

Magnetic feedback optimization by including the dynamic response of the wall in RFX-mod and DIII-D

Lidia Piron, Consorzio RFX

in collaboration with L. Marrelli, P. Martin, P. Piovesan, A. Soppelsa for RFX-mod and J. Hanson, Y. In, M. Okabayashi, E.J. Strait, H. Reimerdes for DIII-D

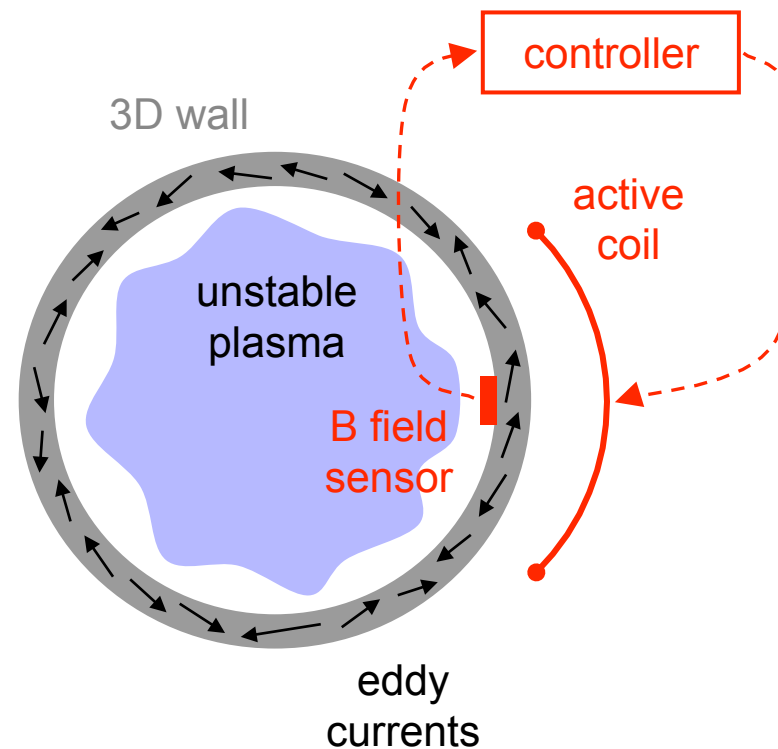
15TH WORKSHOP ON MHD STABILITY CONTROL

Madison, WI, USA, November 15th-17th, 2010



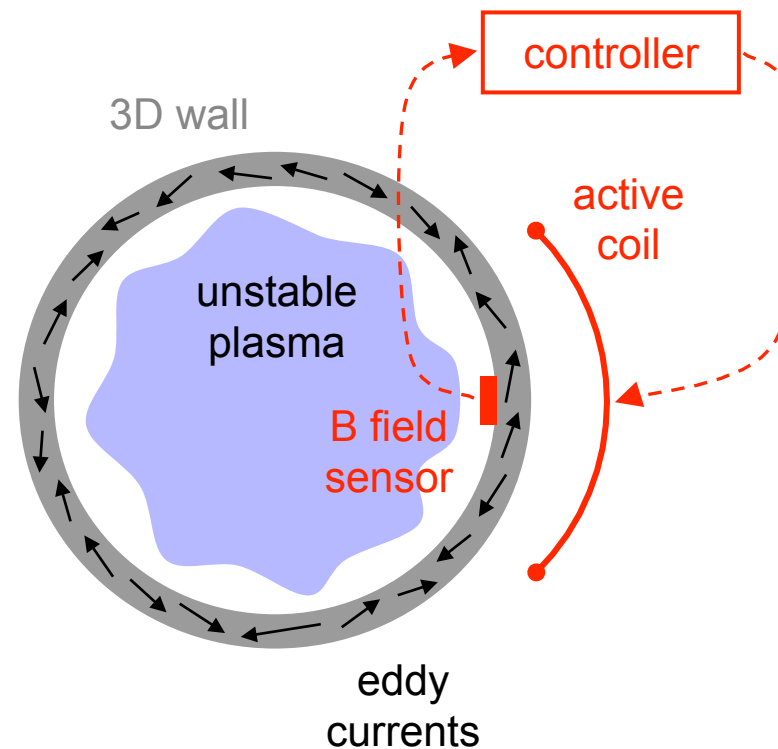
Background and motivation

- Any realistic wall has complicated **3D conducting structures** that modify its **dynamic response (eddy currents)** to any time-varying magnetic field, produced by the plasma or by any external source, such as feedback and axisymmetric coils



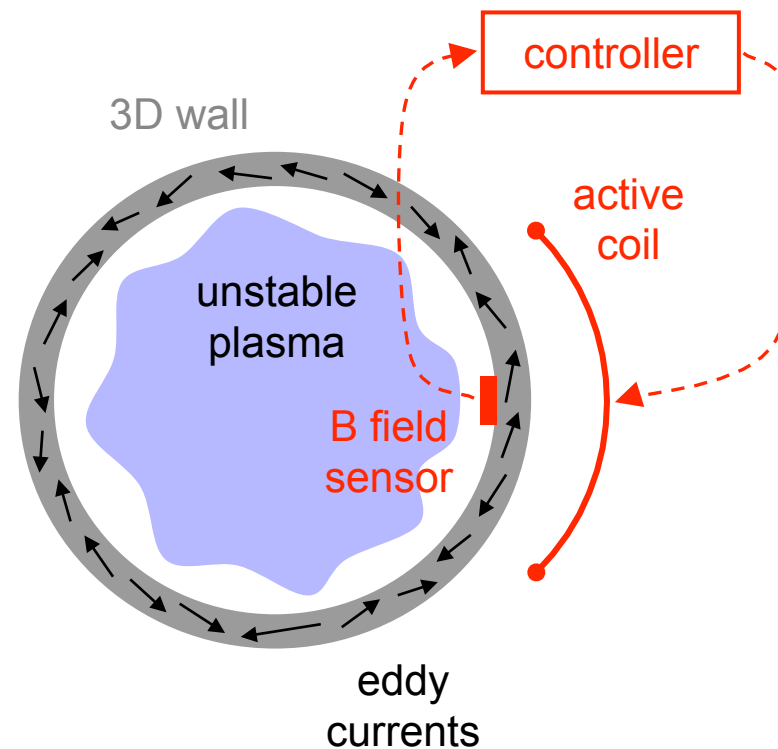
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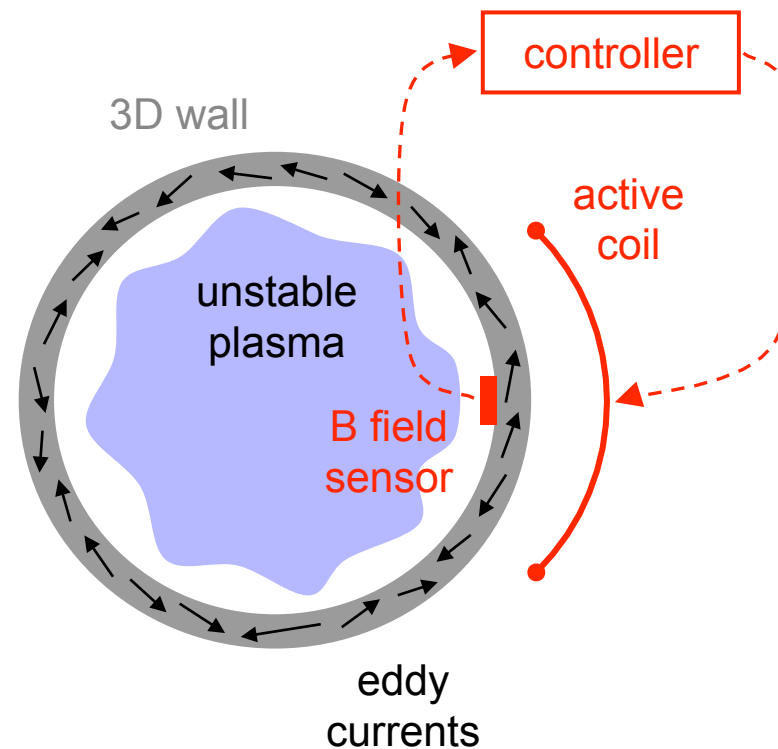
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 - the magnetic field produced by the active coils inside the wall can be modified by the wall response → *error fields may be thus produced*
 - the sensors measure contributions from the plasma and from all external magnetic field sources, but also from the time-dependent eddy currents induced in the wall → *possible need of AC compensation schemes*



- Two different examples of how the **frequency-dependent response of a 3D wall** can affect magnetic feedback in a real experiment will be presented, along with **strategies to compensate for such effects** by including in the feedback algorithm information on the wall response:

RFX-mod

- control of a helical RFP equilibria by imposing helical boundary conditions

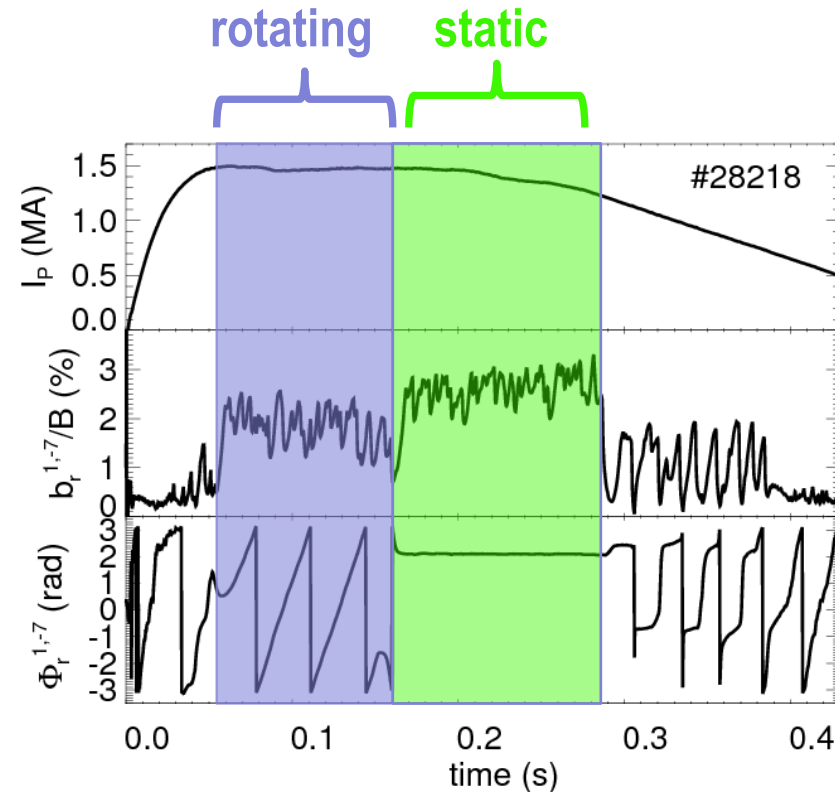
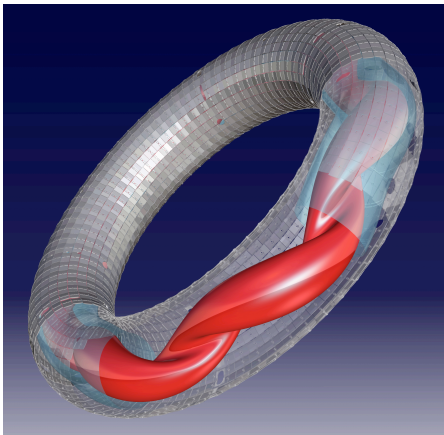
DIII-D

- DC vs AC compensation of sensor signals for $n=1$ RWM suppression both in Ohmic and high- β discharges

RFX-mod

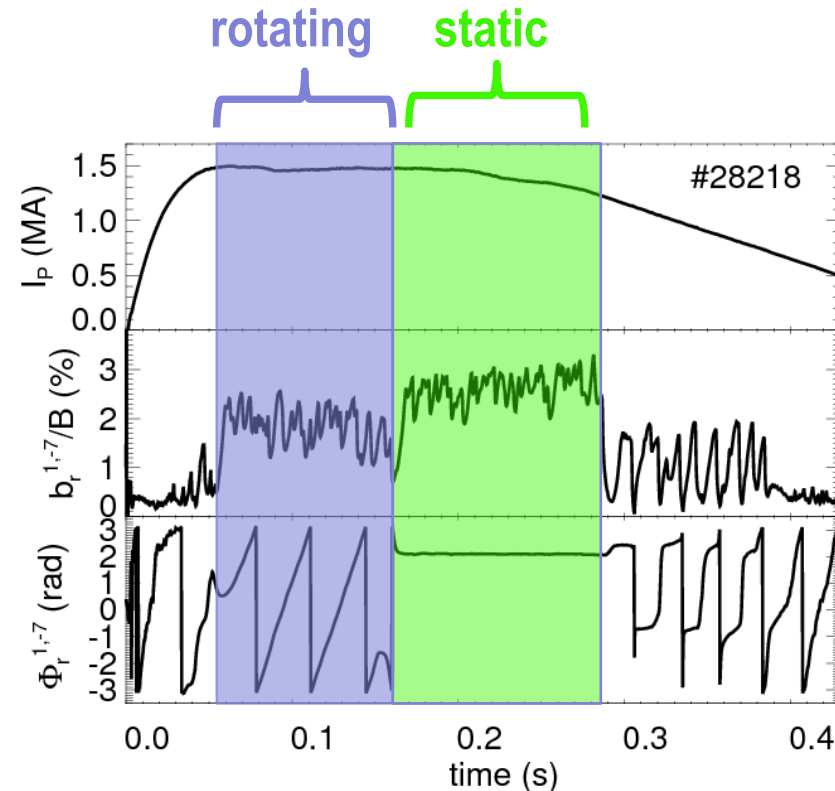
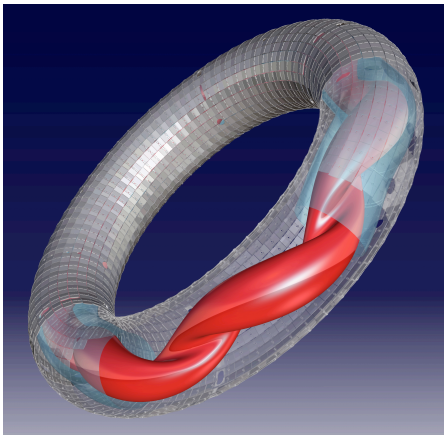
Control of helical equilibria in RFX-mod

- Helical RFP equilibria can be sustained at high I_p by imposing $m=1$, $n=-7$ helical boundary conditions through magnetic feedback → *more in P. Piovesan's talk*



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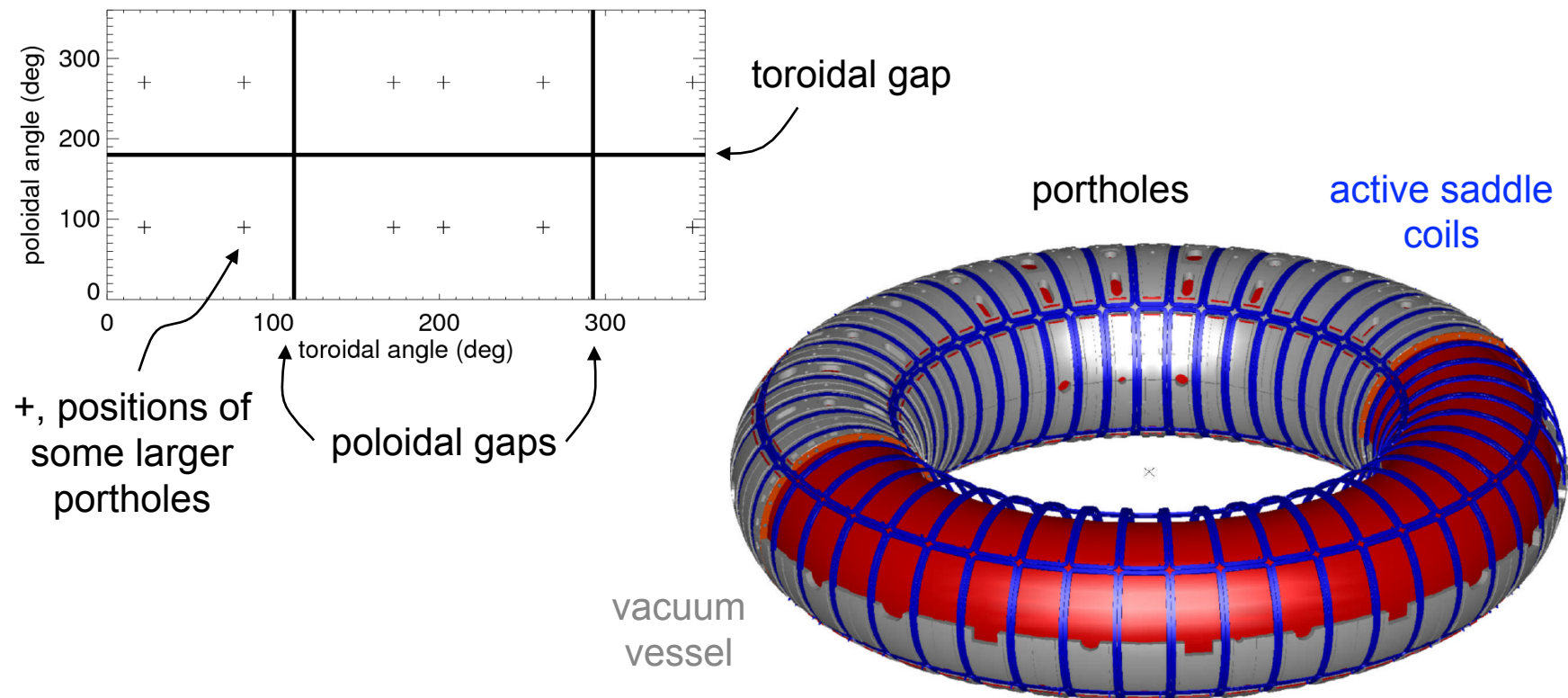
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- Error fields due to 3D wall structure may excite secondary instabilities and affect this improved confinement regime

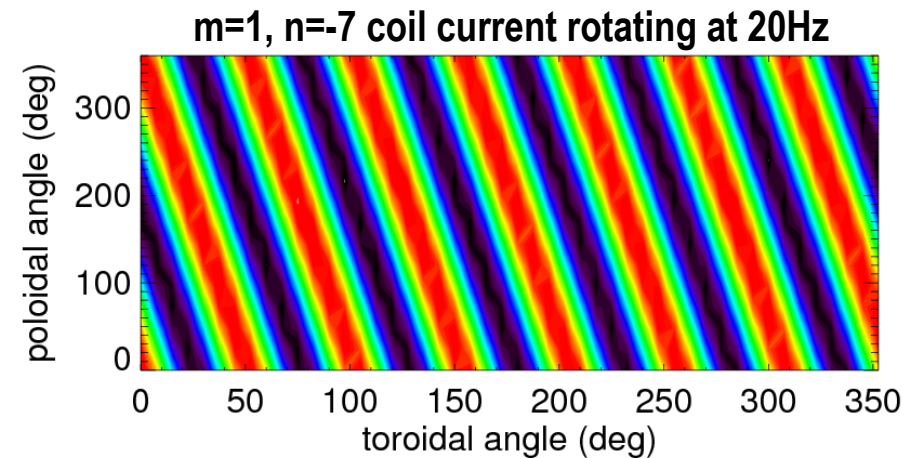
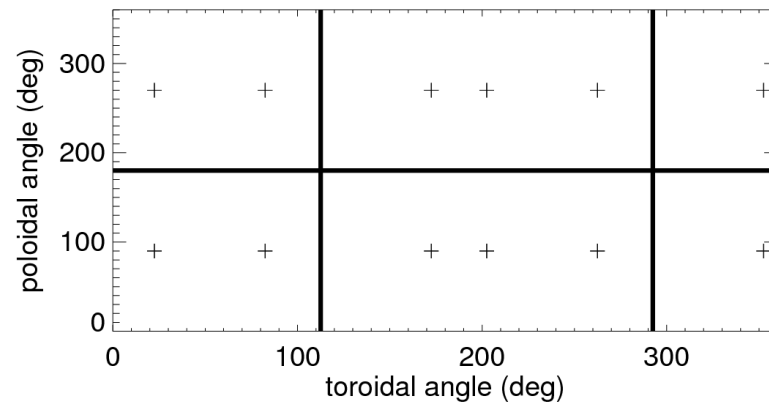
RFX-mod wall layout & control coils scheme

- The RFX-mod wall has **two poloidal gaps** and one **toroidal gap**, plus some **large portholes** that significantly affect its dynamic response to external magnetic fields
- $48 \times 4 = 192$ active coils (external) and respective sensors (internal) fully cover the torus



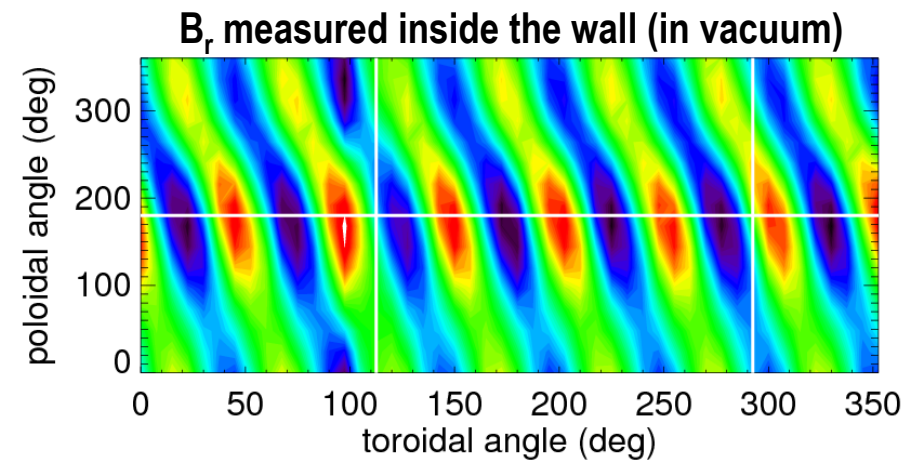
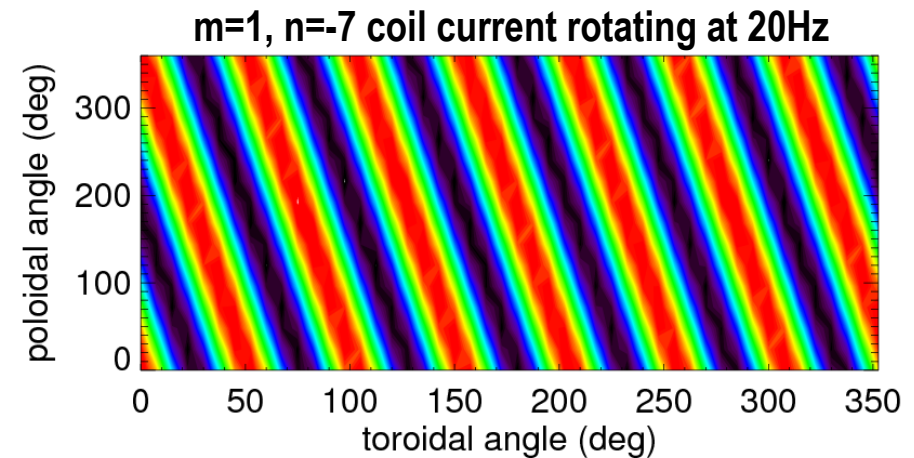
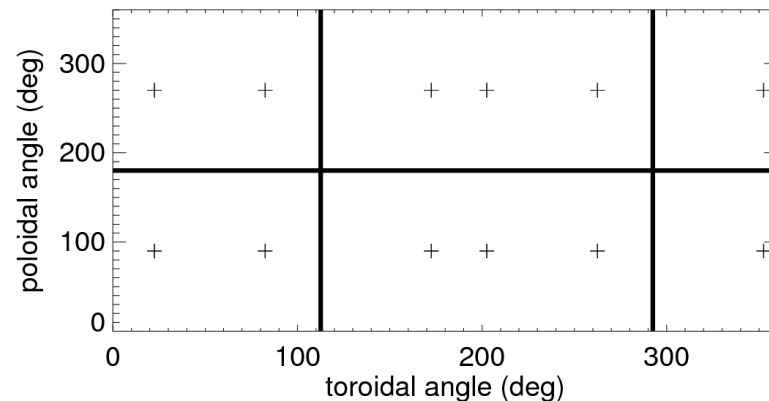
Error fields due to the equatorial gap

- The radial magnetic field produced inside the wall by a rotating $m=1, n=-7$ current in the active coils is affected mostly by the presence of the toroidal gap



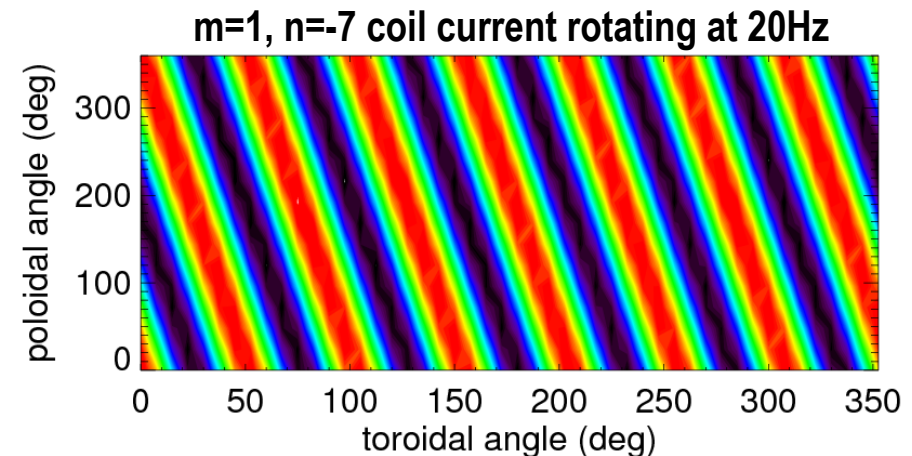
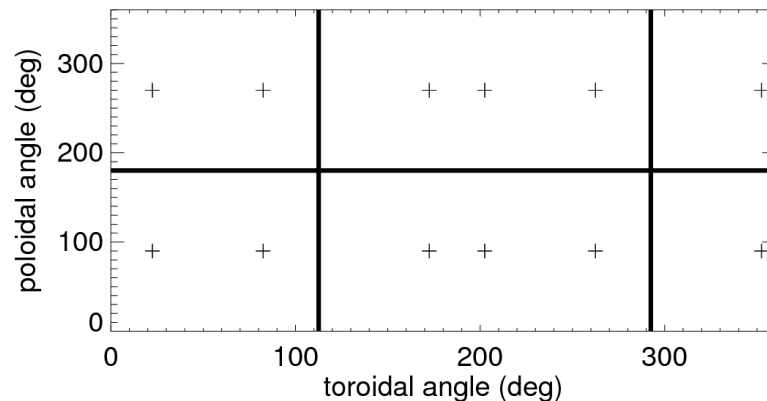
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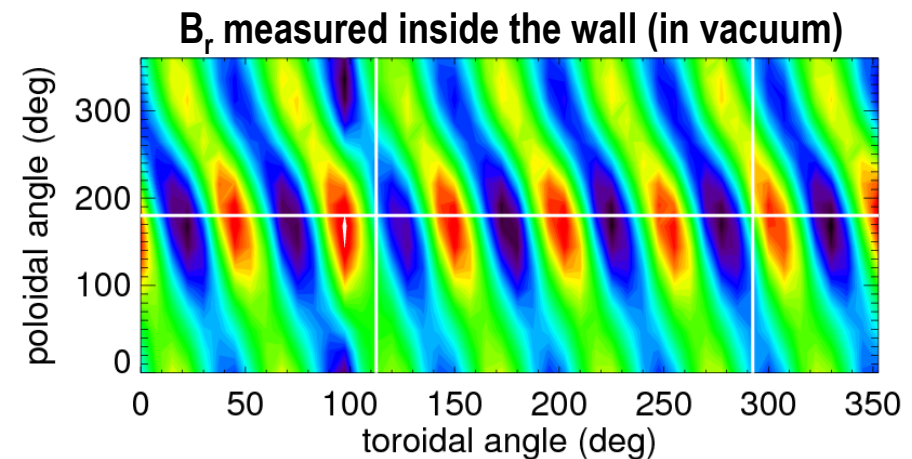


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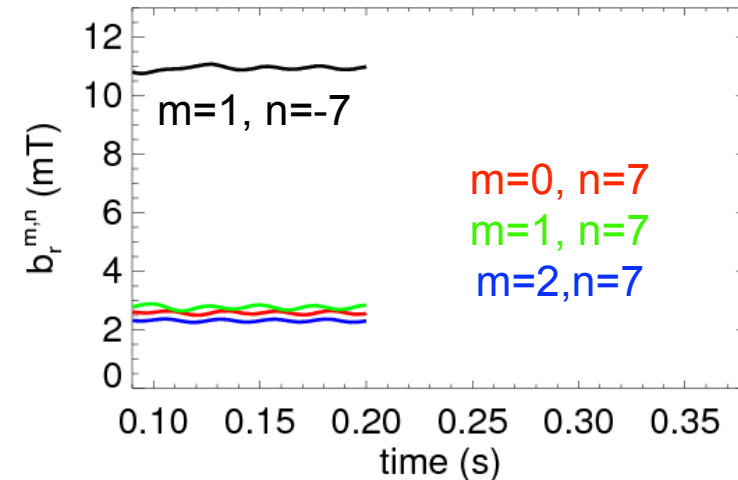
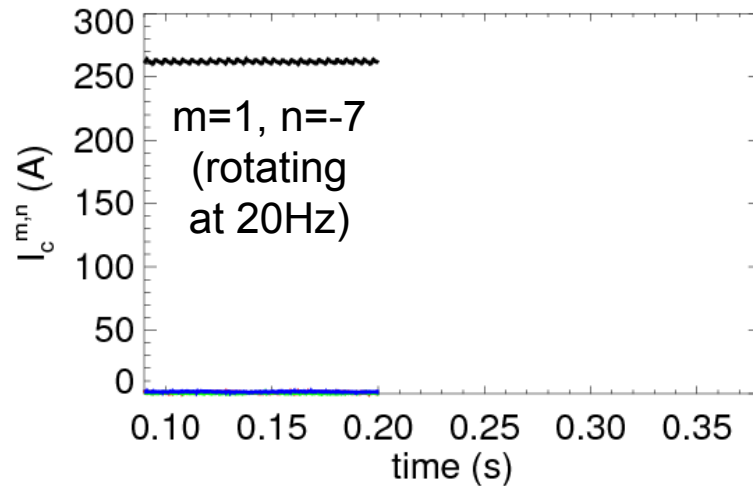


- As a result, the radial magnetic field as a dominant $m=1, n=-7$ harmonic, plus $m=0,1,2, n=-7$ harmonics to represent the radial field perturbation localized at the toroidal gap



Coupling of different n=7 poloidal harmonics

- The toroidal gap introduces a coupling between different poloidal harmonics that can be represented by **complex transfer functions**, which have been measured in vacuum

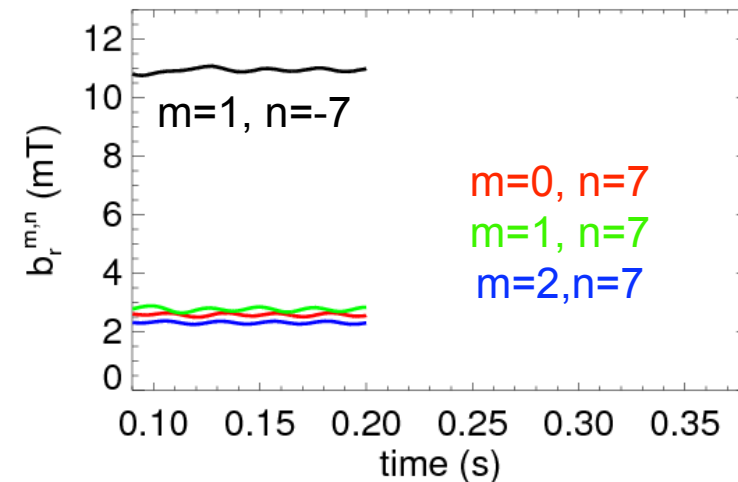
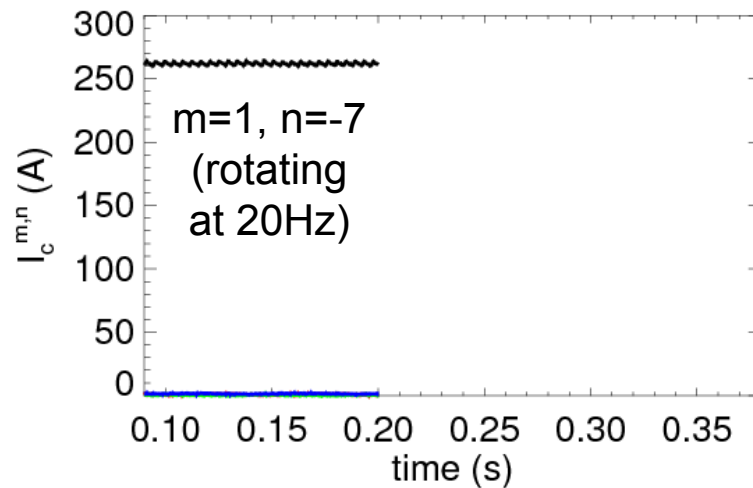


Transfer function matrix

$$\begin{bmatrix} M_{1,-7}^{1,-7} & M_{0,7}^{1,-7} & M_{1,7}^{1,-7} & M_{2,7}^{1,-7} \\ M_{1,-7}^{0,7} & M_{0,7}^{0,7} & M_{1,7}^{0,7} & M_{2,7}^{0,7} \\ M_{1,-7}^{1,7} & M_{0,7}^{1,7} & M_{1,7}^{1,7} & M_{2,7}^{1,7} \\ M_{1,-7}^{2,7} & M_{0,7}^{2,7} & M_{1,7}^{2,7} & M_{2,7}^{2,7} \end{bmatrix} \cdot \begin{bmatrix} I_c^{m=1,n=-7} \\ 0 \\ 0 \\ 0 \end{bmatrix} = \begin{bmatrix} b_r^{m=1,n=-7} \\ b_r^{m=0,n=7} \\ b_r^{m=1,n=7} \\ b_r^{m=2,n=7} \end{bmatrix}$$

Dynamic decoupler for the n=7 harmonic

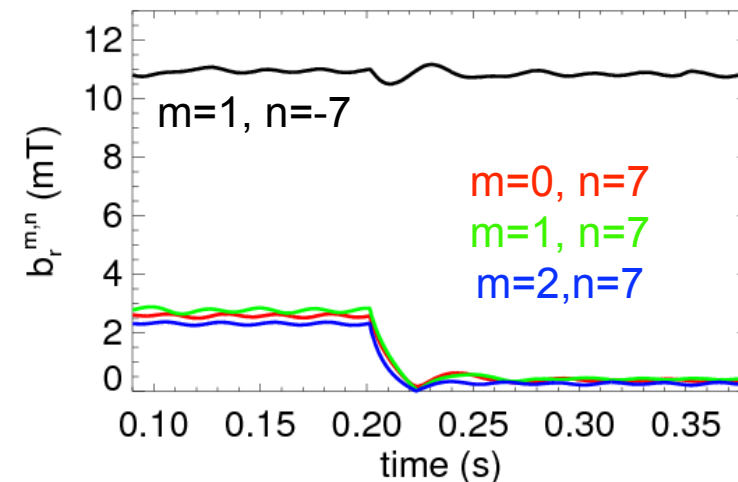
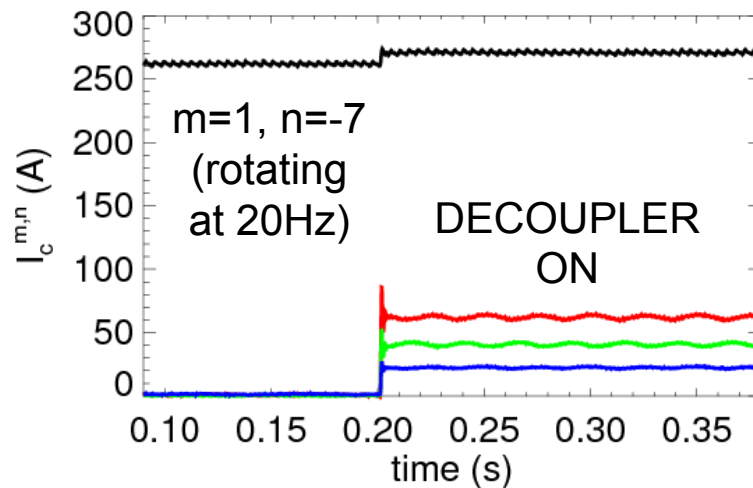
- A dynamic decoupler for the n=7 mode was built, by **inverting the transfer function matrix**, to compute the current distribution that generates a pure 1/-7 B_r inside the wall



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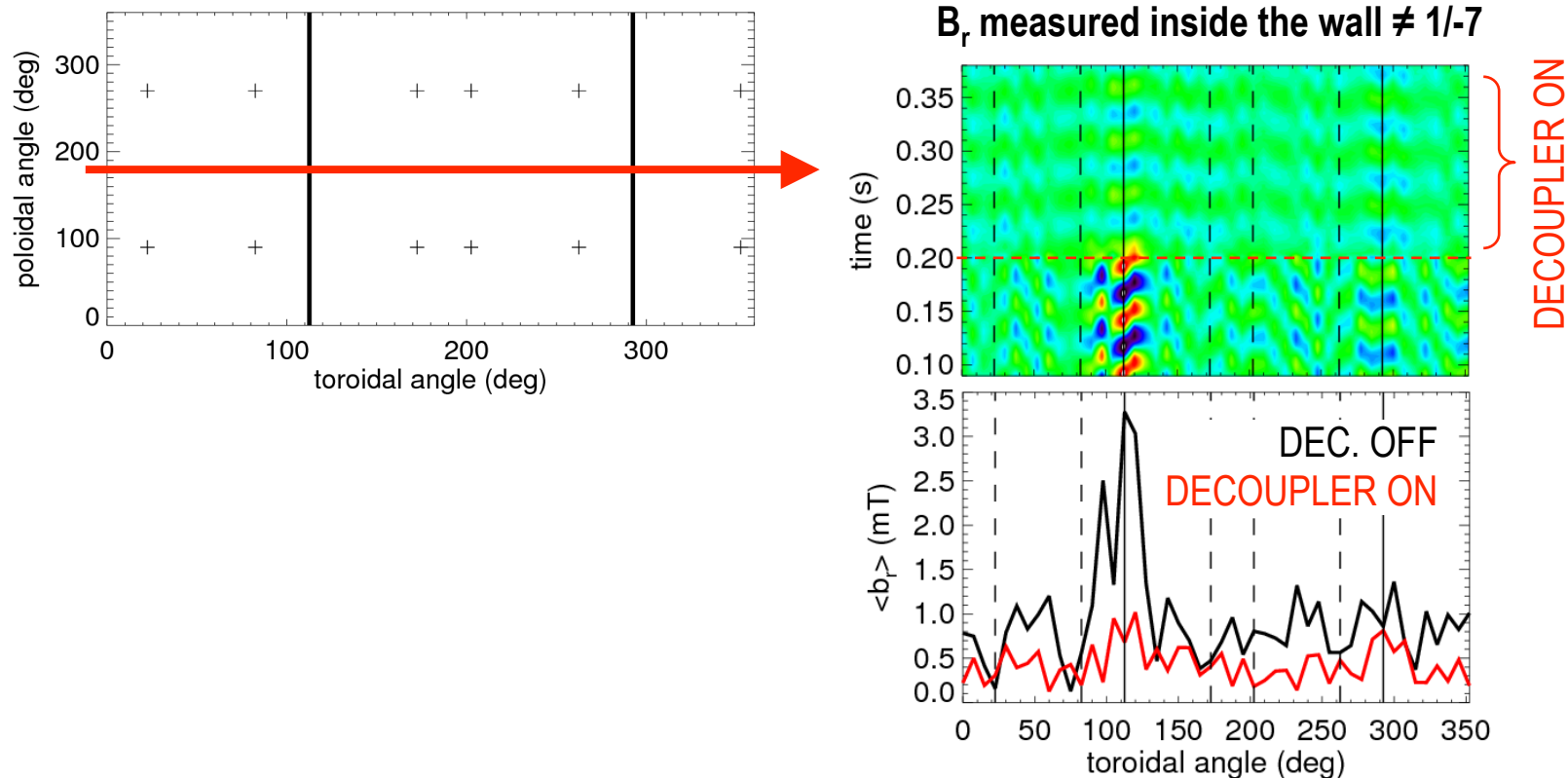
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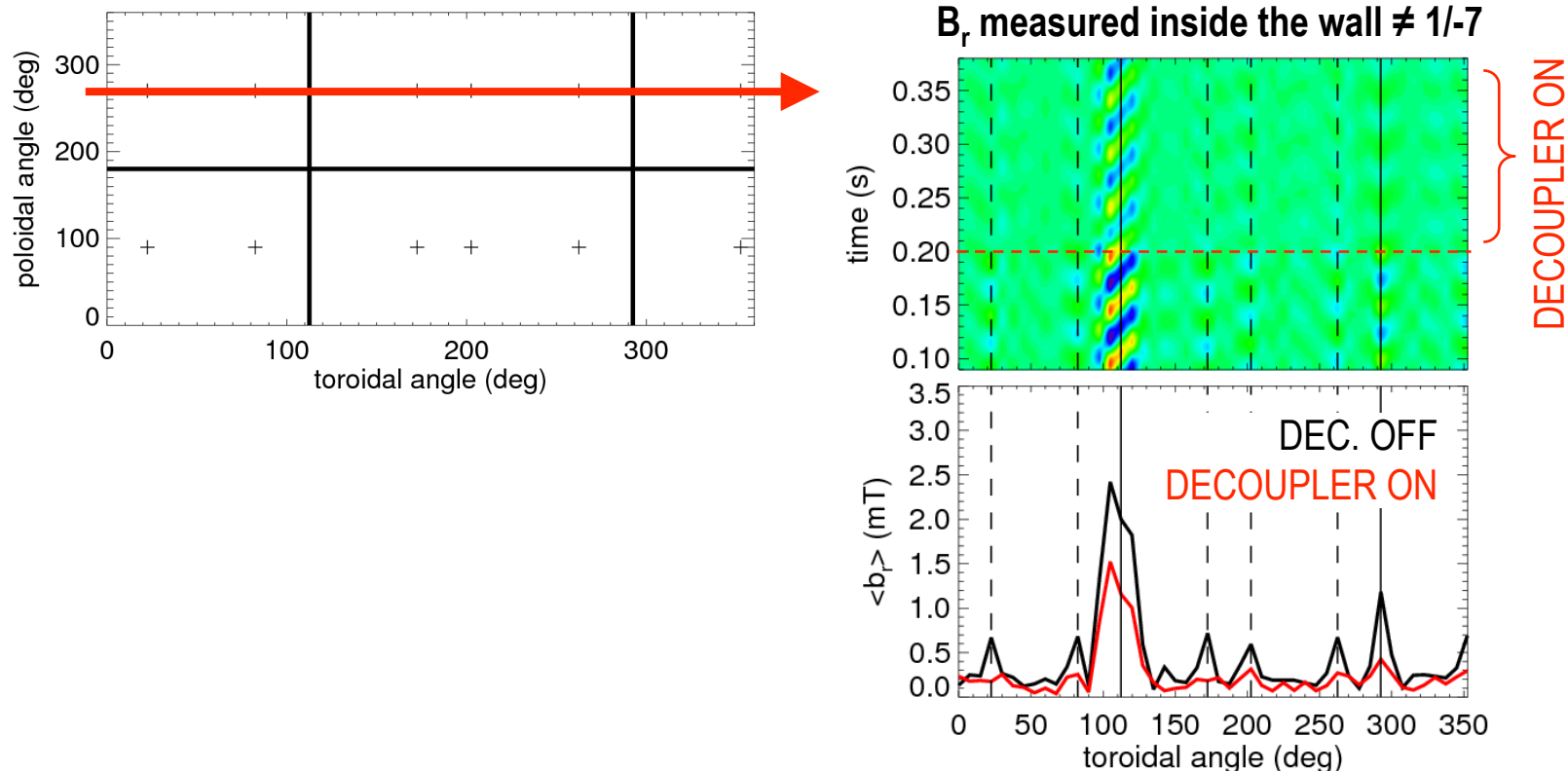
Other error fields due to poloidal gaps and portholes

- Other error fields with lower amplitude are produced by the rotating 1/-7 perturbation near the two poloidal gaps and some large portholes
- A more general decoupler that inverts the matrix of transfer functions among all 192 active coils and 192 B_r sensors was used to compensate for these error fields



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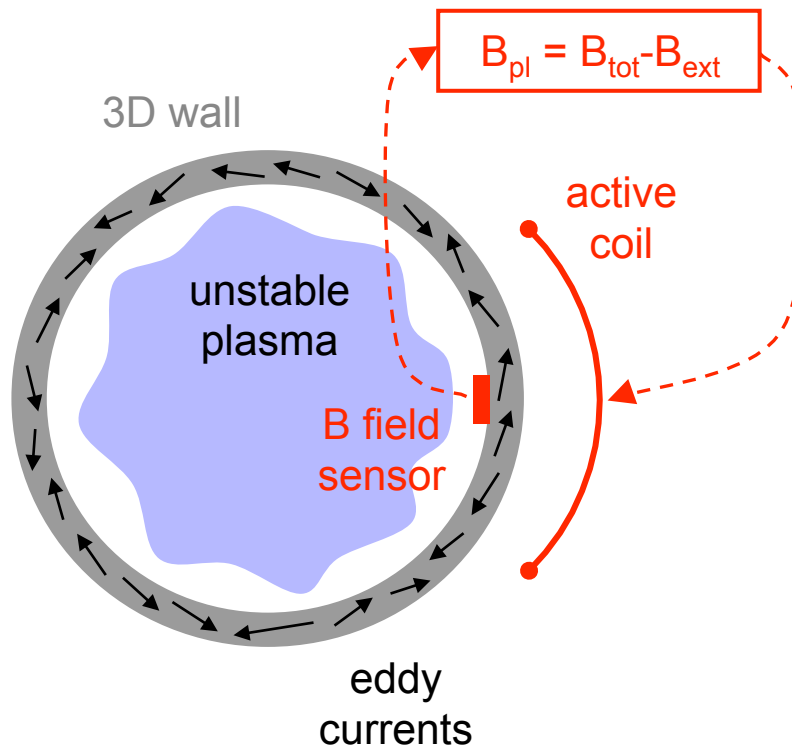
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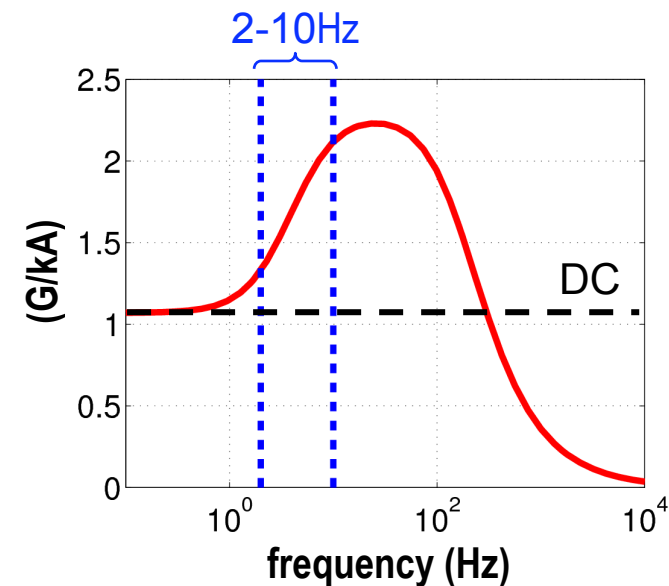
DIII-D

DC vs AC sensor signal compensation in DIII-D

- To suppress a n=1 RWM, or the plasma response to n=1 error fields, the measured magnetic field must be **compensated for spurious n=1 fields** produced by other sources, such as the active coils, slightly misplaced axisymmetric coils, or by **eddy currents induced in the 3D wall structures** by any magnetic field varying in time

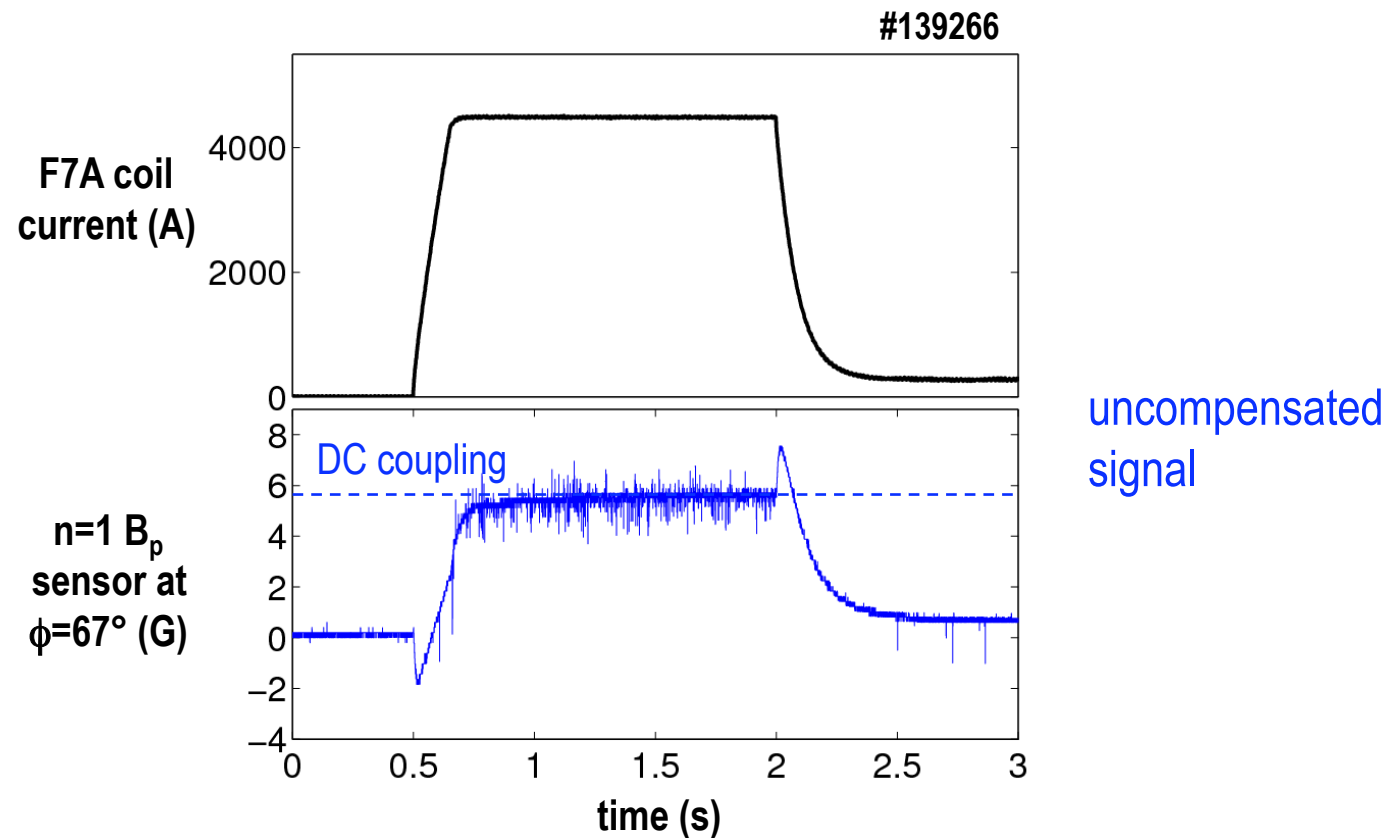


Transfer function among a field shaping coil (F7A) and a n=1 B_p sensor at $\phi=67^\circ$



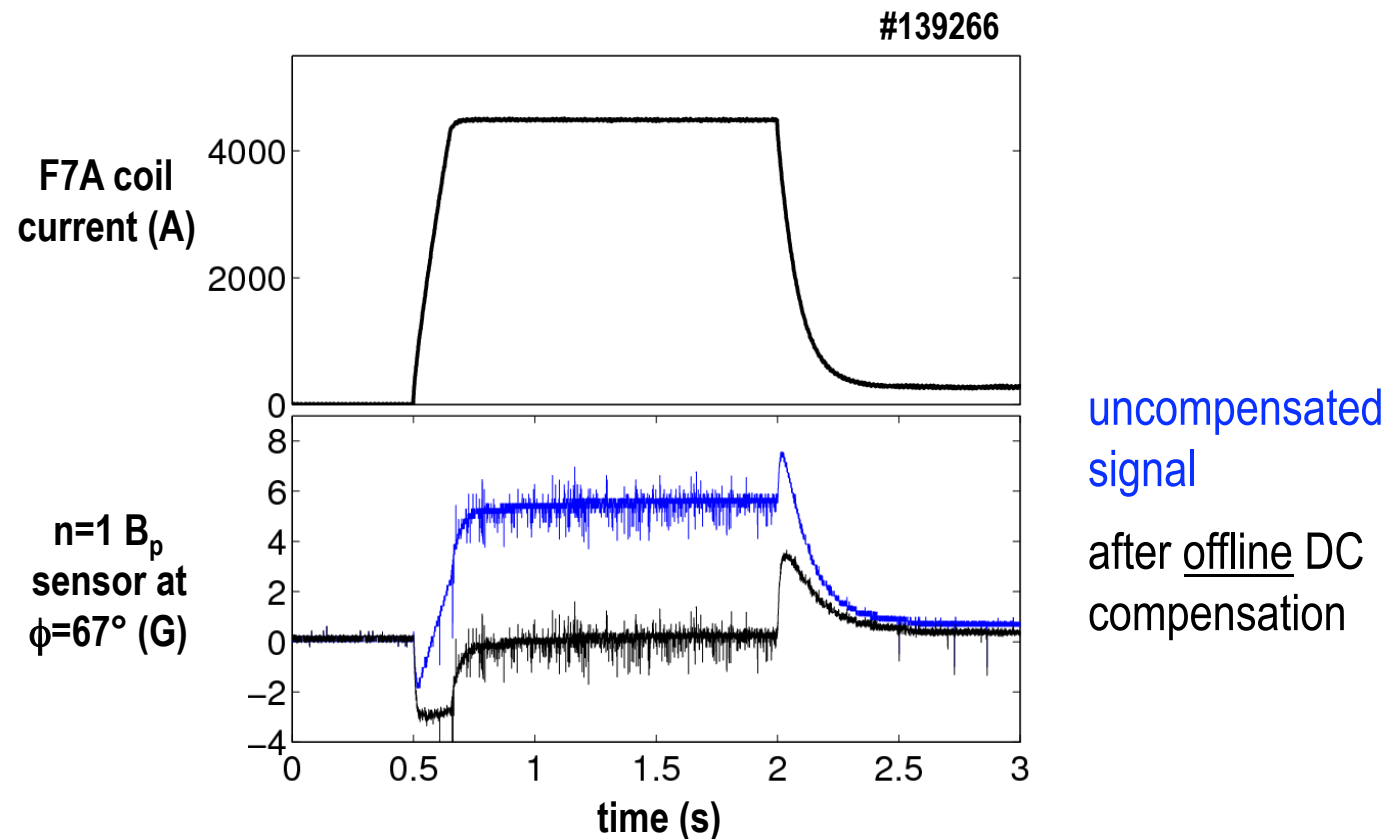
Example: spurious $n=1$ B_p fields from the F coils

- The field shaping coils produce spurious contributions to the $n=1$ B_p signal used for feedback, which need to be compensated for in real-time. This is usually done with a **DC compensation scheme** based on **vacuum measurements of the DC couplings**



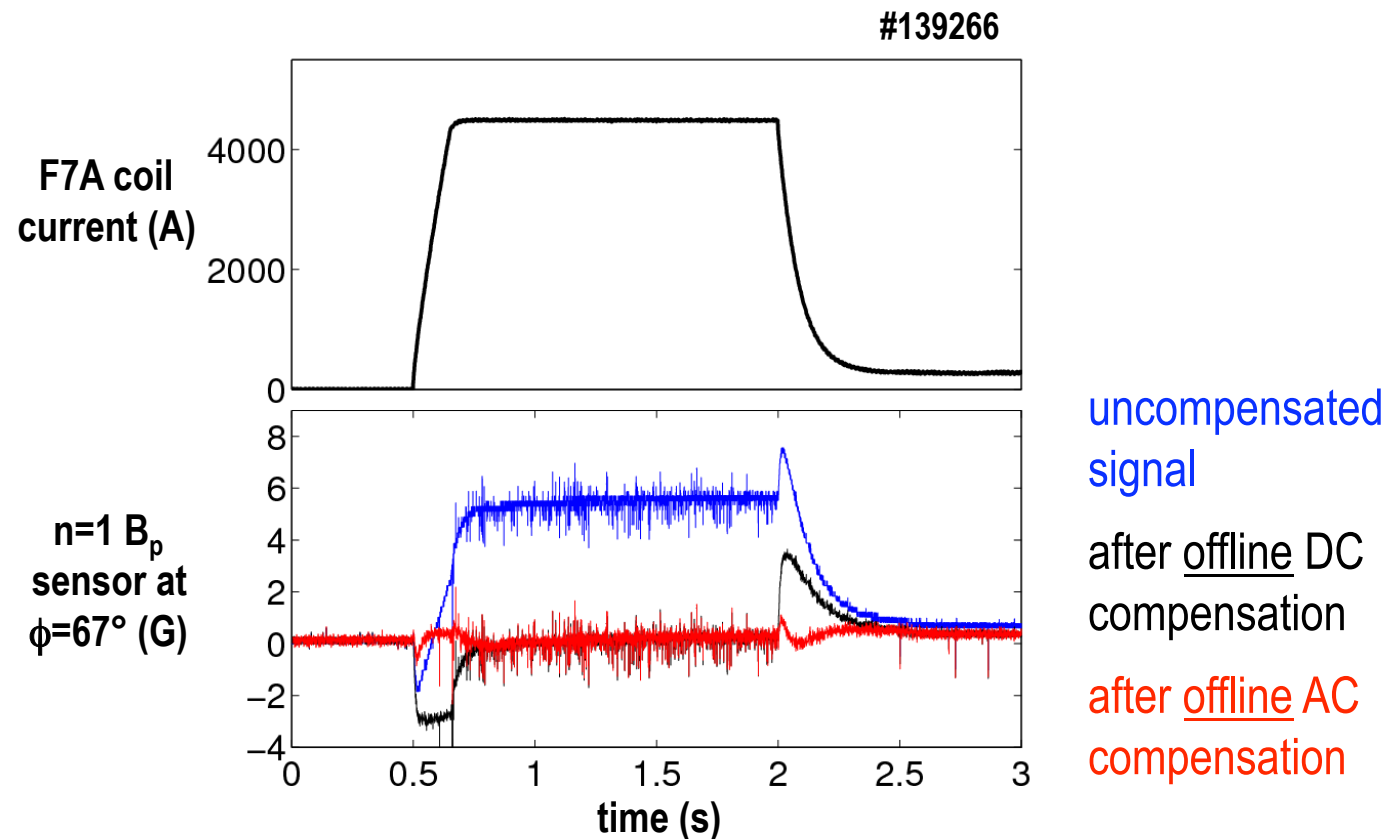
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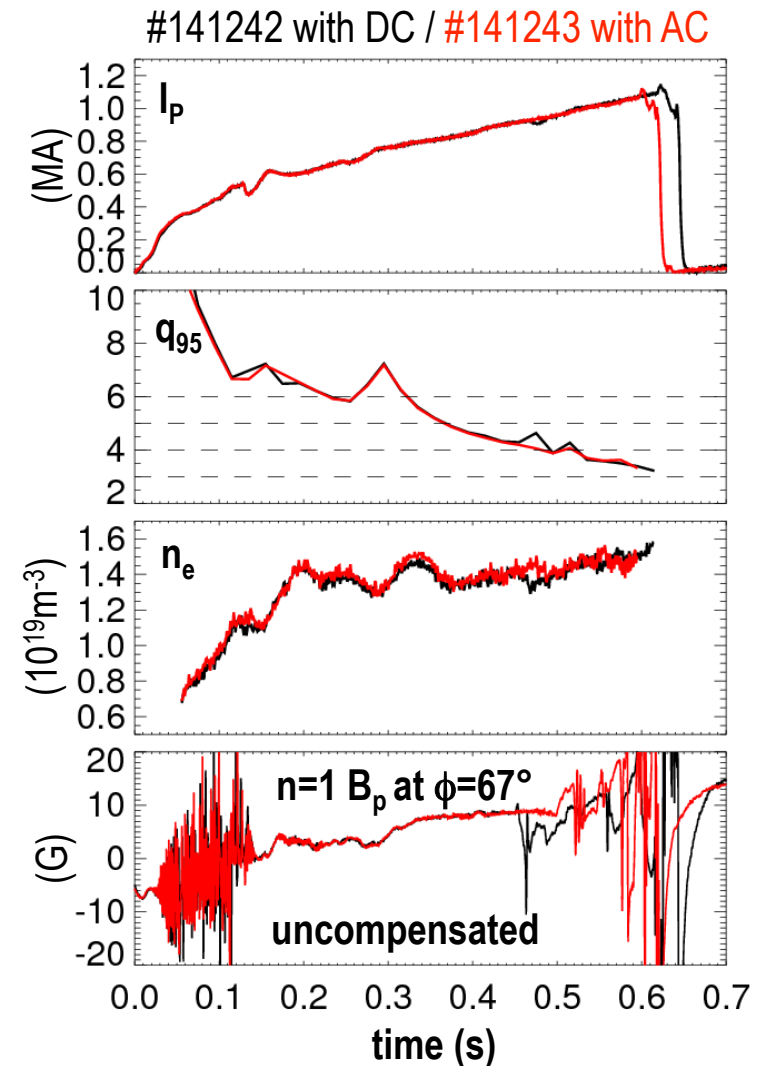
First tests of AC compensation in vacuum

- During transients, the eddy currents induced in the wall produce a significant effect on the $n=1$ B_p signal, which can be compensated with a scheme recently implemented in real-time in DIII-D, based on **transfer functions measured in vacuum**



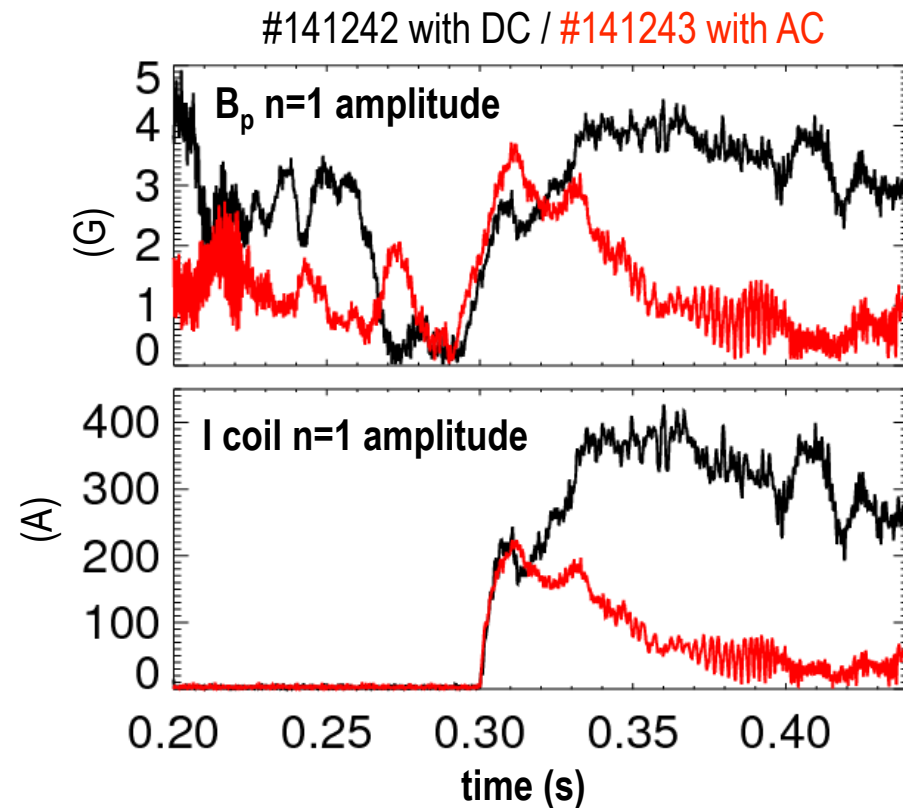
Test of AC compensation in Ohmic plasmas

- DC and AC compensation have been compared in two very similar Ohmic discharges
- A fast current ramp was programmed to destabilize current-driven RWMs, but unfortunately the targeted $n=1$ mode at $q_{95}=4$ turned out not to be unstable
- Nonetheless, we could compare the effect of DC and AC compensation on **dynamic error field correction**



Test of AC compensation in Ohmic plasmas

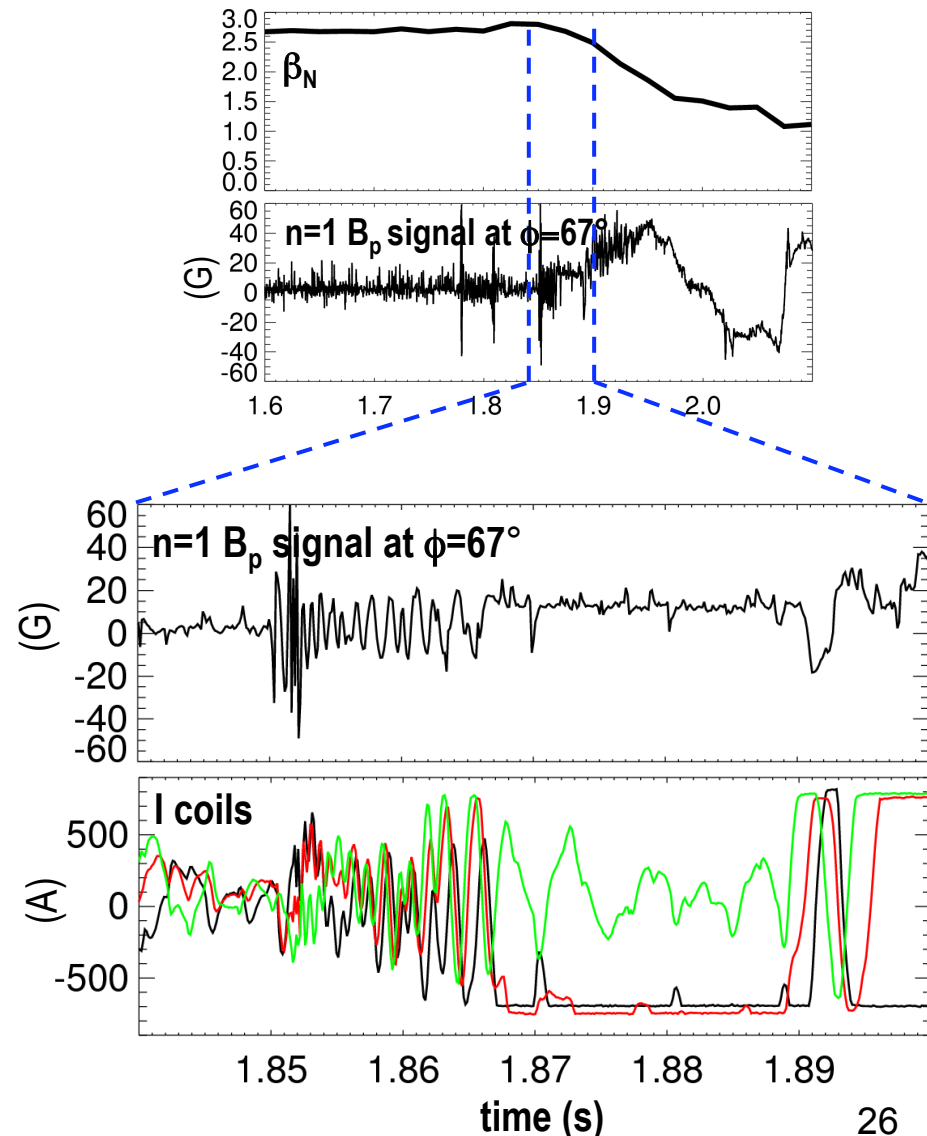
- In these cases, the largest AC effects are associated with the **variation of the F coil currents** during the current ramp-up
- The $n=1$ B_p amplitude estimated with AC compensation is significantly smaller than with DC compensation. This is **completely due to the different algorithm**, not to different plasma parameters
- As a consequence, the I coil current is significantly reduced and the production of additional error fields avoided



Possible relevance for high- β plasmas

- In high- β plasmas, fast feedback is quite effective in suppressing **ELM** or fishbone-driven **RWMs**
[M. Okabayashi *et al.*, Nucl. Fusion **49**, 125003 (2009)]

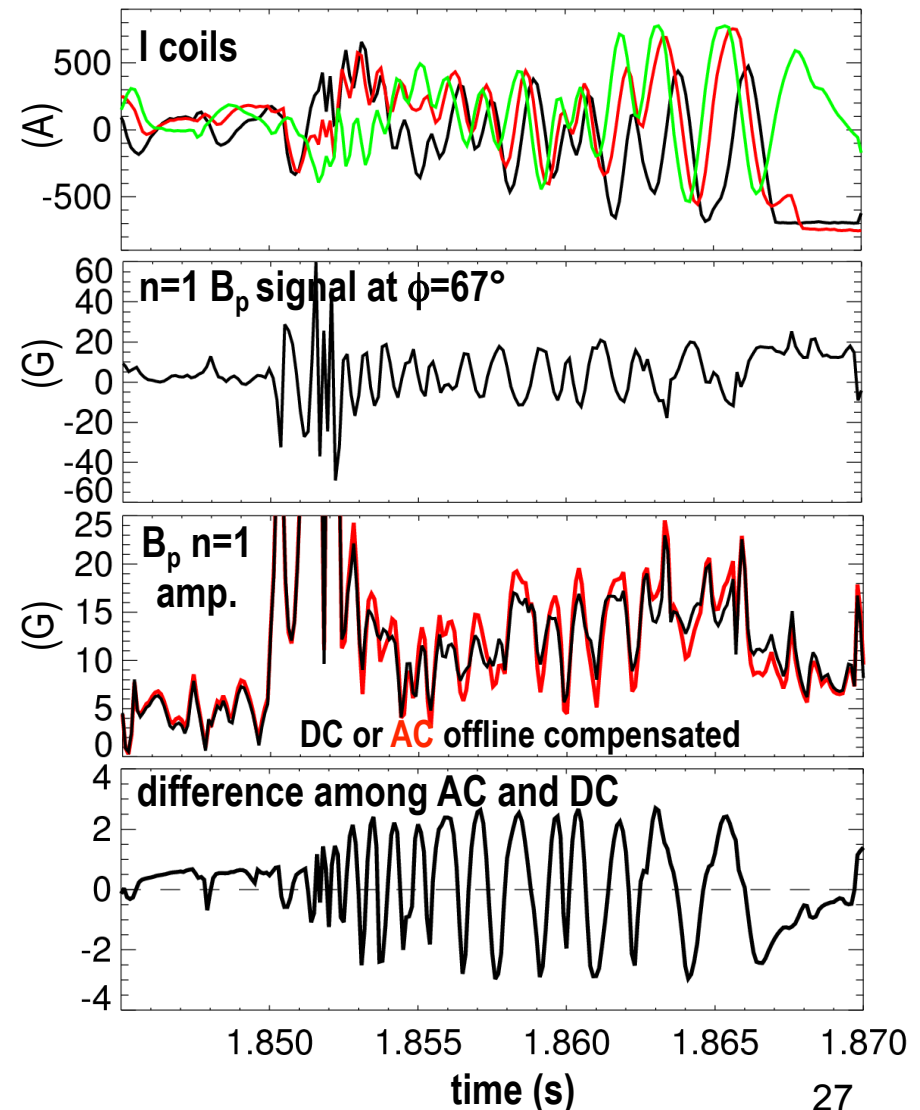
- In some cases a residual oscillation remains, causing the I coil currents to increase and saturate, which brings to a β collapse
- The residual oscillation stops as the I coil currents saturate. This may indicate that this is not an instability, but possibly amplification of the field produced by the I coils.
- But why does this oscillation persists?



Error fields due to AC effects may be destabilizing

#125639

- The $n=1$ mode amplitude computed including AC effects differs from that with DC compensation by an error that oscillates in time
- At some times the feedback overcorrects the mode, adding a small error field; while at others it undercorrects it
- In high- β plasmas, such errors can be amplified and may induce the residual oscillation responsible for the saturation of the I coils and the β collapse



Conclusions and further work



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 - *Future work*: more general decoupling schemes may be developed and implemented in real-time, to assess their relevance especially for high-performance scenarios at high plasma current up to 2MA, but also in low-current 100kA tokamak plasmas

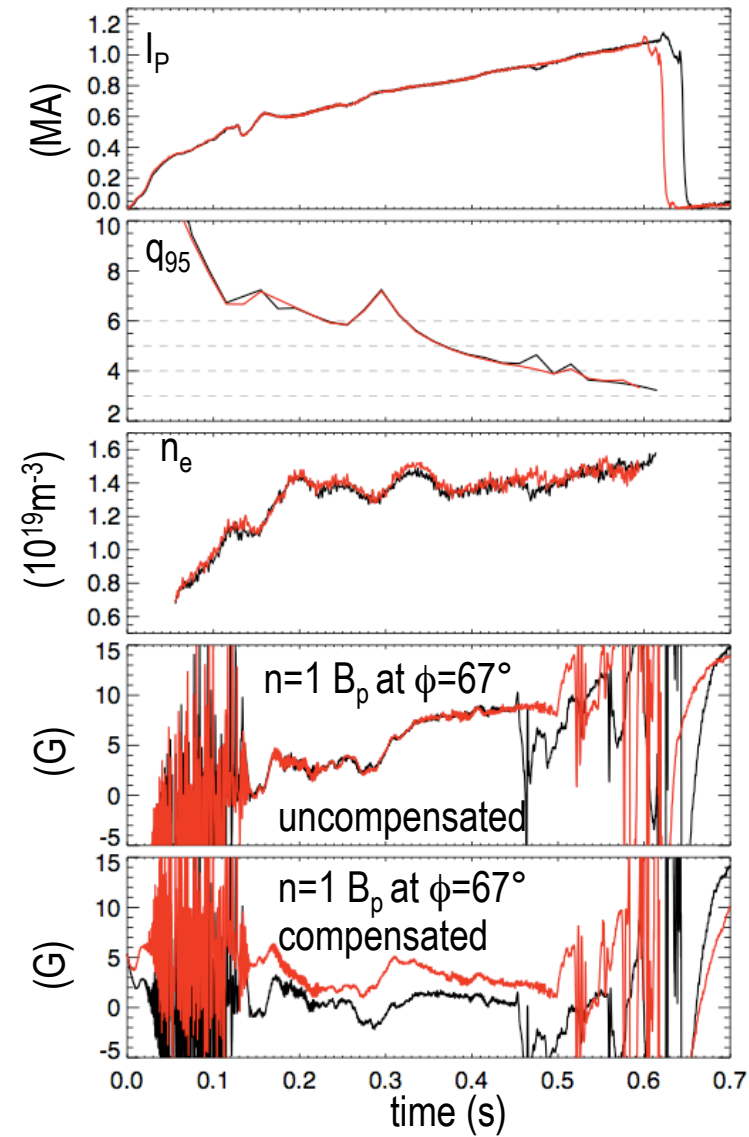
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**Thank you for your
attention**

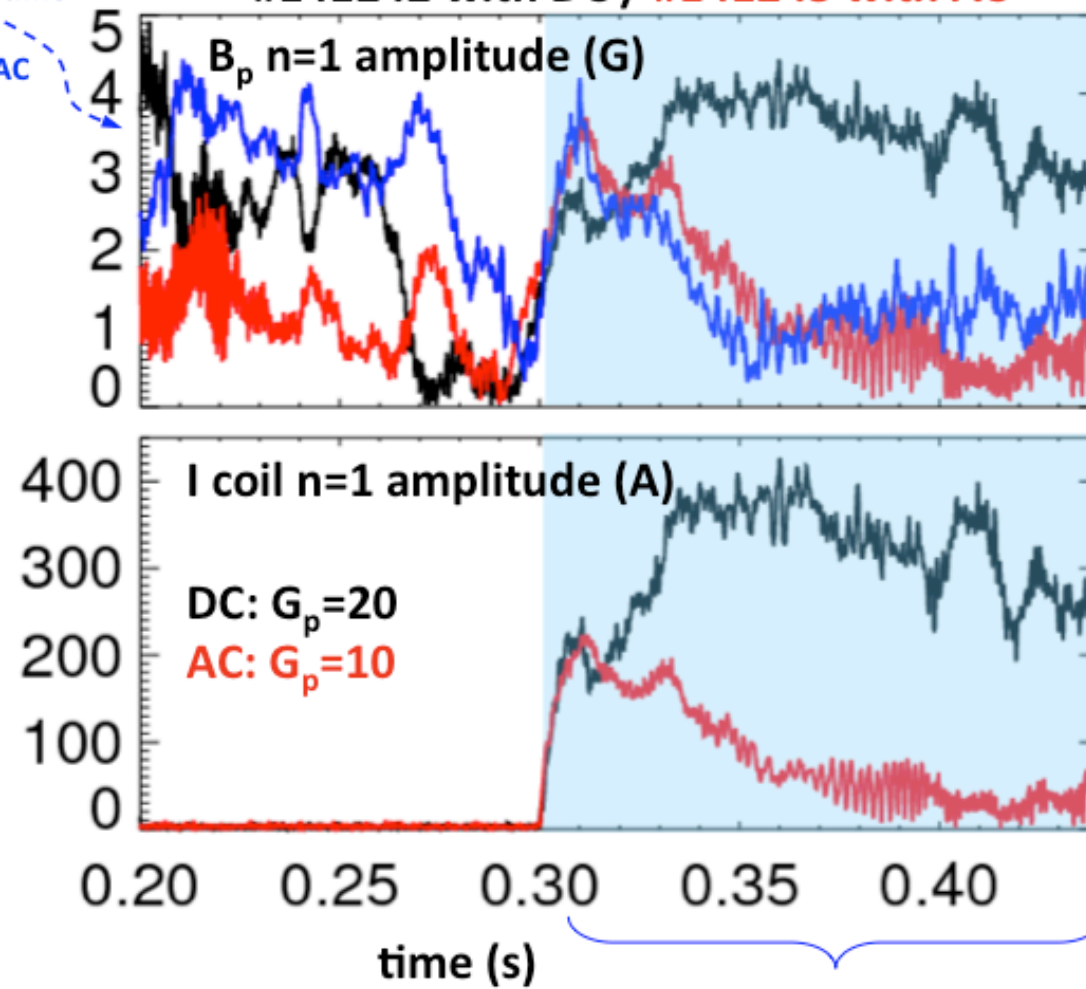
Spare slides

#141242 with DC / #141243 with AC



Case with real-time
DC comp.
compensated AC
OFFLINE.

#141242 with DC / #141243 with AC



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