

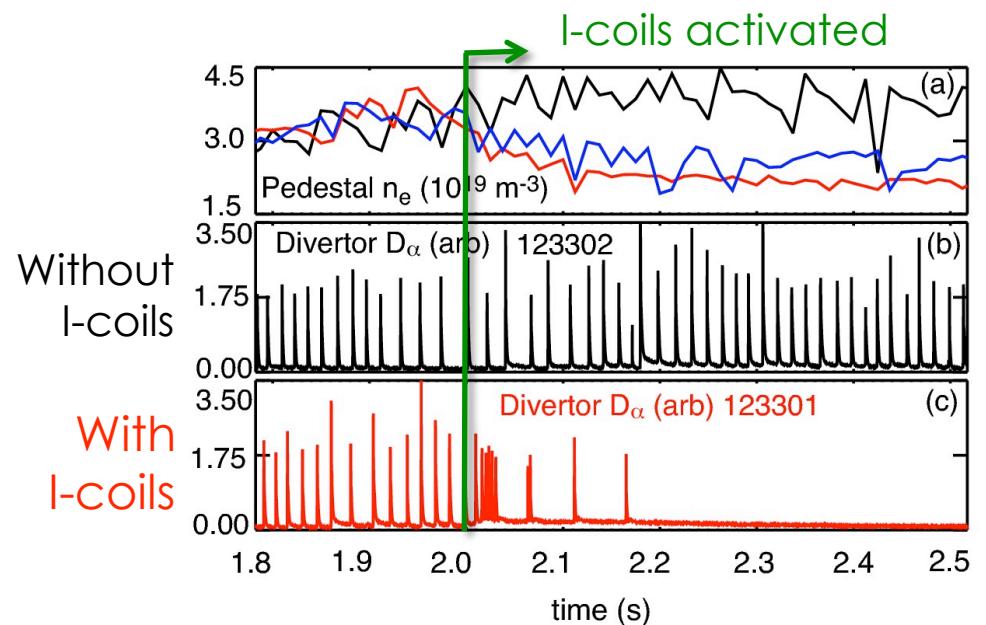
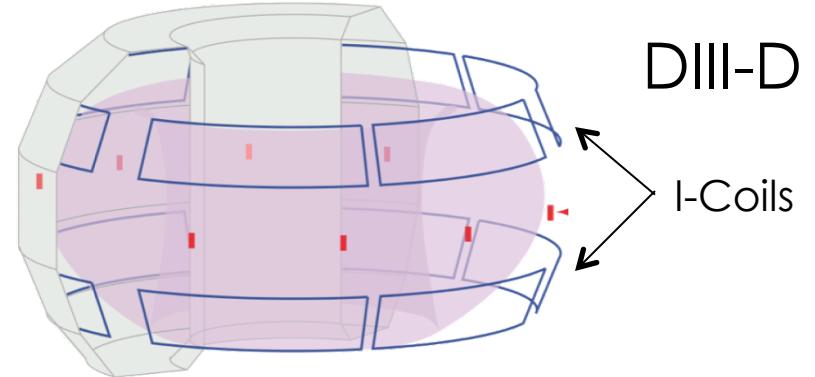
Physics of Tearing Mode Response to 3D Fields

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3D Fields Significantly Affect Tokamak Performance

- **Edge-Localized Modes**
 - Mitigate/suppress ELMs in H-mode
- **Transport**
 - Density pump-out
- **Drive/Brake Rotation**
 - Affects RWM stability
 - Affects tearing mode stability
 - Allow access QH-mode without NBI



Evans, et al. *Phys. Plasmas* **13** (2006)

A Predictive Capability Requires Understanding Plasma Response

- **Predictive capability is challenging because plasma response is complicated**

- Plasma may strongly enhance/suppress spectral components of applied field
- New fields affect transport and rotation
- Rotation strongly affects plasma response

- **New tools are being developed and applied to gain predictive understanding (M3D-C1)**

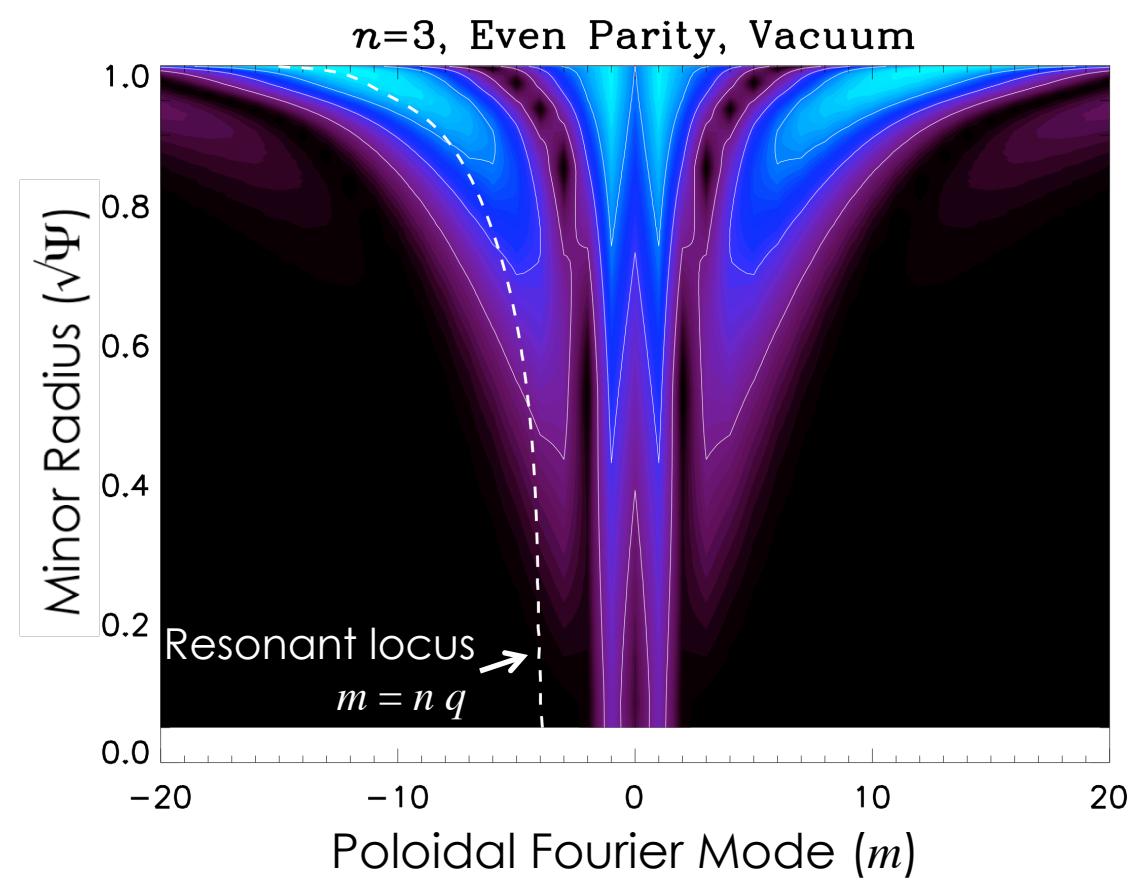
Outline

- **Basics of 3D Response**
- **Introduction to M3D-C1**
- **Linear Results**
 - Influence of rotation
 - Ion rotation vs. Electron rotation
 - When is a linear model appropriate?
- **Implications for Stability Control**

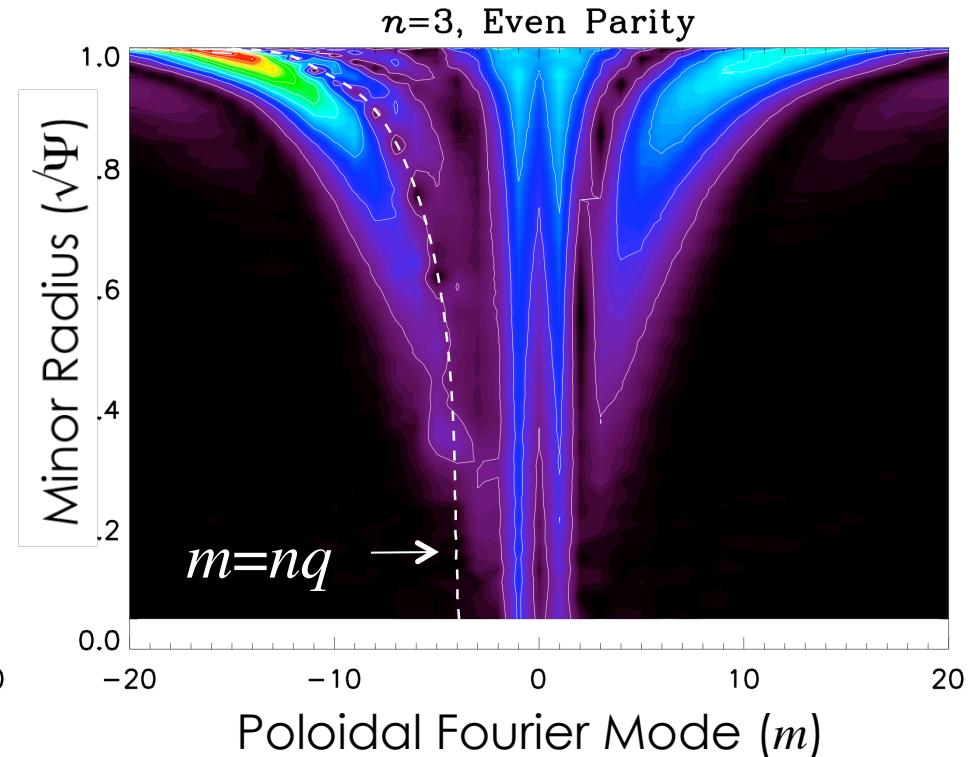
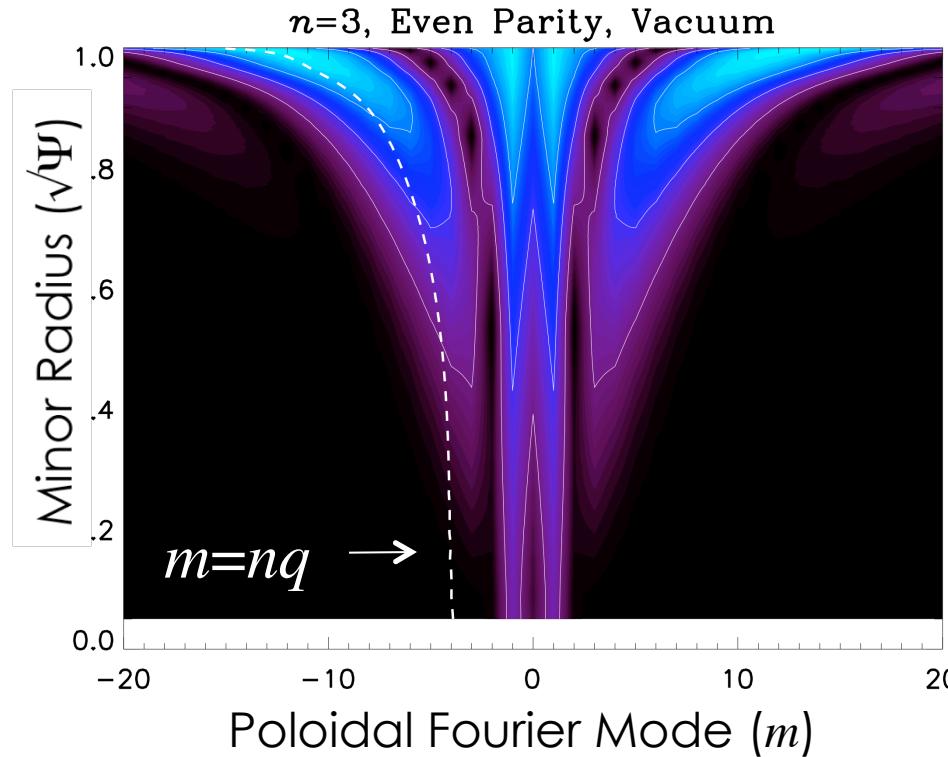
3D Response Basics

Resonant Fields Tear Surfaces; Non-Resonant Fields Bend Surfaces

- Plot shows Fourier spectrum of B_n
- B_n = component of applied field normal to equilibrium magnetic surfaces
- Resonant components (along dashed line) cause islands
- Non-resonant components cause bending of surfaces
- Poloidal spectrum of B_n depends on Ψ



Plasma Response Modifies Spectrum

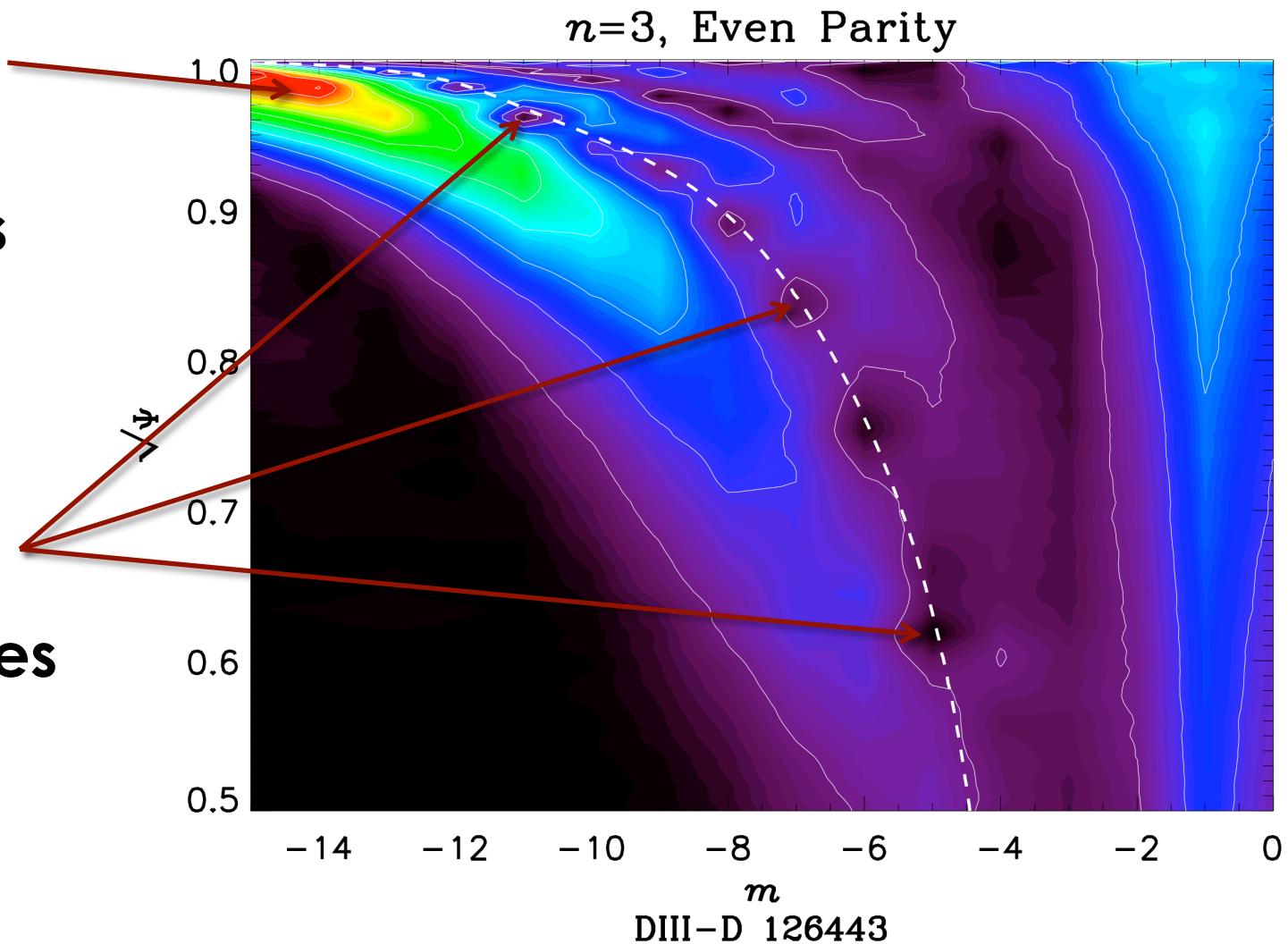


- **Ideal response \rightarrow no islands \rightarrow reduction in resonant components**
- **Excited ideal modes \rightarrow enhancement of non-resonant components**

Plasma Can Kink and Screen

- **“Kinking”**
- **Distorts surfaces**

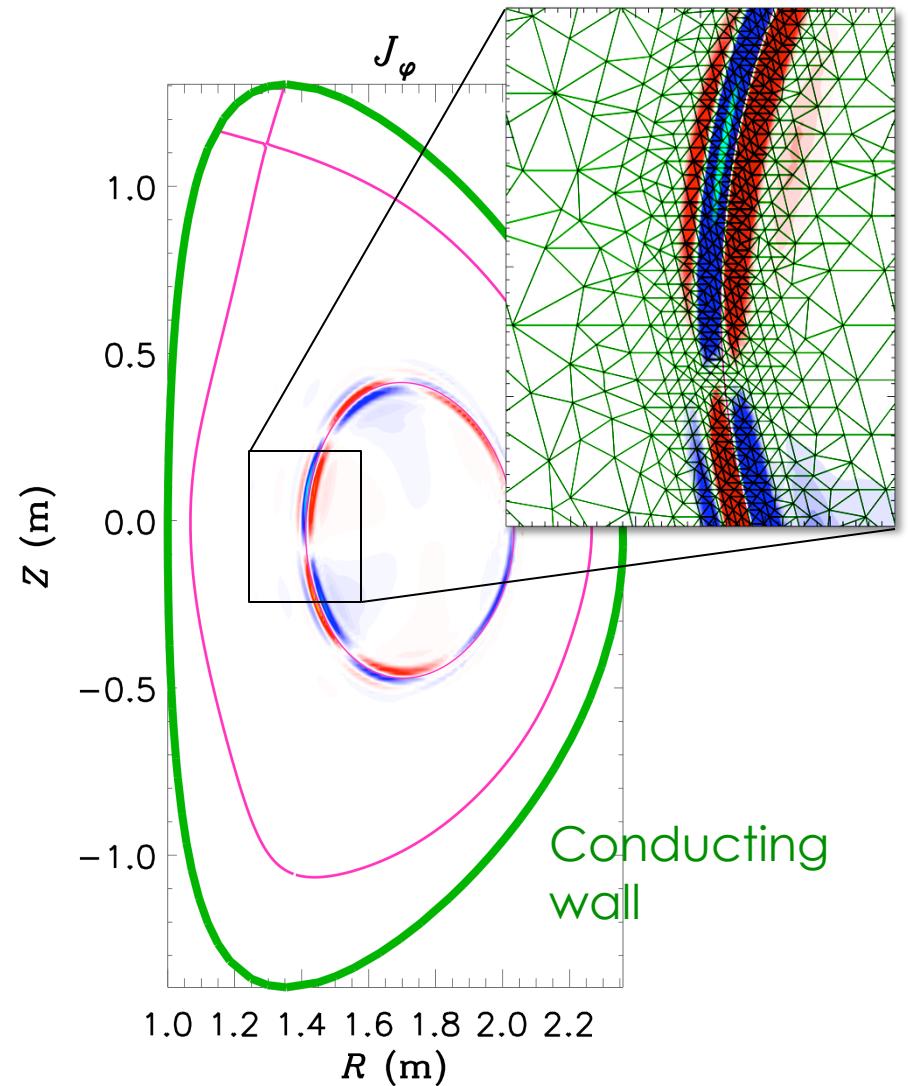
- **Screening**
- **Eliminates islands**



Numerical Methods (M3D-C1)

M3D-C1 Can Calculate Two-Fluid Response

- **M3D-C1 is a two-fluid resistive finite element code**
 - Shares some design principles with M3D
 - (R,Z) coordinates (not spectral in poloidal angle)
- **Computational domain includes plasma, separatrix, and open field-line region**
- **Unstructured mesh allows resolution packing at rational surfaces**
- **Both linear and nonlinear models are implemented**



Two-Fluid Model

$$\frac{\partial n}{\partial t} + \nabla \cdot (n\mathbf{u}) = 0$$

$$n \left(\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} \right) = \mathbf{J} \times \mathbf{B} - \nabla p - \nabla \cdot \Pi$$

$$\frac{\partial p}{\partial t} + \mathbf{u} \cdot \nabla p = -\Gamma p \nabla \cdot \mathbf{u} - \frac{d_i}{n} \mathbf{J} \cdot \left(\Gamma p_e \frac{\nabla n}{n} - \nabla p_e \right) - (\Gamma - 1) \nabla \cdot \mathbf{q}$$

$$\frac{\partial \mathbf{B}}{\partial t} = -\nabla \times \mathbf{E}$$

$$\mathbf{E} = -\mathbf{u} \times \mathbf{B} + \eta \mathbf{J} + \frac{d_i}{n} (\mathbf{J} \times \mathbf{B} - \nabla p_e)$$

$$\Pi = -\mu \left[\nabla \mathbf{u} + (\nabla \mathbf{u})^T \right]$$

$$\mathbf{q} = -\kappa \nabla \left(\frac{p}{n} \right) - \kappa_{\parallel} \mathbf{b} \mathbf{b} \cdot \nabla \left(\frac{p_e}{n} \right)$$

$$\mathbf{J} = \nabla \times \mathbf{B}$$

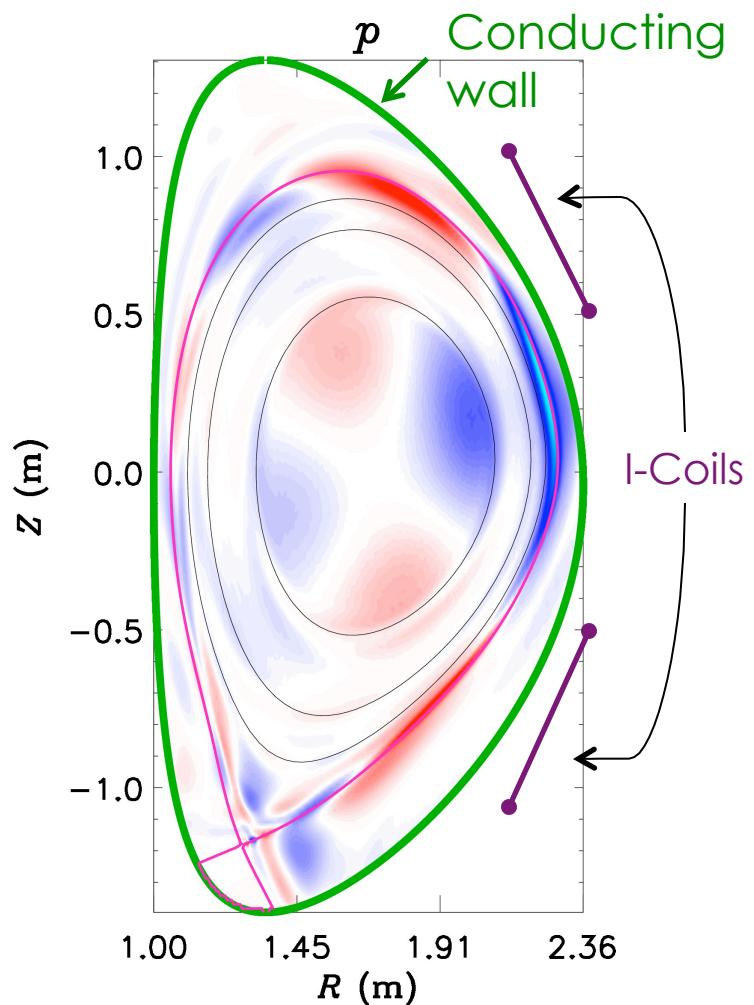
$$\Gamma = 5/3$$

$$p_e = p/2$$

- Complete (not reduced) **two-fluid** model is implemented
- Time-independent equations may be solved directly for linear response

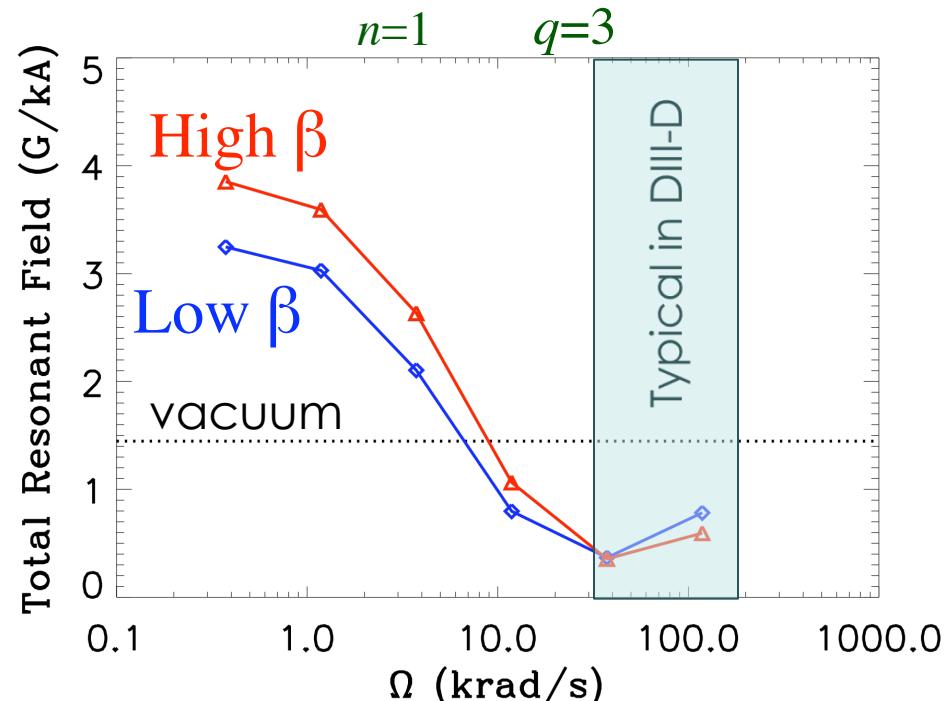
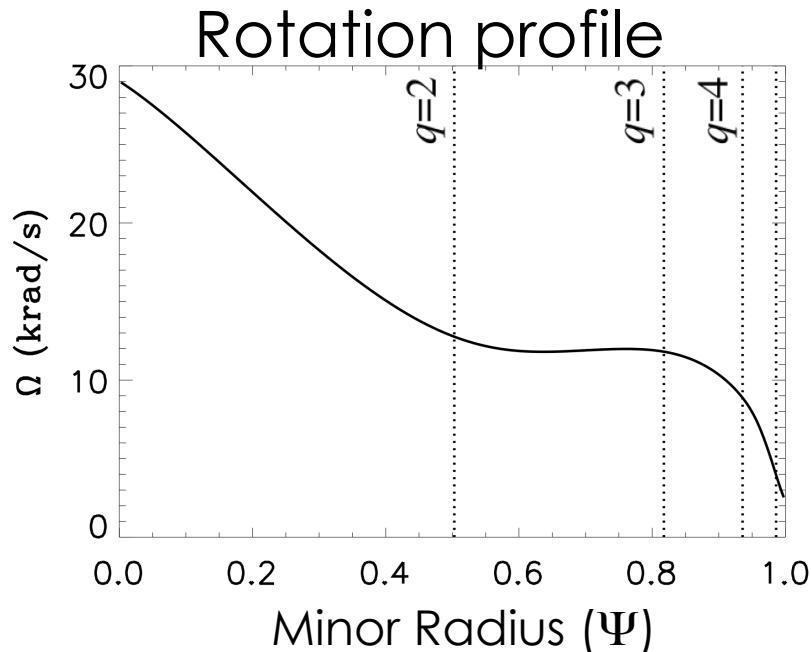
Analysis Considers Reconstructed DIII-D Equilibria

- Vacuum fields generated by DIII-D I-coils
- Boundary conditions:
 - Vacuum B_n is held constant at the boundary
 - No-slip ($v=0$)
- Realistic transport parameters
 - Lundquist number $\sim 10^9$
- Toroidal rotation
 - Rotation is added self-consistently: $p \neq p(\psi)$



Single-Fluid Results

Single-Fluid Result: Rotation (Usually) Improves Screening



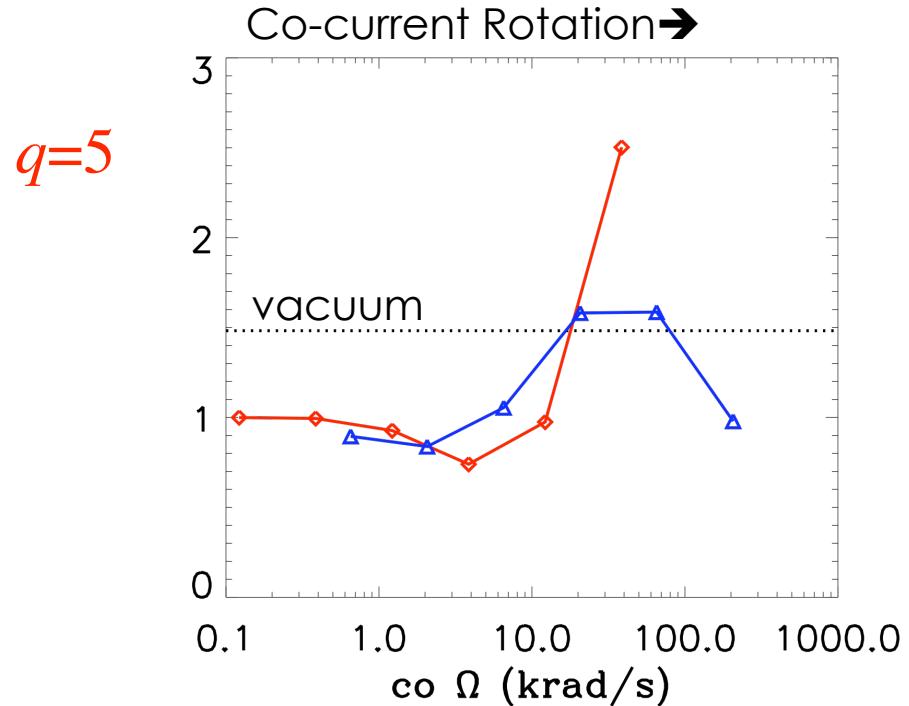
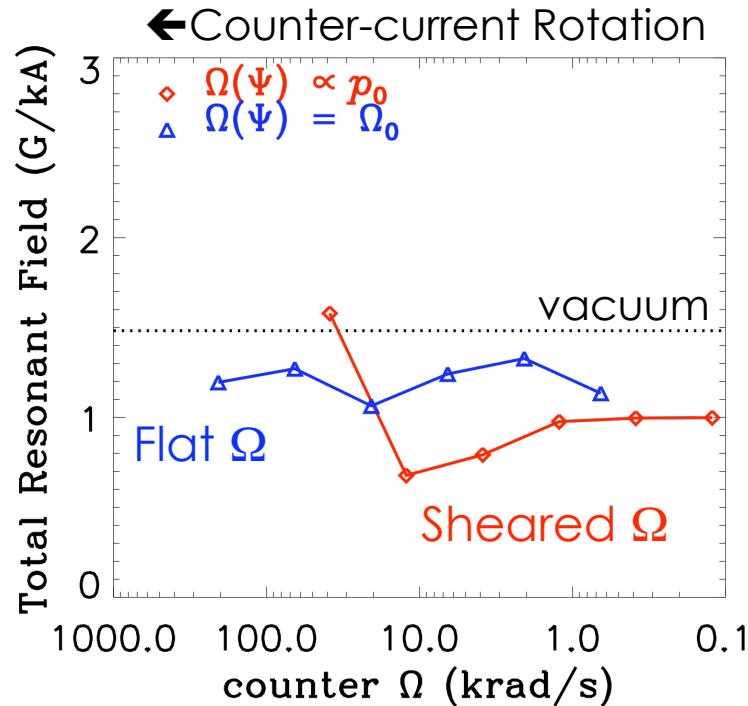
- Plasma may enhance resonant fields at low rotation
- Large rotation screens resonant fields
- Response depends on beta

Why Is Plasma Response Sensitive to Rotation? Why Is It Sensitive To Beta?

- **From a (rotating) plasma's perspective, the static external fields are oscillating**
 - If field is oscillating faster than tearing response, plasma won't tear
- **Rotation drives static tearing modes away from marginal stability**
- **Higher Beta moves modes closer to marginal stability**
 - At marginal stability, an infinitesimal perturbation yields an infinite response

Single-Fluid Result: Rotation Shear Increases Edge Response

DIII-D 135762

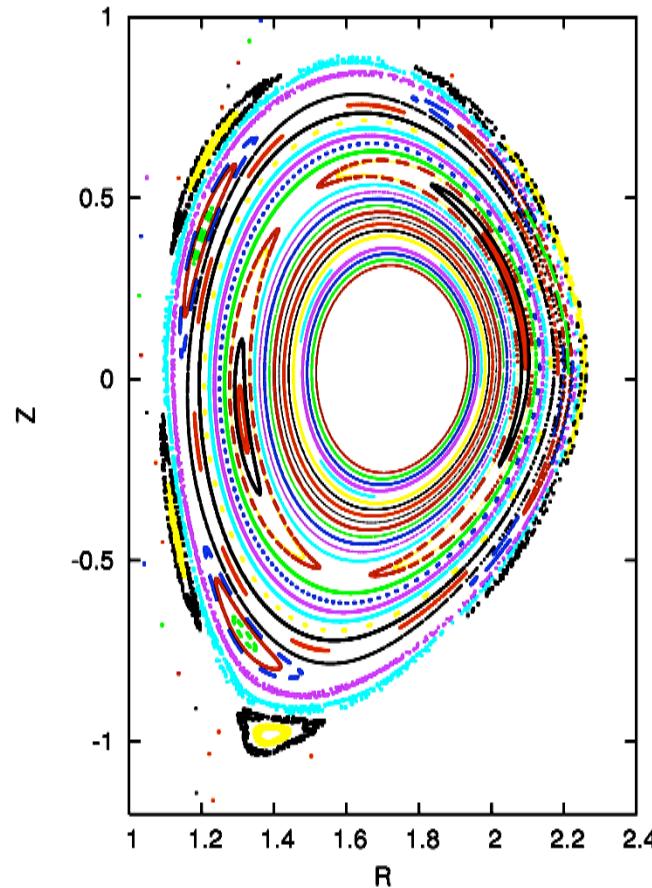


- Large rotation shear seems to increase edge response
- Why? Theory predicts Ω' is destabilizing to low- n peeling-balloonning modes*

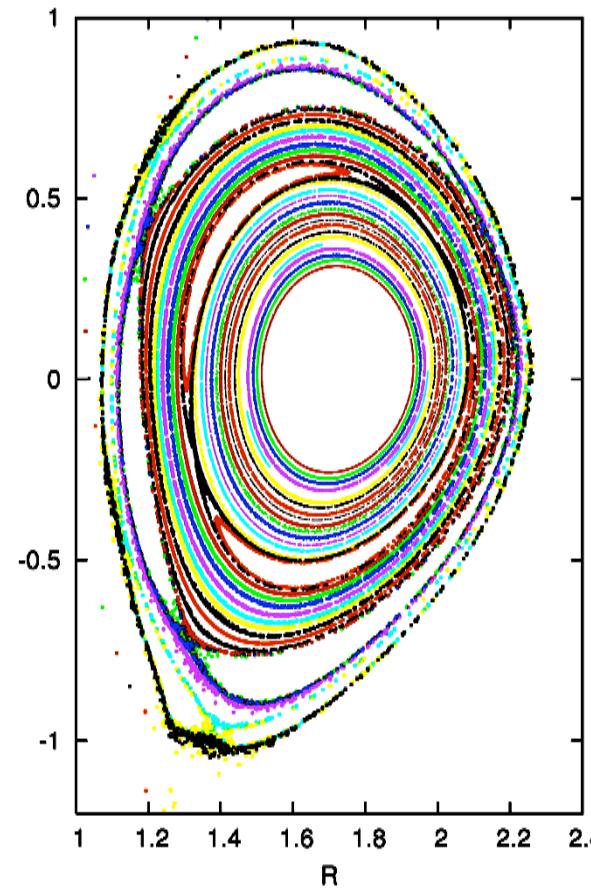
* Snyder, et al. Nucl. Fusion **47** (2007)
 Aiba, et al. Nucl. Fusion **50** (2010)
 Ferraro, et al. Phys. Plasmas **17** (2010)

Rotation Improves Core Screening; But Sheared Rotation Stochasticizes Edge

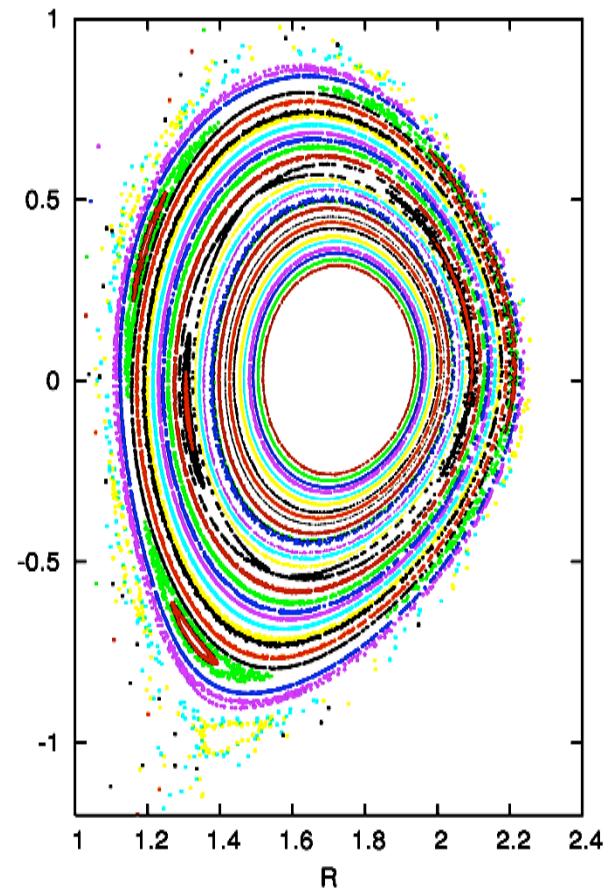
Vacuum



Plasma, Static



Plasma, Rotating



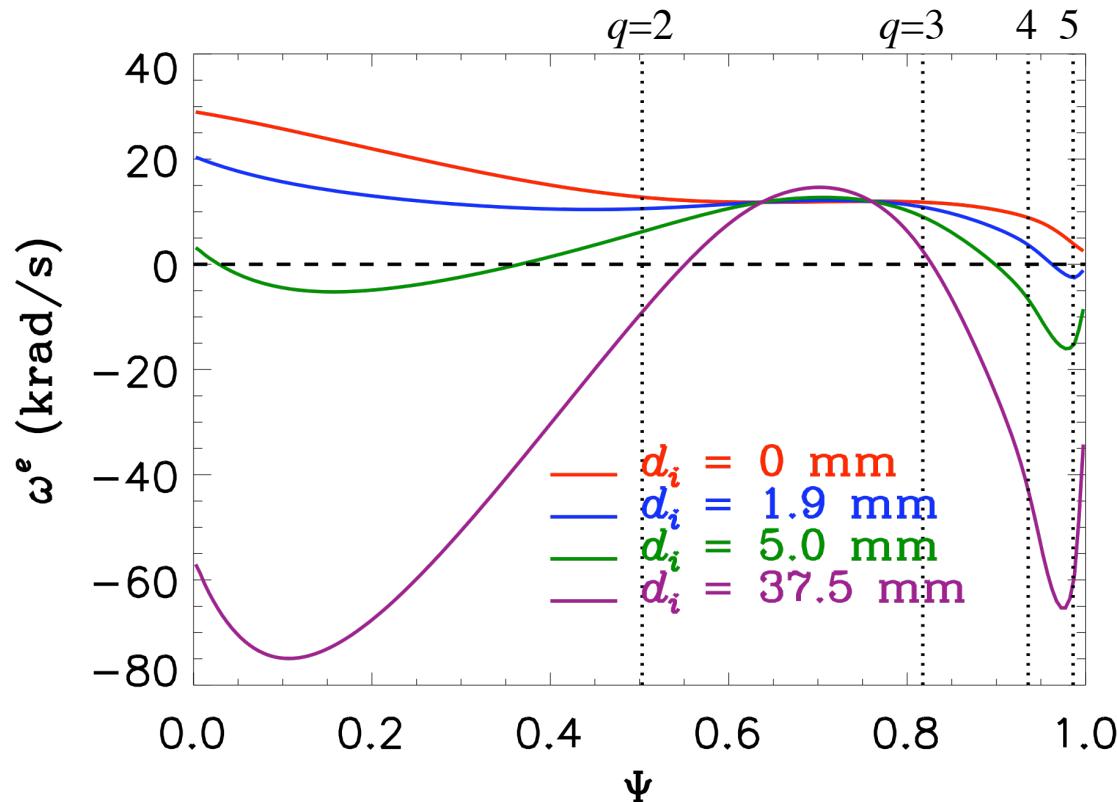
$$\Omega_0 = 0$$

$$\Omega_0 = 300 \text{ krad/s}$$

Two-Fluid Results

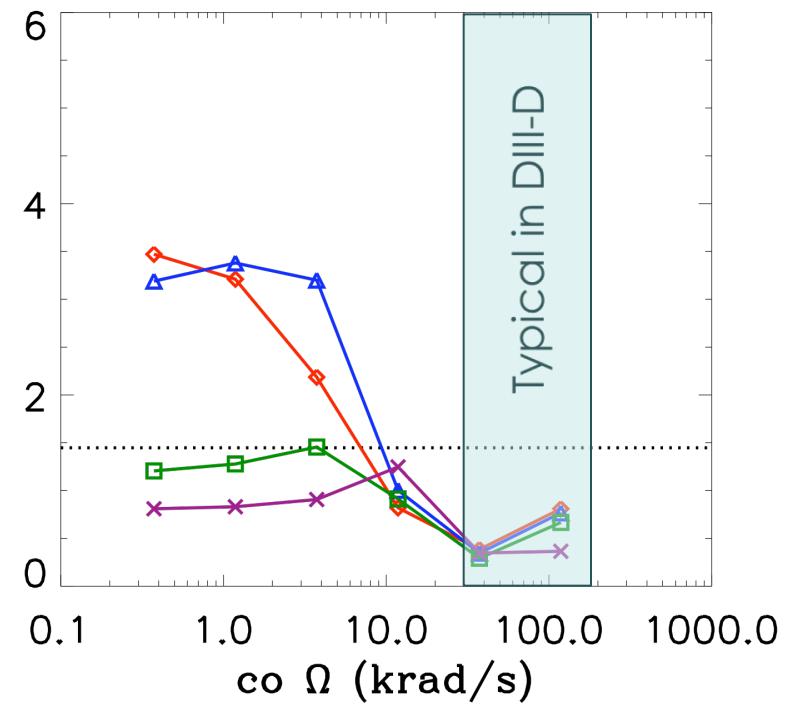
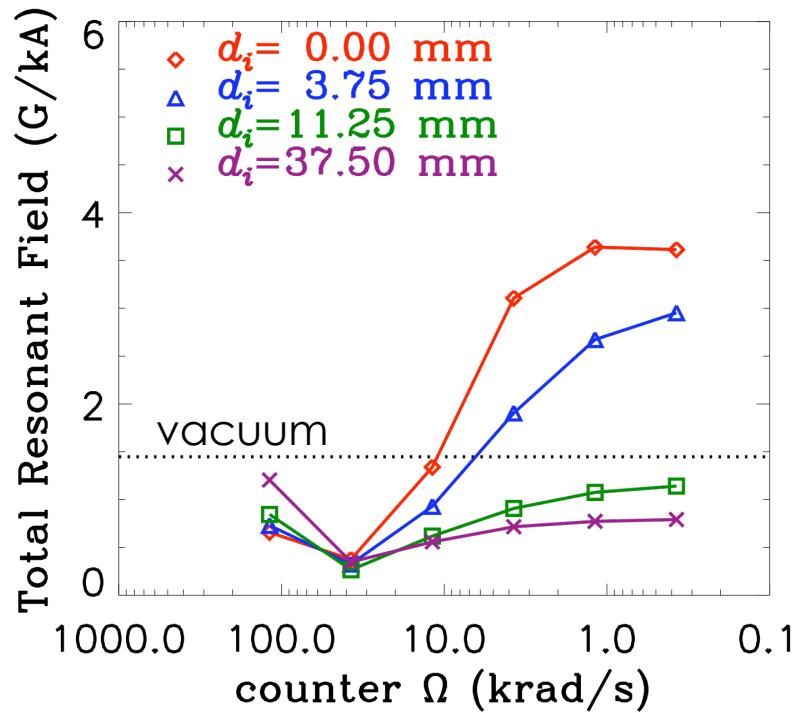
Two-Fluid Results: Ion and Electron Rotations are Distinct

- Given Ω , we can change $\Omega^e = \Omega + \omega_*$ by adjusting $\omega_* = d_i p' / n$



Two-Fluid Effects Shift Resonance

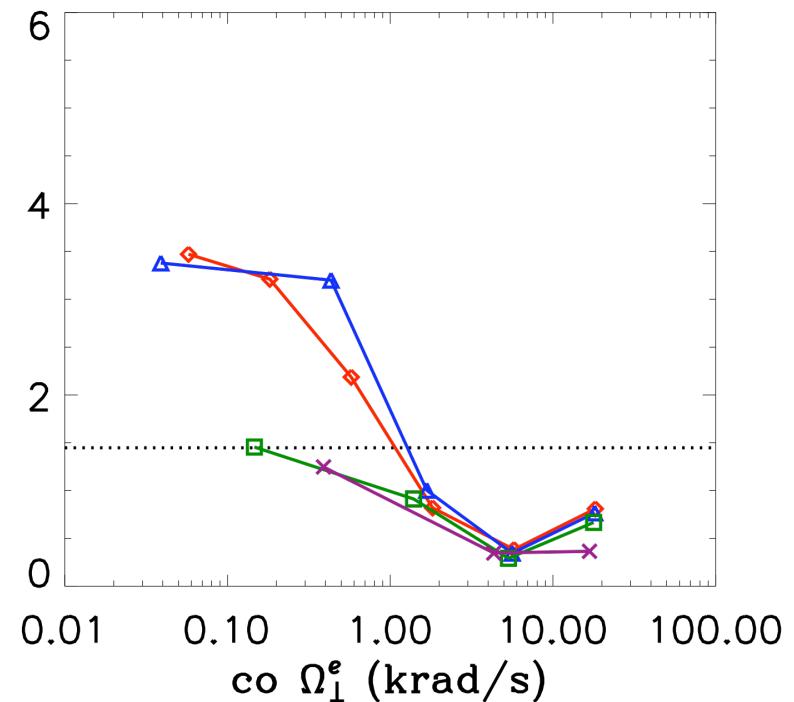
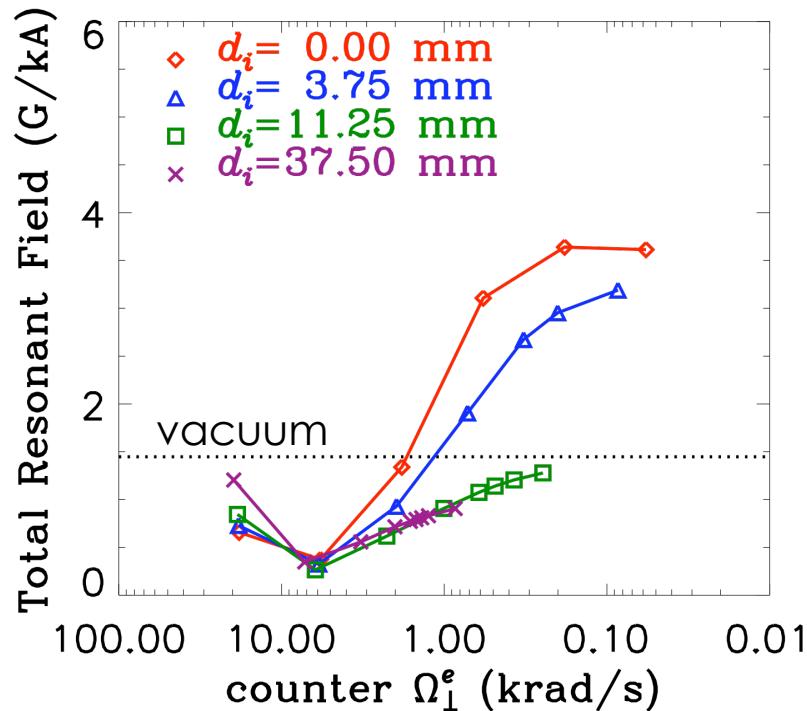
(Mass) rotation at $q=3$



- **Strongest tearing no longer occurs at $\Omega = 0$**

Penetration In Core Depends on Electron Rotation

Perpendicular electron rotation at $q = 3$



- Screening of $q=3$ island clearly depends more on Ω^e than Ω

What is “Perpendicular” Electron Velocity?

- The perpendicular angular velocity is defined as

$$\Omega_{\perp}^{e,i} = \frac{\mathbf{v}^{e,i}}{R} \cdot \frac{\mathbf{B} \times \nabla \psi}{|\mathbf{B} \times \nabla \psi|}$$

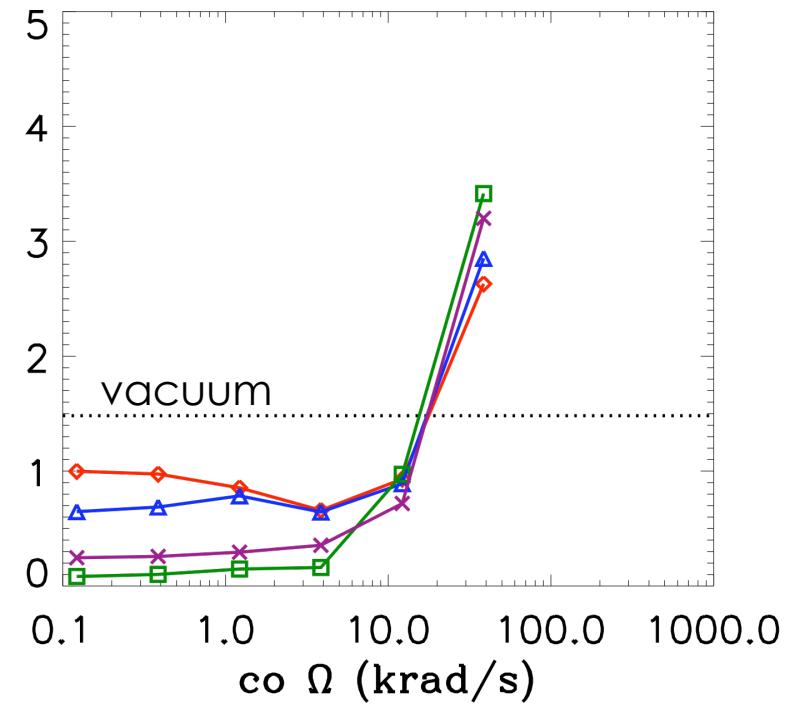
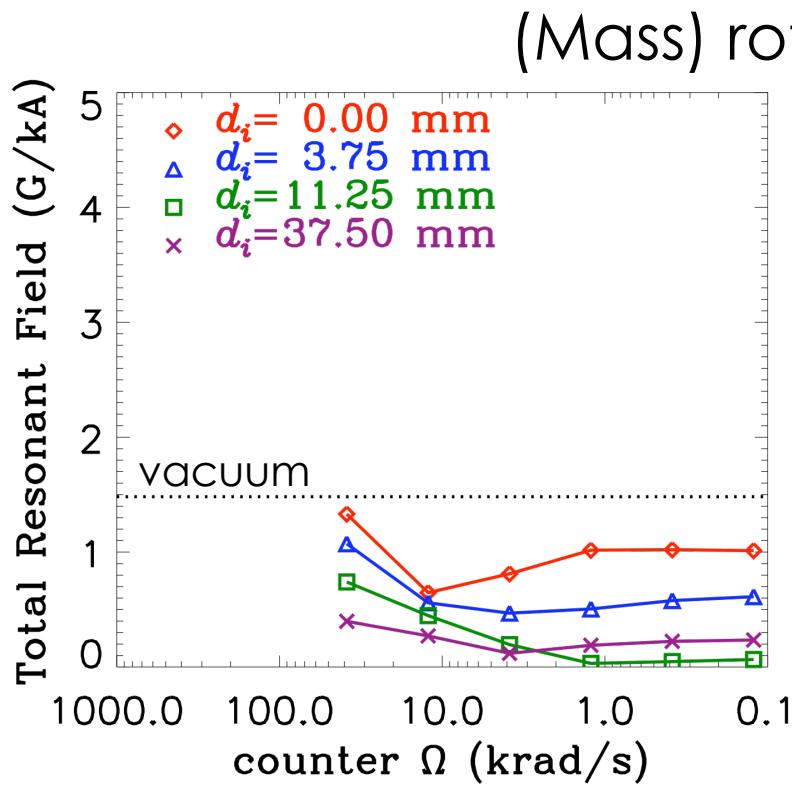
- To lowest order, $\mathbf{v}^{e,i} = R^2 \omega^{e,i}(\psi) \nabla \varphi + \lambda^{e,i} \mathbf{B}$. Thus:

$$\Omega_{\perp}^{e,i} = \frac{|\nabla \psi \times \nabla \varphi|}{|B|} \omega^{e,i}(\psi)$$

- From radial force balance: $\omega^{e,i}(\psi) = \phi'(\psi) + \frac{p_{e,i}'(\psi)}{n_{e,i} q_{e,i}}$

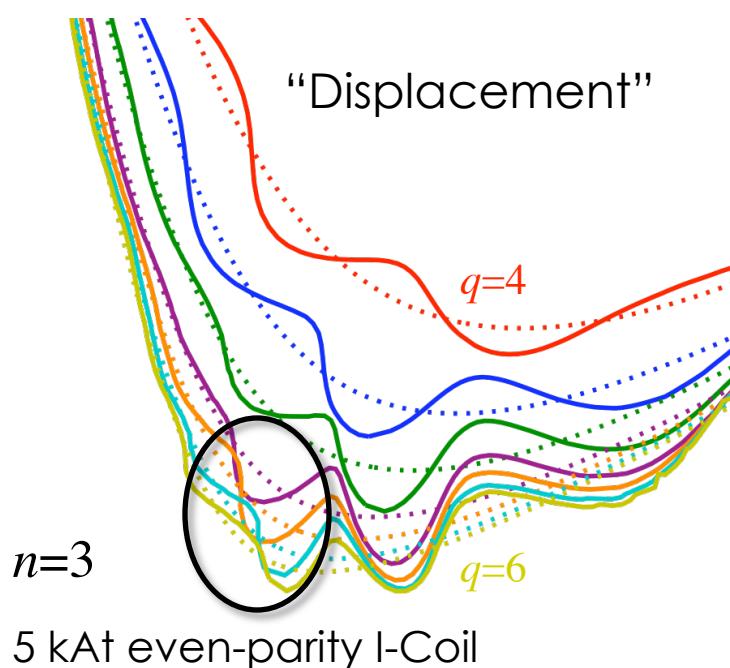
Edge Response Depends on Mass Rotation Shear

- Tearing of edge modes is dependent on ion, not electron, rotation shear



When is Linear Response Appropriate?

- For typical experimental parameters, linear response may not be strictly valid in some regions
 - Large current density near rational surfaces
 - Back-reaction on rotation is important



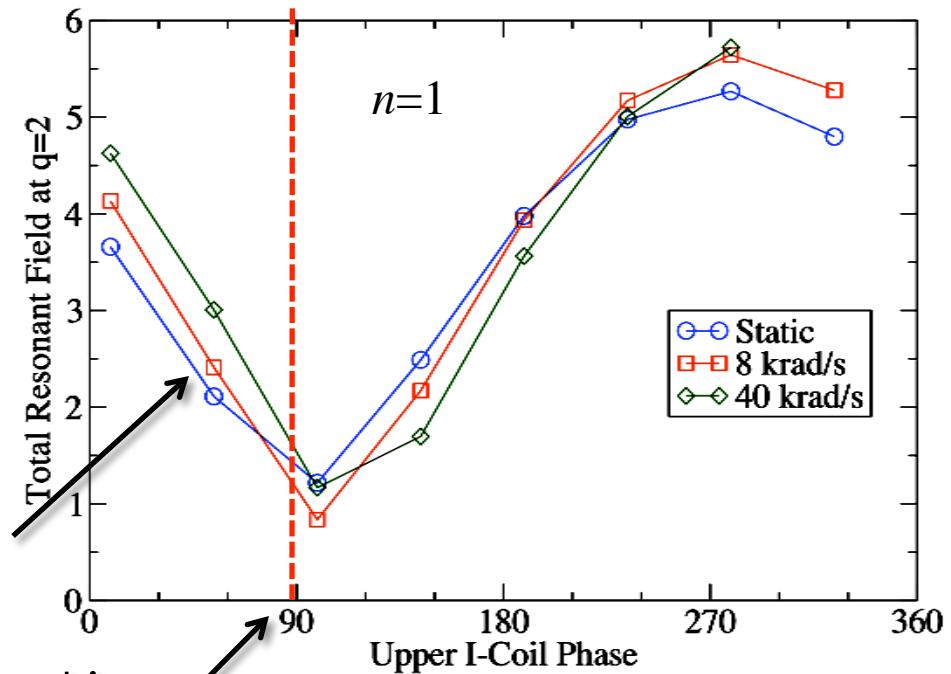
- “Displacement” shows overlapping surfaces near separatrix!
- Quantitative predictions of island size, stochasticity from linear calculations are suspect

Linear Response Gets Some Things Right

- Which modes are most sensitive
- How parameters (rotation, viscosity, etc.) affect sensitivity
- How to optimize coil design

Calculated resonant field
(proxy for resonant torque)

Empirical phase least prone to locking



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Summary of Theory/Modeling Results

- We can now quantify the dependence of plasma response on parameters in experimentally relevant regimes
- Tearing response is closely correlated with rotation
- Core tearing is correlated with electron rotation
- Edge response is correlated with ion (mass) rotation shear

Confidence in Control Methods Require Model Validation

- **Linear modeling is probably sufficient for coil optimization**
- **Nonlinear (or QL) modeling is necessary for other things**
 - Changes to $n=0$ rotation/pressure profiles
 - Locking bifurcation threshold
 - Effect of 3D equilibrium on ELM stability
- **Modeling to test/inform ELM suppression hypotheses is underway**

Implications For ELM Suppression Control

- Empirical q_{95} resonances strongly imply that **current profile control** will be necessary
- It is not yet clear whether tearing is a necessary condition for ELM stabilization
 - If it is, then modeling implies **rotation control** might be necessary
 - ECCD (electron rotation control) can cause/heal tearing
 - Feedback on island size is not feasible
- If transport (in pedestal or at pedestal top) is necessary, are there other ways of driving it?