# **CONTROL OF INTERNAL TRANSPORT BARRIERS IN DIII-D\***

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### OVERVIEW

- $E \times B$  shear is leading effect in creating internal transport barriers (ITB).
  - 1999: ITBs broader with counter-NBI than in similar discharges with co-NBI.
    - Co-NBI: Rotation dominates and opposes pressure gradient contribution to E×B shearing rate; increased or broadened pressure profile is *destabilizing* to turbulence.
    - Counter-NBI: Pressure gradient term dominates; increased or broadened pressure profile is *stabilizing* to turbulence.
  - 2000: Quiescent Double-Barrier regime (with sustained  $\beta_N H^{89} \le 7$ ) combines counter-injected ITB and steady-state ELM-free H-mode edge condition.
    - Separation between core and edge barriers provided by null in E×B shearing rate.
- Electron thermal transport is more difficult to reduce.
  - Strong electron ITB generated with localized direct electron heating (ECH).
  - Believed to require stabilization of both low-k (same as requirement for ion ITB) and high-k (ETG; an additional requirement) turbulence.
    - E×B shear too weak an effect to reduce high-*k* turbulence.
    - Simulations identify  $\alpha$ -stabilization (also known as Shafranov shift stabilization) as trigger mechanism for electron ITB.



## COUNTER-NBI RESULTS IN BROADER PROFILES

- 99849 (1.17s):
  - Counter-NBI
  - $-W_{\rm MHD} = 0.9 \,\rm MJ$
  - $P_{\rm NBI} = 11.2 \,\rm MW$ (6.5 MW absorbed).
- 87031 (1.82s):
  - Co-NBI
  - $-W_{\rm MHD} = 1.2 \text{ MJ}$
  - $-P_{\text{NBI}} = 9.6 \text{ MW}$ (7.6 MW absorbed).





# COUNTER-NBI IS FAVORABLE FOR ITB EXPANSION DUE TO INTERPLAY OF TERMS IN E $\times B$ SHEARING RATE

- Shearing rate  $\omega_{E \times B}$  can be separated into pressure gradient and rotation terms.
  - Total shearing rate is species independent, but individual terms depend on species.
  - Calculation shown is for main (deuterium) ions:
    - Total  $\omega_{\text{E}\times\text{B}}$  from CER carbon impurity measurements.
    - ∇p term uses TRANSP calculation of main ion thermal density.
    - Rotation term by subtraction.
- Counter-NBI is more favorable than co-NBI for barrier expansion:
  - Co: rotation term dominates  $\Rightarrow \omega_{E \times B}$ decreases with increased or broadened pressure profile  $\Rightarrow$  turbulence is destabilized.
  - Counter:  $\nabla p$  term dominates  $\Rightarrow \omega_{E \times B}$ increases with increased or broadened pressure profile  $\Rightarrow$  turbulence is stabilized.





#### DISTINCT CORE AND EDGE BOUNDARIES COEXIST IN THE QUIESCENT DOUBLE-BARRIER REGIME

- Distinct barriers coexist in larger devices, but only with degradation of the ITB in co-injected discharges in DIII-D:
  - With ELMs: ELMs penetrate plasma and impede core barrier.
  - ELM-free: Core and edge barriers merge.
    - Ion thermal transport can become neoclassical throughout entire plasma, but cannot be sustained.
- Double-barrier mode predicted with counter-NBI in DIII-D:
  - $E_r$  strongly negative both in core and at edge (H–mode) barriers.
    - Flattening of  $E_r$  profile inside H-mode edge locally removes E×B shear stabilization.
  - Inherent low confinement region expected to separate core and edge.
- Quiescent H-mode edge allows core barrier to exist without ELM degradation:
  - Quiescent Double-Barrier (QDB) regime allows sustained high performance.



# SUSTAINED HIGH PERFORMANCE IN THE QUIESCENT DOUBLE-BARRIER REGIME

- Quiescent double-barrier (QDB) regime combines:
  - Quiescent H-mode edge barrier:
    - Only observed with counter-NBI.
    - Edge Harmonic Oscillation (EHO): a benign replacement for ELMs.
    - Density control achieved through divertor cryopumping.
      - Particle flux enhanced by EHO.
  - Core barrier:
    - Characteristics similar to L-mode edge ITB with a pedestal.
    - Constant over lifetime of QDB regime.
- Counter-NBCD maintains  $q_{\min} > 1$ .
- Parameters obtained to date (all with  $I_p$ =1.3MA,  $B_T$ =1.8-2.1T):  $\beta_N \leq 2.9, H^{89} \leq 2.5, \beta_N H^{89} \leq 7, S_N \leq 4 \times 10^{15} \text{ neutrons/s}$
- Sustained for length of beam pulse.





( $H^{89}$  corrected for beam ion orbit losses: increased by ~10%)

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## COUNTER-NBI HEATED CORE BARRIER PERSISTS WITH THE ADDITION OF AN H-MODE EDGE

L-mode edge ITB (99849 1.12s)

- Core profiles similar to Lmode edge ITB with additional edge pedestal.
- Flat region in *E*<sub>r</sub> profile corresponds to separation between core and edge barriers.
  - Note that the barriers frequently merge in co-NBI discharges in DIII-D.





QDB (103740 3.3s)

#### CORE TRANSPORT IN QDB REGIME SIMILAR TO L-MODE EDGE ITB

- Core transport with ITB reduced to similar levels regardless of L- or H-mode (QDB) edge.
  - Second barrier appears near edge of QDB discharge.
- Separation between edge and core barriers corresponds to flattening of  $E_r$  profile and null in shearing rate  $\omega_{E \times B}$ .
- All three discharges with counter-NBI.





#### TURBULENCE IS REDUCED THROUGHOUT MOST OF THE QDB PLASMA

• With reduced broadband turbulence, core Alfvén modes are clearly visible in FIR scattering data, as are the low frequency edge harmonic oscillations associated with QH-mode operation.





#### DECREASED CORE TURBULENCE CORRELATION LENGTHS IN QDB REGIME INDICATE REDUCED TURBULENCE TRANSPORT STEP SIZE

- In L-mode, correlation lengths are observed to scale approximately with the poloidal ion gyroradius  $\rho_{\theta,s}$  (or 5-8  $\rho_s$ )
- In QDB discharges, core correlation lengths are significantly different: factor of 2-8 smaller than L-mode.





#### THE ELMLESS "QUIESCENT H-MODE" EDGE IS A KEY FEATURE OF THE QDB REGIME

- ELM-free regime... but particle transport near boundary is sufficient to allow density control via cryopumping.
- Edge gradients similar to ELMing phase.
  Extremely deep *E*<sub>r</sub> well.
- Elimination of ELMs' periodic divertor heat pulses a desirable feature for reactor class devices.
- QH-mode only obtained with counter-NBI.





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### EDGE HARMONIC OSCILLATION BELIEVED RESPONSIBLE FOR DESIRABLE FEATURES OF THE QH-MODE EDGE

- Localized near edge of plasma.
- Can appear at one or more toroidal mode numbers: n=1-10 has been seen.
  - Sometimes shifts mode number during QH-mode with no apparent change to profiles.
- Visible in density, temperature and magnetic fluctuations.
- Drives enhanced particle transport.
  - Allows particle control.





- The EHO has been observed between ELMs in ELMy discharges with both coand counter-NBI.
- ELMs eliminated only with counter-NBI.

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#### ELECTRON ITB PRODUCED BY ECH

- An electron thermal transport barrier has been obtained in a discharge heated with low-power ECH (0.5 MW) and neutral-beam injection (0.5 MW).
  - Central electron temperature approaches 6 keV with  $T_e/T_i \ge 10$ .
  - Electron thermal transport essentially eliminated in a narrow barrier region.
- The electron transport barrier appears under conditions where we would normally expect any barrier to be impeded by turbulence.
  - Low- and high-k turbulence growth rates calculated too large to be suppressed by E×B shear alone.
    - E×B shear may still have an important effect on low-k stability.
  - Growth rates are very sensitive to  $\alpha$  (normalized pressure gradient).
- Predictive simulation with the GLF23 model reproduces the development of the electron transport barrier with sufficiently large  $\alpha$ .
- $\alpha$ -stabilization appears to be a requirement in order to enter this regime.



# DIRECT ELECTRON HEATING GENERATES AN ELECTRON INTERNAL TRANSPORT BARRIER

- Beam blips for MSE and CER, time averaged  $P_{\text{NBI}} \approx 0.5$  MW (counter).
  - Poor beam ion confinement at low current and counter-NBI... most of this power is lost.
- Electron thermal ITB develops rapidly after onset of ECH.
  - Barrier increases in strength while remaining nearly stationary in position.
  - Little or no barrier appears in ion thermal, particle or angular momentum channels.
    - But the sources for these channels are very small.





#### ELECTRON TRANSPORT BARRIER DEVELOPS RAPIDLY FOLLOWING ECH ONSET



- Profiles flat or slightly hollow inside barrier.
- Barrier location expands ahead of ECH heating location.
- Barriers with nearly identical profiles have been observed with co- and counter-ECCD and pure heating (radial launch).
- Smaller response in  $T_i$  profile appears even without ECH.



#### TRANSPORT DECREASES IN BOTH THE ION AND ELECTRON CHANNELS AT THE ECH TURNON



#### REFLECTOMETER INDICATES CORE FLUCTUATION INTENSITY REDUCED WITH ECH

• Low-*k* measurement indicative of ITG and/or TEM stability.

- Reduction appears consistent with hypothesis that both low- <u>and</u> high-k turbulence must be stabilized for formation of e-ITB.
  - Peak at  $\rho \approx 0.37$  may correlate with feature in shearing rate profile; requires further analysis.





# ESTIMATED E×B SHEAR IS NOT LARGE ENOUGH TO SUPPRESS LOW-k MODES



- $E \times B$  shearing rate estimated using neoclassical (NCLASS) poloidal rotation.
  - No measurement available for this discharge, but  $v_{\theta}$  is small contribution.
  - Error bars probably significant: typically ~25% with measured  $v_{\theta}$ .
- Growth rate for low-k modes somewhat exceeds shearing rate .
  - ITG and/or TEM predicted unstable... should prevent ion or electron ITB.
- Growth rate increases when  $\alpha \Rightarrow 0$  ( $\beta$  set to zero in calculation).



## ELECTRON TEMPERATURE GRADIENT IS MARGINAL FOR STABILITY TO ETG MODES IN THE BARRIER

- Electron temperature gradient in barrier region at marginal stability level for ETG mode.
  - Consistent with previous observations with reduced core electron transport [B.W. Stallard, *et al.*, Phys. Plasmas 6, 1978 (1999)]
- α=0 (β set to zero in code) reduces critical gradient below experimental profile.
- Large calculated critical gradient inside barrier suspicious.
  - May be numerical consequence of resistive interchange instability.





# TRAPPED ELECTRON AND ELECTRON TEMPERATURE GRADIENT MODES BOTH HAVE SIGNIFICANT CALCULATED GROWTH RATES

- Spectra shown at point with peak a/L<sub>Te</sub>, where gradient slightly exceeds marginal level for ETG.
  - ETG feature vanishes at critical level.
  - Increases rapidly above critical level.
  - This condition can enforce marginality.
- Estimated E×B shearing rate appears too small to suppress turbulence in either range by itself.
  - May be large enough to have an effect on low-k range.





#### NEGATIVE MAGNETIC SHEAR AND SHAFRANOV SHIFT CAN BE STRONGLY STABILIZING INFLUENCES

- In general, growth rate spectrum is reduced for negative magnetic shear in the collisionless limit [Waltz, *et al.*, Phys. Plasmas 4, 2482 (1997)].
  - Data shown for a typical case.
- Further reductions occur with increasing  $\alpha$ .
  - Normalized pressure gradient  $\alpha = -\mu_0 P'(\Psi) V'(\Psi) (V/4\pi R_0)^{1/2}$
  - $\alpha$  can be *destabilizing* for strong positive shear.



Shafranov shift and  $\alpha$  stabilization, both commonly used terms, are synonymous.



#### BIFURCATION IN ELECTRON TEMPERATURE PROFILE IS PREDICTED BY GLF23 MODEL

- The GLF23 model, including both heat flux and momentum bifurcation mechanisms, dynamically follows bifurcations leading to formation of ITBs.
  - Includes ITG (ion temperature gradient mode), TEM (trapped electron mode), ETG (electron temperature gradient mode).
- *n*, *q*, sources, sinks, and equilibrium from analysis.
- *T* and  $v_{\phi}$  profiles initialized at pre-barrier levels and are evolved including the effects of E×B shear stabilization calculated from predicted profiles.
  - Uses generalized E×B shearing rate [R.E. Waltz, R.L. Miller, Phys. Plasmas 6, 4265 (1999)].
- Boundary conditions enforced at  $\rho\text{=}0.9$  using experimental data.
- Simulated evolution predicts barrier formation in  $T_{\rm e}$  profile.





#### TIME DEPENDENT SIMULATION REPRODUCES ELECTRON THERMAL TRANSPORT BARRIER



- Dynamic simulation begins with experimental profiles prior to barrier formation.
- Electron ITB forms in simulation with sufficiently large  $\alpha.$ 
  - $\alpha = c_{\alpha} \alpha^{calc}$ , where  $\alpha^{calc}$  is the value of  $\alpha$  calculated from the profiles used by the code, and  $c_{\alpha}$  is an arbitrary coefficient.
  - No barrier forms when  $c_{\alpha} < 1.35$ .

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- Without E×B shear,  $c_{\alpha} \ge 1.7$  required for barrier formation.
  - 35% is well within experimental or numerical uncertainty.

# SIMULATED BARRIER ONLY MAINTAINED WHEN EFFECTS OF $\alpha\mbox{-}STABILIZATION$ ARE INCLUDED



 "Backwards" simulations start with experimental profiles of fully developed e-ITB and run to steady-state.



# A WORKING MODEL FOR TRANSPORT AND ITS CONTROL

- Ion ITB can occur from stabilization of low-*k* modes alone.
  - Either E×B shear or  $\alpha$ -stabilization can be effective by themselves.
- Electron ITB requires stabilization of turbulence at both low- and high-*k*.
  - E×B shear alone insufficient.
  - α-stabilization appears able to explain electron ITBs in DIII–D.
  - ETG streamer formation may complicate the picture, but streamers not expected with negative shear [F. Jenko, *et al.*, Phys. Plasmas 7, 1904 (2000)].
- Since the requirements for electron ITB formation are a superset of those for ion ITB formation, additional power to the ion channel in an electron ITB may trigger simultaneous electron and ion barriers.





#### SUMMARY

- Addition of a Quiescent H-mode edge to a counter-injected ITB plasma results in the sustainable high performance Quiescent Double-Barrier regime.
  - ELMs replaced by the more benign Harmonic Edge Oscillation.
    - Little or no impact on core barrier.
    - Allows particle control.
    - Eliminates pulsed divertor heat load characteristic of ELMs.
  - Separation between core and edge barriers provided by null in  $E\infty B$  shearing rate.
  - Parameters obtained to date (all with  $I_P$ =1.3MA,  $B_T$ =1.8-2.1T):  $\beta_N \leq 2.9$ ,  $H^{89} \leq 2.5$ ,  $\beta_N H^{89} \leq 7$ ,  $S_N \leq 4 \times 10^{15}$  neutrons/s.
- Direct, localized electron heating (ECH) can trigger formation of a strong electron ITB.
  - Formation of e-ITB reproduced by theory-based simulations.
    - Requires  $\alpha$ -stabilization; E×B shear not effective stabilizer of short scale modes.
  - Current understanding suggests simultaneous ion and electron barriers can be obtained by heating ions in the e-ITB target plasma.
- These results taken together support a working model for transport in which E×B shear and α-stabilization both stabilize turbulence at different scales, resulting in the frequent observation of ion ITBs and less frequent observation of electron ITBs.

