

Parallel and Perpendicular Flows driven by Stochastic Magnetic Fields in MST

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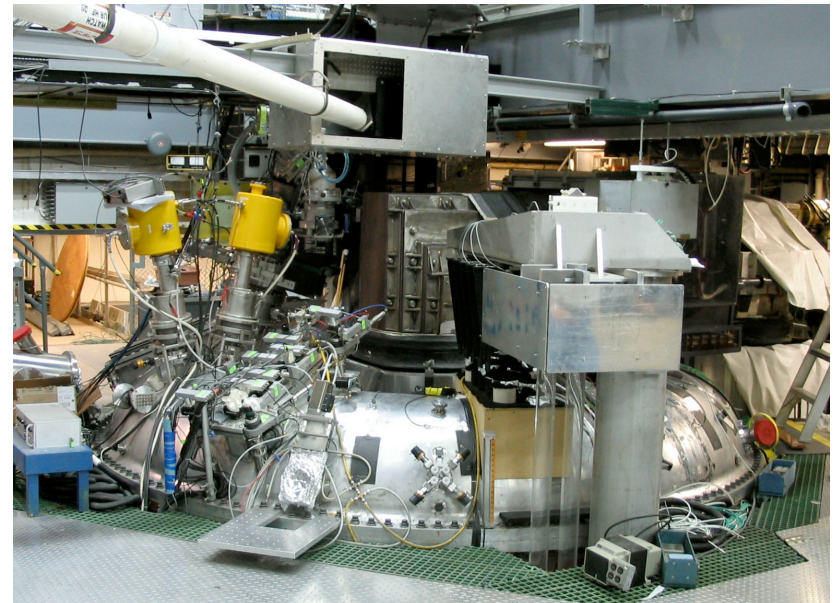
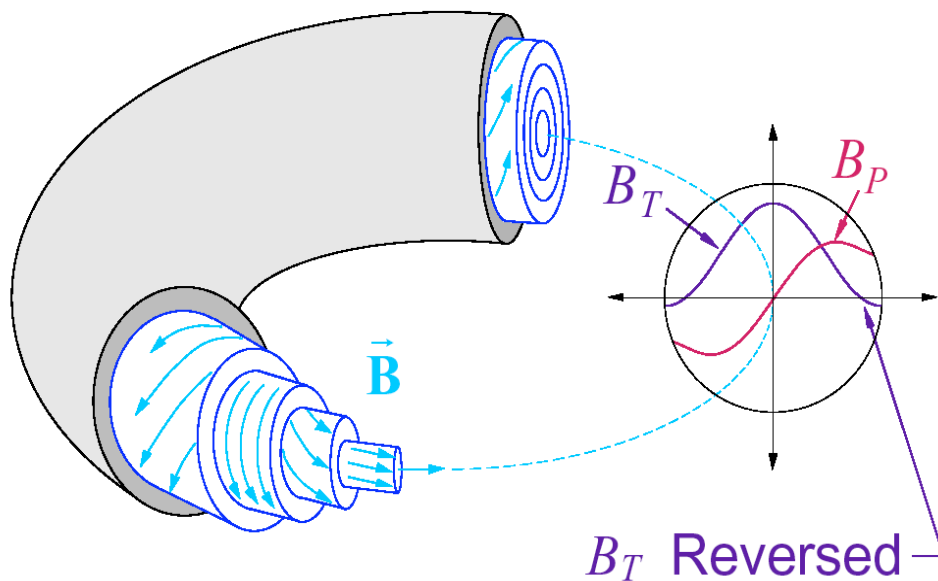
GOAL: Identify the role of stochastic magnetic fields in momentum transport and flow generation during reconnection

Results:

- (1) transport of parallel momentum from stochastic fields is measured and found to be significant in the plasma core
.....matching parallel velocity relaxation
- (2) Stochastic field driven charge transport produces localized perpendicular flow and radial electric field (shear)
.....to be presented at poster

Madison Symmetric Torus - MST

MST Reversed-Field Pinch (RFP) is toroidal configuration with relatively weak toroidal magnetic field B_T (i.e., $B_T \sim B_p$)

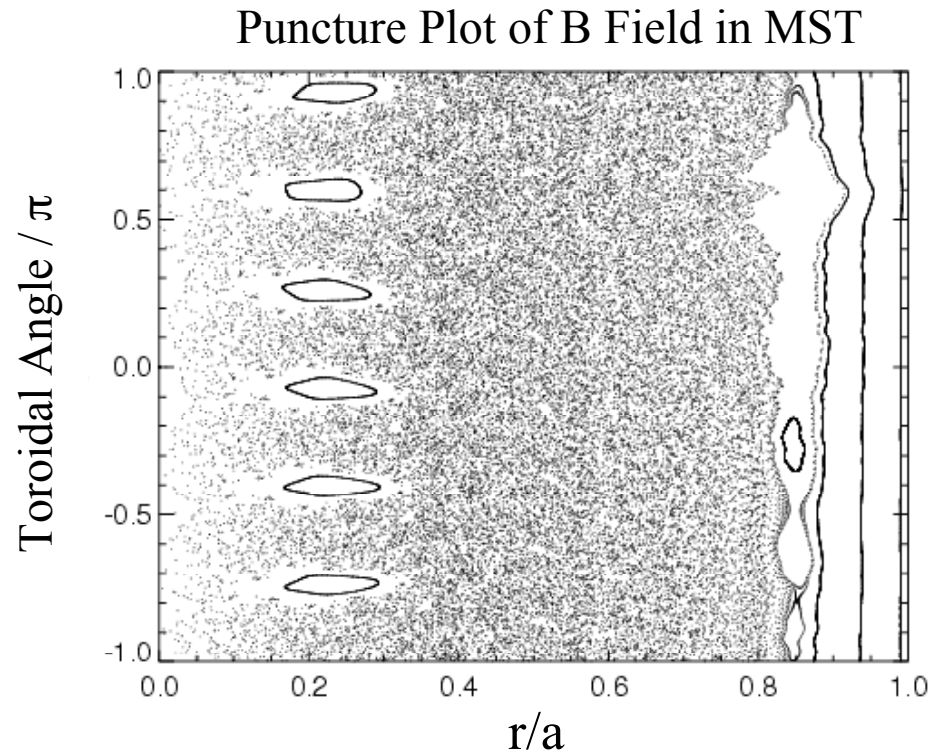


$$R_0 = 1.5 \text{ m}, a = 0.51 \text{ m}, I_p < 600 \text{ kA}$$

$$B_T \sim 3\text{-}4 \text{ kG}, n_e \sim 10^{19} \text{ m}^{-3}, T_{e0} < 1.3 \text{ keV}$$

$$\tau_E \sim 10 \text{ ms}, \beta = \langle p \rangle / B^2(a) = 15\%$$

magnetic field in MST is typically stochastic and should contribute to momentum transport



- J. Finn (1991?) suggests momentum diffusion with $D \sim D_m c_s$
- Plugging in MST parameters: $\tau_m \sim 1$ ms in ambient case
 $\tau_m \sim 100$ μ s during relaxation event

momentum transport by stochastic B field

- Momentum density in z direction being transferred in x direction is

$$\int f(\mathbf{x}, \mathbf{v}) m v_z v_x d^3 v$$

- Assuming a strong B field with small fluctuations and taking z along B, we can write the contribution from parallel streaming particles using $v_z \cong v_{\parallel}$ and $v_x \cong v_{\parallel} (b_x/B)$ as

$$\int f(\mathbf{x}, \mathbf{v}) m v_{\parallel}^2 \frac{b_x}{B} d^3 v = \frac{p_{\parallel} b_x}{B}$$

- Averaging over flux surface (z and y directions) we get the average momentum flux.

$$\frac{\langle p_{\parallel} b_x \rangle}{B}$$

- If $\langle \mathbf{b}_x \rangle = 0$, then we are left with

$$Flux: \quad \Gamma_r^{mom} = \frac{\langle \tilde{p}_{\parallel} \tilde{b}_x \rangle}{B}$$

- In order for the momentum at x to change, need more flux coming in from $-x$ than that leaving toward $+x$.

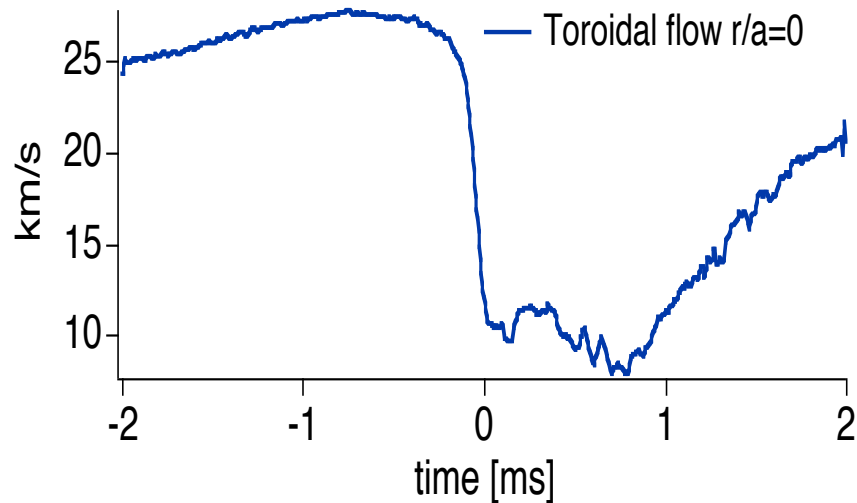
$$\frac{\partial \langle \rho v_{\parallel} \rangle}{\partial t} \sim -\nabla \left(\frac{\langle \tilde{p}_{\parallel} \tilde{b}_x \rangle}{B} \right)$$

- above relation provides a direct measure of the transport of momentum along the mean magnetic field direction due to stochastic magnetic field.

- Substituting $\tilde{p}_{\parallel} = T_{\parallel} \tilde{n} + n \tilde{T}_{\parallel}$, we get two terms

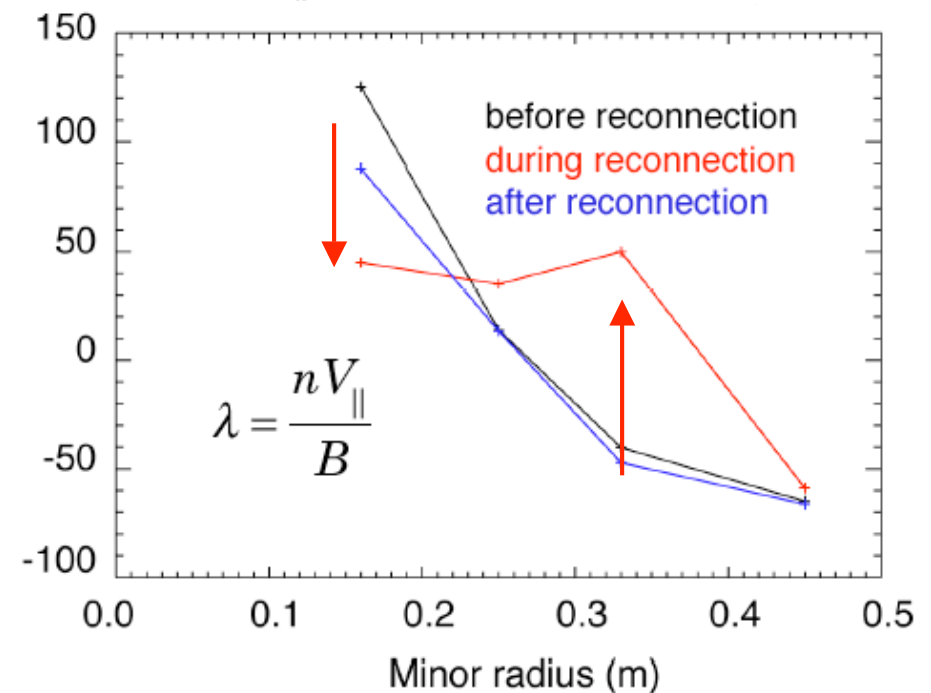
$$\frac{\partial \langle \rho v_{\parallel} \rangle}{\partial t} \sim -T_{\parallel} \nabla \left(\frac{\langle \tilde{n} \tilde{b}_x \rangle}{B} \right) - n \nabla \left(\frac{\langle \tilde{T}_{\parallel} \tilde{b}_x \rangle}{B} \right)$$

Momentum transported from core to the edge during reconnection



Parallel momentum profile

$$nV_{\parallel} / B \text{ (} 10^{19} \text{ m}^{-3} \cdot \text{km/s} \cdot \text{T}^{-1}\text{)}$$



During reconnection, parallel momentum redistributes within $\sim 100 \mu\text{s}$, much faster than classical collision time.

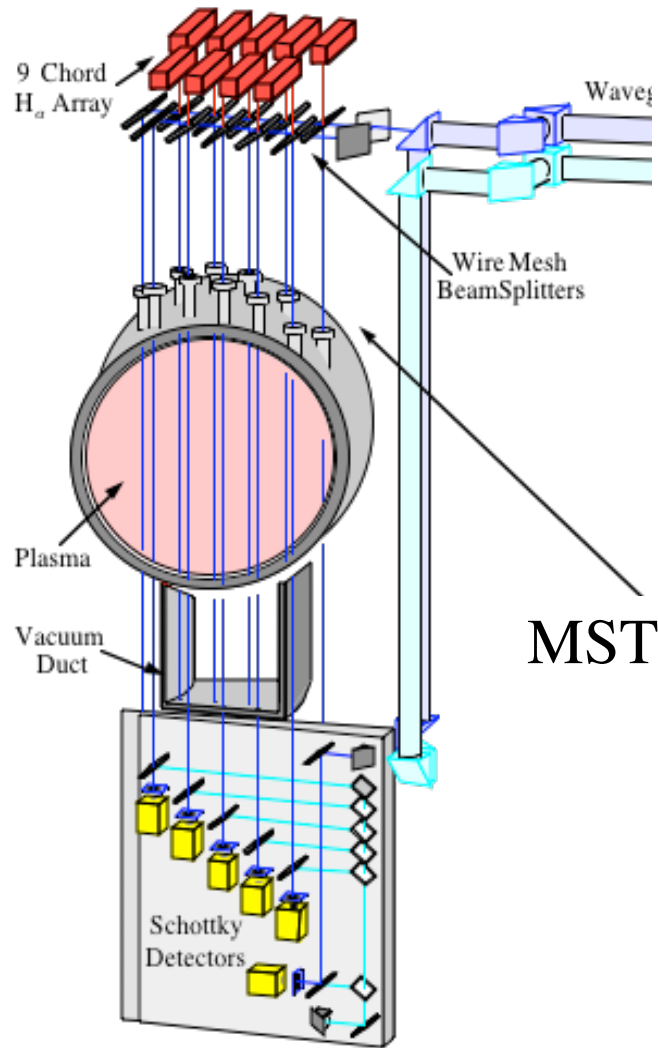
Stochastic field-driven transport is directly measured (first term only)

$$\frac{\partial \langle \rho v_{\parallel} \rangle}{\partial t} \sim \boxed{-T_{\parallel} \nabla \left(\frac{\langle \tilde{n} \tilde{b}_r \rangle}{B} \right)} - n \nabla \left(\frac{\langle \tilde{T}_{\parallel} \tilde{b}_r \rangle}{B} \right)$$

Convection

- Measure $T_{//,\text{ion}}$ with Rutherford Scattering (or CHERS)
- Measure b_r and B with FIR laser Faraday Rotation
- Measure \tilde{n} and $\nabla \tilde{n}$ with differential FIR interferometer
- *Measurement details provided at poster....*

FIR Polarimeter-Interferometer System



Interferometer $\rightarrow \tilde{n}$

$$\phi \sim \int n dl + \int \tilde{n} dl$$

Differential Interferometer $\rightarrow \nabla \tilde{n}$

$$\frac{\Delta \tilde{\phi}}{\Delta x} \quad \Delta x \sim 1 \text{ mm}$$

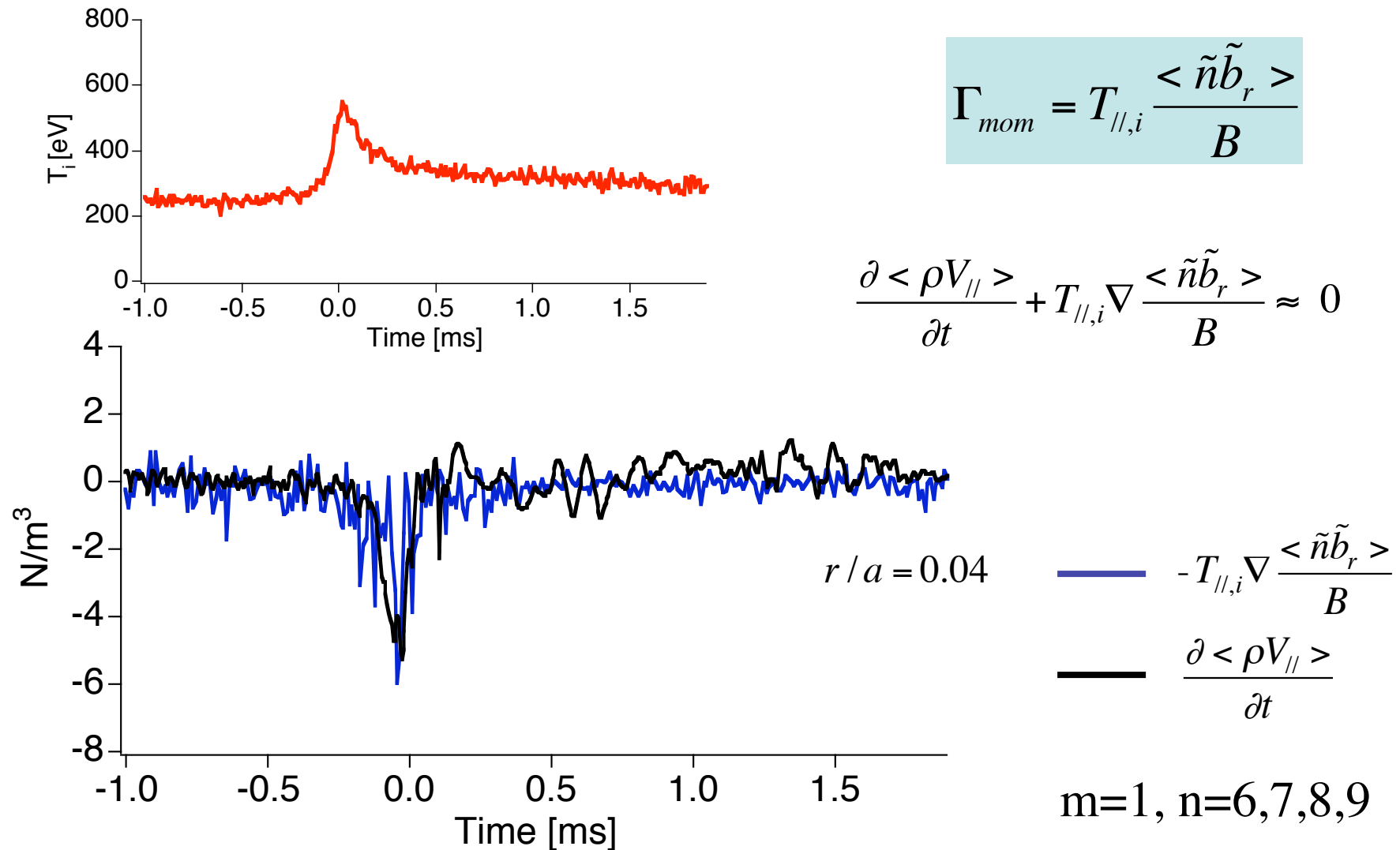
Faraday rotation $\rightarrow \tilde{b}$ and \tilde{j}

$$\Psi \sim \int n \vec{B} \cdot d\vec{l} + \int n \vec{b} \cdot d\vec{l} + \int \tilde{n} \vec{B} \cdot d\vec{l}$$

$$\oint_L \vec{b} \cdot d\vec{l} = \mu_0 \tilde{I}_\phi \approx \frac{\tilde{\Psi}_1 - \tilde{\Psi}_2}{c_F \bar{n}_0} \Rightarrow \tilde{j}_\phi$$

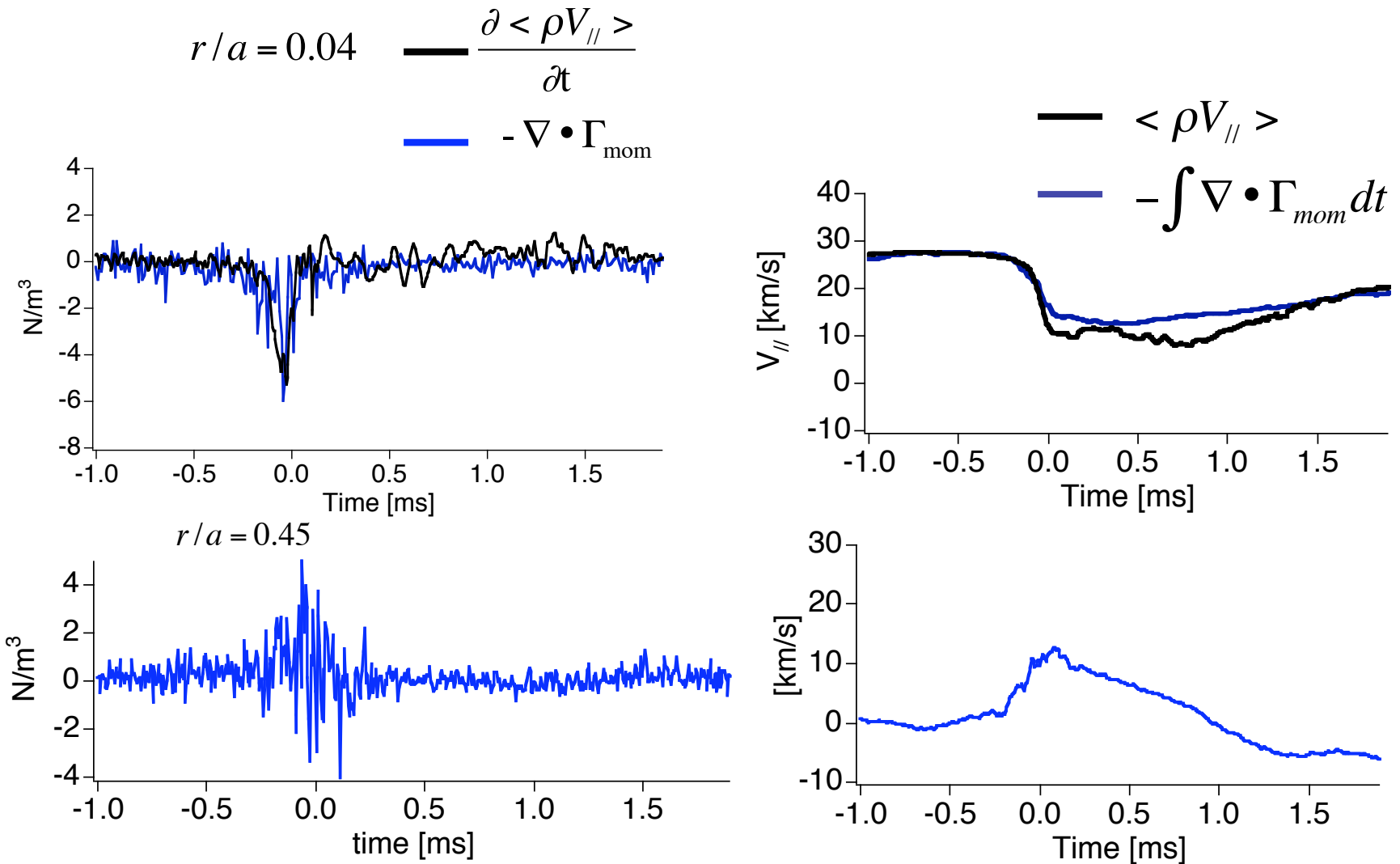
11 chords, $\Delta x \sim 8 \text{ cm}$, $\Delta \phi \sim 0.05^\circ$, $\Delta t \sim 1 \mu\text{s}$

Momentum flux and parallel momentum change



Magnetic fluctuation-induced (convective) flux drives parallel momentum drop

Momentum flux changes direction away from magnetic axis



*Direction change consistent with momentum transport
 (not total momentum loss)*

Summary: Stochastic magnetic fields drive parallel and perpendicular flows

Measurements indicate :

