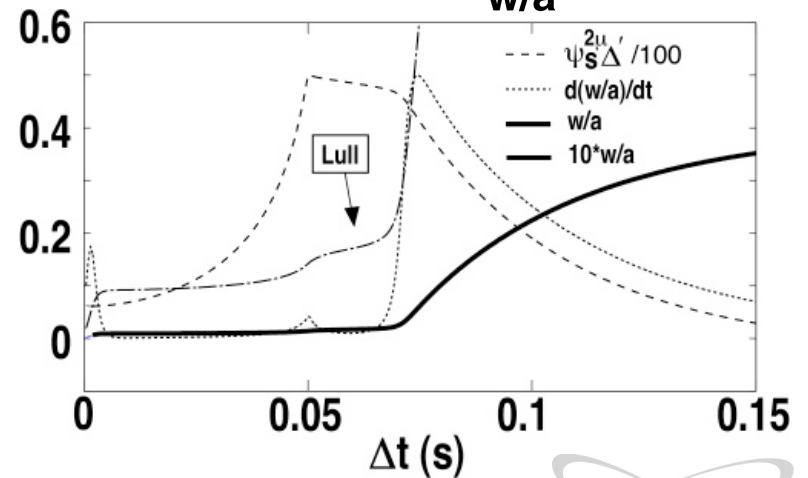
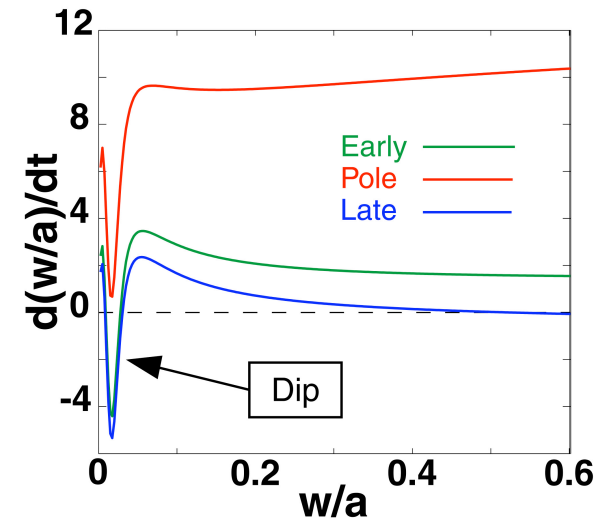
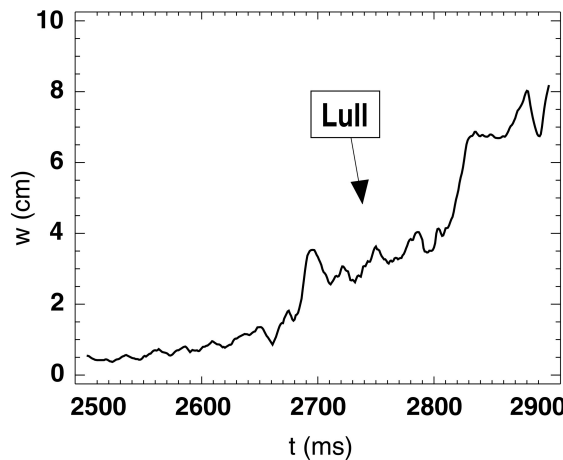
Modeling of β_N Pole Experiment Reproduces dw/dt Lull Shortly After Onset

D_{pol} term negligible at small island width. A model for D_{pol} is assumed which decays at $w < w_b$.

$$\frac{dw}{dt} = \frac{\eta^*}{k_0} \left(\Delta^* [\beta_N(t)] + \frac{w D_{tot}}{w_d^2 + w^2} + \frac{w D_{pol}}{w_b^4 + w^4} \right)$$

Resulting phase space plots have a Dip in growth rates.

$\beta_N(t)$ driving function reproduces Lull in growth



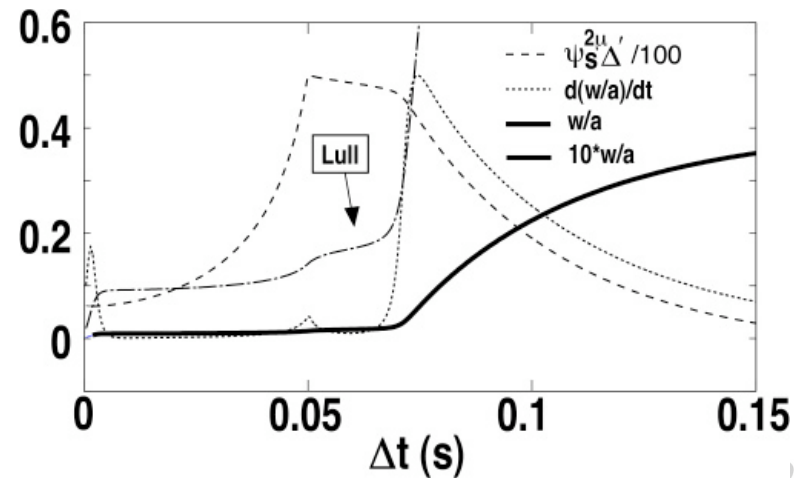
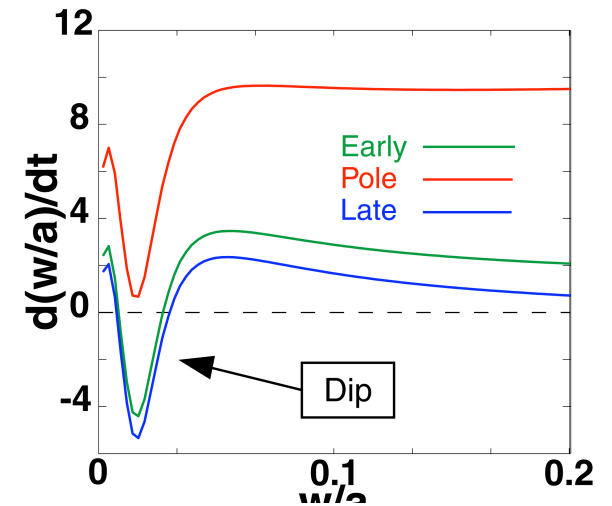
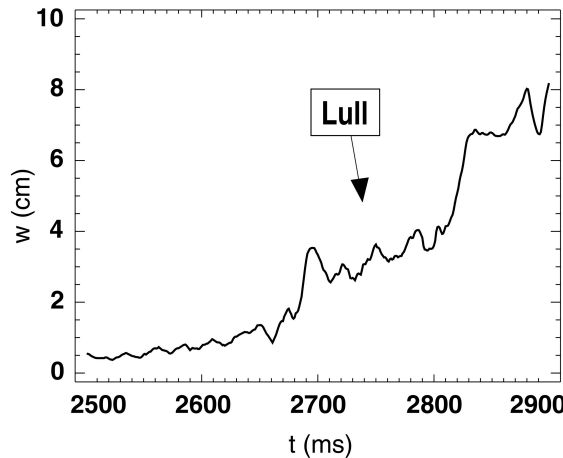
Modeling of β_N Pole Experiment Reproduces dw/dt Lull Shortly After Onset

D_{pol} term negligible at small island width. A model for D_{pol} is assumed which decays at $w < w_b$.

$$\frac{dw}{dt} = \frac{\eta^*}{k_0} \left(\Delta^* [\beta_N(t)] + \frac{w D_{tot}}{w_d^2 + w^2} + \frac{w D_{pol}}{w_b^4 + w^4} \right)$$

Resulting phase space plots have a Dip in growth rates.

$\beta_N(t)$ driving function reproduces Lull in growth



Summary of Results

- Increasing β sharply increases β_{crit} near ideal limit, due to approaching a pole, destabilizing NTMs
 - In a sawtooth seeding a NTM case, eventual onset was caused by increased β_{crit} due to approach of a pole
 - Nonlinear NIMROD simulations confirm that axisymmetric β_{crit} is meaningful for nonlinearly coupled finite islands
- DIII-D experiment confirmed effect of poles in β_{crit} on NTM stability with varying rates of $d\beta/dt$ causing spontaneous NTMs
 - Data from early evolution are in agreement with prediction, both as function of $d\beta_N/dt$ and detailed time evolution.