

# MHD STABILITY ISSUES IN A BURNING PLASMA

by  
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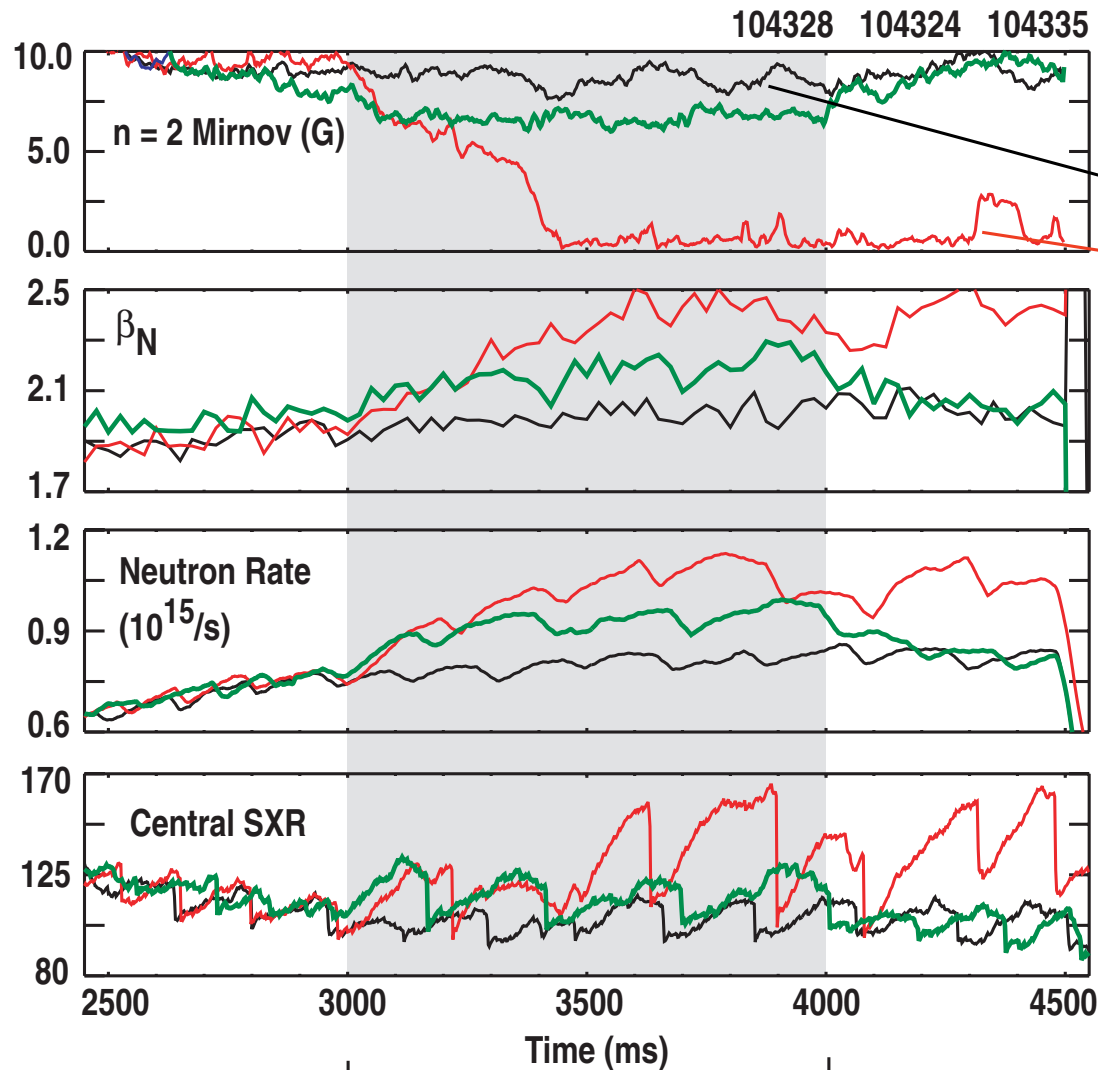
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# PRESENT UNDERSTANDING OF MHD STABILITY LIMITS IS SUFFICIENT TO DESIGN A BURNING PLASMA EXPERIMENT

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- **Ideal MHD stability limits are well understood and predictable**
  - Upper limit to plasma stability
  - Credible foundation for design of next-step devices
- **Non-ideal effects introduce greater uncertainty**
  - Resistivity, finite Larmor radius, energetic ions, ...
- **Resistive instabilities are less predictable but may be avoidable**
  - Neoclassical tearing modes can be avoided transiently by profile modification
  - Recent experiments have suppressed NTMs with localized current drive
- **Steady operation very near stability limits has been demonstrated**
- **Burning plasma experiments go beyond present experience with MHD stability, and present new scientific challenges**

# FULL STABILIZATION OF NTM OBTAINED WITH MODEST ECH POWER

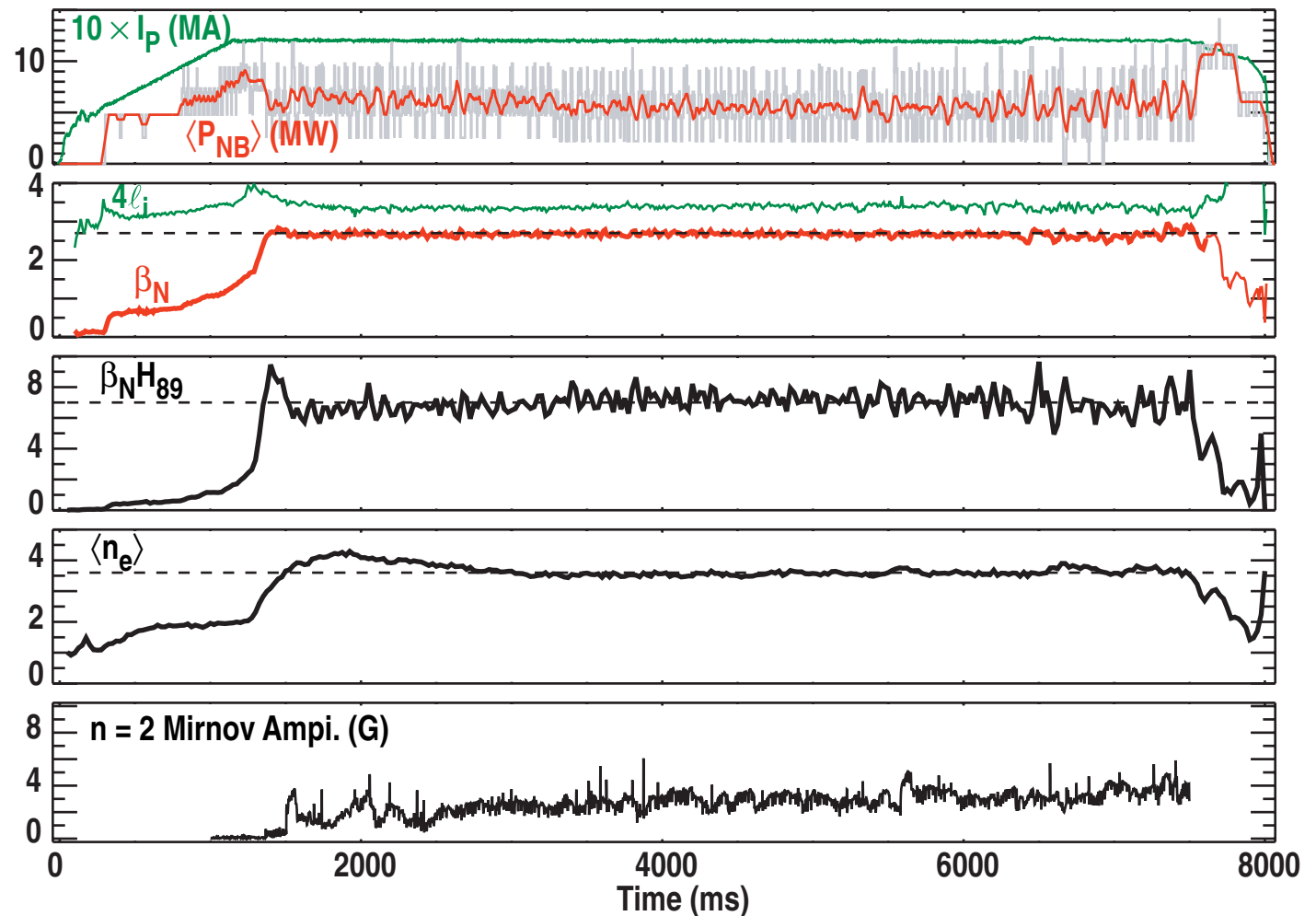


Resonance moved 2 cm outward  
No ECCD  
Full Stabilization

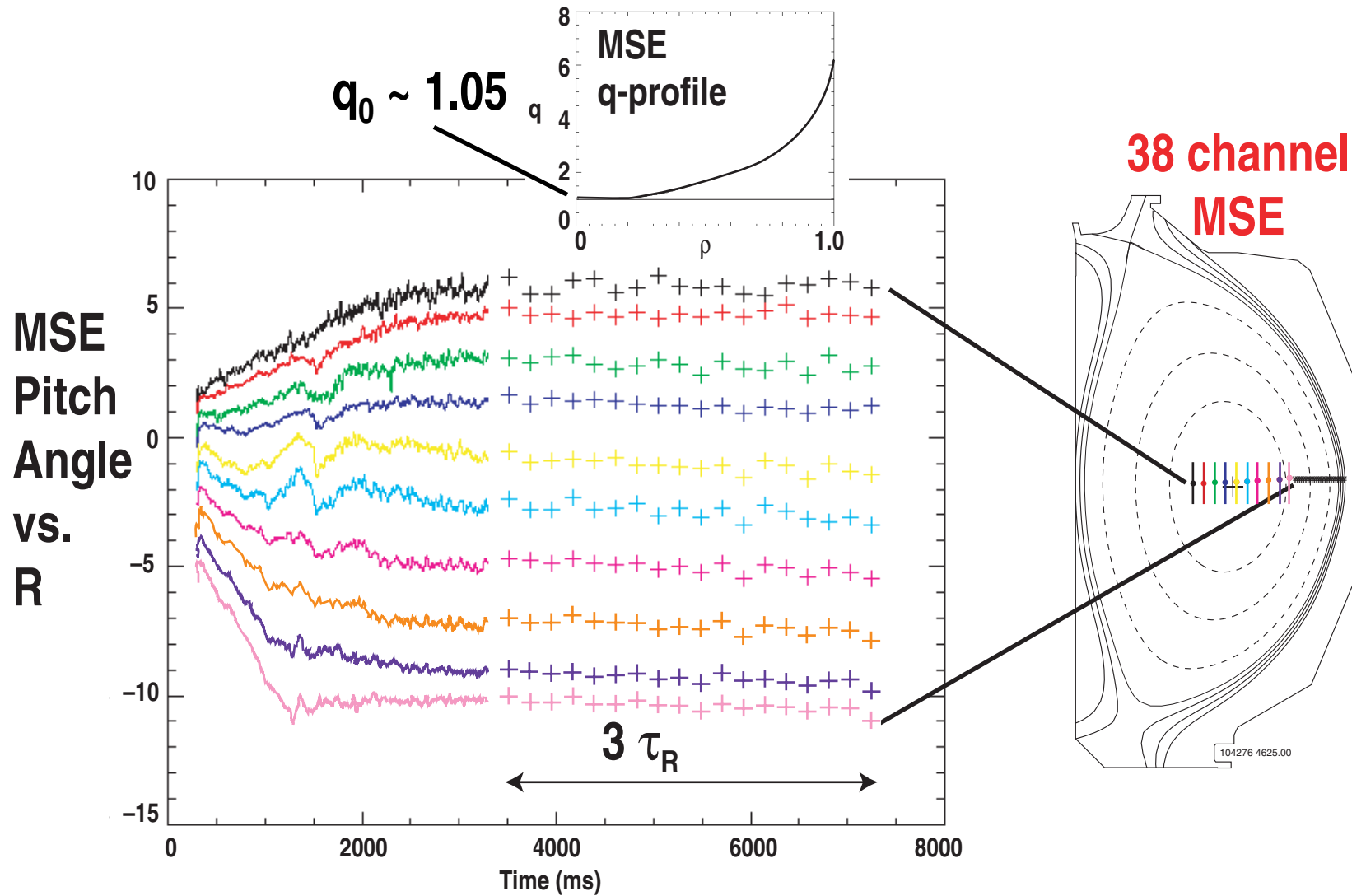
- After reaching the seed size, the stabilization is rapid because the mode growth rate is negative
- $\beta_N$  increases during stabilized phase
- Even in presence of large sawteeth the mode doesn't grow

# STEADY STATE HIGH PERFORMANCE DISCHARGES CAN BE ACHIEVED USING UNDERSTANDING OF STABILITY LIMITS AND DISCHARGE CONTROL

- $\beta$  controlled to remain ~20% below predicted RWM limit
  - $\beta$  also kept 5% below experimental 2/1 NTM  $\beta$  limit
- Discharge continued in steady state until beam termination
- No sawteeth
  - $q_0 \gtrsim 1$



# MSE shows J(r) profile has reached resistive equilibrium with $q_0 \sim 1.05$



# WHAT DISTINGUISHES A BURNING PLASMA FROM EXISTING EXPERIMENTS?

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- **Self-heating**
  - Less external control over profiles ( $p, j, \Omega$ )
- **Energetic particle effects**
  - Large isotropic population of fast ions
- **New ranges of dimensionless parameters**
  - $\rho_i^* = \rho_i/a \sim T^{1/2}/aB$
  - $S = \tau_A/\tau_R \sim aBT^{3/2}/n^{1/2}Z_{\text{eff}}$
  - $v^* = v_i/\epsilon\omega_{bi} \sim nqRZ_{\text{eff}}/\epsilon^{3/2}T^2$

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	DIII-D	C-MOD	JT-60U	JET	FIRE	IGNITOR	ARIES-RS	ITER-FEAT	ITER-FDR
$aB$ (m-T)	1.3	1.7	3.5	4.3	5.3	6.1	10	11	16

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# EXISTING EXPERIMENTS ARE SUFFICIENT TO INVESTIGATE MANY ISSUES OF MHD STABILITY

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- **Ideal MHD stability limits**
  - Profile dependence
  - Shape dependence
  - Aspect ratio dependence
- **Feedback stabilization of RWM**
- **ECCD stabilization of NTM**
- **Edge-driven instabilities**
  - Identification of instability
  - Dependence on bootstrap current
- **Stability with non-inductively driven current profiles**

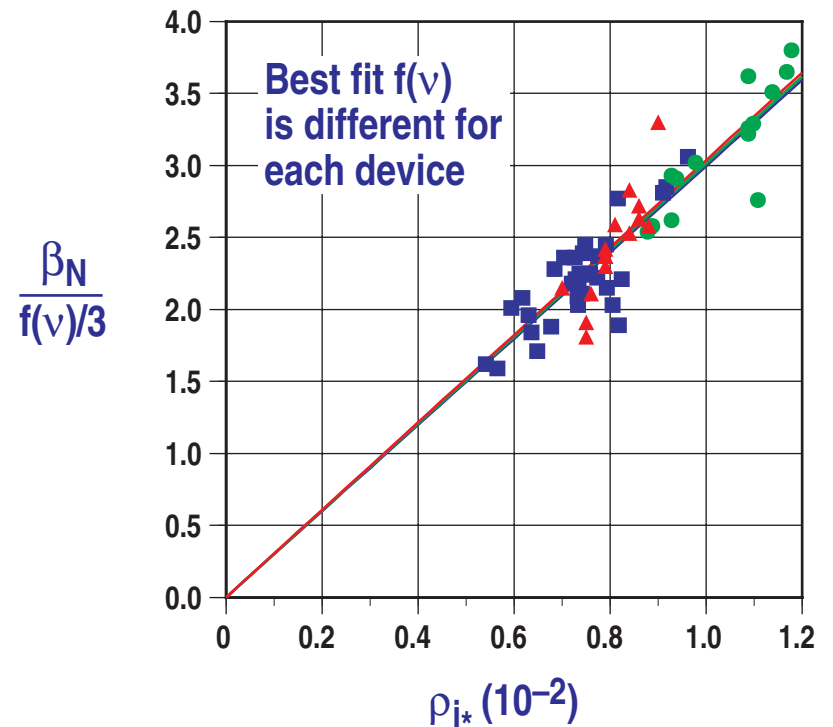
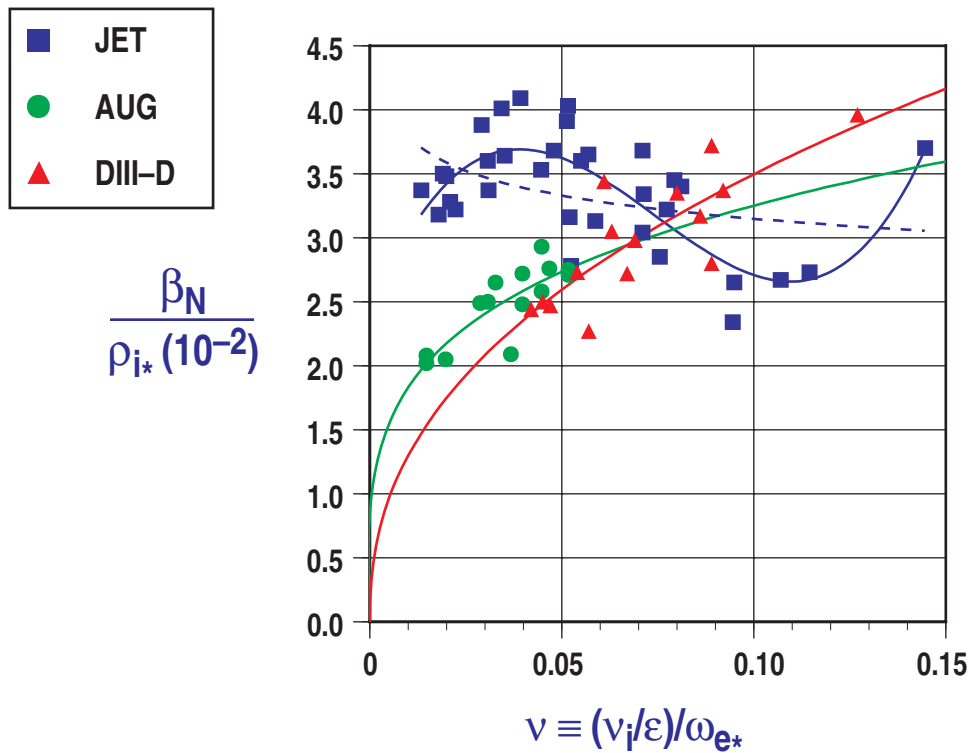
# BURNING PLASMA-SIZE EXPERIMENTS (WITHOUT ALPHA HEATING) ARE REQUIRED TO INVESTIGATE SCALING OF MHD STABILITY PHYSICS

- **NTM beta limit scaling**
  - Threshold island size decreases with decreasing  $\rho_i^*$
  - Seed island size decreases with increasing S
- **Edge-driven instabilities**
  - Edge gradients determine stability limit
  - Pedestal width determines coupling to core
  - Scaling of edge parameters is not well understood
- **Resistive wall mode stability**
  - Rotation frequency required for stabilization may increase with S ( $\Omega \tau_A \sim 0.05$ )
- **Runaway avalanche during disruption**
  - Number of e-foldings increases with plasma current
  - Runaway electron current multiplication
    - ★  $\gtrsim 10^2$  at  $I_p = 2$  MA
    - ★  $\gtrsim 10^6$  at  $I_p = 5$  MA



# NTM THRESHOLD SCALES LINEARLY WITH NORMALIZED ION LARMOR RADIUS

- But scaling of  $\beta_N/\rho_{i*}$  with collisionality is not consistent between machines
  - Possible additional dependence on  $\rho_{i*}$  or  $S$
- $\beta_N \propto \rho_{i*} f(v)$  is consistent with polarization/inertial model of Wilson et al.

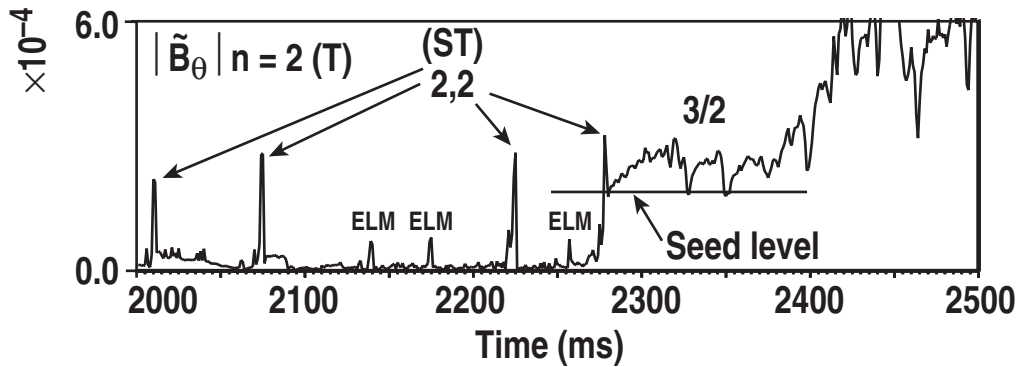


- Sawtooth-induced 3/2 NTM, ELMing H-mode

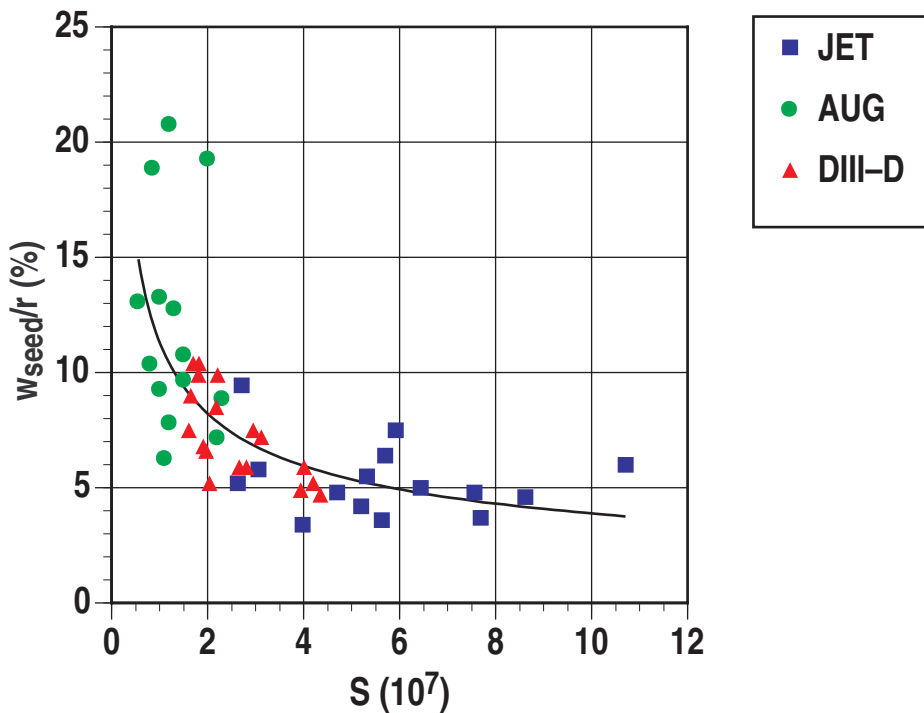
# SAWTOOTH INDUCED SEED ISLANDS SCALE INVERSELY WITH MAGNETIC REYNOLD'S NUMBER

- Seed islands estimated from  $m/n = 3/2$  Mirnov level upon excitation

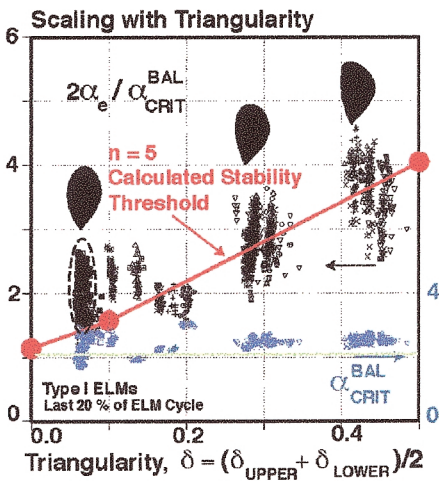
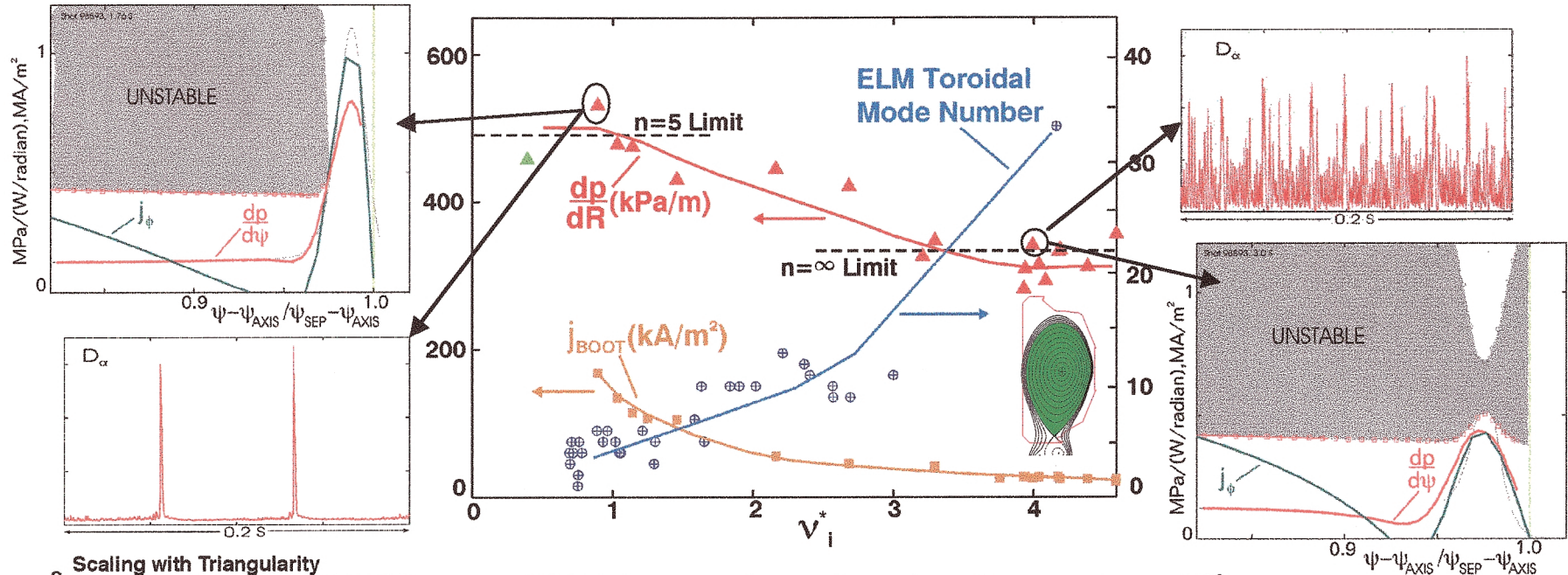
★  $w_s \approx \left( \frac{16rR |\tilde{B}_r|}{3s B_T} \right)^{1/2}$  with  $|\tilde{B}_r| \approx \frac{1}{2} \left( \frac{b}{r} \right)^4 |\tilde{B}_\theta|_{\text{wall}}$



- Best fit has  $w_{\text{seed}}/r \propto S^{-0.46 \pm 0.05}$ , correl  $r = -0.74$  consistent with dynamical coupling model of Hegna et al.



# EDGE STABILITY AND ELM CHARACTER DEPEND CRITICALLY ON COLLISIONALITY



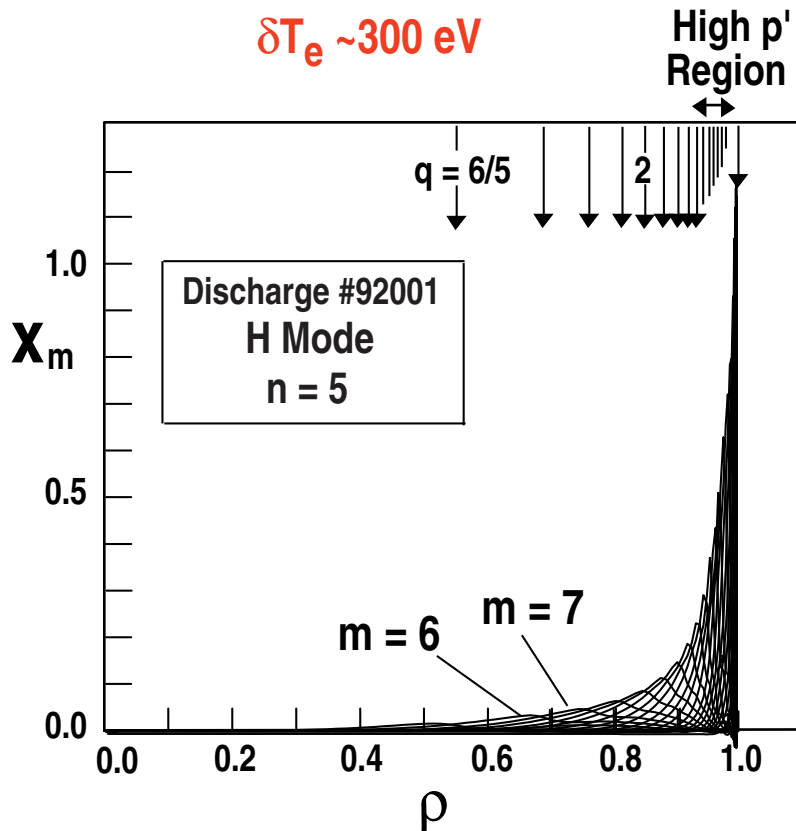
- ◆ With increasing edge density or  $v_i^* \propto n/T^2$ .
  - Calculated  $j_{BOOT}$  decreases  $\Rightarrow$  edge magnetic shear increases,  $S \approx S_0 - 2 \left\langle \frac{j_{TOR}^{EDGE}}{j_{TOR}} \right\rangle \Rightarrow$  SS access lost
  - ELM modes increase in n.
  - Pressure gradient is reduced from calculated limit for n=5 edge localized ideal kink/ballooning (GATO) to ideal high n ballooning mode limit (BALOO).

# ELM SIZE CORRELATES WITH RADIAL WIDTH OF PREDICTED UNSTABLE INTERMEDIATE $n$ KINK MODE

- Highly localized instability computed from GATO

⇒ Type I ELM has little effect

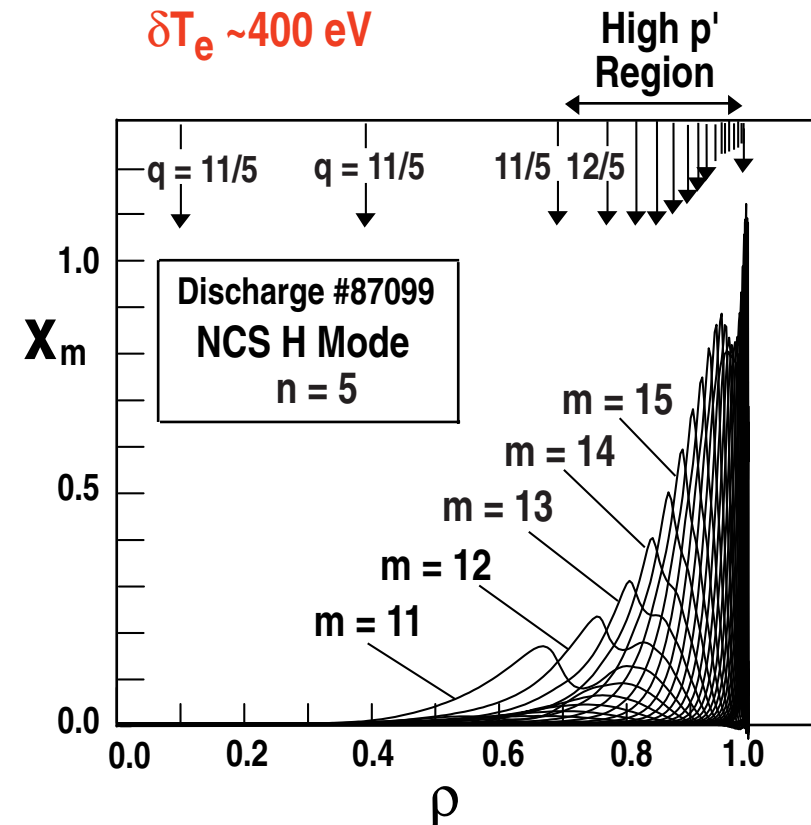
$\delta T_e \sim 300$  eV



- Predicted instability computed from GATO code penetrates into core

⇒ High performance is lost

$\delta T_e \sim 400$  eV



# A BURNING PLASMA (STRONG ALPHA HEATING) IS NEEDED TO INVESTIGATE KEY ISSUES OF MHD STABILITY

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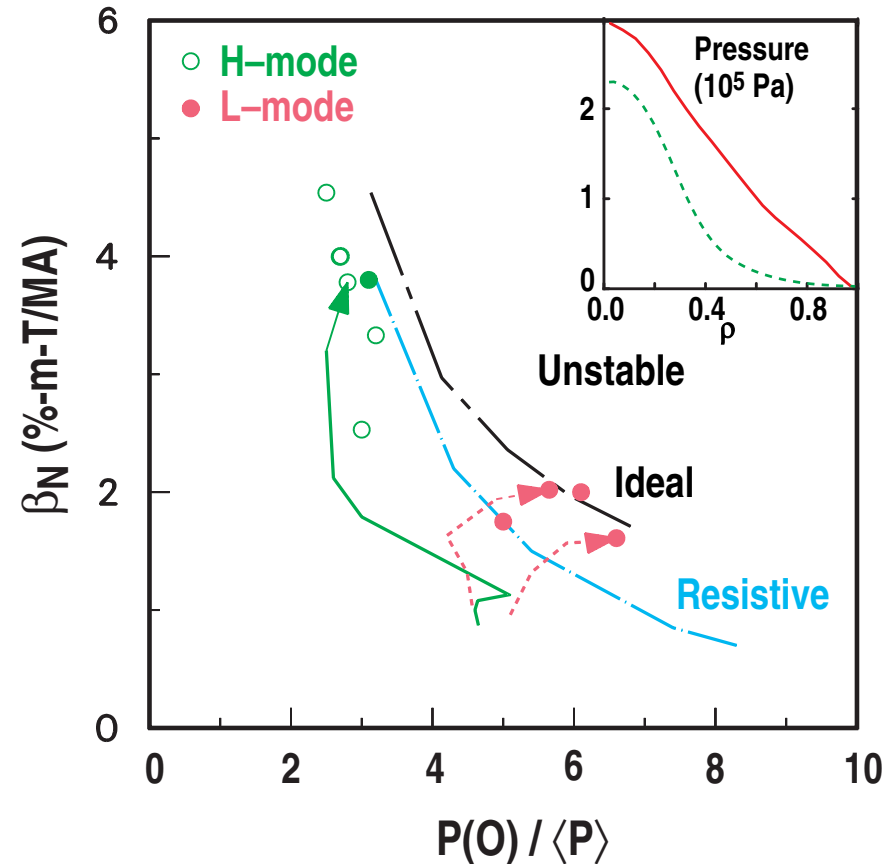
- Energetic particle interactions with MHD modes (sawteeth, fishbones, TAE, ballooning modes, etc.)
  - Stabilization or destabilization of MHD modes by alphas
  - Enhanced transport of alphas by MHD modes
- Self-heating ( $P_\alpha \gg P_{\text{external}} \Rightarrow Q \geq 10$ )
  - Stability limits with pressure profiles determined by alpha heating
  - Plasma rotation with little or no external momentum input (RWM stability, mode locking, error field sensitivity)

$$\Omega \sim \omega^* \sim T/a^2B \quad ?$$

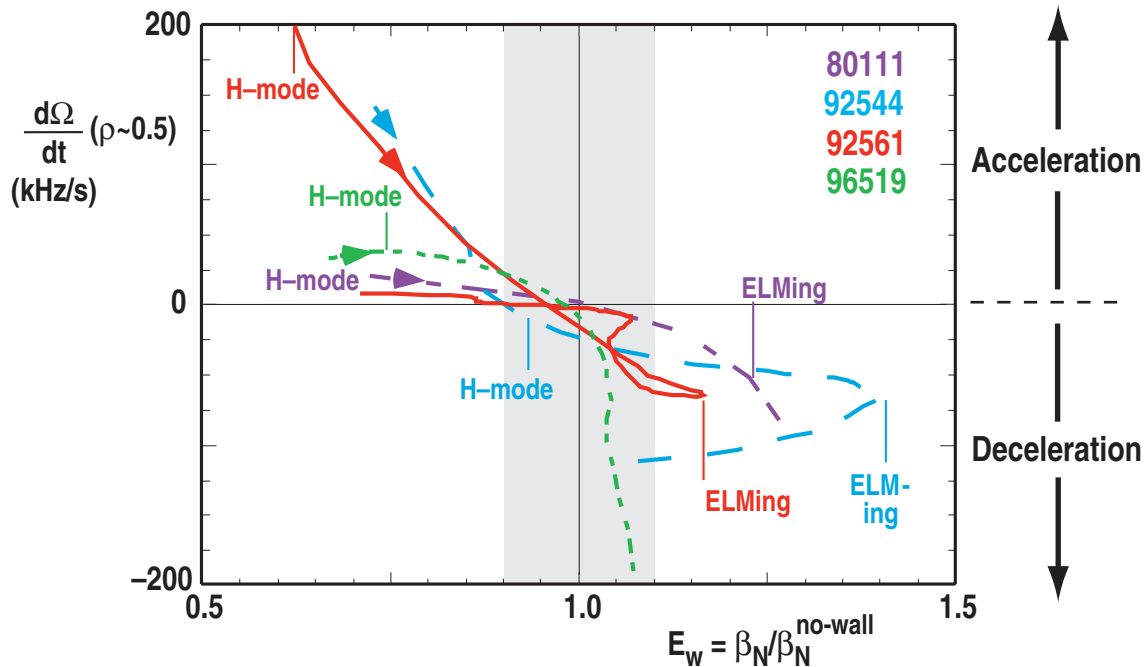
- Steady-state operation ( $\tau > \tau_{\text{CR}} \sim a^2 T^{3/2}/Z_{\text{eff}}$ )
  - Stability limits with self-consistent current density and pressure profiles

# STABILITY LIMIT DEPENDS STRONGLY ON THE FORM OF THE PRESSURE PROFILE

- DIII-D high  $p_0/\langle p \rangle \sim 6.0$  (L-mode):  
 $\beta_N \lesssim 2.5$   
 — Limited by fast  $n = 1$  disruption
- TFTR high  $p_0/\langle p \rangle \sim 6.0$  (ERS-mode):  
 $\beta_N \lesssim 2$   
 — Limited by fast  $n = 1$  disruption
- DIII-D low  $p_0/\langle p \rangle \sim 2.5$  (H-mode):  
 $\beta_N \lesssim 4$   
 — No disruption  
 limited by ELM-like activity from finite edge pressure gradients



# ROTATION DECELERATES ABOVE THE NO-WALL $\beta$ LIMIT (EVEN WITH LARGE TORQUE)



- Two competing models are being investigated
  - Gimblett and Hastie torque balance model with marginally unstable RWM predicts qualitative behavior
  - New data is consistent with resonant amplification of static error fields by marginally stable RWM

# CONCLUSIONS

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- Some issues of MHD stability require burning-plasma parameters to investigate
  - NTM beta limit scaling
  - Edge-driven instabilities
  - Resistive wall stabilization
  - Disruption scaling (runaway avalanche)
- Some key issues of MHD stability can only be addressed with strong alpha heating
  - Energetic alpha interactions with MHD modes
  - Stability with profiles determined by self-heating ( $t \gg \tau_E$ )
  - Stability with self-heating and relaxed current density profile ( $t \gg \tau_{CR}$ )
- Many of the issues requiring a burning plasma are not purely MHD stability issues but issues of integration (transport, profile control, burn control, etc.)



# INTEGRATION OF SEPARATE ELEMENTS MAY BE THE MOST IMPORTANT MISSION FOR A BURNING PLASMA EXPERIMENT

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- Strong coupling of transport, heating, and stability leads to a more “self-organized” plasma than in a short-pulse, externally heated tokamak
  - Pressure profile → Fusion rate → Alpha heat deposition → Thermal transport → Pressure profile
  - Pressure profile → Bootstrap current → Current profile → Thermal transport → Pressure profile
- MHD instabilities can intervene in these loops:
  - Pressure, current density, and fast ion profiles → Instabilities → Modification of profiles
- Investigation of such a complex, non-linear system represents a scientific challenge, and may yield some surprises

**RECOMMENDATION: A “next step” burning plasma experiment is needed as the only way to address this challenge**