

# Study of Critical Pressure Gradient at the Edge of Type I ELMy H-mode plasmas

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# Objectives

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- ❑ Examine approaches for determining magnetic shear and safety factor used in a prediction of  $T_{\text{ped}}$
- ❑ Examine sensitivity of  $T_{\text{ped}}$  on the functional form assumed for the elongation and triangularity
- ❑ Investigate conditions for possible access to 2<sup>nd</sup> stability of the ballooning mode instability
- ❑ Simulation of a  $\rho^*$  scan utilizing a predictive pedestal temperature model in the BALDUR integrated predictive transport code

# Influence of $s$ and $q$ on $T_{\text{ped}}$

- The pedestal width, based on magnetic and flow shear stabilization ( $\Delta \propto \rho s^2$ ), yields:

$$T_{\text{ped}} [\text{keV}] = C_w^2 (3.231 \times 10^{37}) \left( \frac{B^2}{q^4} \right) \left( \frac{A_H}{R^2} \right) \left( \frac{\alpha_c^2}{n_{\text{ped}}^2} \right) s^4$$

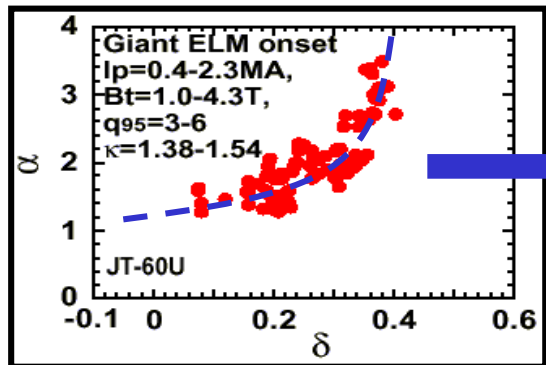
- $C_w$  is a constant in the pedestal width model
- Pedestal temperature is strongly dependent on magnetic shear,  $s$ , and safety factor,  $q$
- Equation for  $T_{\text{ped}}$  may be nonlinear depending on how  $s$  and  $q$  are computed

# Scaling of $\alpha_c$

□ Normalized critical pressure gradient,  $\alpha_c$ , depends not only on magnetic shear,  $s$ , but also on a geometrical factor,  $f(\kappa, \delta)$

We express  $\alpha_c$  in terms of  $s$  and  $f$  using:  $\alpha_c = 0.8 s f(\kappa, \delta)$

## Experimental Suggestion for Geometrical Factor



(Y.Kamada, et al., IAEA 1996)

One approximation considered for the geometrical factor  $f$  has the following dependence on elongation and triangularity:

$$f(\kappa, \delta) = 0.5 (1 + \kappa_{95}^2 (1 + C_0 \delta_{95}^2))$$

$$\text{We use : } \alpha_c = 0.4 s (1 + \kappa_{95}^2 (1 + C_0 \delta_{95}^2))$$

# Calculation of $s$ and $q$ using Approach 1

- Several approaches are proposed to calculate magnetic shear ( $s$ ) and safety factor ( $q$ )

**Approach 1:** Magnetic shear is assumed to be a constant and safety factor is calculated at the 95% flux surface

$$s = 2$$

$q_{95}$  obtained from D. P. Stoler, et al., Comp. Phys. Comm. 81, 261 (1994)

$$q_{95} = \frac{5a^2 B (1 + \kappa_{95}^2 (1 + 2\delta_{95}^2 - 1.2\delta_{95}^3)) (1.17 - 0.65\varepsilon)}{IR \cdot 2(1 - \varepsilon^2)^2}$$

# Calculation of $s$ and $q$ using Approach 2

**Approach 2:** Magnetic shear ( $s$ ) is calculated with bootstrap current effect included. Safety factor is calculated at 95% flux surface.

$$s = \frac{2\sqrt{\varepsilon}}{\sqrt{\varepsilon} + 0.2b(v^*, \varepsilon)f(\kappa, \delta)}$$

$$q_{95} = \frac{5a^2 B (1 + \kappa_{95}^2 (1 + 2\delta_{95}^2 - 1.2\delta_{95}^3))(1.17 - 0.65\varepsilon)}{IR 2(1 - \varepsilon^2)^2}$$

**Derivation of  $s$  is shown in A. H. Kritz's poster**

# Calculation of $s$ and $q$ using Approach 3

**Approach 3:** Magnetic shear and safety factor are calculated at one pedestal width away from the separatrix  
(M. Sugihara, Nuclear Fusion, 40 (2000) 1743)

$$\Delta = a(1 - x) = \text{Pedestal width}$$

$$q(x) = \frac{q_{95}}{2.933} \left[ \left( 1 + \left[ \frac{x}{1.4} \right]^2 \right)^2 + 0.267 |\ln(1 - x)| \right]$$

$$s = \frac{x}{q} \frac{\partial q}{\partial x}$$

# Calculation of $s$ and $q$ using Approach 4

**Approach 4:** Magnetic shear with bootstrap current included and safety factor calculated one pedestal width away from separatrix

$$\Delta = a(1 - x) = \text{Pedestal width}$$

$$q(x) = \frac{q_{95}}{2.933} \left[ \left( 1 + \left[ \frac{x}{1.4} \right]^2 \right)^2 + 0.267 |\ln(1-x)| \right]$$

$$s = \frac{s_0 \sqrt{\varepsilon}}{\sqrt{\varepsilon} + 0.1 s_0 b(v^*, \varepsilon) f(\kappa, \delta)} \quad s_0 = \frac{x}{q} \frac{\partial q}{\partial x}$$

This Approach used in A.H. Kritz and G. Bateman's poster

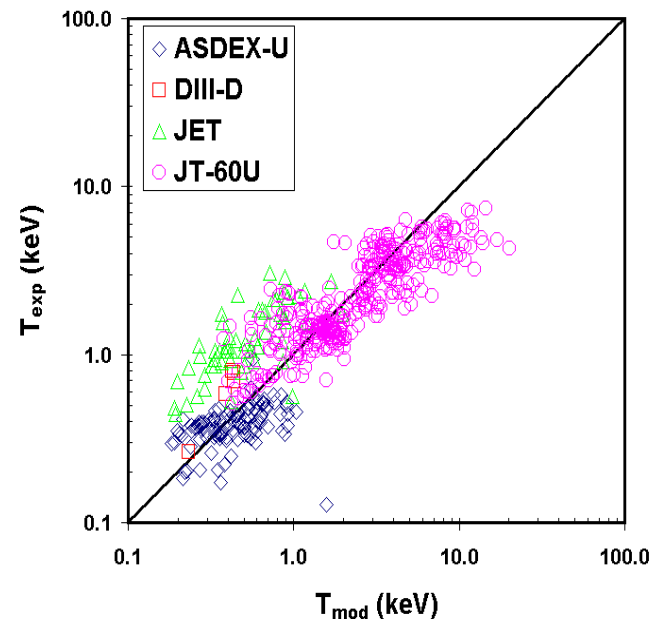


# Result of Approach 1

□ The predictions of  $T_{\text{ped}}$  using Approach 1 compared with 533 experimental data points.

➤ Shear is constant ( $s = 2$ ) and  $q = q_{95}$

$C_0$	RMSE(%)
2.0	57.8
<b>5.0</b>	<b>52.0</b>
10.0	50.7
20.0	60.5

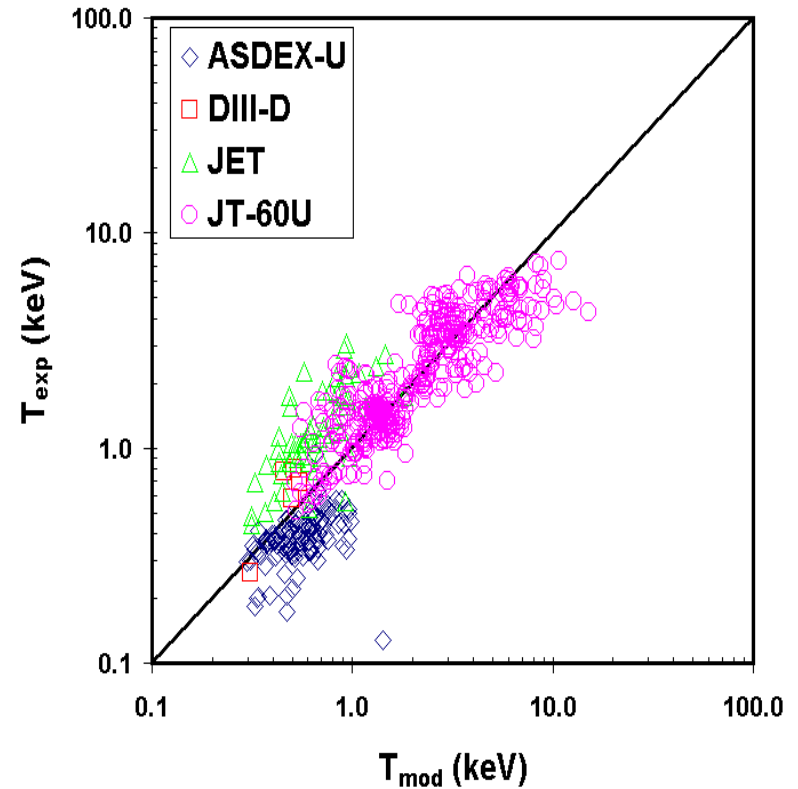


$$f(\kappa, \delta) = 0.5(1 + \kappa_{95}^2(1 + C_0 \delta_{95}^2))$$

# Result of Approach 2

➤ Magnetic shear with bootstrap current and  $q = q_{95}$

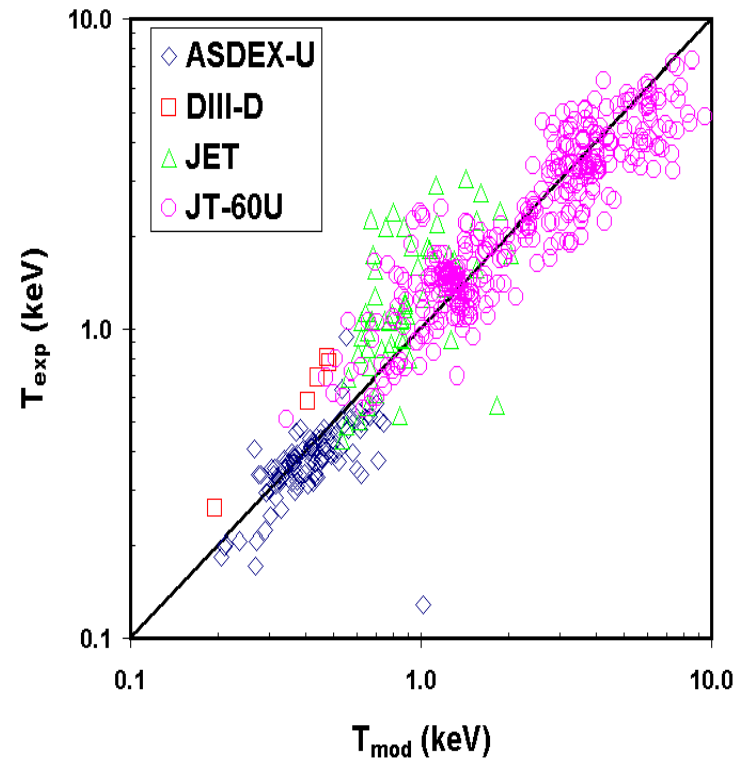
$C_0$	RMSE(%)
2.0	47.5
<b>5.0</b>	<b>43.3</b>
10.0	40.2
20.0	38.7



# Result of Approach 3

➤  $s$  and  $q$  calculated one pedestal width away from the separatrix and bootstrap current effect not included

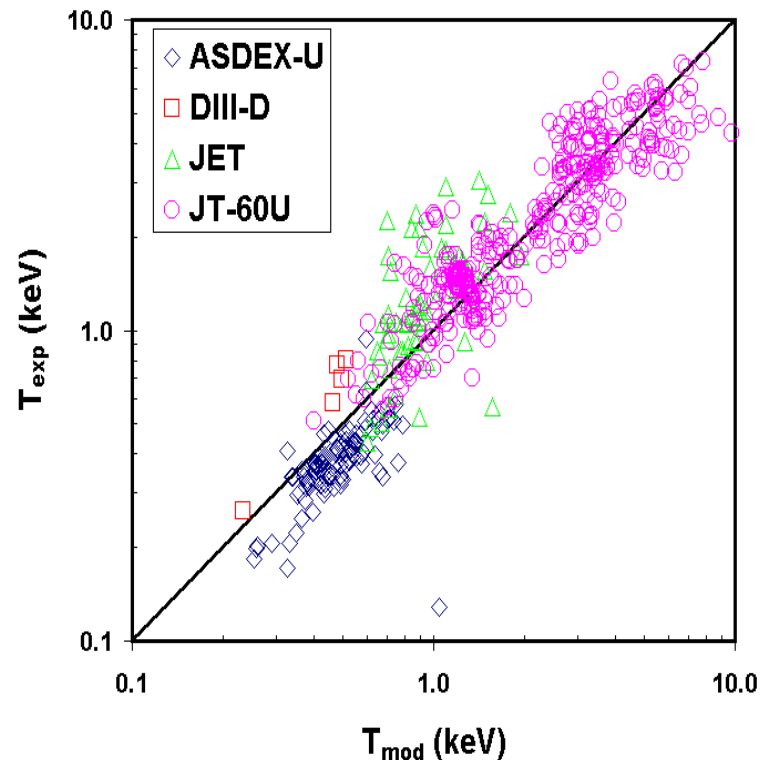
$C_0$	RMSE(%)
2.0	36.7
<b>5.0</b>	<b>32.2</b>
10.0	29.7
20.0	32.9



# Result of Approach 4

➤  $s$  with bootstrap current effect included and  $q$  are calculated one pedestal away from separatrix

$C_0$	RMSE(%)
2.0	35.4
<b>5.0</b>	<b>32.0</b>
10.0	29.6
20.0	29.0



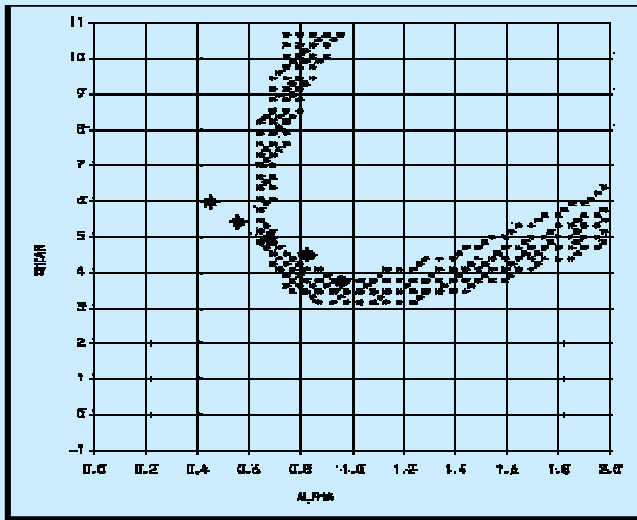
# RMS of different approaches for $s$ and $q$

□ Predictions of pedestal temperature based on different approaches for  $s$  and  $q$  is compared with 533 experimental data points

	$C_0 = 2$	$C_0 = 5$	$C_0 = 10$	$C_0 = 20$
Approach 1	57.8	52.0	50.7	60.5
Approach 2	47.5	43.3	40.2	38.7
Approach 3	36.7	32.2	29.7	32.9
Approach 4	35.4	32.0	29.6	29.0

# Access to 2<sup>nd</sup> stability

Evolution of pressure gradient at 95% flux surface



□ JET Discharge 53186 (low triangularity) does not have access to 2<sup>nd</sup> stability

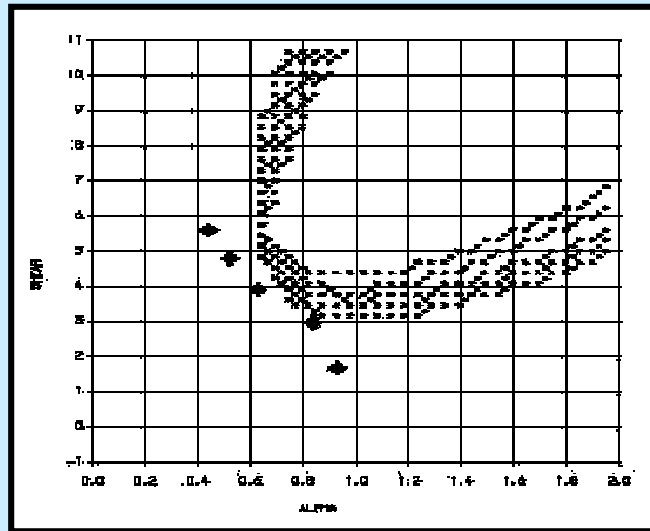
- Enter unstable regime of ballooning mode
- Remains at the transition between 1<sup>st</sup> and 2<sup>nd</sup> stability (at the nose of unstable curve)

□ Possible access to 2<sup>nd</sup> stability can be achieved by

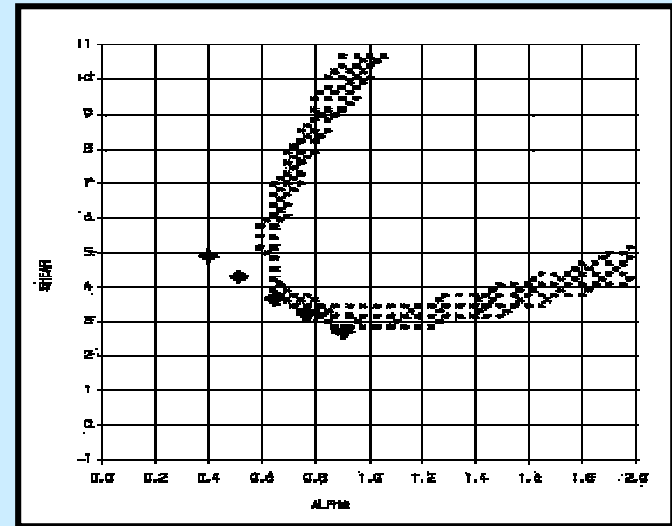
- Strong shaping (high triangularity)
- Strong edge current
  - achieved by increasing bootstrap current or
  - achieved by increasing plasma current

# Result of Increasing of Edge Current

- Bootstrap current is artificially increased by a factor of 2.



- Ramping current up from 2.5 MA to 3.0 MA during 1 sec.



- With a small change in the edge conditions, the edge pressure gradient can change dramatically and impact the performance of the plasma in H-mode

- For example, access to 2<sup>nd</sup> stability can be achieved

# Core and Edge model

□ Pedestal temperature model is implemented in the predictive integrated transport code, BALDUR, to provide boundary conditions

## Core transport model

- Core transport is calculated using MMM95 model
- Agrees with experimental results (within 15% RMS deviation)

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## Pedestal temperature model

- Based on ballooning mode for pressure gradient (approach 4) and flow & magnetic shear stabilization for width
- Approximately 32% RMS deviation from experimental results

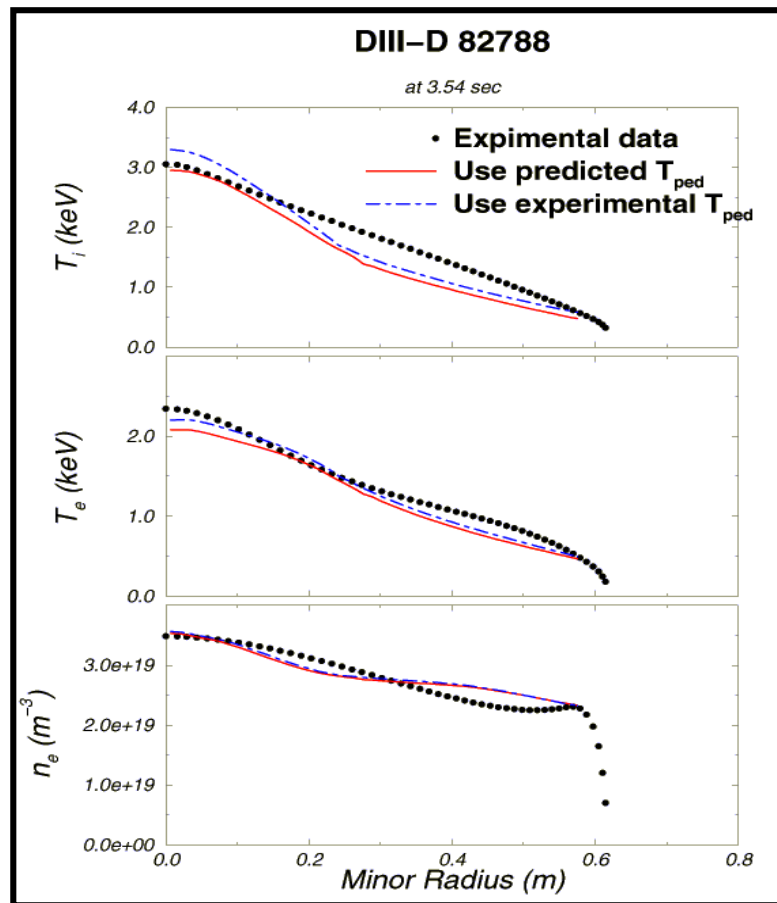
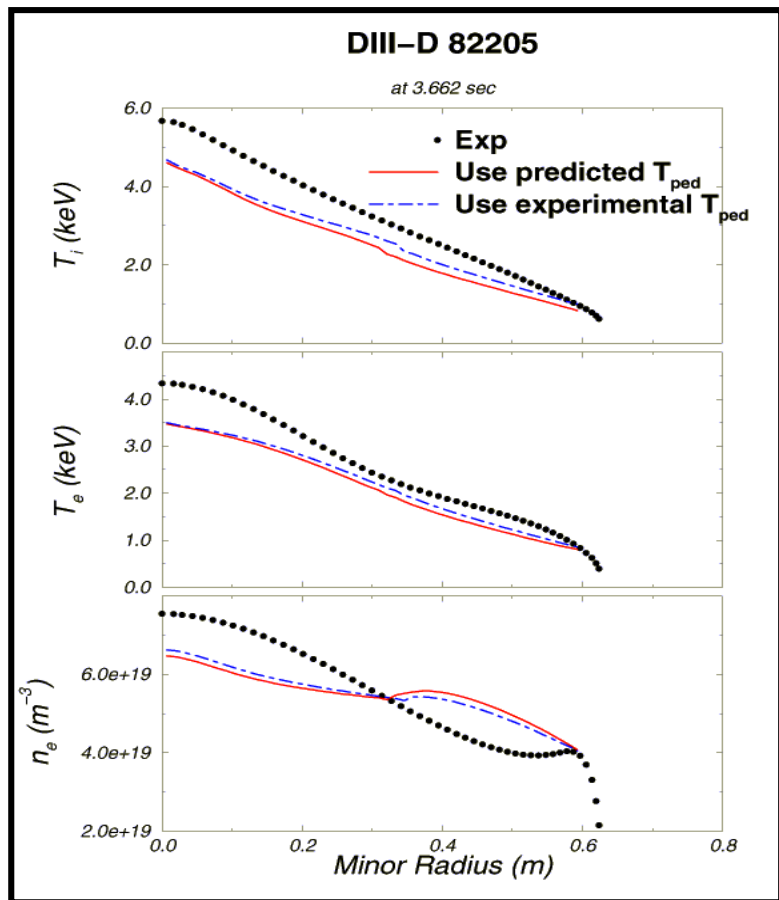


# $\rho^*$ scan

□ Simulations of  $\rho^*$  scans in DIII-D and JET have been carried out using BALDUR code

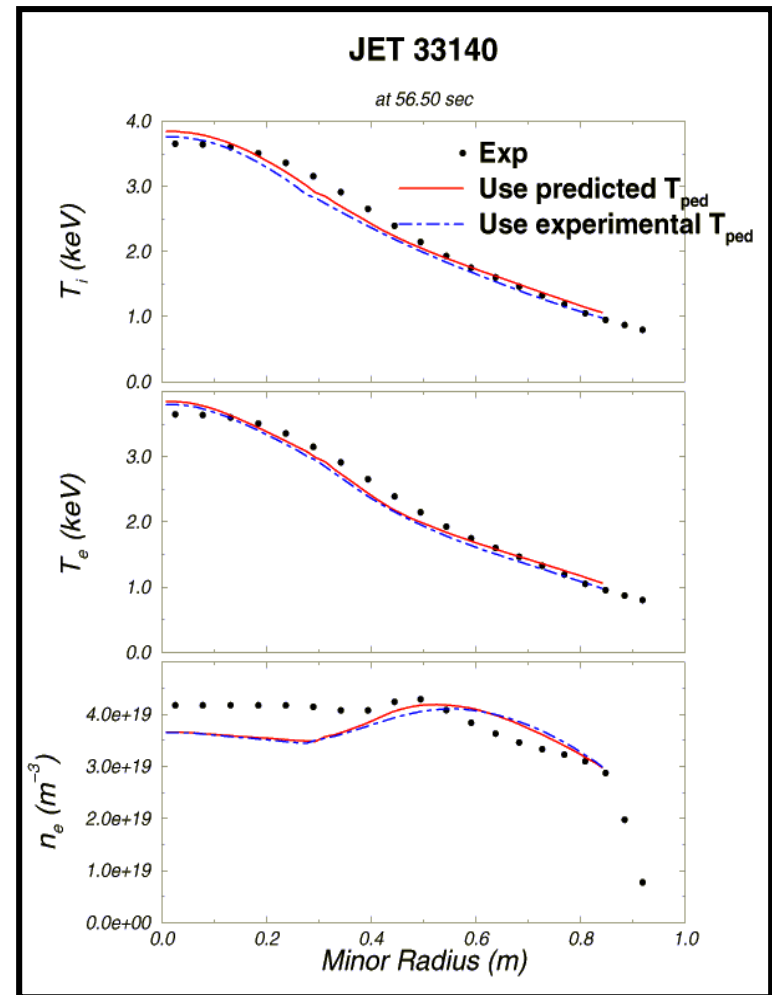
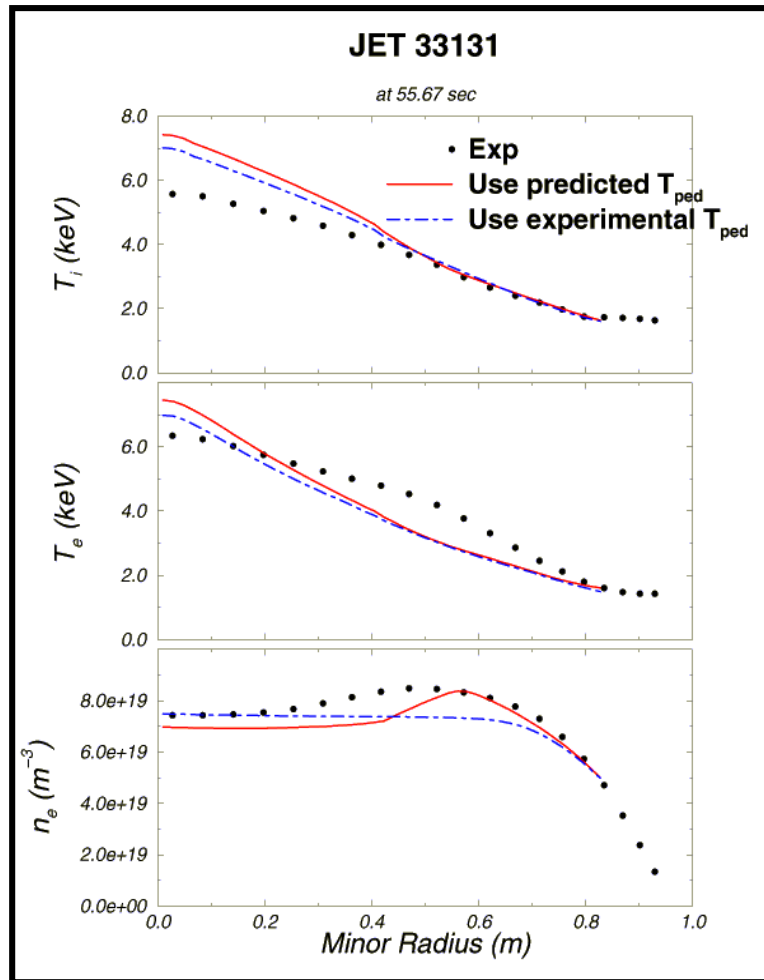
Discharge	D3D 82205	D3D 82788	JET 33131	JET 33140	JET 35156	JET35171
Type	Low $\rho^*$	High $\rho^*$	Low $\rho^*$	High $\rho^*$	Low $\rho^*$	High $\rho^*$
R (m)	1.69	1.68	2.94	2.93	2.87	2.88
a (m)	0.63	0.62	0.94	0.92	0.93	0.94
$I_p$ (MA)	1.34	0.66	2.83	1.61	2.05	1.01
B (T)	1.87	0.94	3.13	1.77	2.17	1.09
$\kappa$	1.71	1.67	1.70	1.56	1.56	1.58
$\delta$	0.37	0.35	0.28	0.26	0.11	0.24
$\rho^*(0)$	0.013	0.019	0.005	0.008	0.005	0.010

# Simulations of $\rho^*$ scan in DIII-D

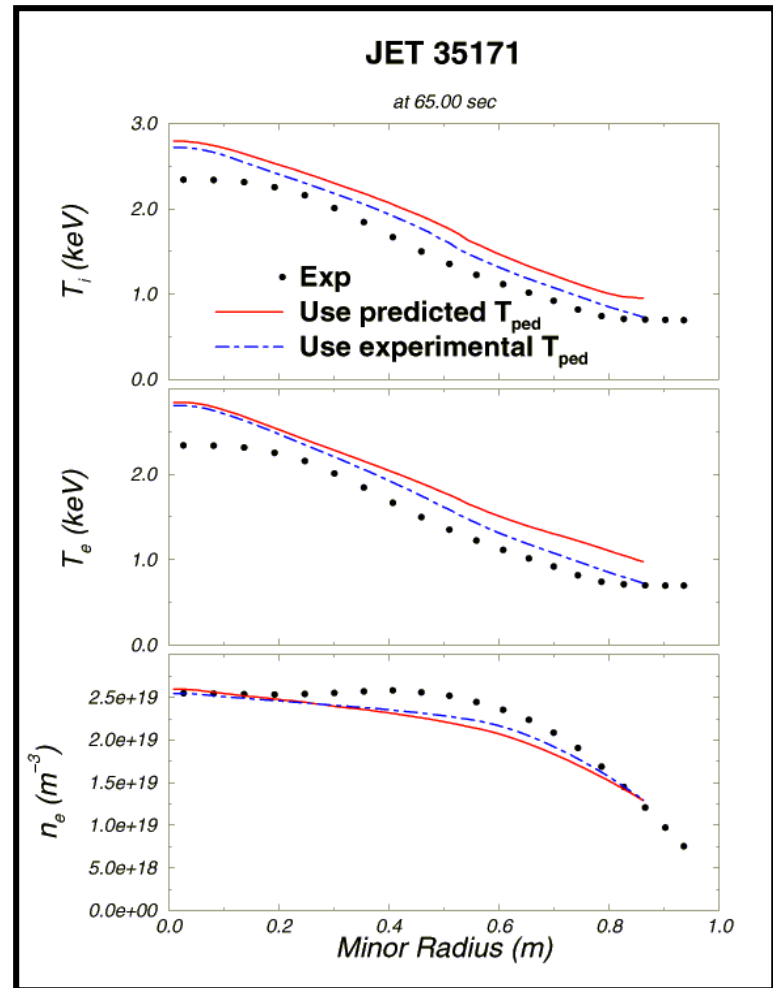
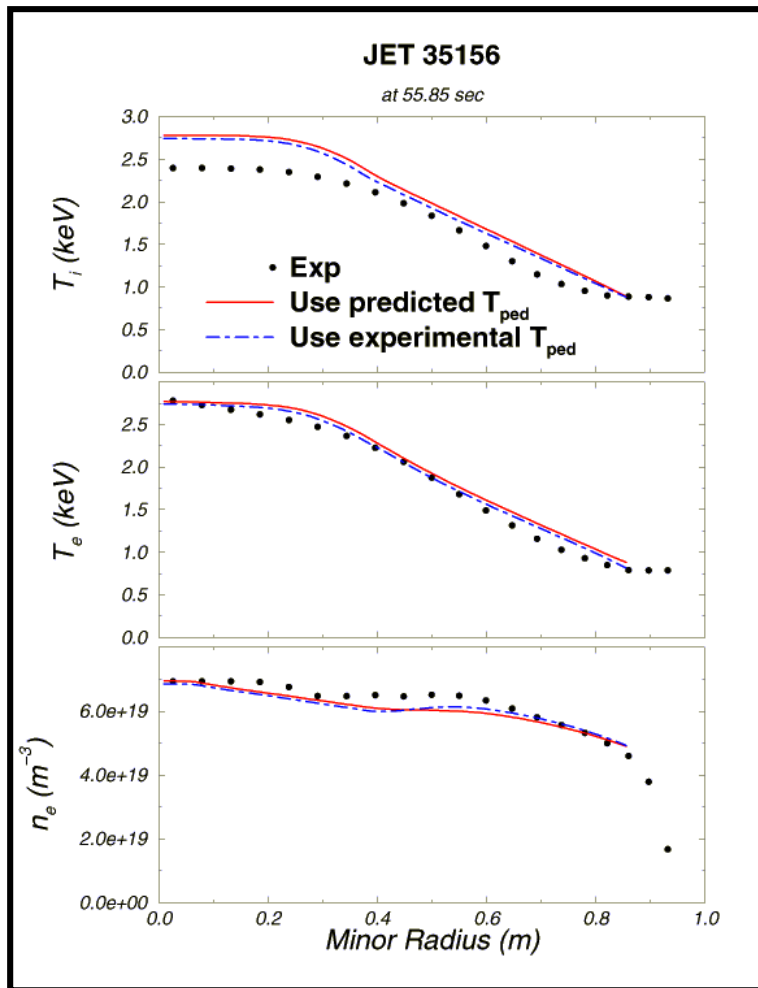


Simulations use **MMM95** +  $\Delta \propto \rho s^2$  + **Approach 4** for  $s$  and  $q$

# Simulations of $\rho^*$ scan in JET



# Simulations of $\rho^*$ scan in JET



# Pedestal scaling in $\rho^*$ scan

□ In  $\rho^*$  scan,  $\beta$  and  $v^*$  are kept constant

➤ This results in the relationship  $T_{\text{core}} \propto B^{0.67}$

Discharge	B (T)	$T_{\text{ped}}$ (keV)
DIII-D 82205	1.87	0.79
DIII-D 82788	0.94	0.46
JET 33131	3.13	1.59
JET 33140	1.77	1.04

$$\frac{T_{\text{ped},82205}}{T_{\text{ped},82788}} = \left( \frac{B_{82205}}{B_{82788}} \right)^{0.79}$$

$$\frac{T_{\text{ped},33131}}{T_{\text{ped},33140}} = \left( \frac{B_{33131}}{B_{33140}} \right)^{0.74}$$

➤ In the pedestal region,  $T_{\text{ped}} \propto B^{0.77}$

➤ Stronger temperature dependence on magnetic field in the pedestal than in the core

# Summary-1

- Using  $s$  at one pedestal width away from the separatrix and modified with bootstrap current (Approach 4) yields lowest RMS error
  - **29.0 %** with the geometrical factor of
$$f(\kappa, \delta) = 0.4 s (1 + \kappa_{95}^2 (1 + 20\delta_{95}^2))$$
- Increasing effect of triangularity yields better agreement with experimental data
- Bootstrap current has limited effect in improving the agreement with data
  - **Negligible effect** when shear is calculated at one pedestal width away from separatrix

# Summary-2

- Plasma can access 2<sup>nd</sup> stability with large plasma edge current
  - Increase bootstrap current or plasma current
- Simulations of  $\rho^*$  scans discharges carried out using the BALDUR integrated transport code with a predictive boundary model for  $T_e$  and  $T_i$ 
  - Errors in the prediction of  $T_{ped}$  do not amplify the error in the core predictions
  - $T_{ped}$  scales as  $B^{0.77}$  in the pedestal region
    - $T_{ped}$  has stronger dependence on  $B$  than does  $T_{core}$