

NIMROD Applications for MHD Control

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for the

NIMROD Team (<https://nimrodteam.org>)

16th Workshop on MHD Stability Control

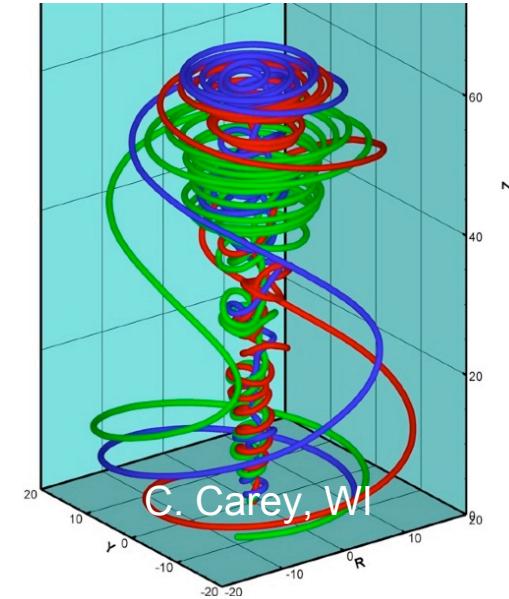
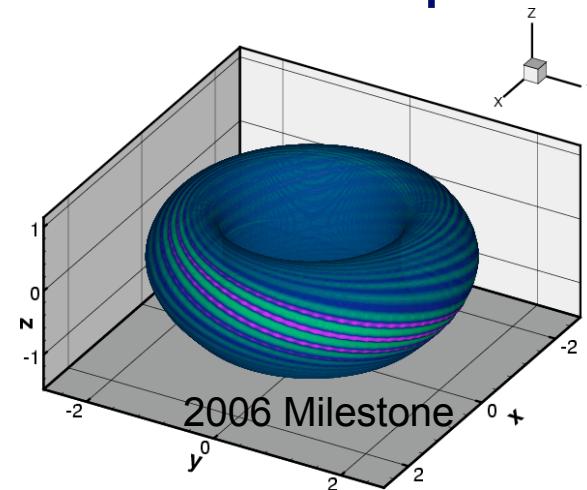
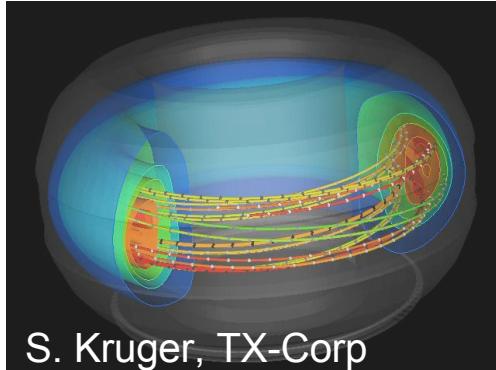
San Diego, CA Nov. 20-22, 2011



Outline

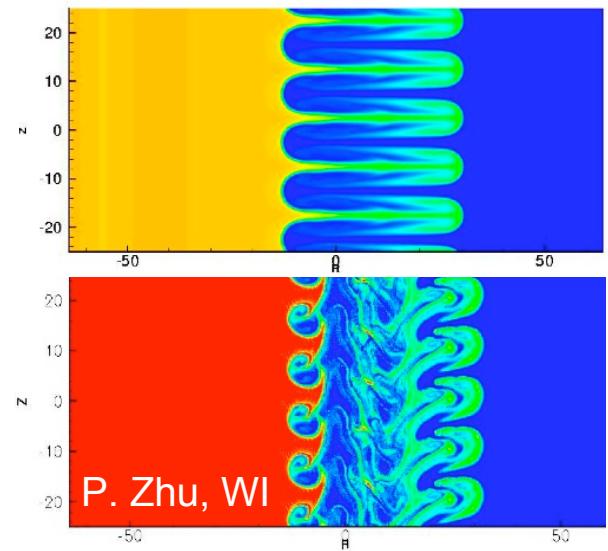
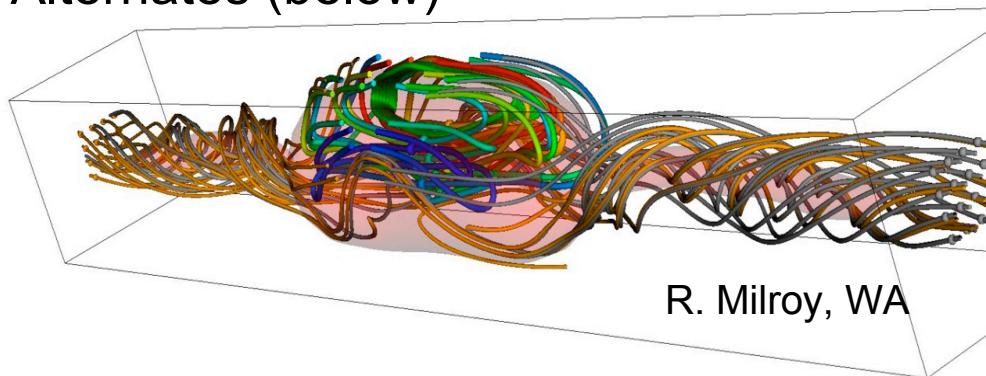
- Introduction of NIMROD
- Control and control-related applications
 - RF/MHD modeling
 - Disruption mitigation
 - Imposed asymmetry
- Relevant model development
- Concluding remarks

Introduction: NIMROD is a nonlinear macroscopic dynamics code for multi-scale plasma studies.



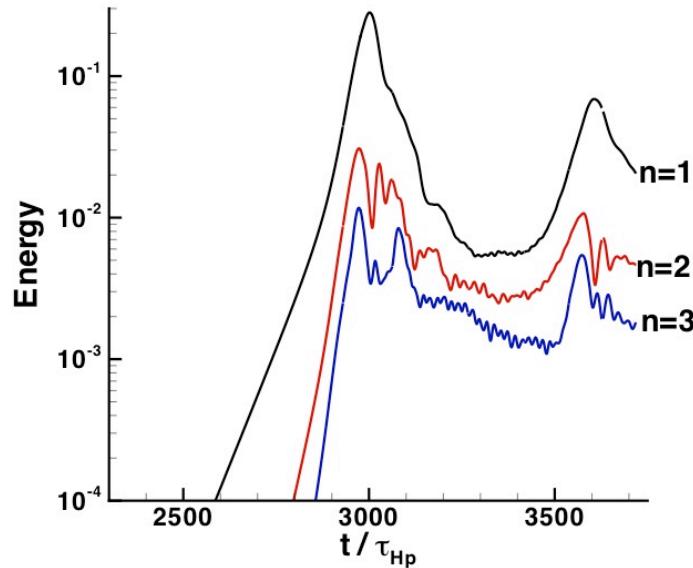
It has been applied to many configurations including:

- Tokamak tearing and ELMS (above)
- MHD jets (above right) and space physics
- Basic plasma phenomena (below right)
- Alternates (below)

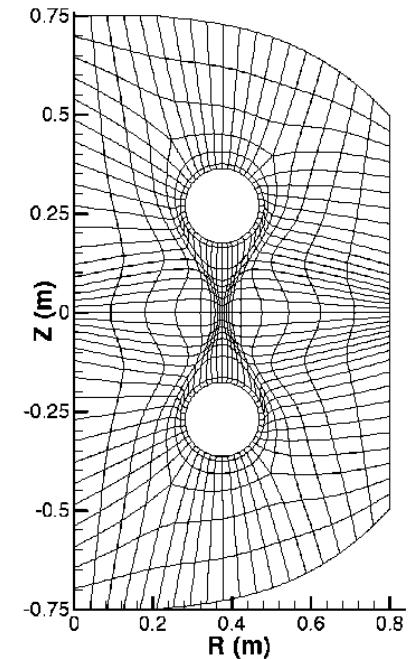
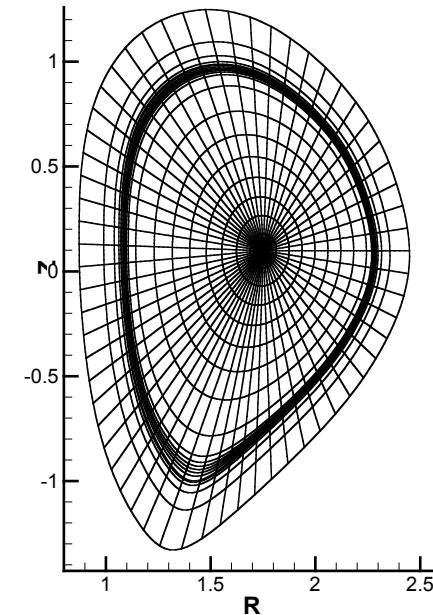


The implementation considers 3D nonlinear dynamics within a geometrically 2D domain.

- The fundamental dependent fields, \mathbf{V} , \mathbf{B} , n , and T (or T_i and T_e), are evolved in time from initial conditions using fluid-based models.
- Variations over the poloidal plane are represented on a mesh of spectral elements.



Two-fluid sawtooth evolution of kinetic fluctuation energies. (Comp. has $0 \leq n \leq 42$.)

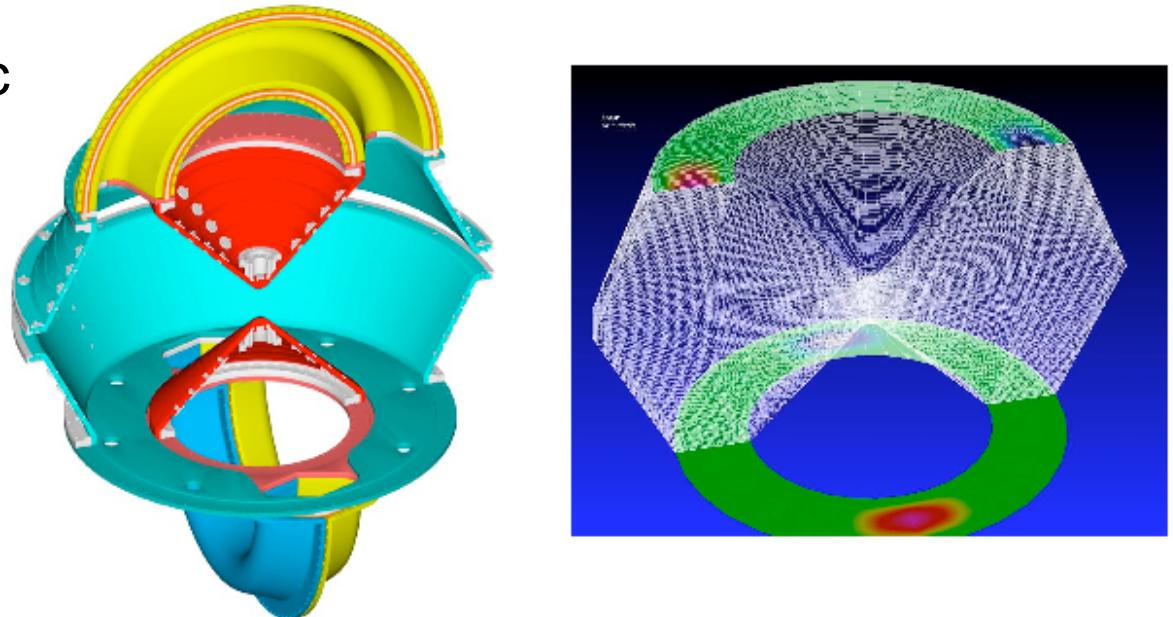


Examples of meshes used for DIII-D tearing (left) and for MRX reconnection (right).

- Variations over the periodic (toroidal or straight) coordinate are represented by finite Fourier series.

Computations related to control include sources or boundary conditions that model the physical system.

- Volumetric sources of momentum density (fluid or electron) and energy density have been used.
- Fluxes and fields applied at the wall may be symmetric or asymmetric.
- A common symmetric flux is the application of E_ϕ from loop voltage.
- An example of asymmetric fields and fluxes are the time-dependent conditions used to model HIT-SI injectors.

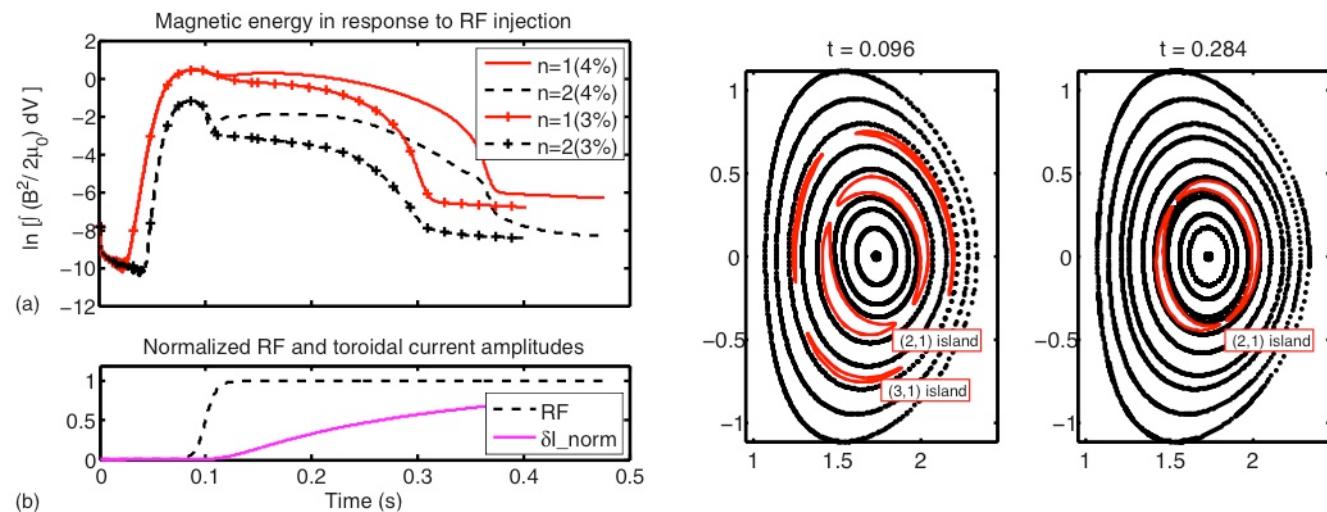


CAD drawing of the HIT-SI spheromak chamber (left) and NIMROD mesh and injector B_z (in color) used in simulation (right). [Izzo and Akcay, Univ. WA]

Control-related applications: To date, the two applications most closely tied to plasma control are RF/MHD coupling and disruption mitigation.

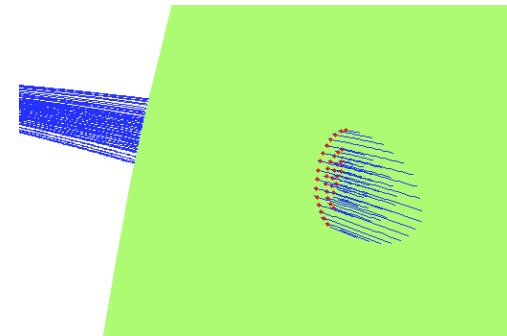
- The RF/MHD effort seeks to model ECCD stabilization of tearing modes as a part of comprehensive simulation capability.
- The project is led by Tom Jenkins of Tech-X in collaboration with Dalton Schnack, Eric Held, and Bob Harvey as part of SWIM.
- Initial work demonstrated island suppression with an ad hoc, toroidally symmetric current drive. [Jenkins, PoP 17, 12502, 2010]

Application of symmetric source affects magnetic islands over the resistive time-scale.

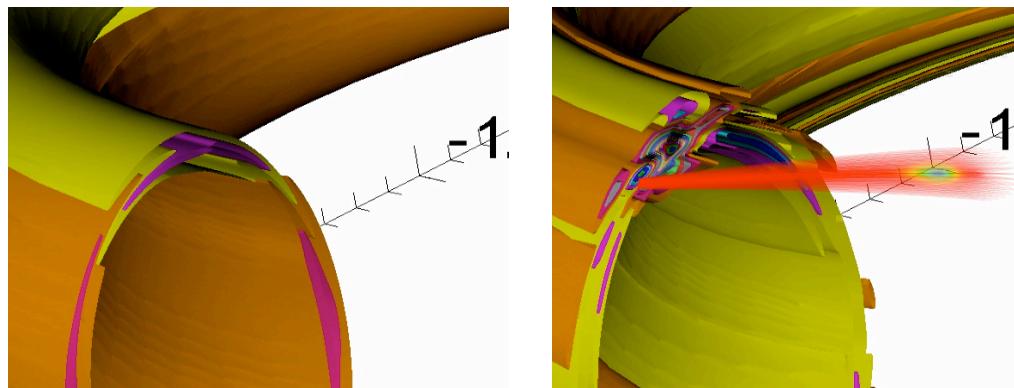


Recent development couples GENRAY ray-tracing and QLCALC diffusion calculation for localized deposition.

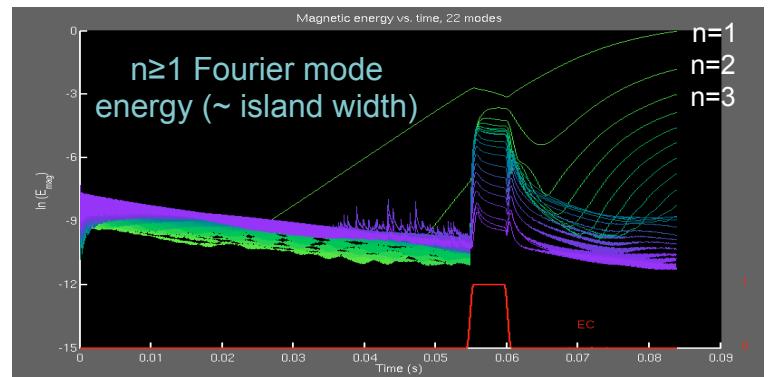
- SWIM's Integrated Plasma Simulator (IPS) is used to coordinate computations.
- Present logic is intended to model an ideal system that targets the island O-point.
- Future work will implement more realistic control logic.
- Predicting power requirements for ITER is the primary physics objective.



GENRAY ray bundle intercepting NIMROD data plane.



Temperature contours before (left) and just after (right) ECCD feedback on 2/1 island.



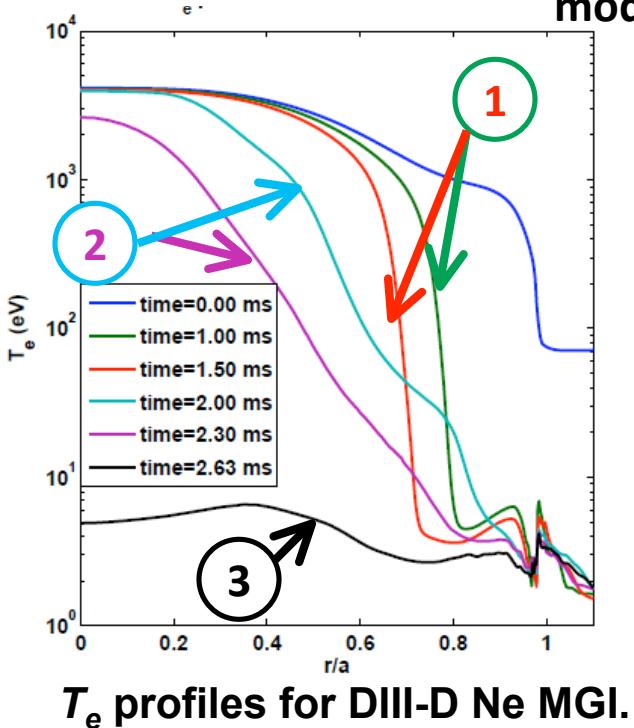
NIMROD magnetic fluctuation energies respond to simulated control.

Disruption mitigation: Simulations of MGI track impurity density and incorporate radiative cooling.

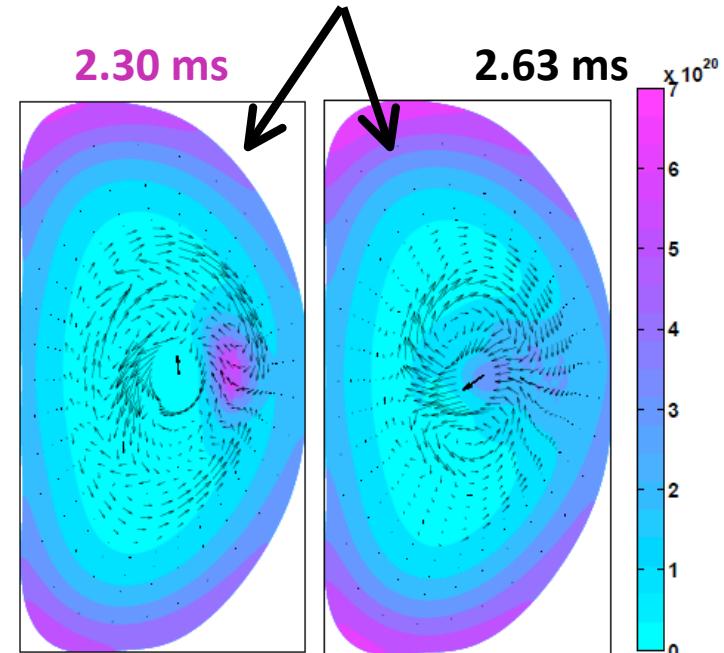
- Val Izzo uses 'NIMRAD' to investigate MHD mixing and resulting efficiency.
- The computations couple NIMROD and the KPRAD radiation code.
- Validation studies have been conducted using T_e and n_e profiles from DIII-D and C-MOD. [Izzo, et al., PoP 15, 56109, 2008]

MGI simulation of DIII-D with Ne injection:

- 1) Injected Ne cools edge, triggers MHD
- 2) Destruction of flux surfaces reduces Te gradient
- 3) Core Te drops rapidly due to $m=1/n=1$ mode



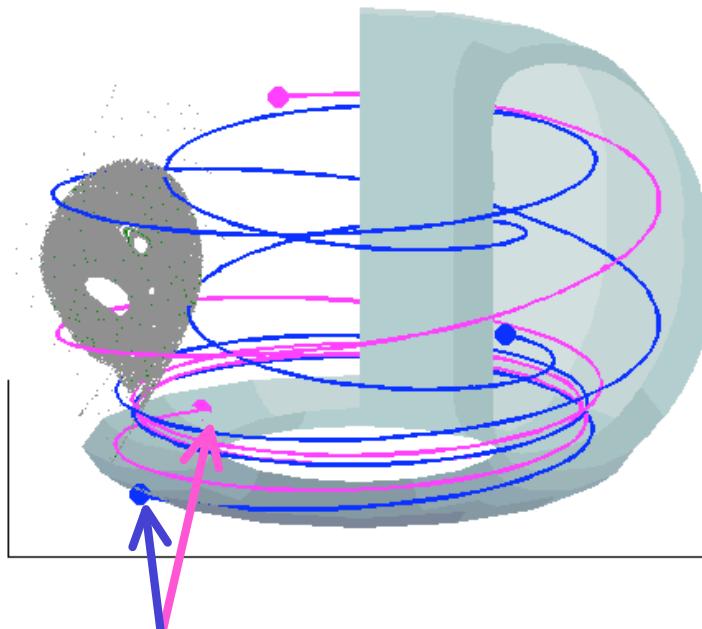
Ne density shows efficient mixing into core due to flows associated with $1/1$ mode, further enhancing core cooling.



Test-particle model calculates runaway electron (RE) drift orbits as MHD fields evolve.

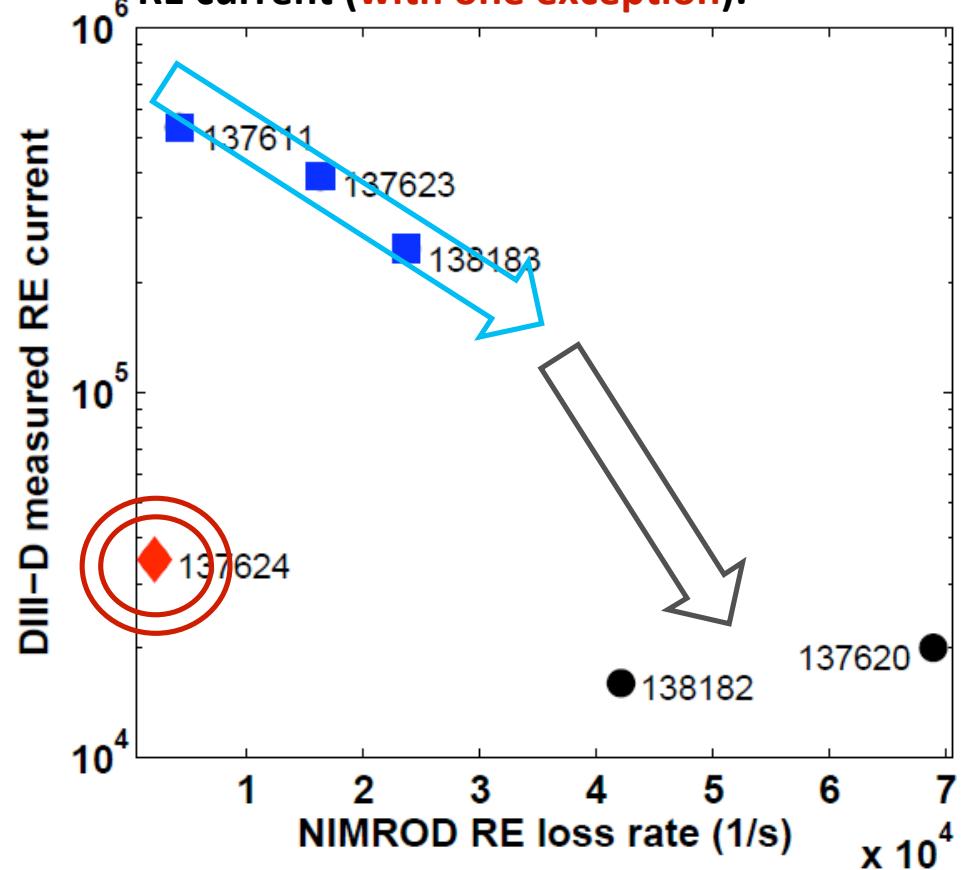
- Tracer electrons evolved with collisions, ∇B and curvature drifts, bremsstrahlung and synchrotron radiation predict confinement and strike points.

Example RE orbits calculated during NIMROD run; typically thousands per simulation.

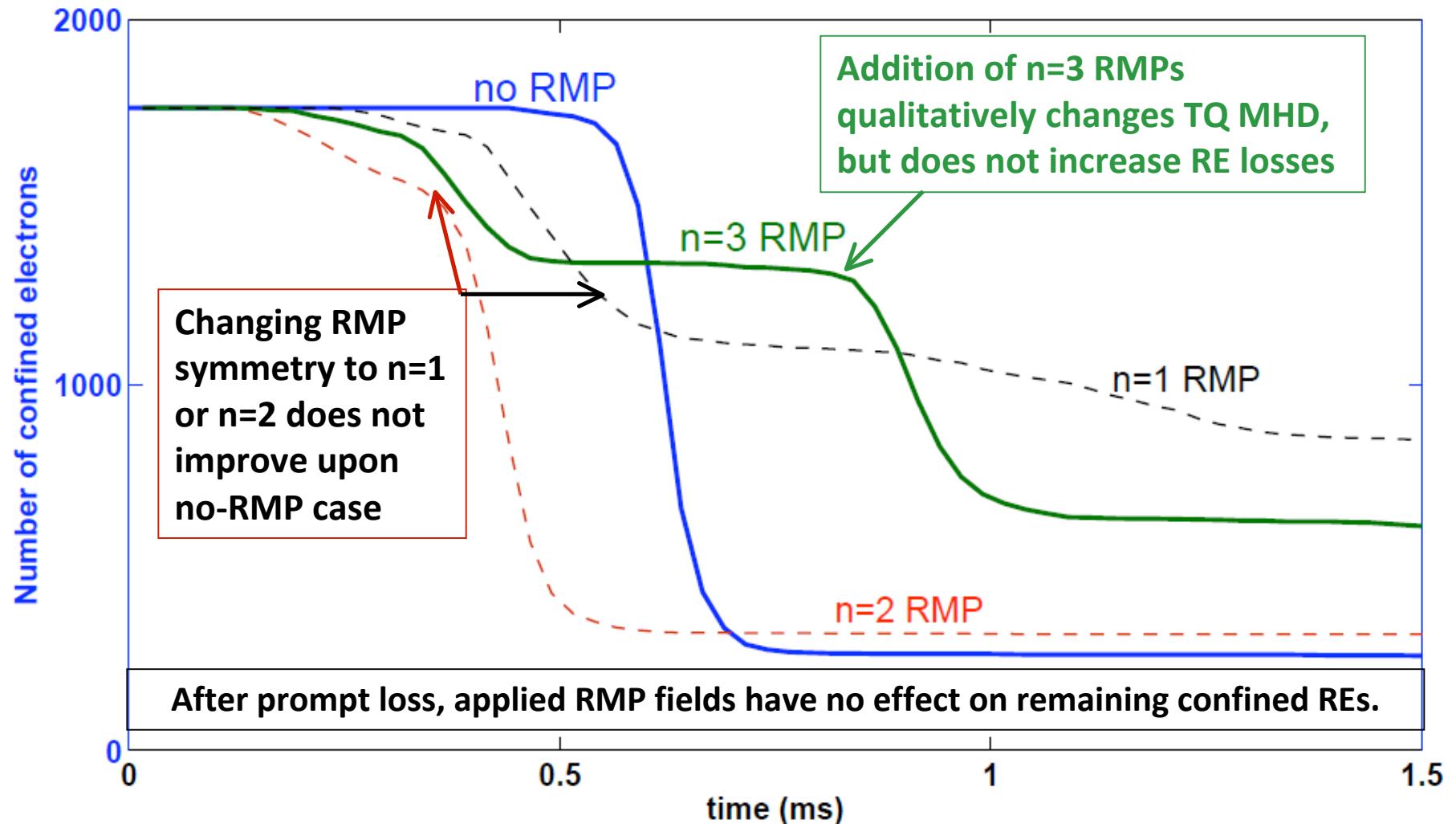


When magnetic fields become stochastic, REs escape, striking outer divertor.

NIMROD predicted RE loss rates show expected relationship to DIII-D measured RE current (with one exception).



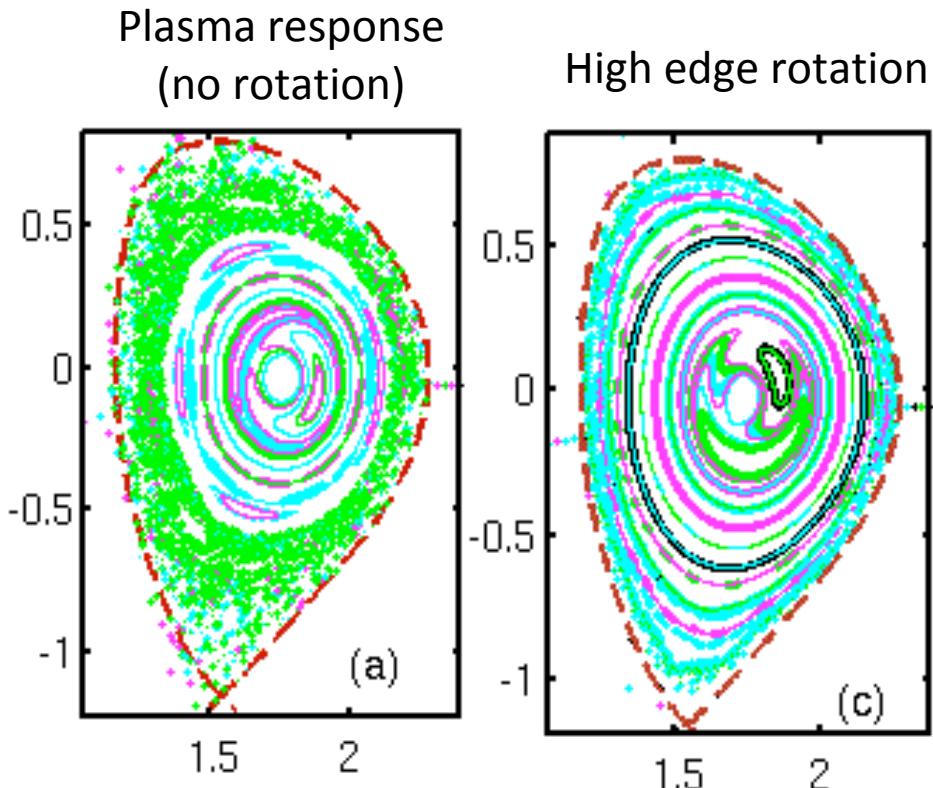
Nonlinear computations with RE tracing and RMP predict little enhancement of losses during thermal quench in DIII-D.



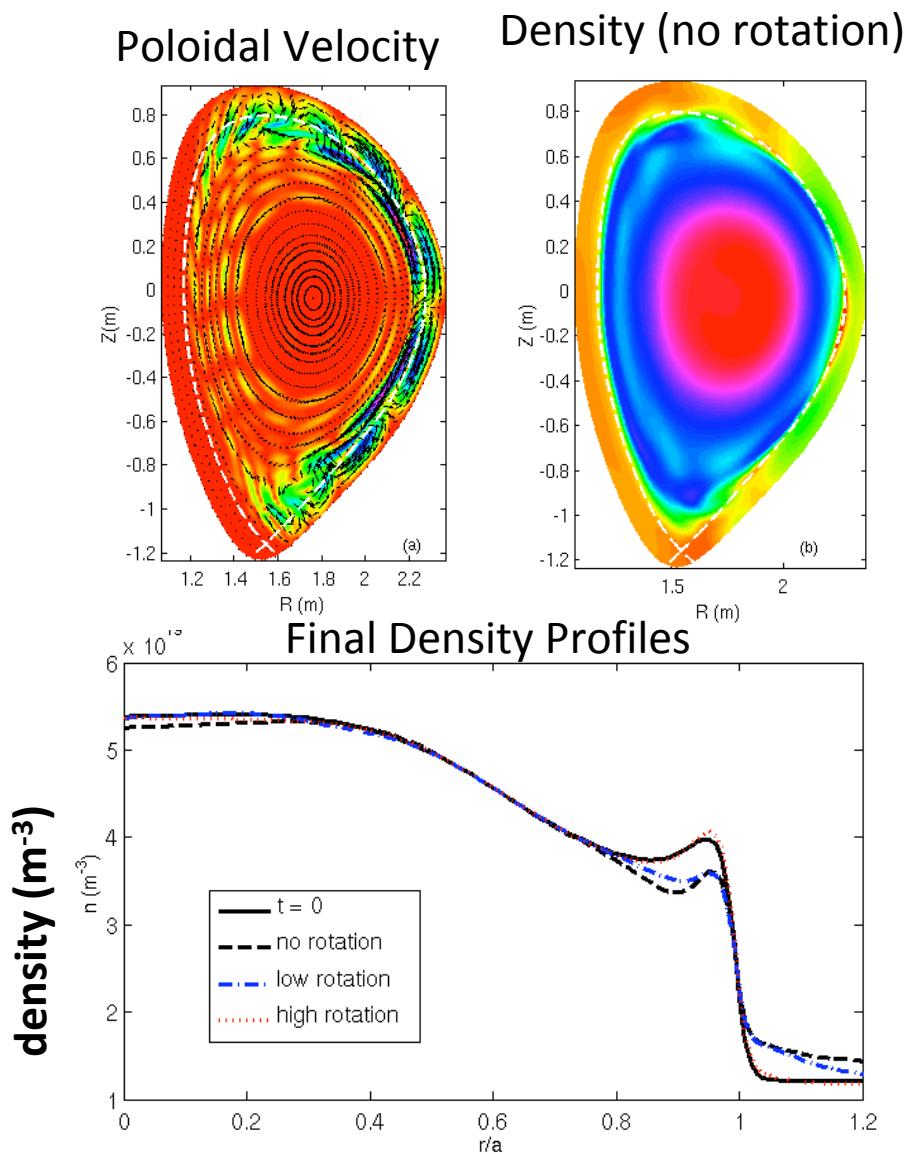
DIII-D experiments showed ambiguous effects of applying $n=3$ RMPs before TQ, i.e. no clear benefit. As in NIMROD, post-TQ RMPs had no effect.

Imposed asymmetry: Simulations are being used to investigate possible benefits of imposing 3D fields.

- Izzo's DIII-D simulation with $n=3$ RMP investigates MHD density pump-out.

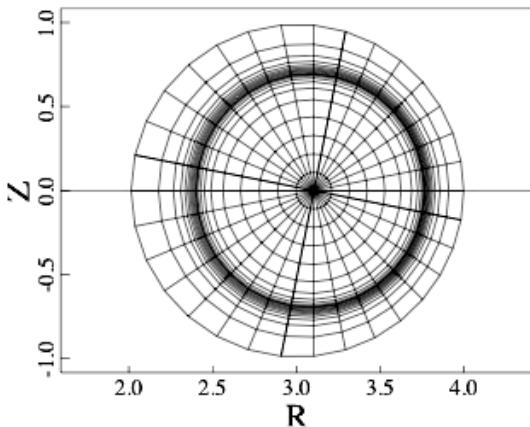


Large toroidal rotation near the edge screens the RMP fields and eliminates the density pump-out effect.



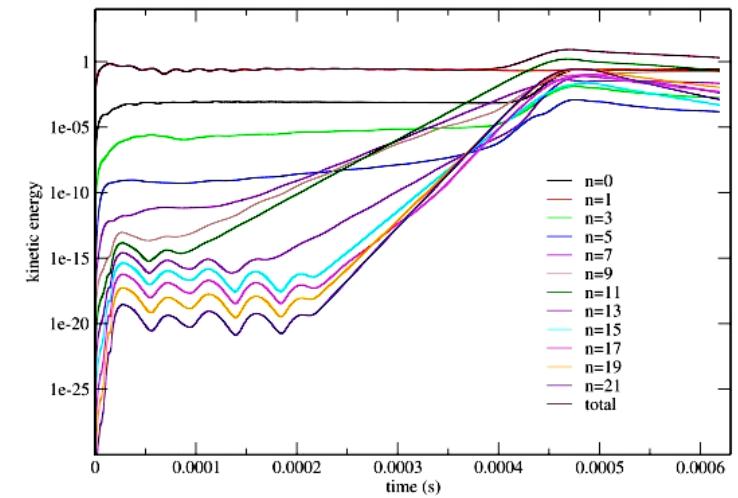
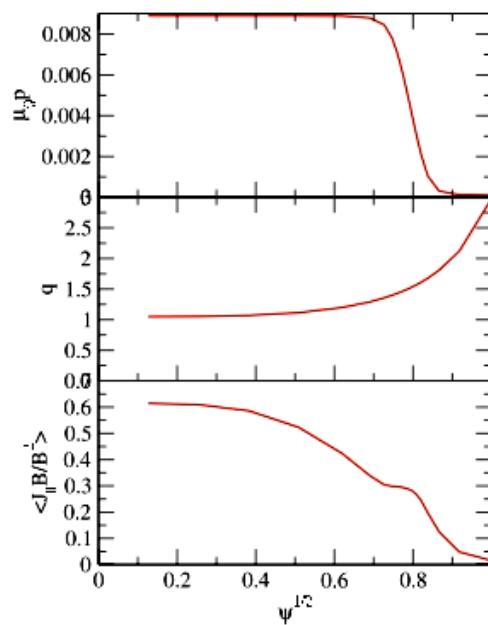
A new RMP study uses circular cross-section toroidal geometry to facilitate analytical comparisons for ELMs.

- Ping Zhu previously performed analytics and computations for an intermediate ballooning regime (no RMP). [Zhu, NF 49, 95009, 2009]
- The new study adds RMP perturbations.



- ▶ $\beta \sim 1\%$ at pedestal top
- ▶ $n = 7 - 13$ modes weakly unstable

The problem setup has uniform pressure within a pedestal region.

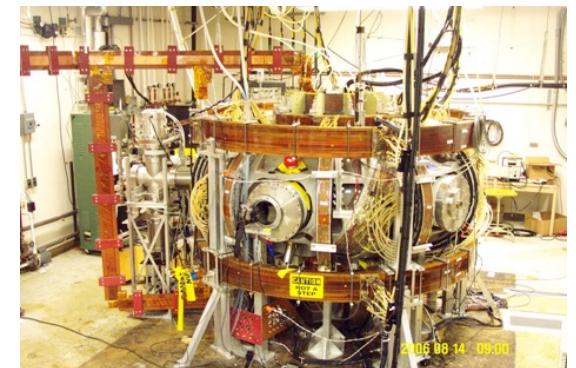


Evolution of perturbed kinetic energy shows RMP effects prior to ballooning saturation.

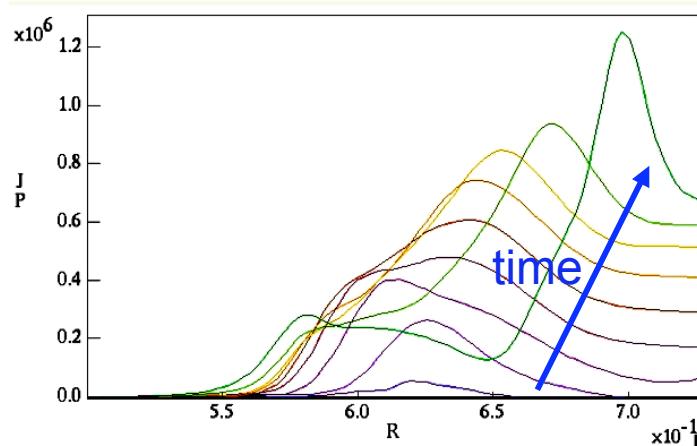
- Poincaré surfaces (not shown) have islands before ballooning saturation, a stochastic layer after.

Computations of the Compact Toroidal Hybrid consider flux-surface evolution with large 3D shaping.

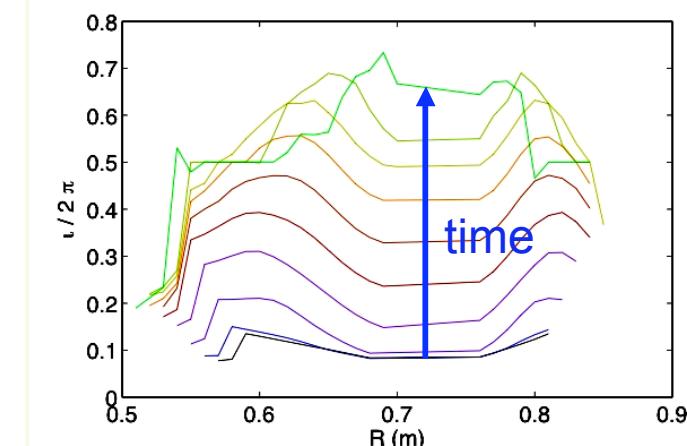
- Auburn's CTH is a heliotron that can have significant rotational transform from plasma current.
- While imposed magnetic fields are asymmetric, helical coils lie outside a toroidally symmetric surface.
- Mark Schlutt (WI) has incorporated vacuum fields provided by Jonathan Hebert (Auburn) into NIMROD boundary conditions.



Auburn CTH experiment.



Evolution of J_ϕ profile shows current penetration for applied $V_{loop} = 4$ V.



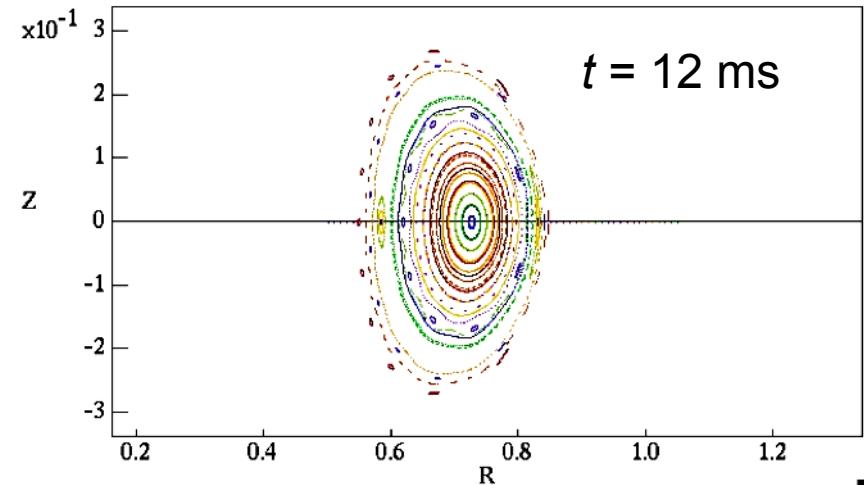
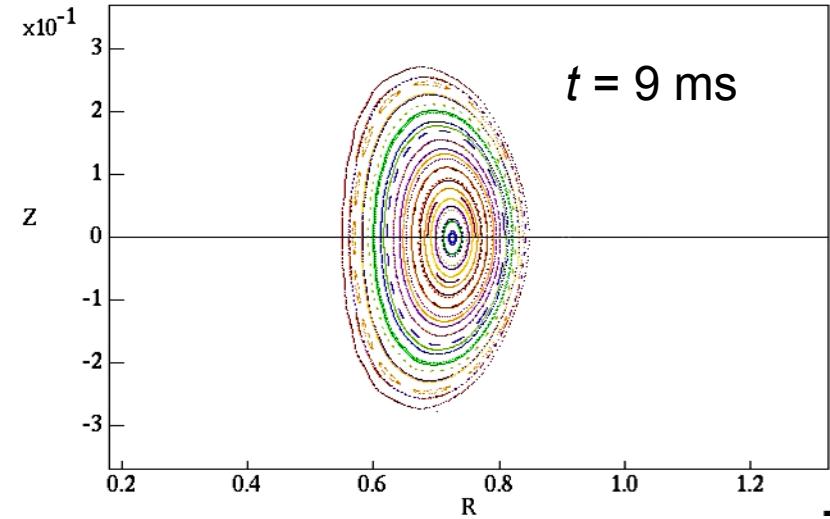
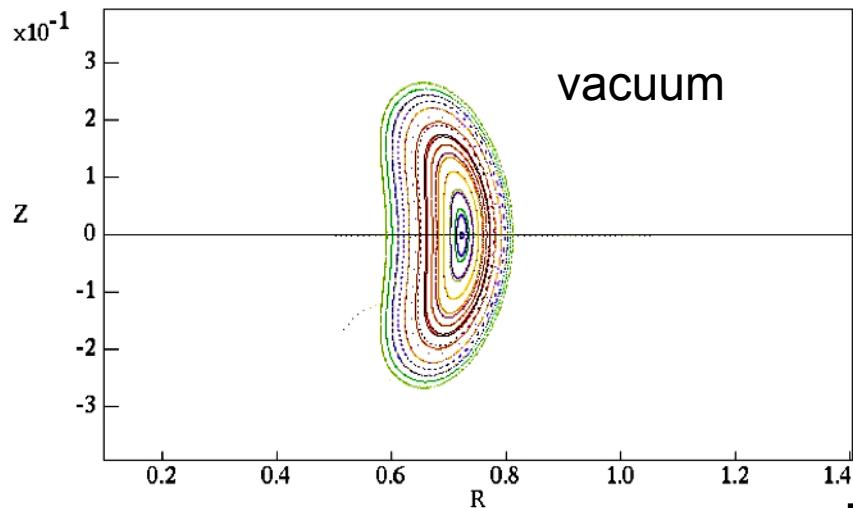
Evolution of rotational transform indicates 2/1 activity late in time.



AUBURN UNIVERSITY



Magnetic topology evolution shows 2/1 island formation prior to loss of flux surfaces.

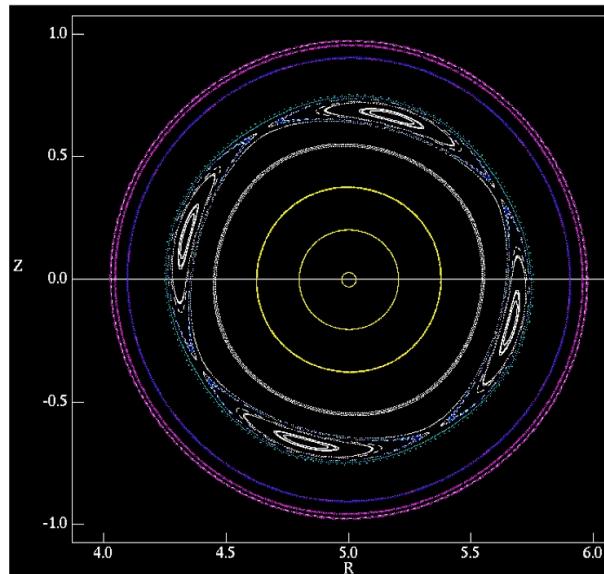


- These 0- β , fixed-resistivity computations show behavior that appears qualitatively similar to the experiment.
- After convergence testing, next steps include temperature evolution to model realistic resistivity profiles that change in time.

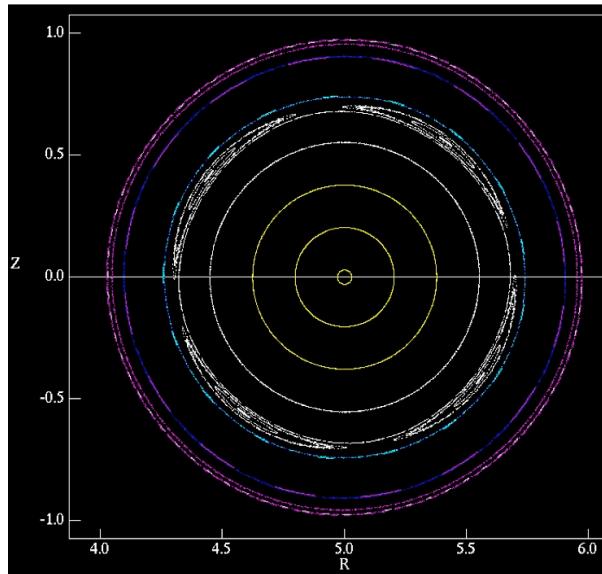
Poincaré surfaces show growth of closed-flux volume then island development.

A study of field-error penetration uses straight cylindrical geometry for verification.

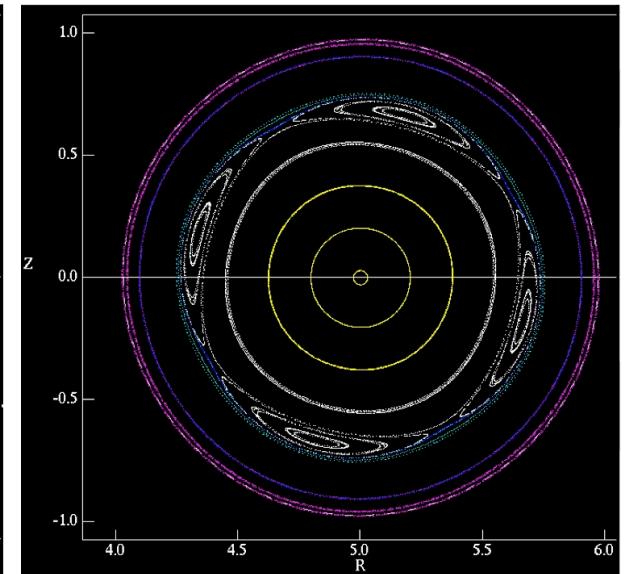
- Andrea Montgomery (WI) has implemented boundary conditions for a thin resistive wall with a helical coil located outside the resistive wall.
- Simple analytical profiles of the form $q(r) = q_0 [1 + (r/a)^{2\lambda}]^{1/\lambda}$ are used to select a stable resonant mode.



Poincaré surface from a linear computation without flow indicates perturbation amplitude.



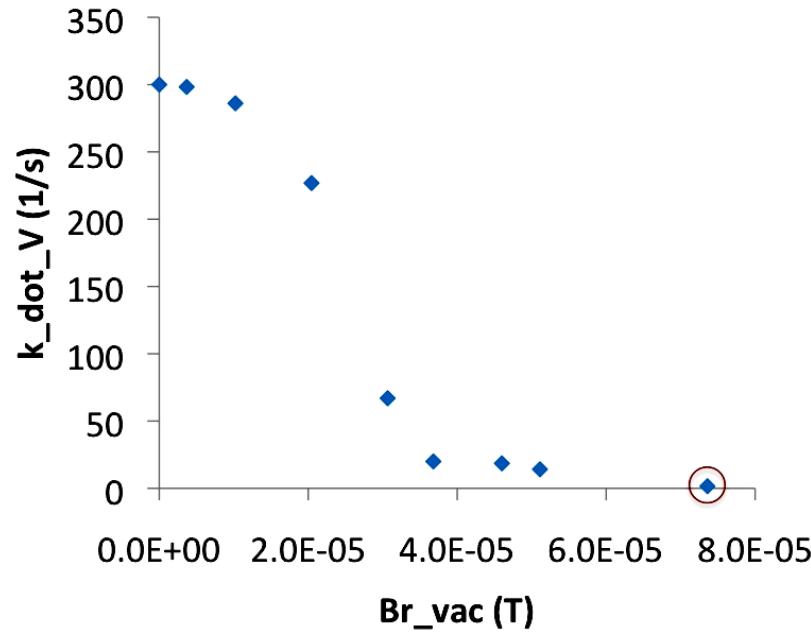
A linear computation with flow shows significant shielding.



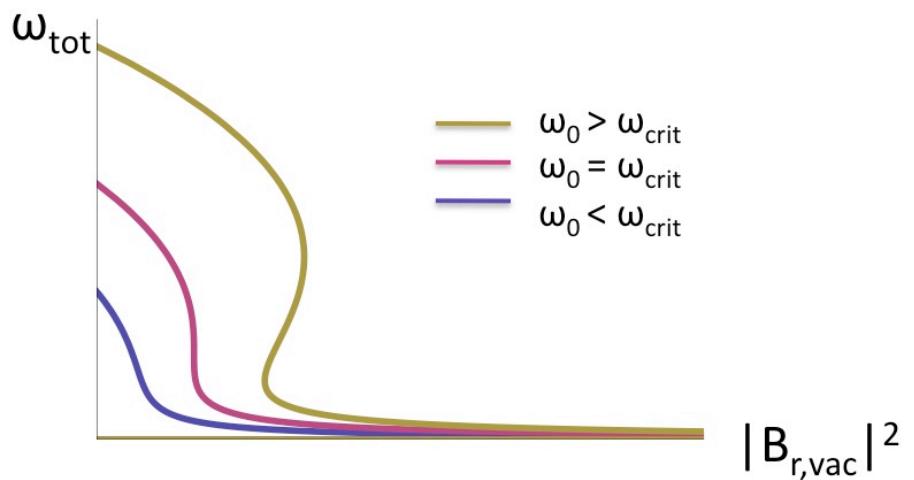
Nonlinear effects with the same flow and perturbation induce locking.



A series of nonlinear computations demonstrates ‘sub-critical’ locking behavior.



Nonlinear results on $k \cdot \langle V \rangle$ as the perturbation amplitude is varied.



Sketch of different classes of locking behavior for flows that are above and below the critical rate.

- Locking information from Waelbroeck, PFB 1, 1989 and Fitzpatrick PoP 5, 1998 will be adapted for verifying these large-island cylindrical results.
- Computations for less-stable, larger-S (7×10^7) conditions maintain smaller islands and will be used for verifying small-island behavior.
- Testing the response in toroidal geometry is the next computational step.

Relevant model development: Our computations are based on single- and two-fluid plasma models.

$$\frac{\partial n}{\partial t} + \nabla \cdot (n\mathbf{V}) = 0$$

particle continuity

$$mn\left(\frac{\partial}{\partial t} + \mathbf{V} \cdot \nabla\right)\mathbf{V} = \mathbf{J} \times \mathbf{B} - \nabla \sum_{\alpha} nT_{\alpha} - \nabla \cdot \underline{\Pi}$$

flow evolution

$$\frac{3}{2}n\left(\frac{\partial}{\partial t} + \mathbf{V}_{\alpha} \cdot \nabla\right)T_{\alpha} = -nT_{\alpha} \nabla \cdot \mathbf{V}_{\alpha} - \nabla \cdot \mathbf{q}_{\alpha} + Q_{\alpha}$$

temperature evolution

$$\frac{\partial \mathbf{B}}{\partial t} = -\nabla \times \left[\eta \mathbf{J} - \mathbf{V} \times \mathbf{B} + \frac{1}{ne} \mathbf{J} \times \mathbf{B} - \frac{T_e}{ne} \nabla n + \frac{m_e}{ne^2} \frac{\partial}{\partial t} \mathbf{J} \right]$$

Faraday's / Ohm's law

$$\mu_0 \mathbf{J} = \nabla \times \mathbf{B}$$

low- ω Ampere's law

$$\nabla \cdot \mathbf{B} = 0$$

divergence constraint

- Closures for stress ($\underline{\Pi}$) and heat flux (\mathbf{q}) often use relations for collisional plasma for practical modeling.

$\underline{\Pi}$ for collisional plasma is a combination of $\underline{\Pi}_{gv}$, $\underline{\Pi}_{\parallel}$, and $\underline{\Pi}_{\perp}$.

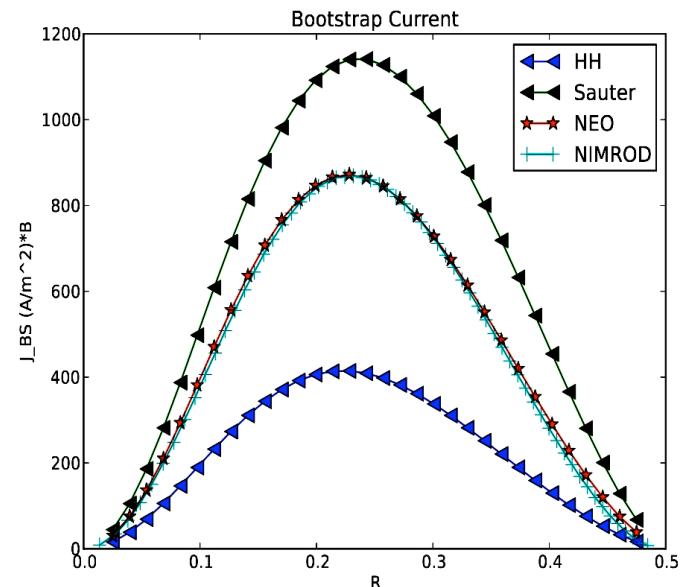
$$\underline{\Pi}_{gv} = \frac{m_i p_i}{4eB} \left[\hat{\mathbf{b}} \times \underline{\mathbf{W}} \cdot (\underline{\mathbf{I}} + 3\hat{\mathbf{b}}\hat{\mathbf{b}}) - (\underline{\mathbf{I}} + 3\hat{\mathbf{b}}\hat{\mathbf{b}}) \cdot \underline{\mathbf{W}} \times \hat{\mathbf{b}} \right], \quad \left(\underline{\mathbf{W}} \equiv \nabla \mathbf{V} + \nabla \mathbf{V}^T - \frac{2}{3} \underline{\mathbf{I}} \nabla \cdot \mathbf{V} \right)$$

$$\underline{\Pi}_{\parallel} = \frac{p_i \tau_i}{2} (\hat{\mathbf{b}} \cdot \underline{\mathbf{W}} \cdot \hat{\mathbf{b}}) (\underline{\mathbf{I}} - 3\hat{\mathbf{b}}\hat{\mathbf{b}})$$

$$\underline{\Pi}_{\perp} \sim -\frac{3p_i m_i^2}{10e^2 B^2 \tau_i} \underline{\mathbf{W}} \text{ has been treated as } -nm_i v_{iso} \underline{\mathbf{W}} \text{ or } -nm_i v_{kin} \nabla \mathbf{V}$$

Kinetic modeling of energetic ions is available, and majority-species drift kinetics is under development.

- Charlson Kim (WA) implemented and verified PIC-based δf computation for minority energetic particles; drift and full-Lorentz options are available.
- A relatively recent application is Dylan Brennan's (U-Tulsa) study of energetic-ion effects on 2/1 tearing in large tokamaks.
- Eric Held is developing drift-kinetic modeling with a complete linearized Coulomb collision operator for neoclassical effects in 3D evolution.
- Comparison of bootstrap current over a NIMROD 2D spatial domain with results from E. Belli's 1-spatial-D NEO code provides an important verification.
- Full RF/MHD coupling will incorporate DK modeling a la Hegna&Callen, PoP 16, 112501, 2009.



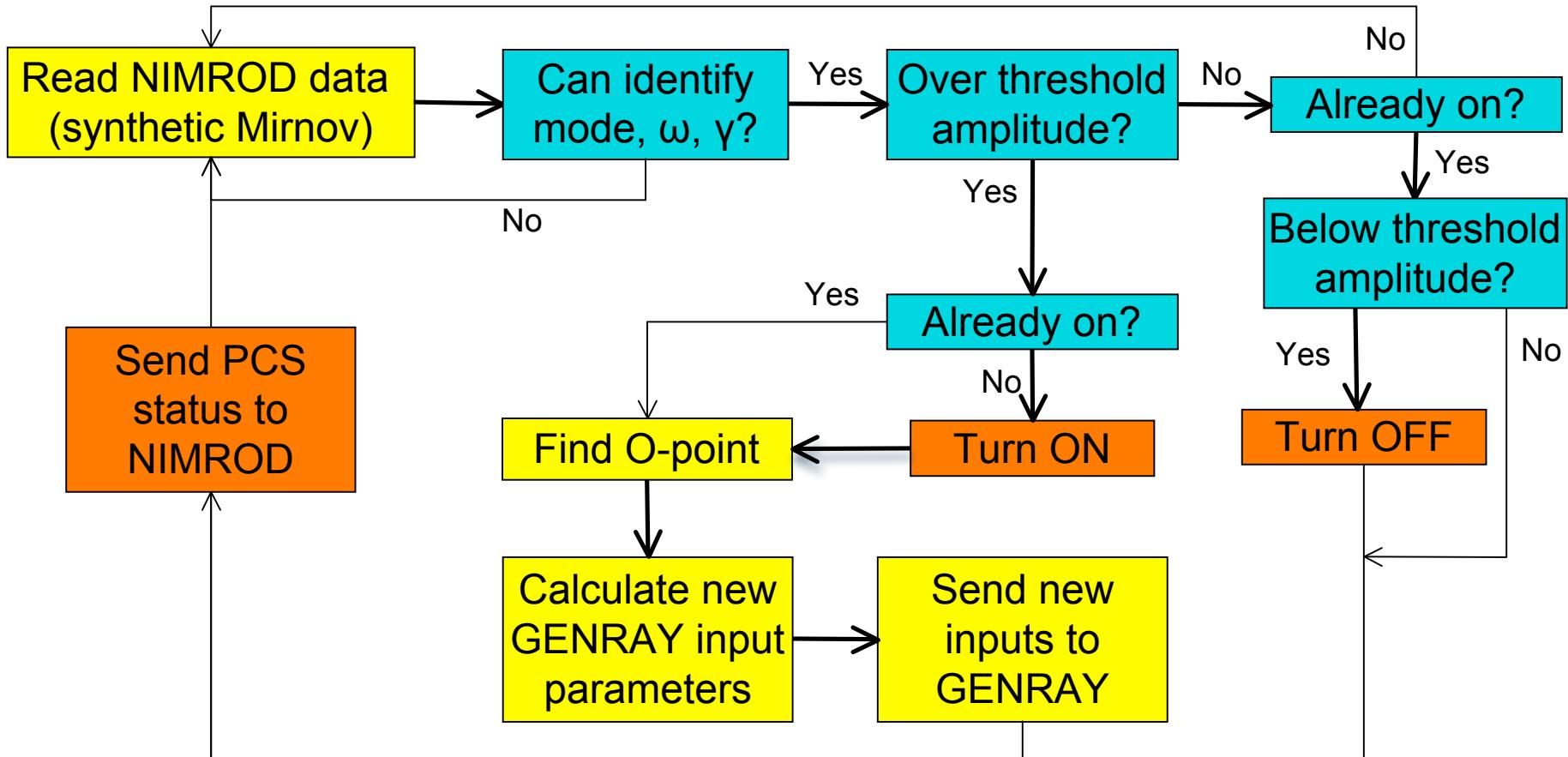
Comparisons of NEO's flux-surface averaged bootstrap current with results from NIMROD's DKE solution.



Concluding Remarks

- Many current NIMROD modeling applications are related to plasma control, either directly or indirectly.
- The applications include boundary conditions or sources that model how external controls influence magnetized plasma over long time-scales.
- Many NIMROD applications are conducted by university faculty, staff, and students.
- To benefit all applications, significant active development of NIMROD continues after more than 15 years.

Rough outline of control system logic



Hierarchy of control systems desired:

“Dumb system”: Mirnov coils only (for detection)

Experimental mimicking system: Implement system similar to DIII-D

Optimized control system: What we want to provide

“Perfect system”: Hit the O-point exactly (what we are doing in this talk)