

Dependence of Achievable β_N on Discharge Shape and Edge Safety Factor in DIII-D Steady-State Scenario Discharges

by
J.R. Ferron

with T.C. Luce, P.A. Politzer, R. Jayakumar,^{*}
and M.R. Wade[†]

^{*}Lawrence Livermore National Laboratory.
[†]Oak Ridge National Laboratory

Presented at
the American Physical Society
Division of Plasma Physics Meeting
Long Beach, California

October 29 through November 2, 2001



Motivation

Motivation: high β_N is an important component of the self-consistent parameter set for a steady-state, advanced tokamak discharge

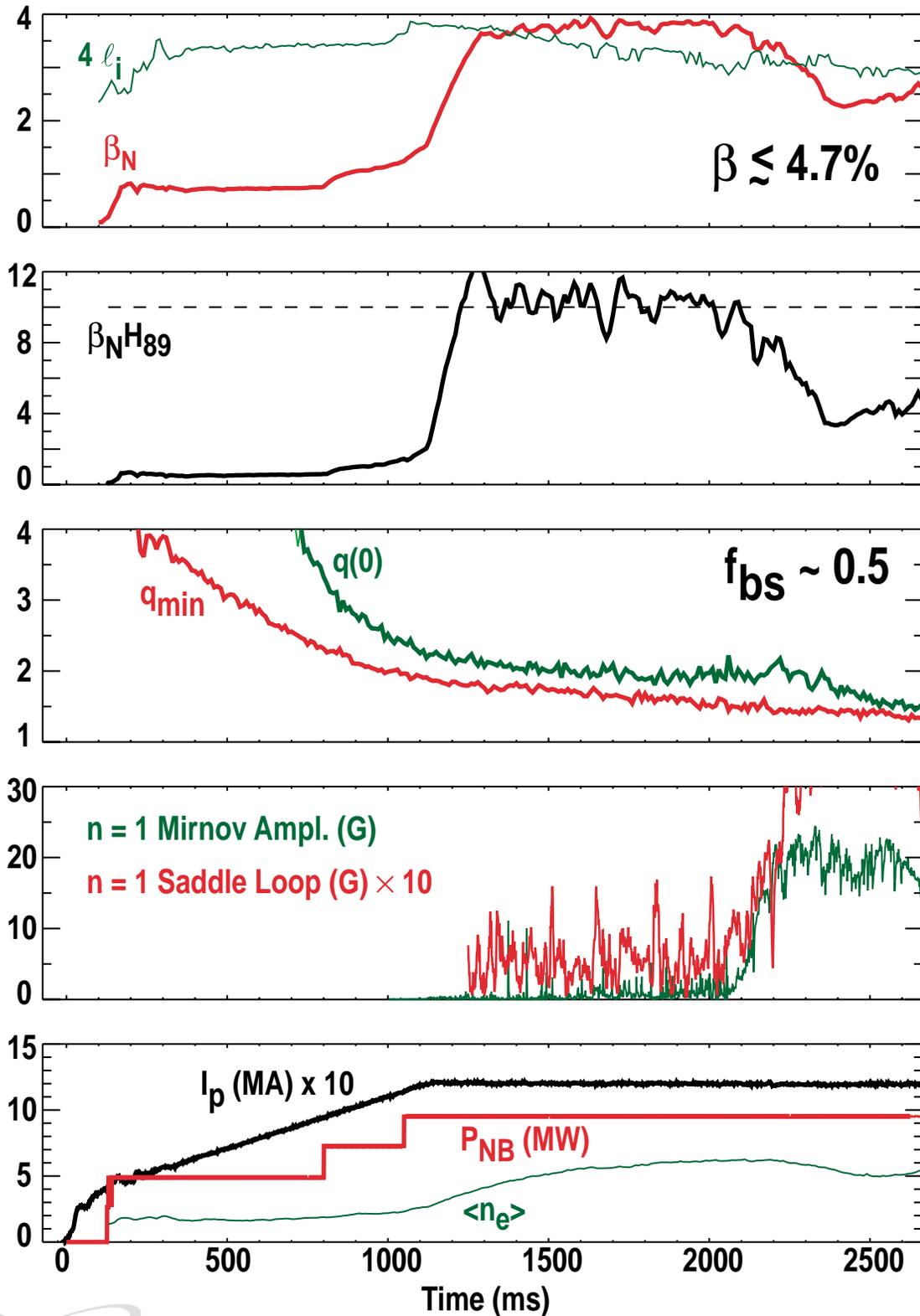
- High $\beta\tau_E \propto (\beta_N/q)(H_{89}/q^\alpha)$ for fusion power output.
- High bootstrap fraction $f_{BS} \propto \beta_P \propto q\beta_N$ for steady state.
- Optimum discharge has relatively high q (including $q(0)$, q_{min} , q_{95}) for f_{BS} , high β_N for power output.
- High β_N enhances off-axis ECCD efficiency.
- Low density required for efficient electron cyclotron current drive (ECCD).
- Low density requires pumping of H-mode divertor exhaust.

Background

Optimization of β_N was sought within the constraints of the advanced tokamak scenario

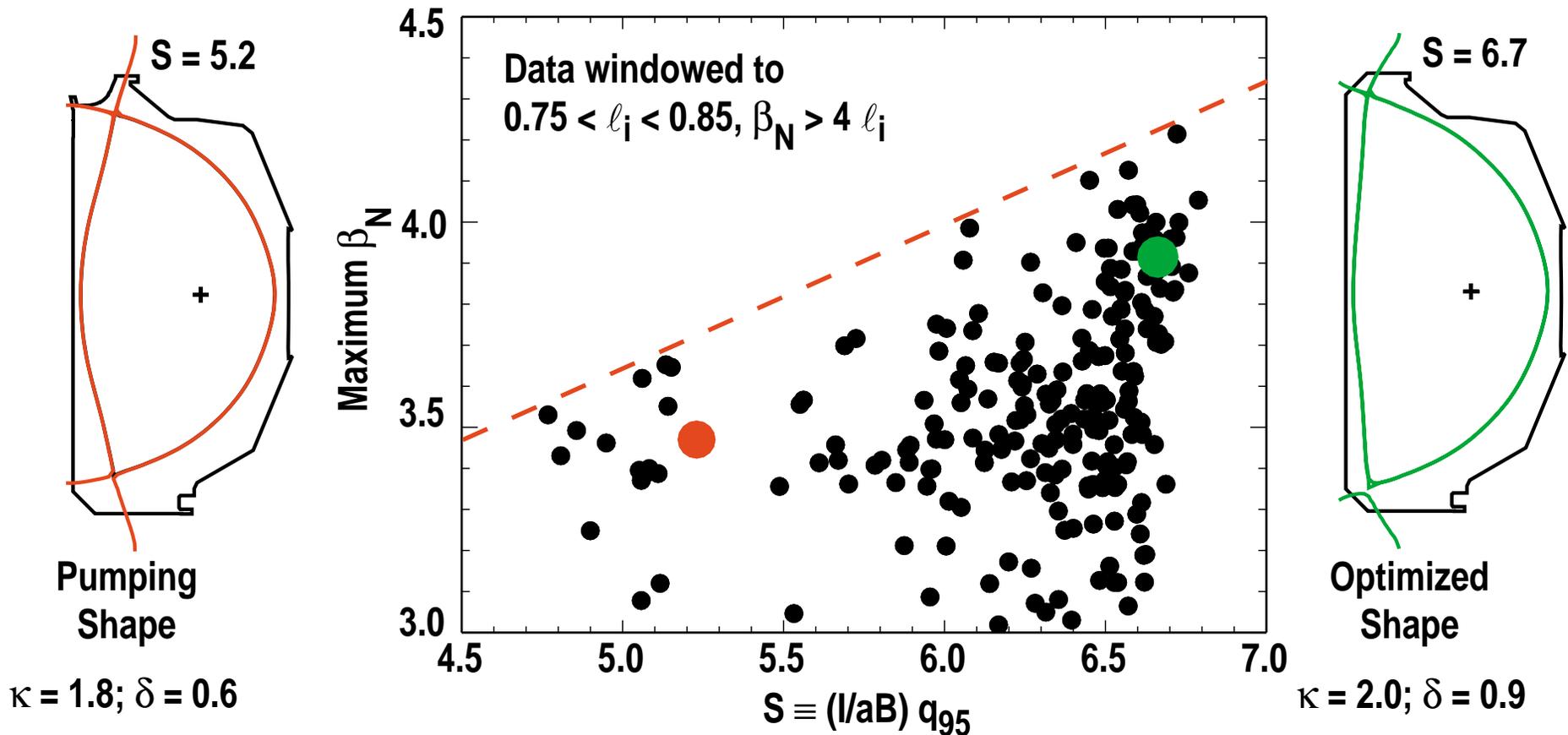
- Prior to upper divertor baffle installation (1999), $\beta_N H_{89} \approx 10$, $\beta_N \approx 3.7$ sustained in $\kappa = 2$, $\delta = 0.9$ shape.
- Achievable β_N (≈ 3.4) reduced in lower elongation, triangularity ($\kappa = 1.7$, $\delta = 0.7$) shape compatible with pumping after divertor baffle installation (2000).
- 1999-2000 operation at $B_T = 1.6T$ for ECCD.
- Data set shows a dependence of the achievable β_N on the shape parameter $S = (I/aB)q_{95}$.
- Leads to 2001 investigation of the effect of variation in shape and q_{95} (by changing B_T) on achievable β_N .

HIGH NORMALIZED PERFORMANCE (~ 10) SUSTAINED FOR $5\tau_E$ IN $\kappa = 2, \delta = 0.9$ SHAPE



OBSERVED β_N LIMIT HAS A SIGNIFICANT DEPENDENCE ON PLASMA SHAPE

- Comparison between shapes available before and after divertor baffle installation



- Experimental scans at constant I/aB

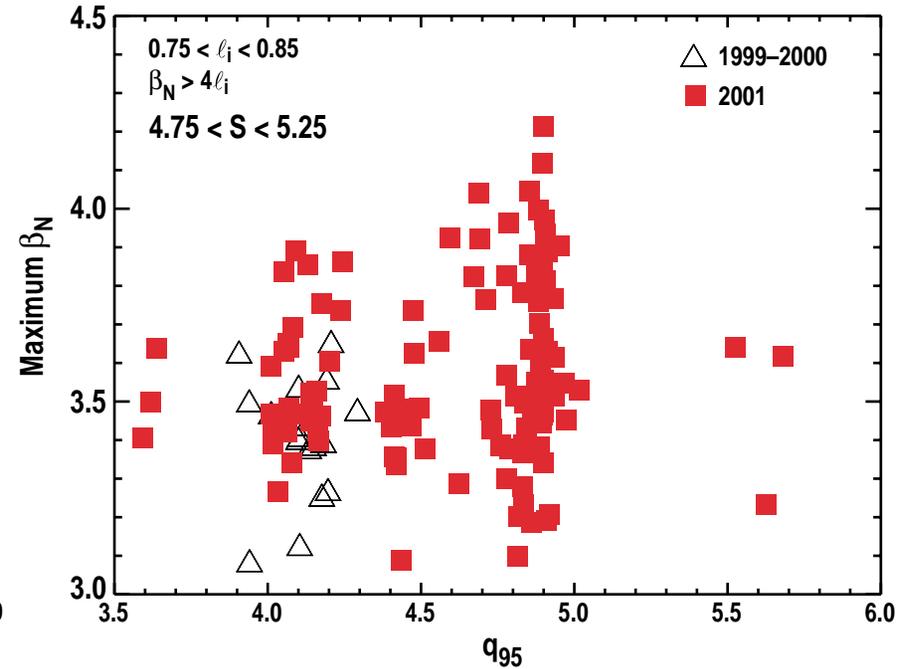
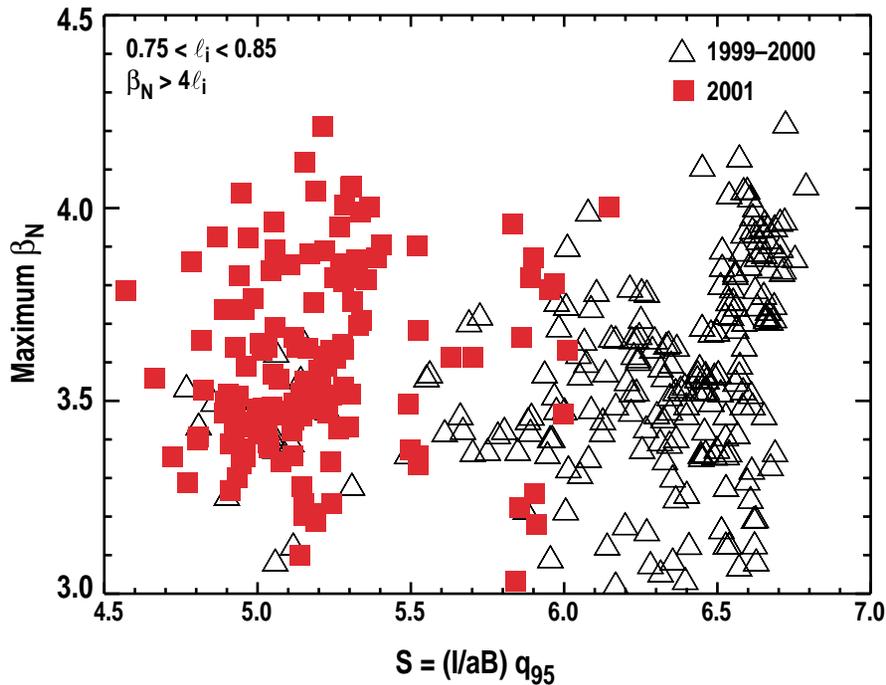
Experimental results

Experimentally achievable β_N increased at higher q_{95}

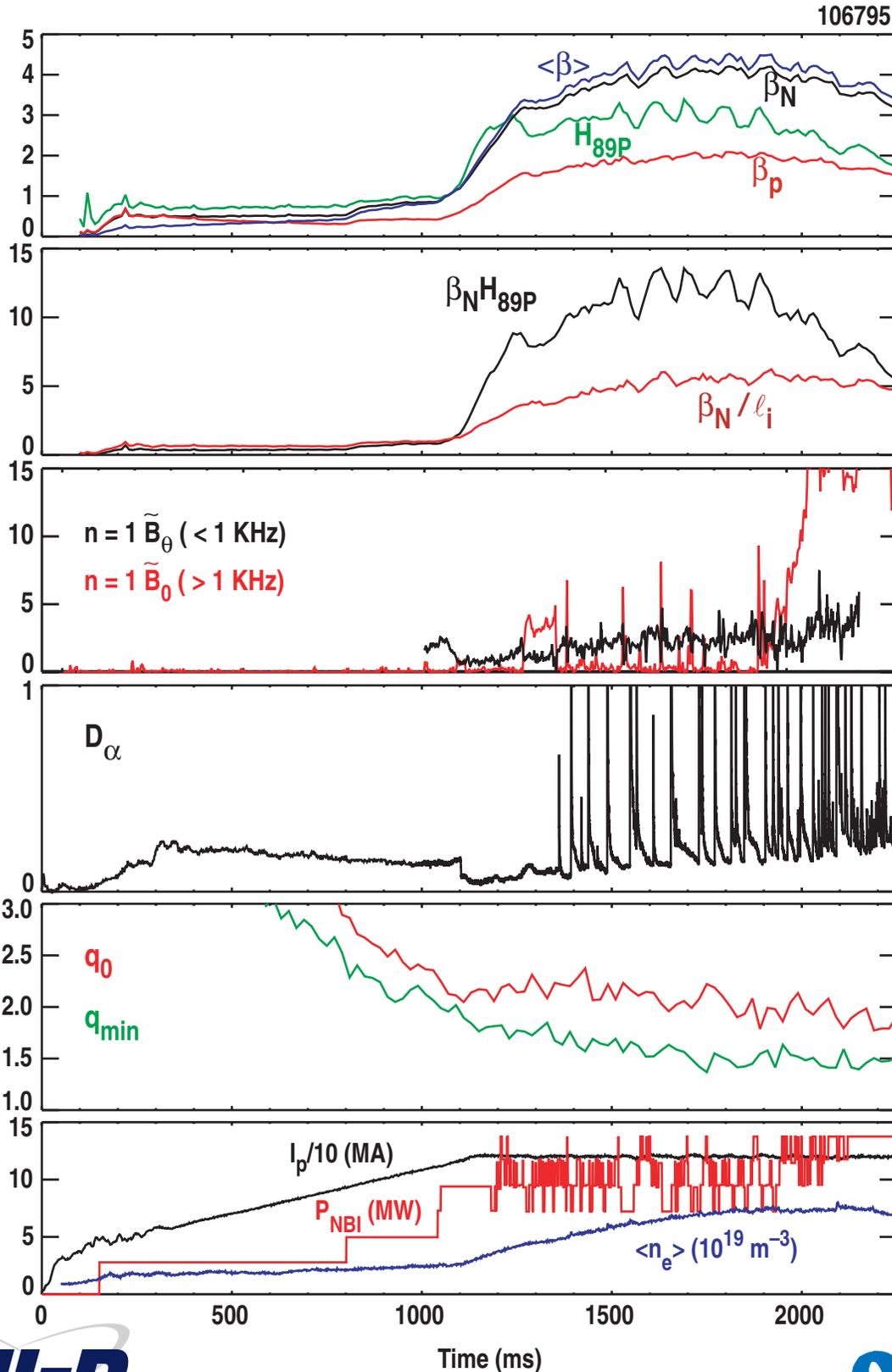
- $\beta_N = 4.1$ sustained at $B_T = 1.85T$ in $\kappa = 1.8$, $\delta = 0.7$ shape (compared to $\beta_N = 3.4$ at $B_T = 1.6T$).
- Improvement over β_N in the older higher elongation, triangularity shape.
- Provides a satisfactory β_N value for an ECCD target.
- No improvement in β_N was found by varying the discharge shape near the pumping geometry.
- β_N limited, high performance phase terminated by instability:
 - Resistive wall modes (RWM).
 - ELMs
 - Neoclassical tearing modes (NTM) at $m/n = 2/1$.
- Maximum β_N :
 - Up to $5.5\ell_i$.
 - Above the ideal, no-wall $n = 1$ limit.
 - In some cases, close to the ideal-wall limit.

EXPERIMENTAL STUDIES INDICATE β_N INCREASES WITH INCREASING q_{95}

- Increasing DRSEP causes a drop in the shape parameter $S = (I/aB) q_{95}$ and q_{95} itself
- 1999-2000 studies indicated variation of RWM β limit with shape parameter and q_{95}

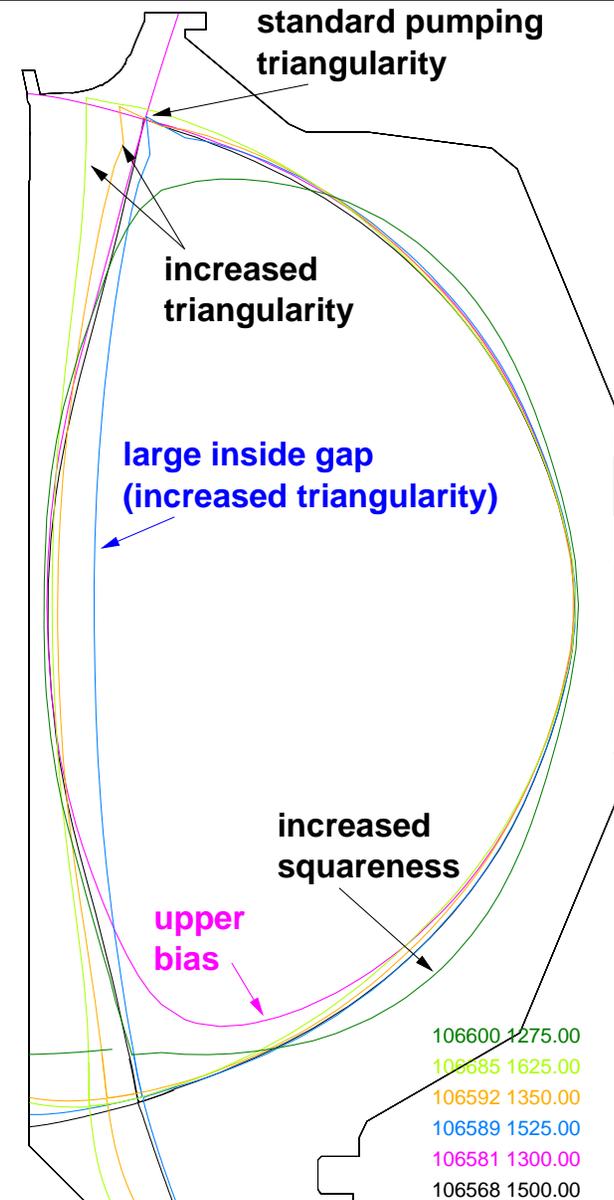
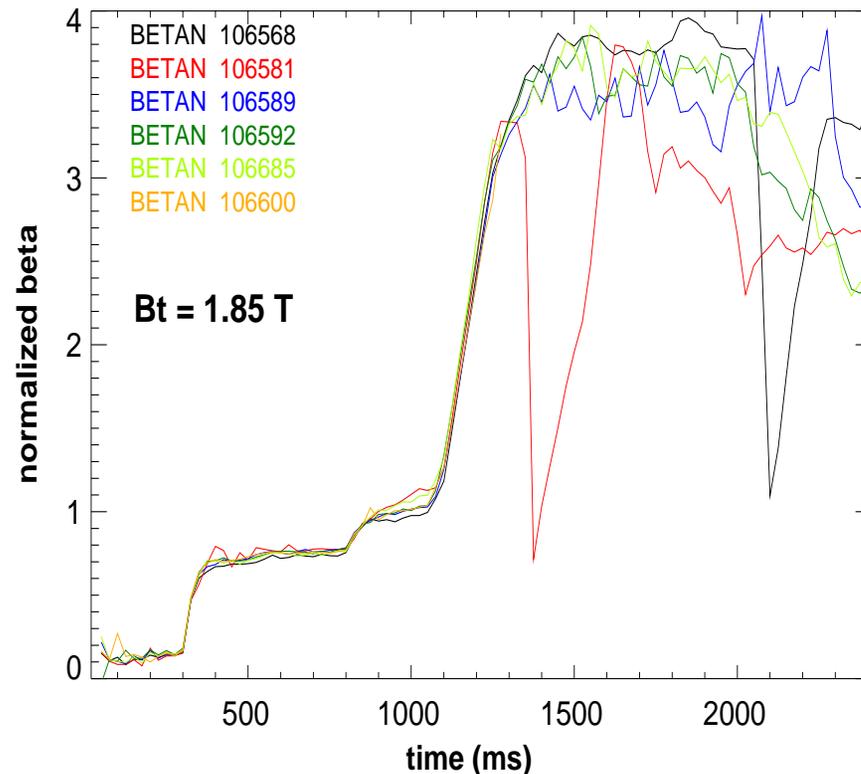


CONDITIONS CONDUCTIVE TO HIGH f_{BS} AND HIGH $\beta\tau$ ($\beta_N = 4$, $H_{89} = 3$ WITH $q_{min} > 1.5$) SUSTAINED FOR $5\tau_E$



No systematic variation in β_N resulted from many small variations about the standard pumping shape

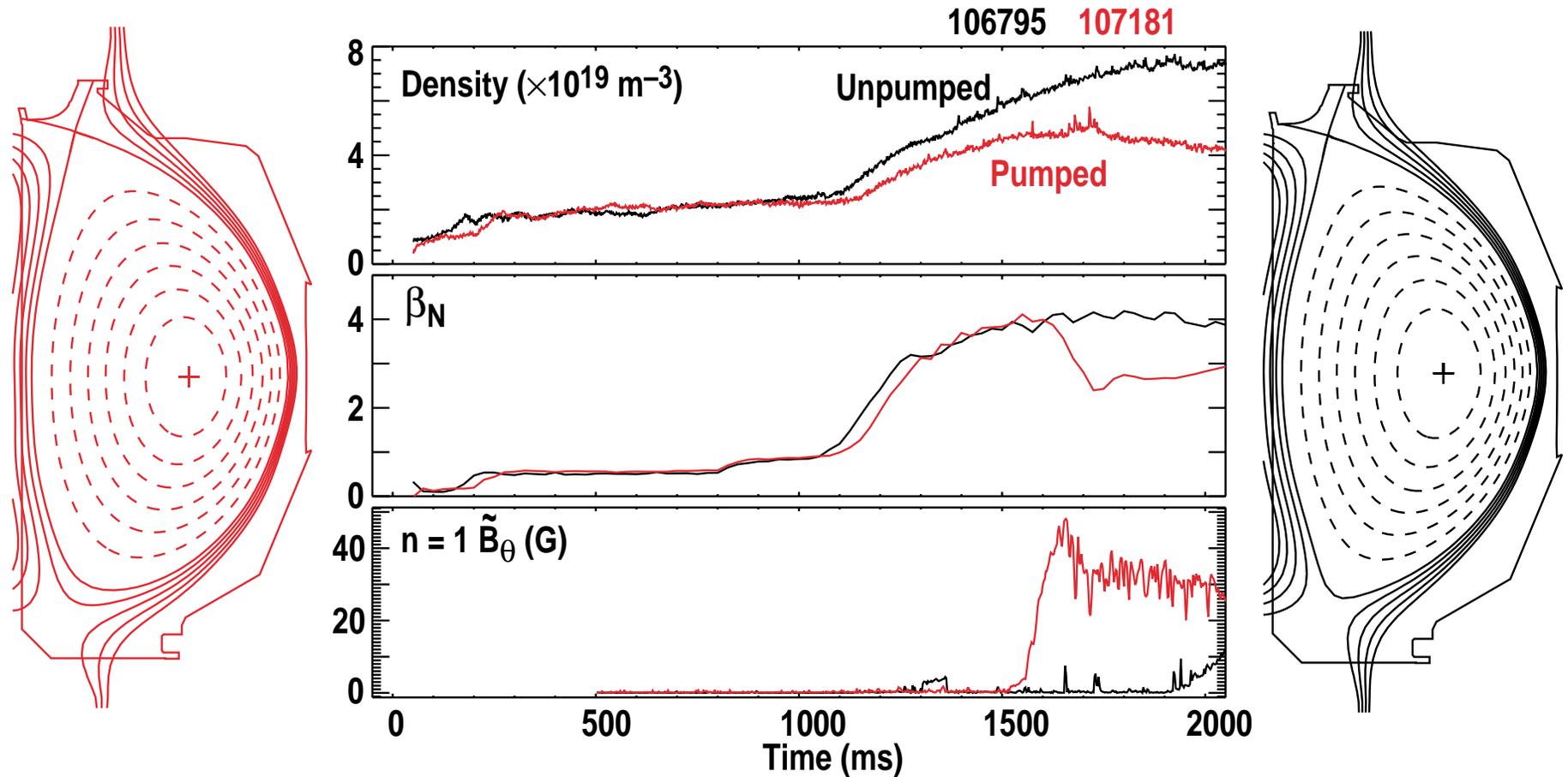
- No opportunity to optimize most shapes.
- Further analysis and modeling comparison required.



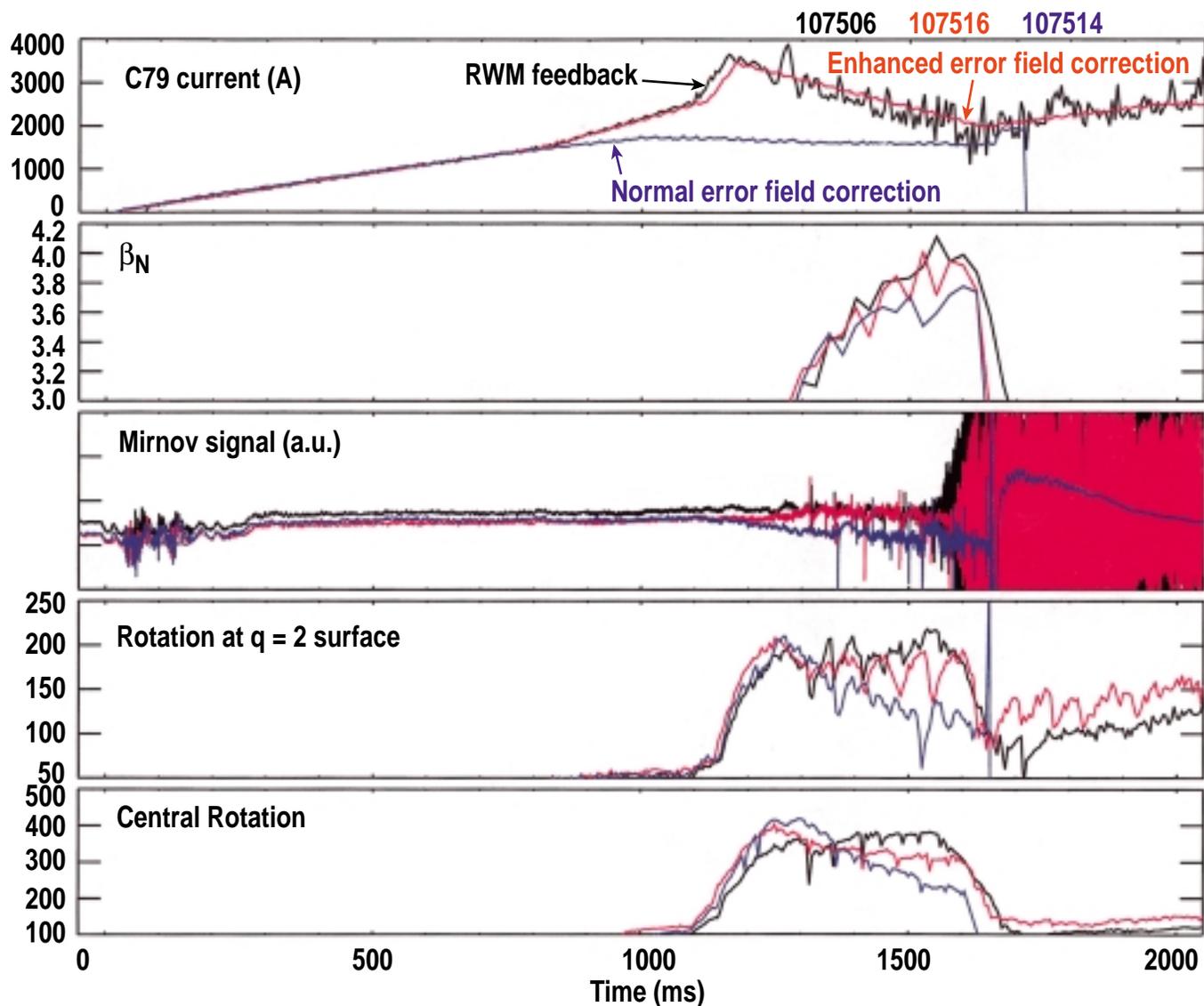
$\beta_N \sim 4$ HAS BEEN ACHIEVED WITH THE UPWARD SHAPE BIAS REQUIRED TO MAINTAIN DENSITY CONTROL

Pumped Shape

Unpumped Shape



ERROR FIELD CORRECTION AND RWM FEEDBACK STABILIZATION HAVE ALSO BEEN SUCCESSFUL IN IMPROVING THE ACHIEVABLE β_N



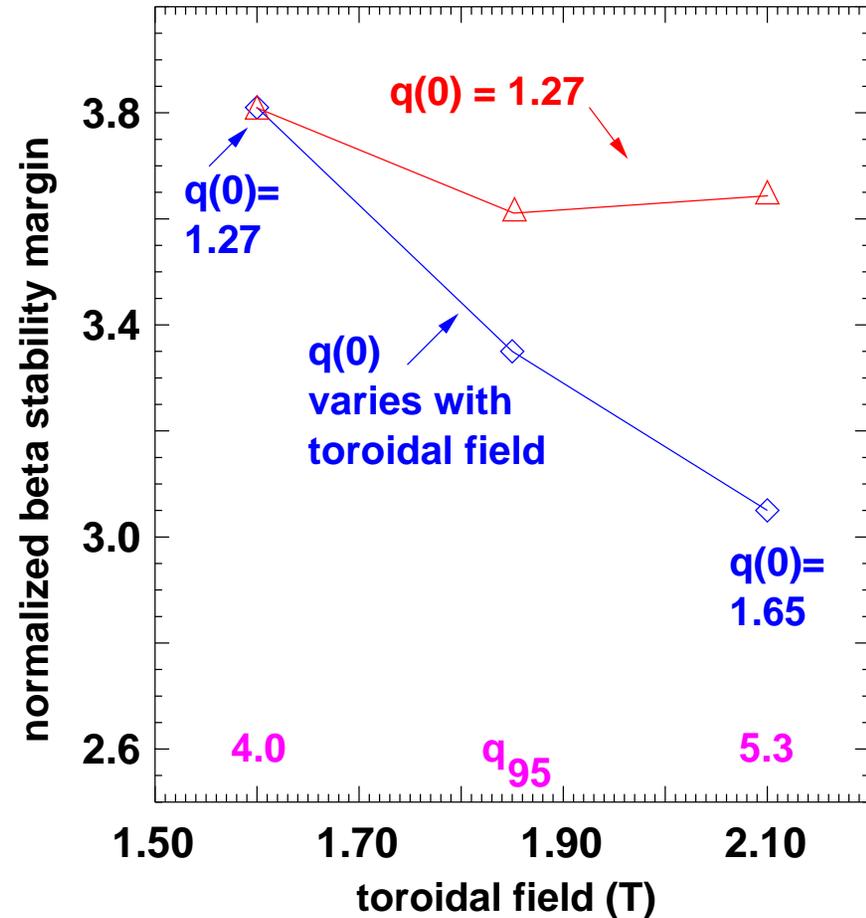
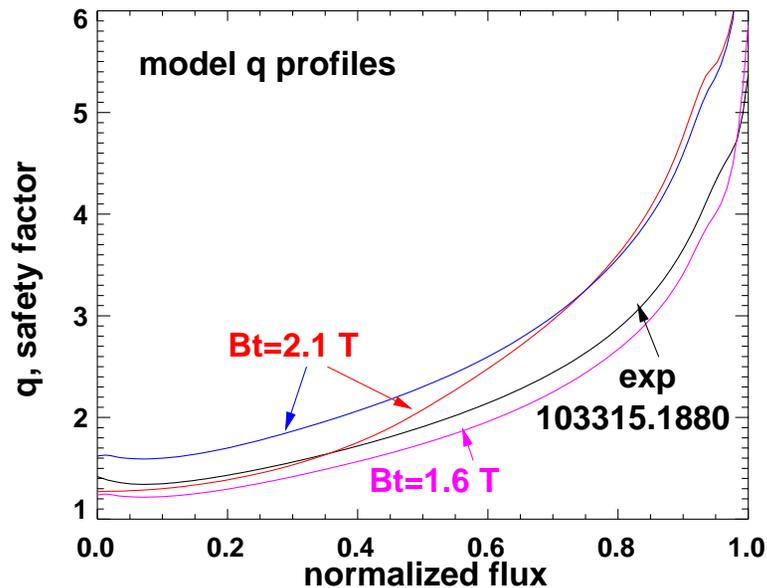
Modeling of no-wall limit

Modeling based on ideal $n = 1$ no-wall theory cannot account for the changes in β_N observed in the experiment

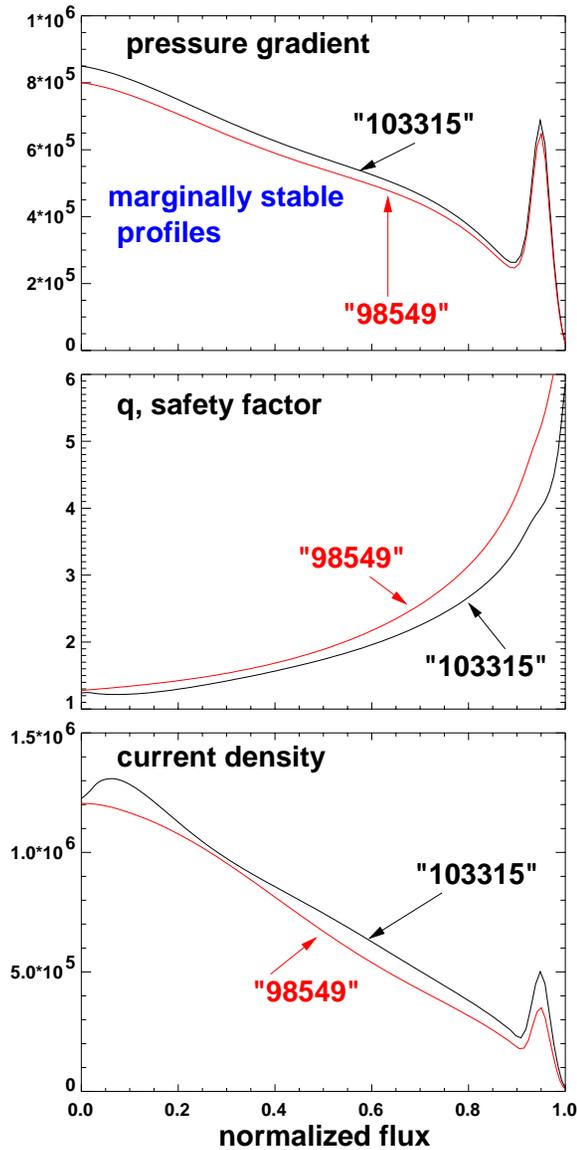
- In contrast to the experiment, the predicted no-wall $n = 1$ limit:
 - Decreases as q_{95} increases.
 - * Scan of B_T at constant shape.
 - * Change in I_p at constant shape.
 - Is lower in the higher elongation, triangularity shape.
- Modeling also predicts that the no-wall $n = 1$ limit:
 - Is optimized at lower $q(0)$.
 - Is not strongly dependent on details of the pressure gradient profile shape.
- Modeling:
 - Equilibria computed with the code TOQ.
 - Current and pressure gradient profiles modeled from measured experimental profiles.
 - Ideal stability computed with GATO with or without a conformal wall or a wall at the DIII-D vessel position.

Stability modeling results show that the marginal value of β_N is constant or decreasing as B_T and q_{95} are increased

- No-wall β_N limit for $n = 1$ from GATO/TOQQ.

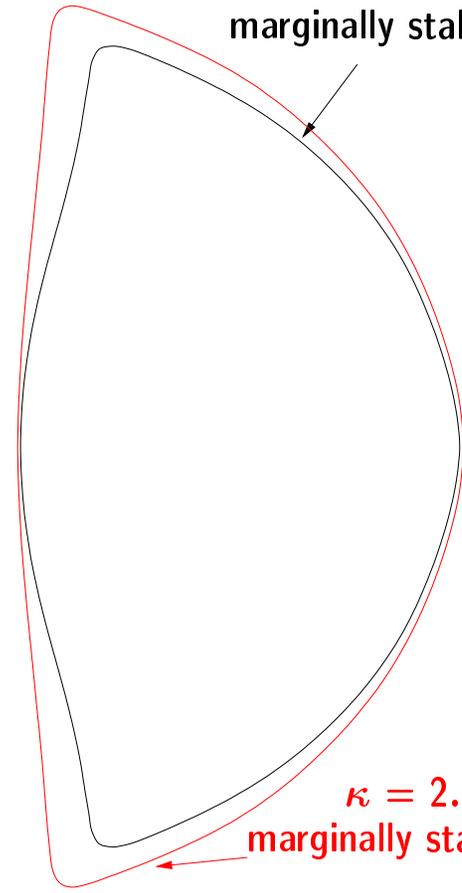


No increase in the ideal, no-wall, $n = 1$ stability limit with extra shaping has been found in initial modeling results



TOQ model of 103315.01880

$\kappa = 1.8, \delta = 0.6, q_{95} = 4$
marginally stable at $\beta_N = 3.74$



Reducing q_{95}
(increasing I_p)
increases
normalized beta
limit

$\kappa = 2.0, \delta = 0.8, q_{95} = 5.3$
marginally stable at $\beta_N = 3.48$

($\beta_N = 3.66$ at $q_{95} = 4$)

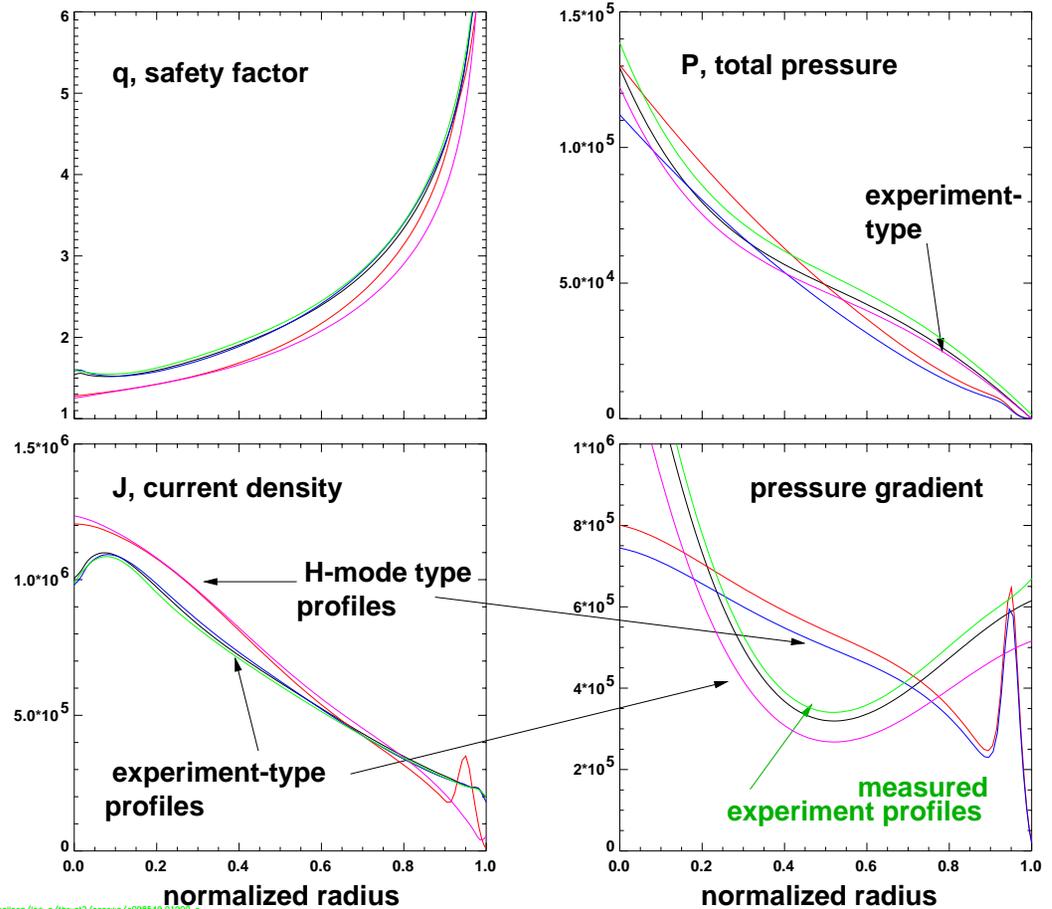
model of 98549.01900

/u/terron/balloon/low_n/thrust2/case#a/g800001.00112
/u/terron/balloon/low_n/thrust2/case#a/g800001.00076

Sensitivity of the β_N limit to the model profiles was studied for the more strongly shaped case

	H-mode type J profile ($q(0) \approx 1.27$)	experiment-type J profile ($q(0) \approx 1.6$)
H-mode P' profile	3.52	3.15
Experiment-type P' profile	3.34	3.62

- β_N limit varies by at most 0.5.
- Strongly shaped discharge had an unusually broad pressure profile.



/u/ferran/balloon/low_n/thrust2/casewa/g98549.015002

/u/ferran/balloon/low_n/thrust2/casewa/g800001.00116

/u/ferran/balloon/low_n/thrust2/casewa/g800001.00010

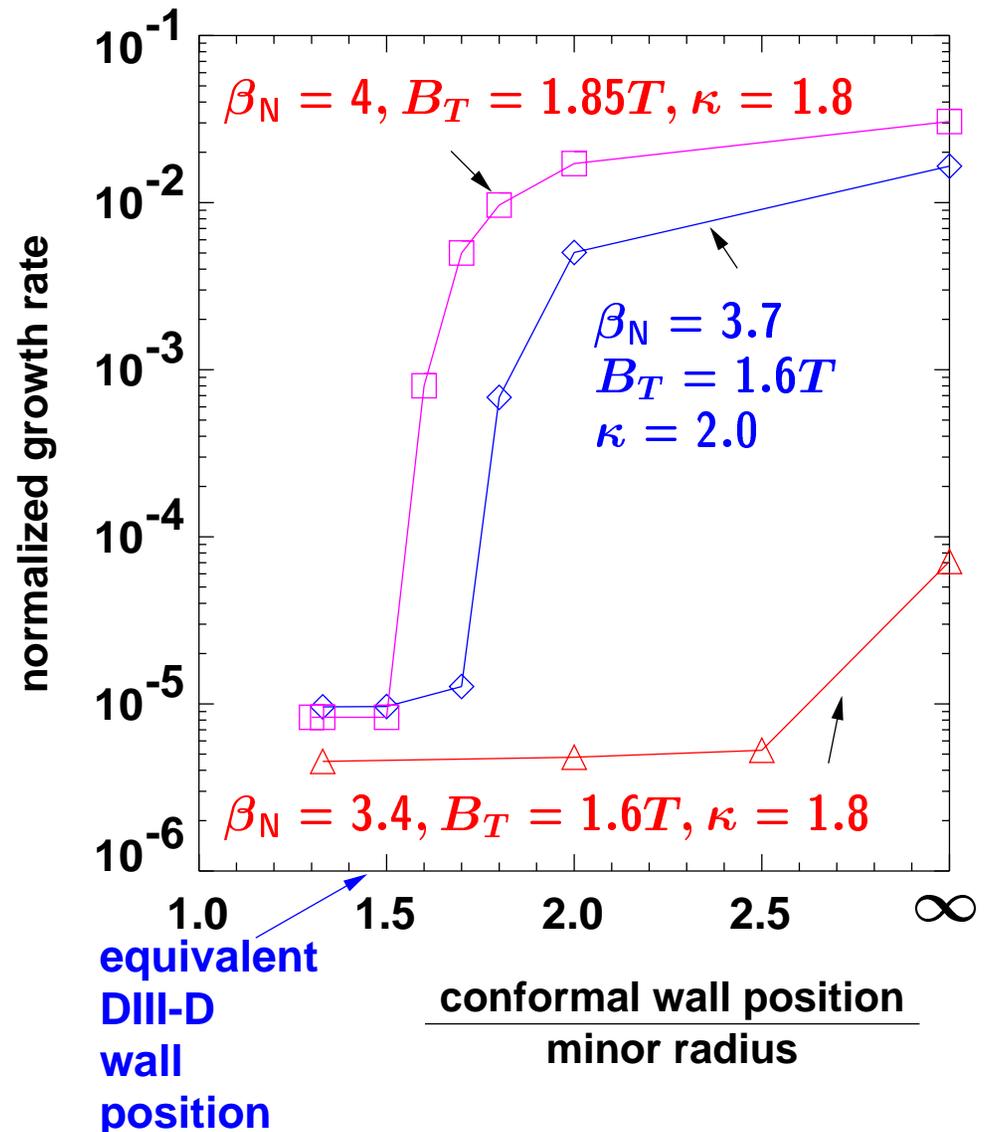
/u/ferran/balloon/low_n/thrust2/casewa/g800001.00128

/u/ferran/balloon/low_n/thrust2/casewa/g800001.00112

Effect of an ideal wall

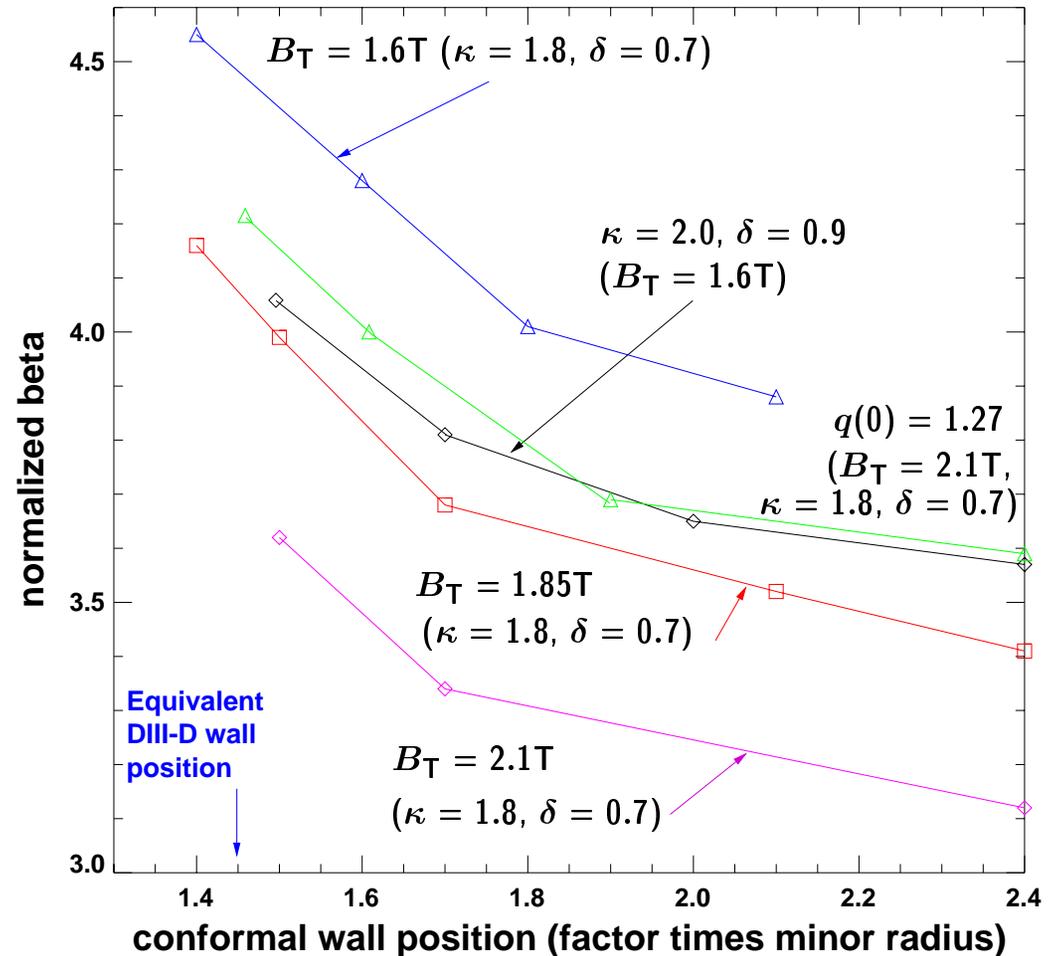
Increases in β_N achieved by exceeding the no-wall limit by a wider margin

- The no-wall limit is exceeded if the ideal wall must be present for stability.
- The more the no-wall limit is exceeded, the closer the stabilizing wall must be to the plasma.



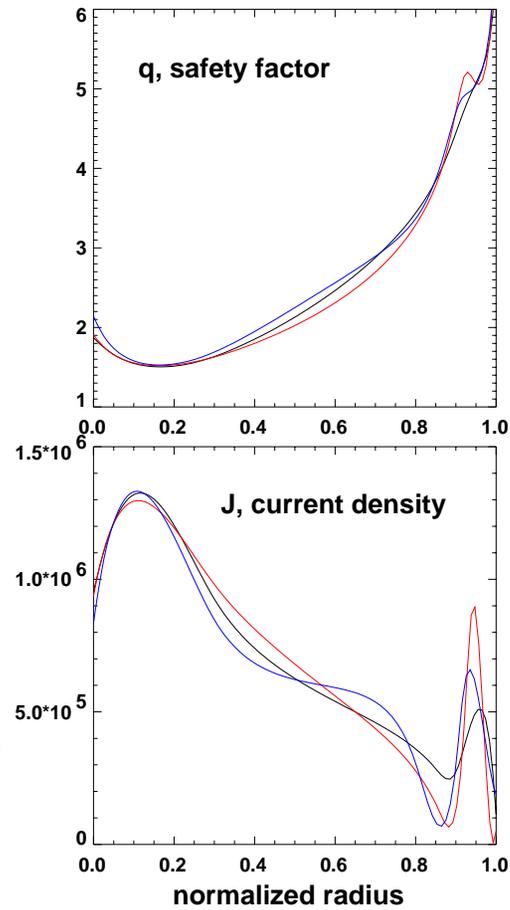
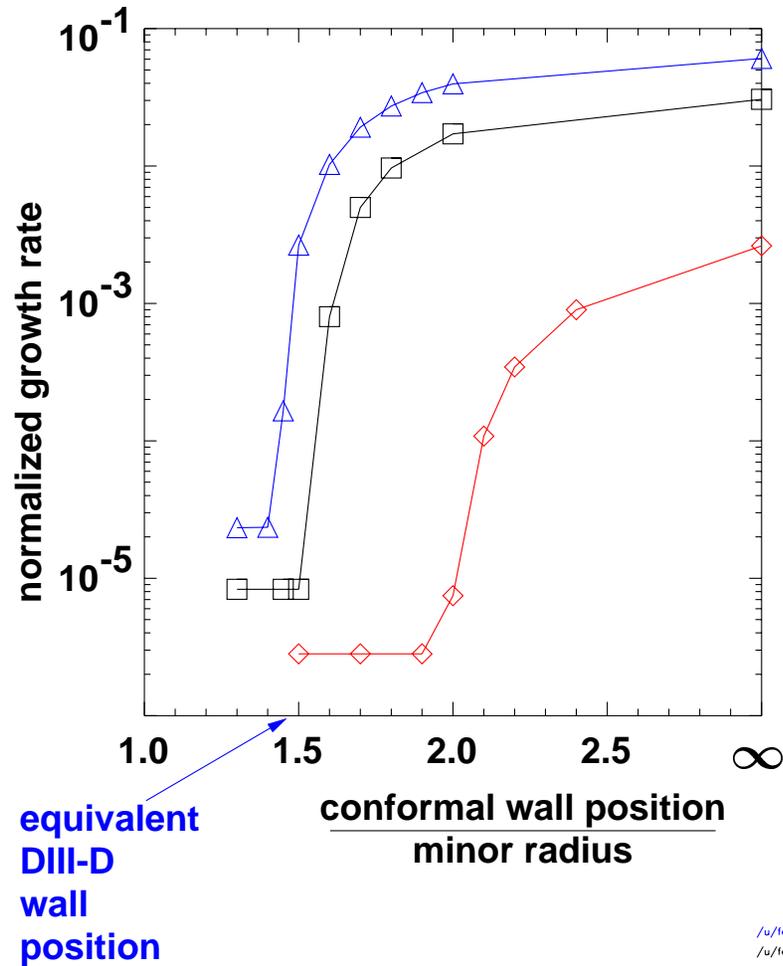
Maximum β_N achievable with an ideal wall determined by modeling stability versus wall position

- Little difference in response to change in wall position as shape and B_T (i.e. q_{95}) are varied.

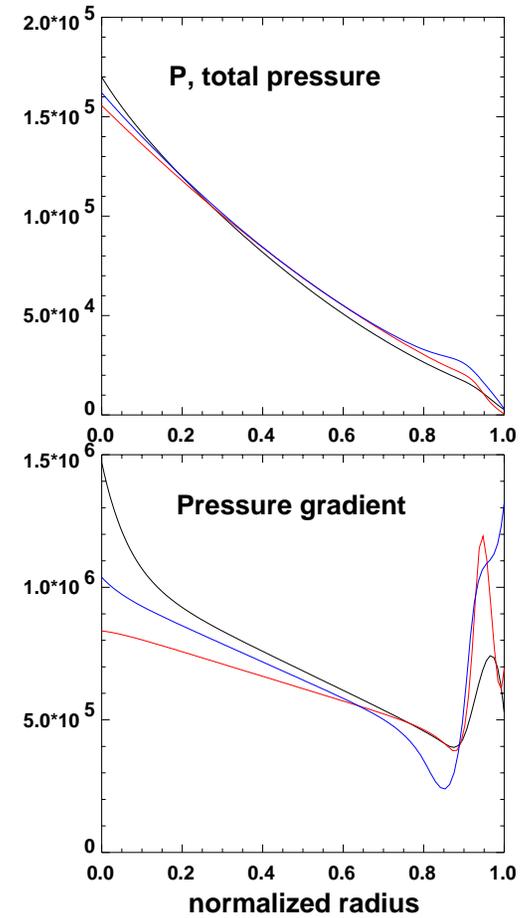


Some experimental discharges with $\beta_N \geq 4$ are near the ideal wall $n = 1$ stability limit

- Effect of wall depends on details of the equilibrium in 3 examples with $\beta_N \geq 4$.



/u/terron/balloon/low_n/thrust2/kinetic/106696/kinetic/g106696.01694_e
/u/terron/balloon/low_n/thrust2/kinetic/106795/Wade/g106795.01800



/u/terron/balloon/low_n/thrust2/kinetic/105981/kinetic/g105981.01584_j

Summary

Summary

- The advanced tokamak steady-state scenario required a method to increase achievable β_N in the lower κ , δ pumping shape.
- An increase to $\beta_N > 4$ was achieved by a 15% increase in B_T .
 - q_{95} is thought to be the key parameter.
 - Improved error field correction may be an enabling factor.
- The higher β_N values cannot be accounted for by an increase in the no-wall β_N limit.
- Increase in β_N likely results from a closer approach to the ideal wall β_N limit.
- Discharges at $q_{95} \approx 5$, $\beta_N \approx 4.1$ appear to be close to the ideal wall limit.
- Further increases in the ideal wall limit appear possible through changes in profiles.