

# **3-D self-organization in RFPs with connections to stellarator MHD**

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# Background

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**Self-organization to helical state is rather general trend in the RFPs in certain operational regimes.**

**A dominant tearing mode grows spontaneously, then transition to QSH occurs with reduction of amplitudes of the secondary modes.**

**Optimum regime for self-organization lies in shallow-reversal regime.**

**Higher plasma current  $I_p$  or higher values of  $S$  may be favorable to longer and purer QSH (or SH) state.**

**Growth time of the single dominant mode may have machine (or some other parameter(s) ) dependence – fast growth in RFX-mod and slow growth in MST.**

**The self-organized helical RFP is accompanied by formation of the internal (electron) transport barrier.**

# Motivation

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**Helically deformed RFP is also realized in our low-aspect-ratio machine RELAX.**

**Current dynamics of self-organization in RFPs is dominated by current driven tearing modes.**

**What would be the role of pressure driven modes in high beta RFP plasmas?**



**Studies on pressure driven modes have been carried out in LHD.**

# Outline

**A brief review of 3-D self-organization in RFPs**

**Operational regimes in low-aspect-ratio RFP plasmas in RELAX**

**Comparison of QSH in  $A=2$  RFP with 3-D MHD computation results**

**Internal structure of pressure driven MHD instability and its role in profile change in LHD**

# Conventional self-organization theory for RFP

## - Taylor relaxation -

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In low-beta plasmas, minimum energy state for a specified value of magnetic helicity  $K$  can be described by the force-free equilibrium.

$$\nabla \times \mathbf{B} = \mu \mathbf{B}$$

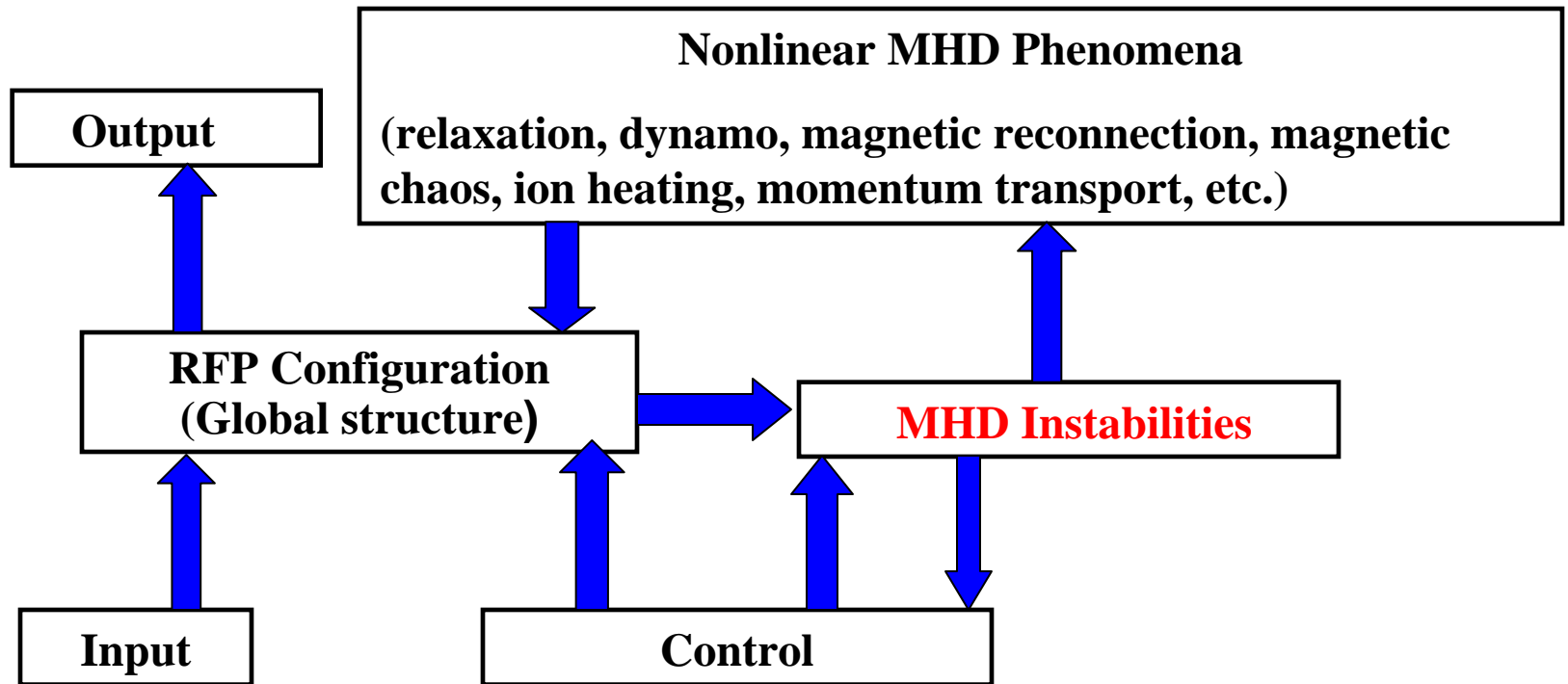
In a cylindrical approximation, the minimum energy solution to the force-free equilibrium bifurcates to symmetric and (symmetric + m=1 helical) states when the magnetic helicity is increased.

For  $K$  larger than a critical value, the minimum energy state is a mixture of symmetric and m=1 helical state, where helix of the m=1 component coincides with that of the external field – externally nonresonant helical structure.

**There exists the major difference between m=1 Taylor state and the present experimental helical RFP state – externally nonresonant helix (Taylor relaxation) or internally resonant helix (experimental self-organized state).**

# Concept of self-organization in the RFP

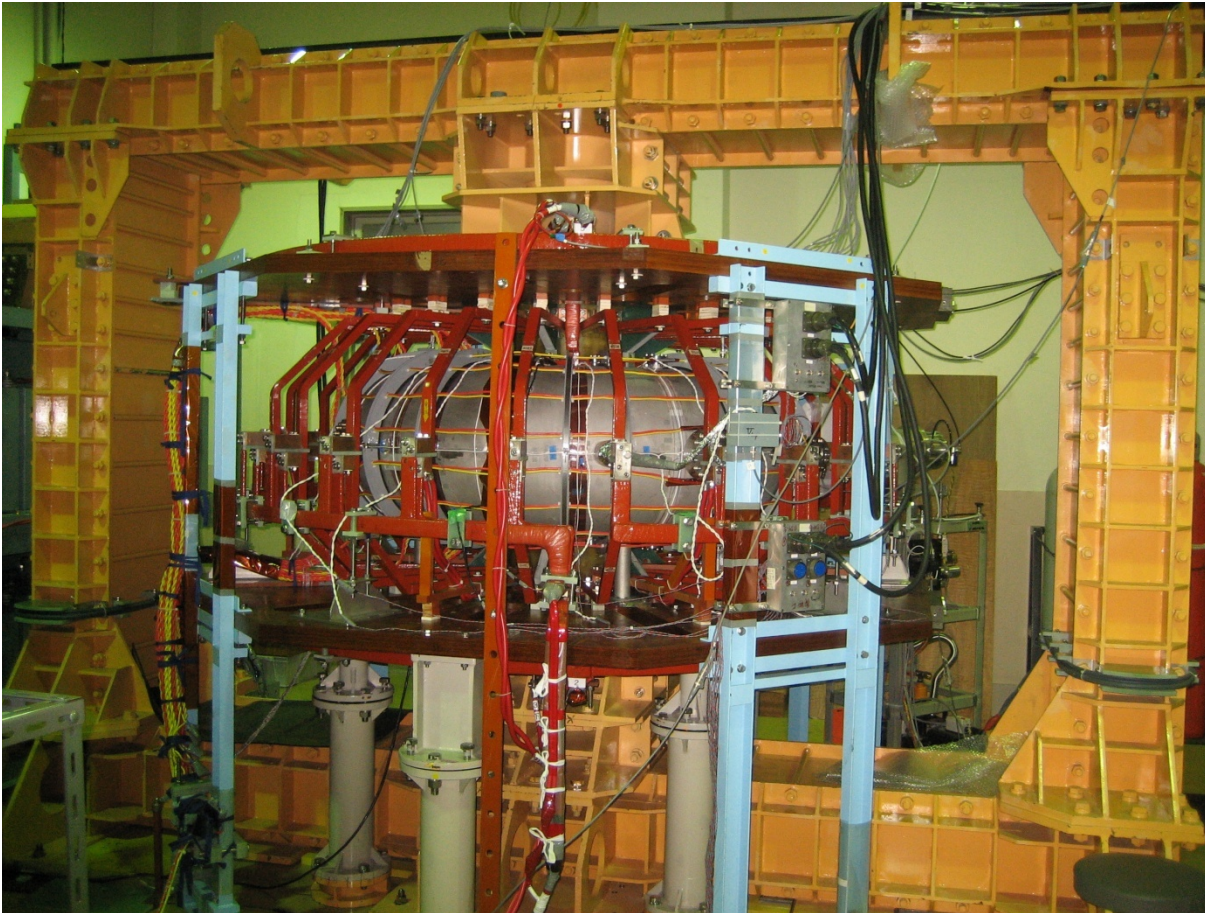
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- formation and sustainment of the configuration through nonlinear MHD
- of general interest as a control problem of highly nonlinear system

# REversed field pinch of Low-Aspect-ratio eXperiment

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•  $R/a = A = 2$   
(51 cm/25 cm)

resistive wall:

$\tau_W \sim 1.5 \text{ ms}$

$I_p < 100 \text{ kA}$

$n_e = 10^{18} \sim$   
 $2 \times 10^{19} \text{ m}^{-3}$

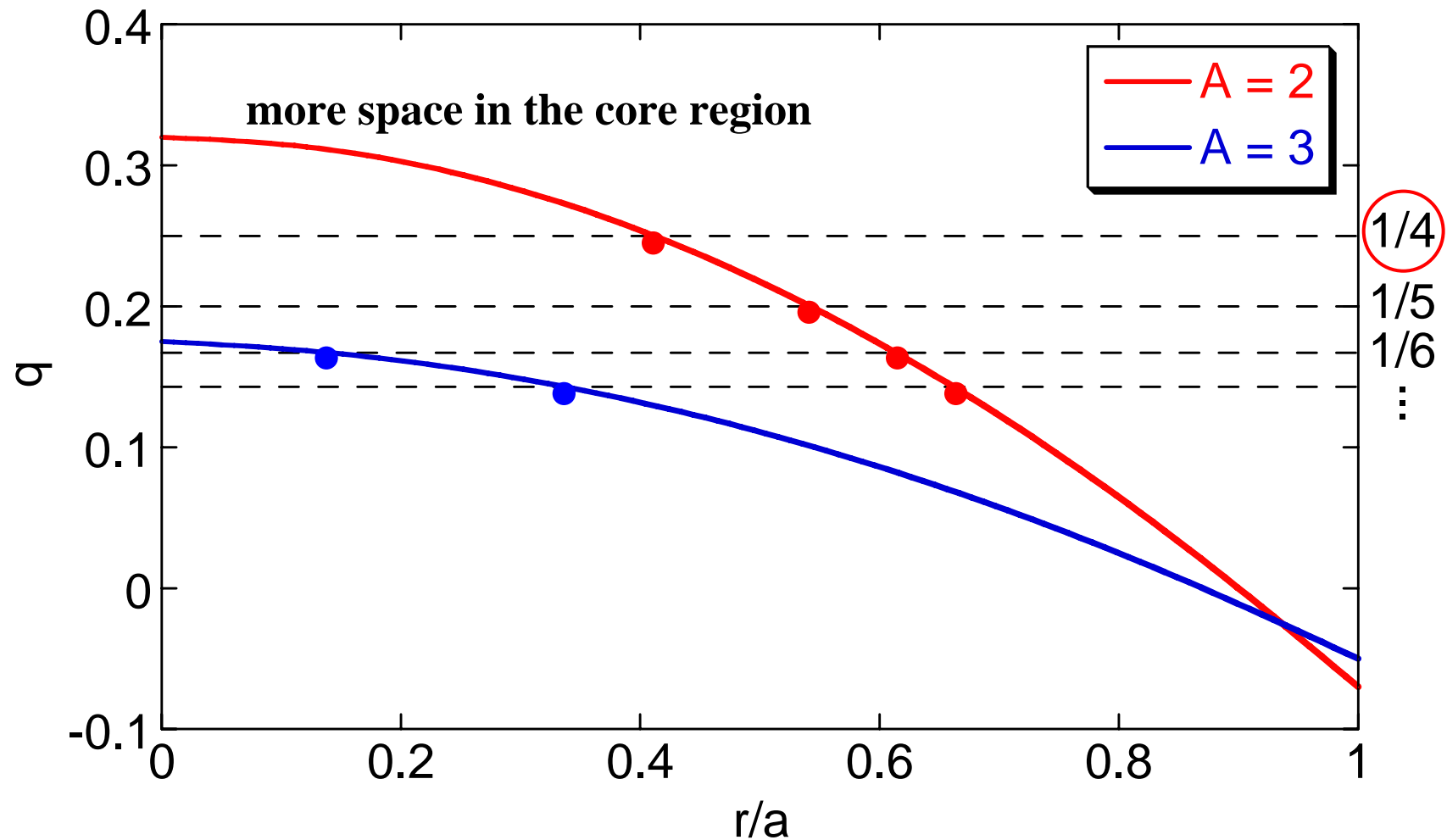
$T_e < 100 \text{ eV}$

$\tau_D \sim 2 \text{ ms}$

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## Lower $A$ means lower $n$ for dominant $m = 1$ modes

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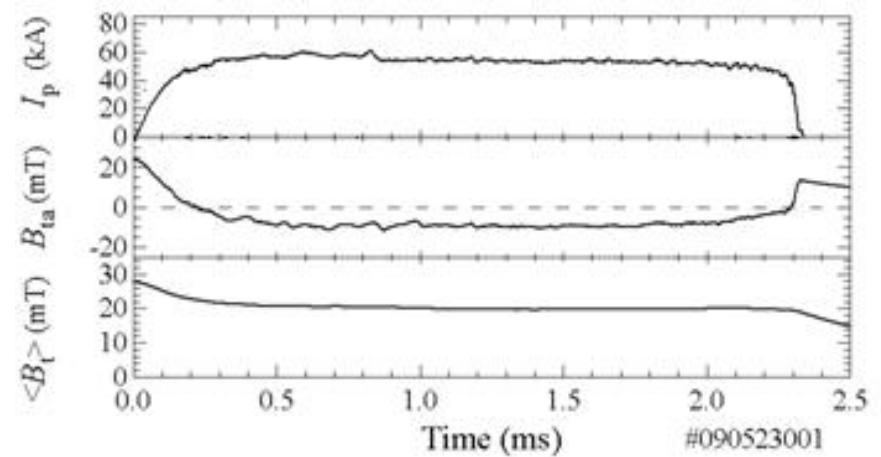
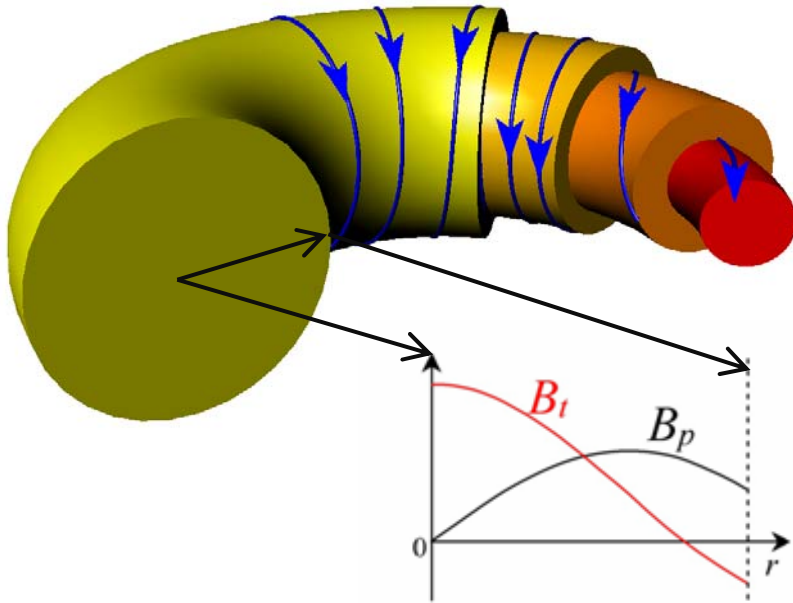


# Motivation to explore the low-aspect-ratio RFP

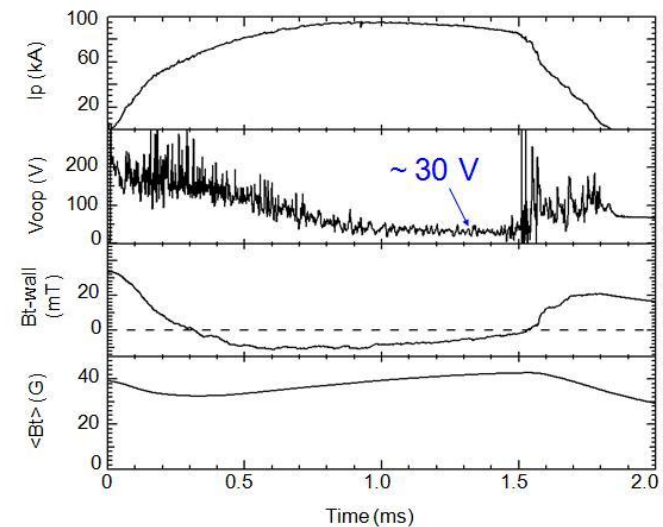
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- With an adequate choice of  $q_0$ , more space would be available in the core region without major resonant surface
  - easy growth of magnetic island associated with core resonant dominant mode
  - Good confinement with QSH for achieving high beta would be expected
- The bootstrap current would become sizable if very high beta could be achievable  
(target parameters:  
 $T_e \sim 300\text{eV}$ ,  $n_e \sim 4 \times 10^{19}\text{m}^{-3}$  at  $I_p \sim 100\text{kA}$ )

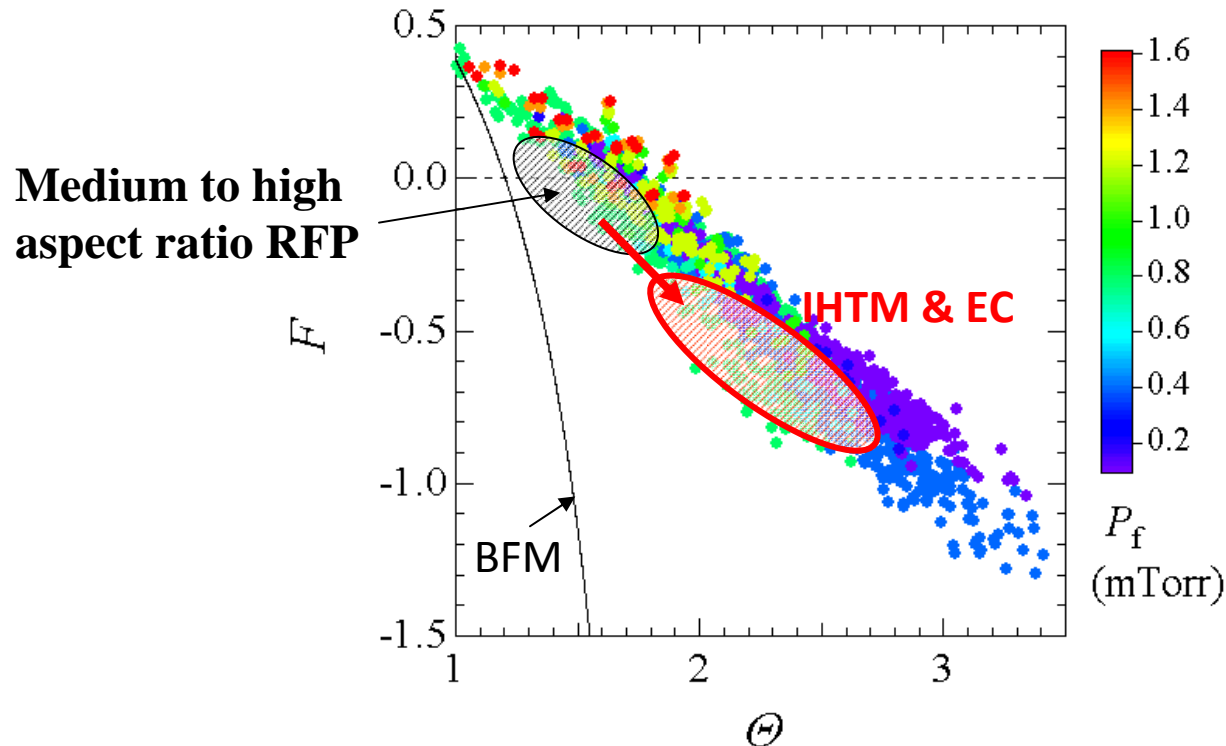
# RFP plasmas in RELAX



- **Low-aspect-ratio RFP plasmas:**  
 $I_p < 100 \text{ kA}$  ,  $n_e = 10^{18} \sim 2 \times 10^{19} \text{ m}^{-3}$  ,  
 $T_e < 100 \text{ eV}$  ,  $\tau_D \sim 2 \text{ ms}$
- MHD properties have been studied in discharges with  $40 \text{ kA} < I_p < 100 \text{ kA}$  during flat-topped phase longer than 0.5ms



# $F$ and $\Theta$ keep some relation over wide range of parameters



$$\Theta = \frac{B_p(a)}{\langle B_t \rangle}$$

$$F = \frac{B_t(a)}{\langle B_t \rangle}$$

- In shallow reversal region,
  - Periodic QSH or Helical Ohmic RFP state tends to be realized
- In deep reversal, high- $\Theta$  region,
  - Amplitudes of resonant modes are suppressed significantly
  - SXR emission increases, indicating good plasma performance

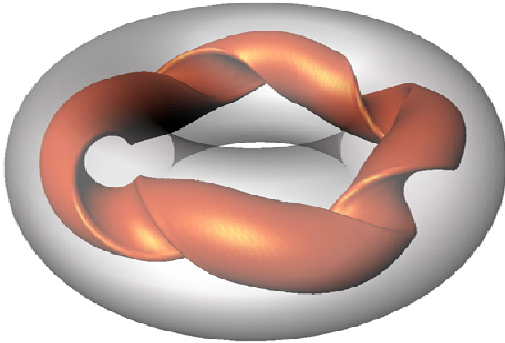
# Quasi-single helicity state in RELAX

$m=1/n=4$  mode behavior:

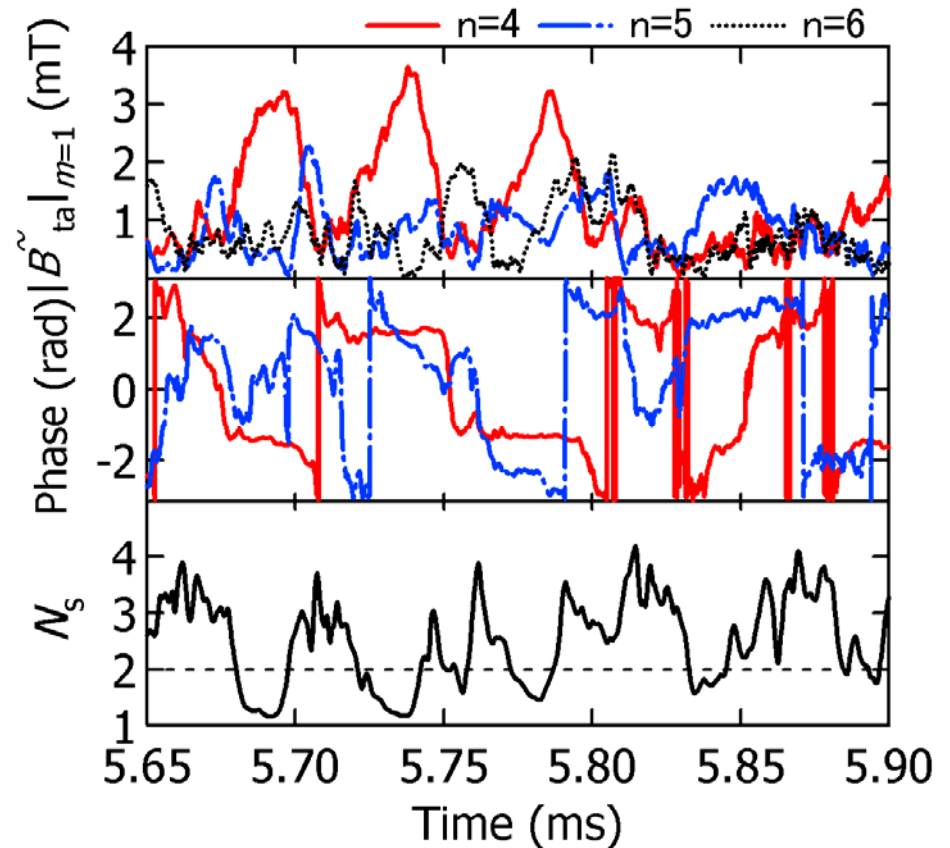
- Longer QSH period for slower rotation
- Shorter QSH period with **higher spectral index  $N_s$  for faster rotation**

Spectral index:

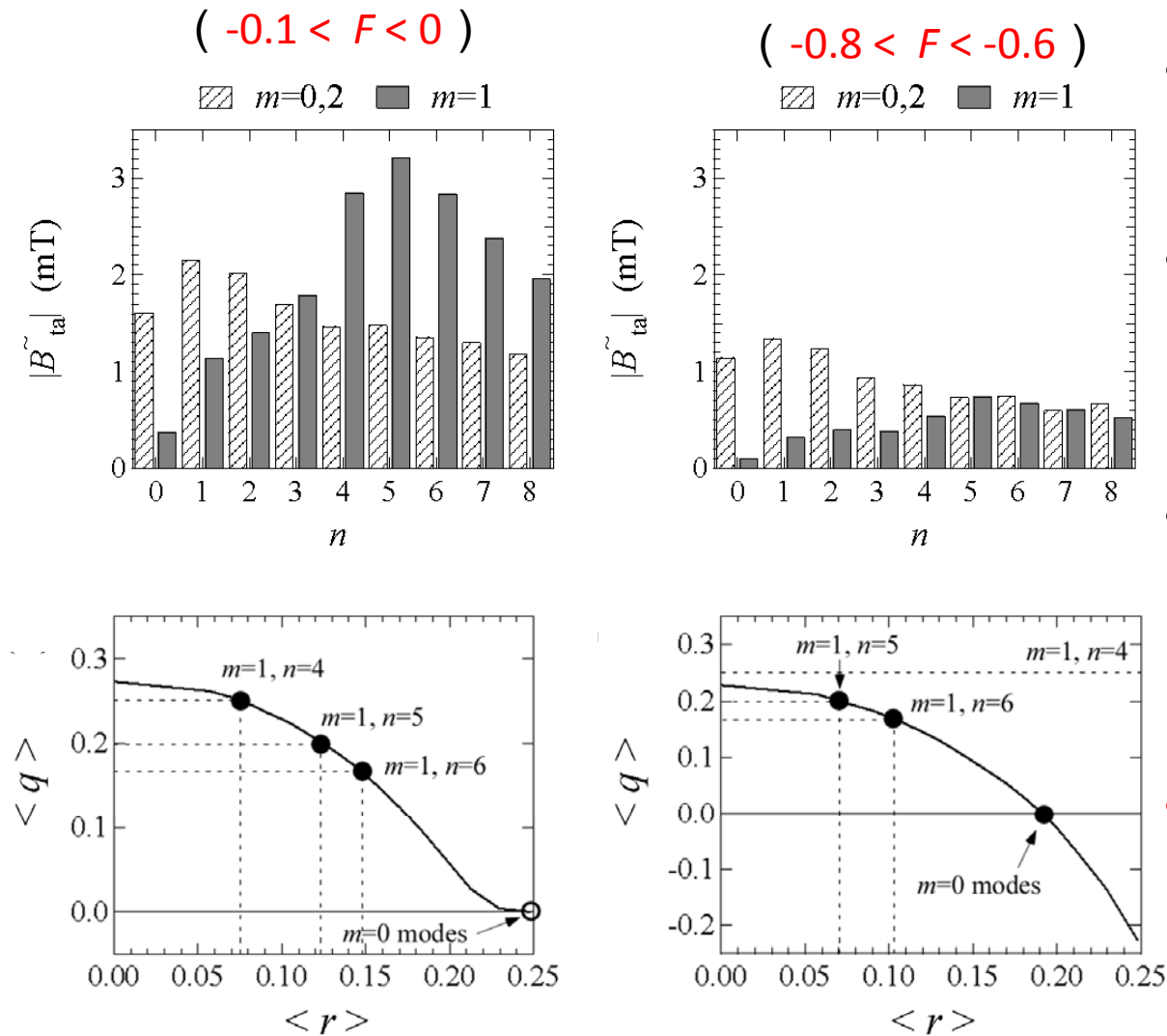
$$N_s = \left[ \sum_{n=n_{\min}}^{n_{\max}} \left( \frac{b_{1,n}^2}{\sum_n b_{1,n}^2} \right)^2 \right]^{-1}$$



Helical hot core is suggested from SXR pin-hole camera



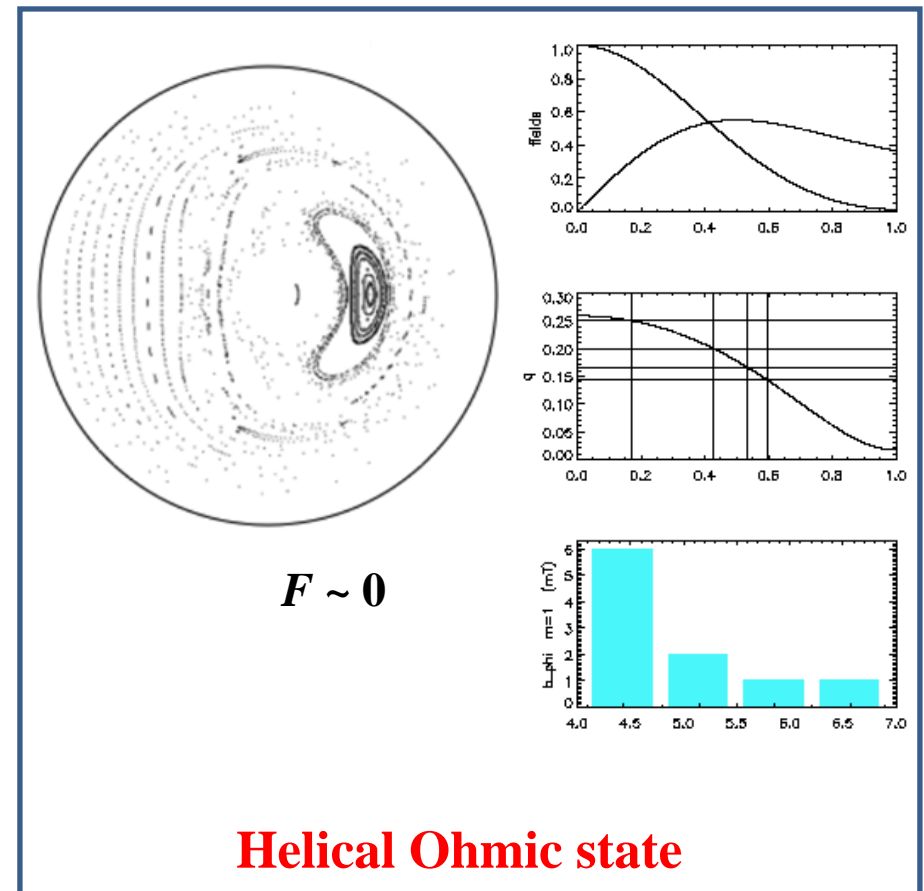
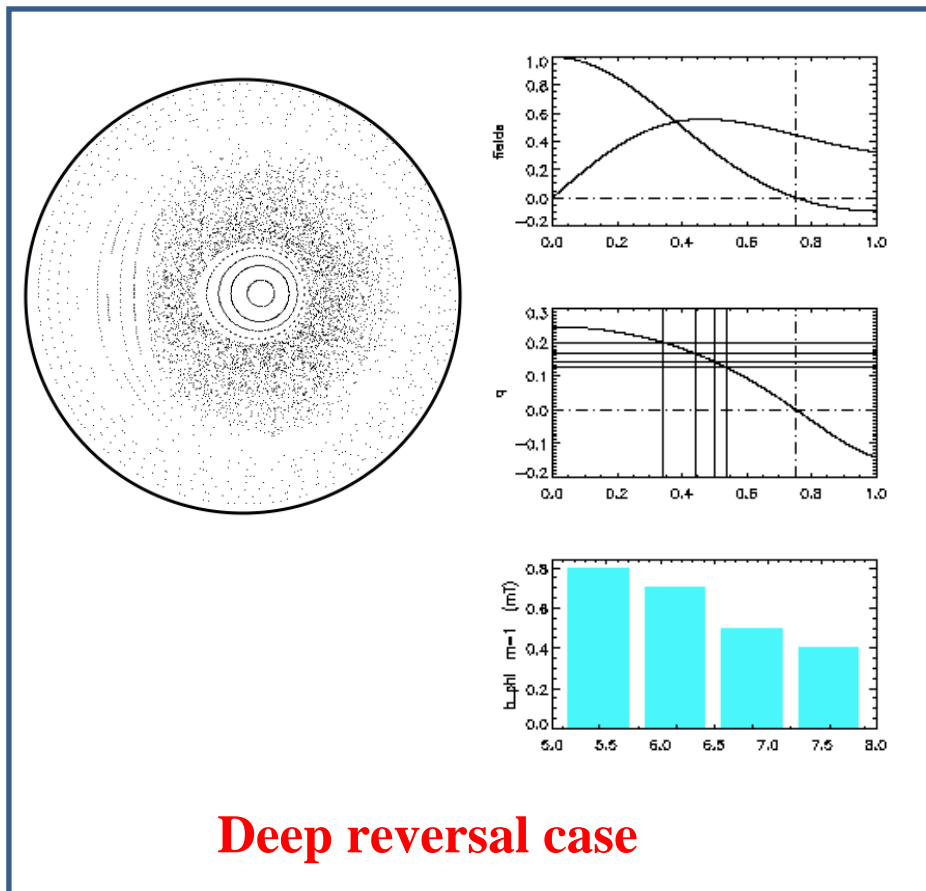
# $q$ profiles and toroidal mode spectrum of $m=1$ modes



- **Shallow reversal region:**  
Dominant modes are core resonant  $m=1/n=4,5,6$ .
- **Deep reversal region:**  
Amplitudes of  $m=1$  modes are reduced to  $\sim 1/4$ .  
Broader spectrum
- Increased magnetic shear may contribute to the lower fluctuation amplitudes in deep reversal case.
- **No evidence of externally resonant mode?**

$q$  profiles from equilibrium reconstruction code RELAXFit

# Field line trace using ORBIT suggests recovery of flux surfaces

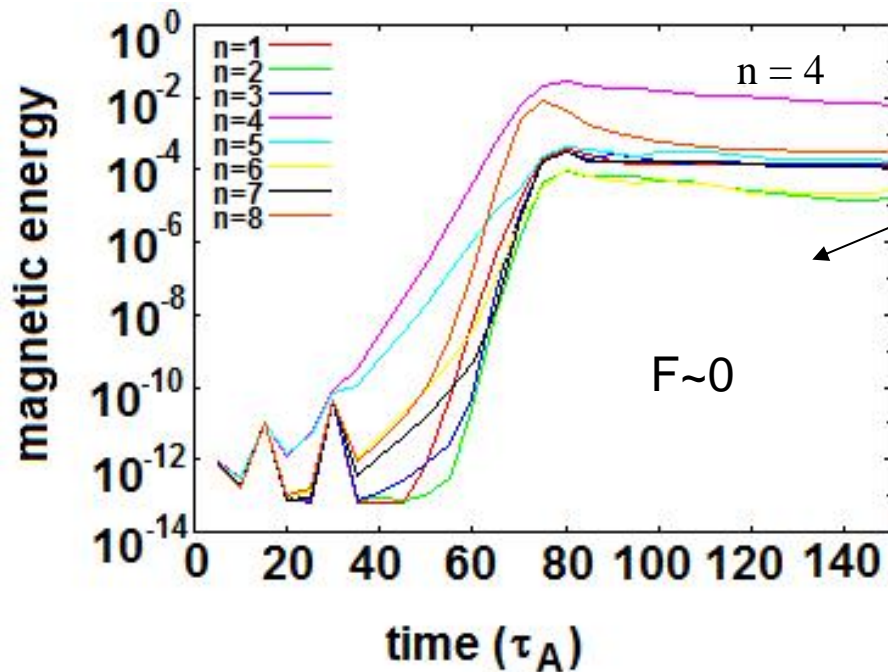


1-D equilibrium , linear  
eigenfunction in cylindrical geometry  
in ORBIT

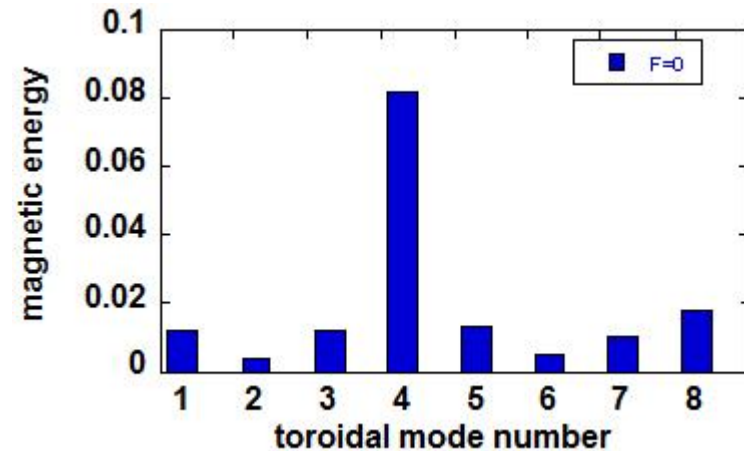
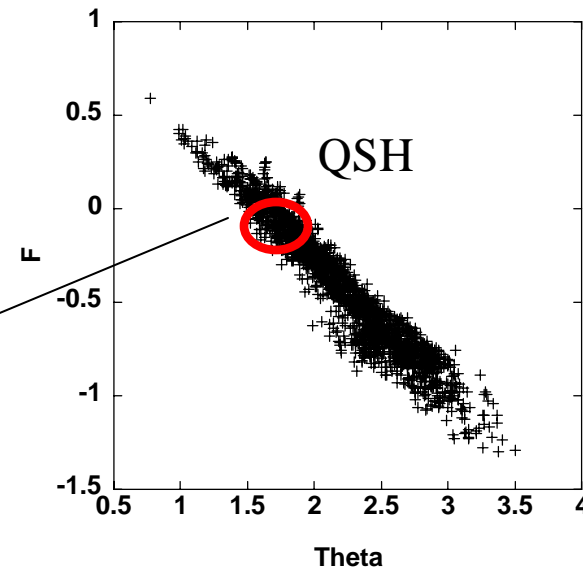


3-D equilibrium and 3-D  
eigenfunctions are needed

# 3-D MHD computation results for $A=2$ RFP



Linear growth of each mode up to  $t \sim 60 \tau_A$ , followed by rapid self-organization to helical equilibrium in  $\sim 20 \tau_A$ .



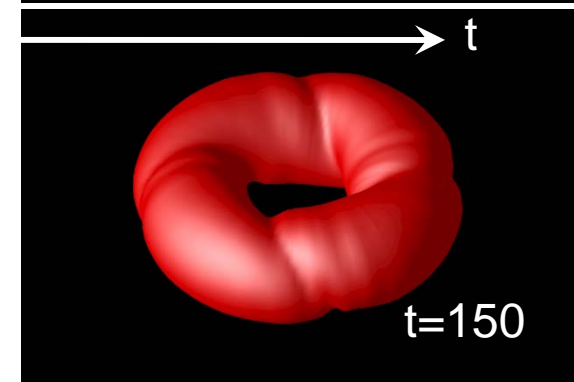
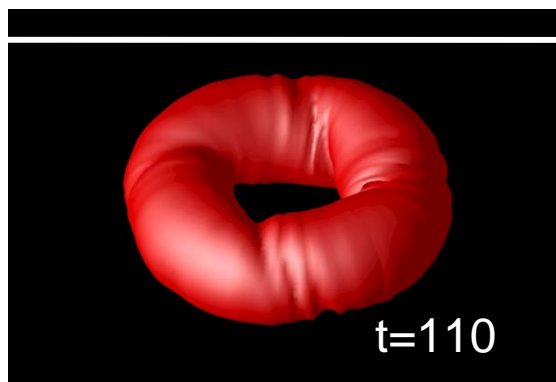
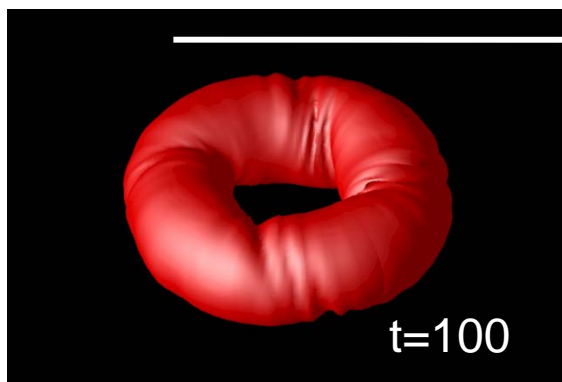
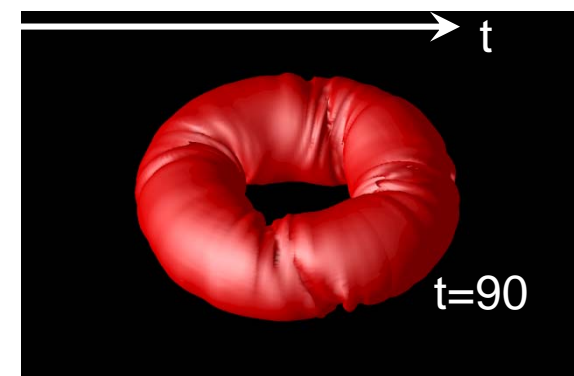
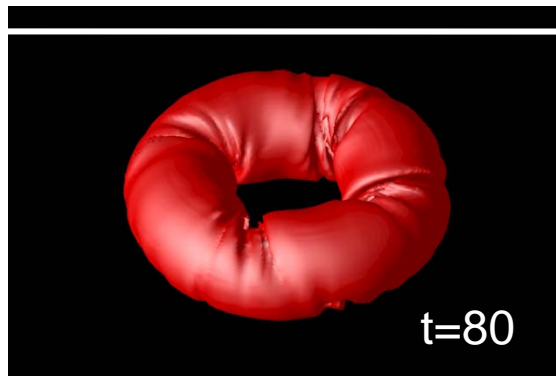
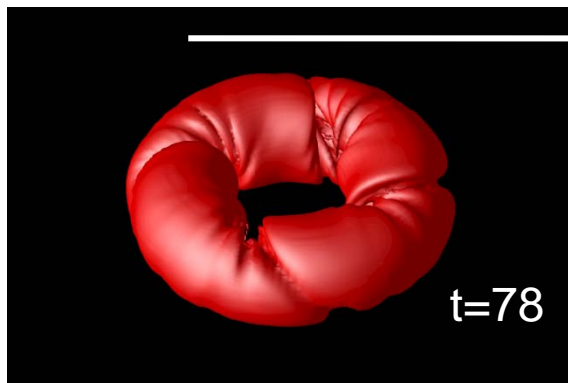
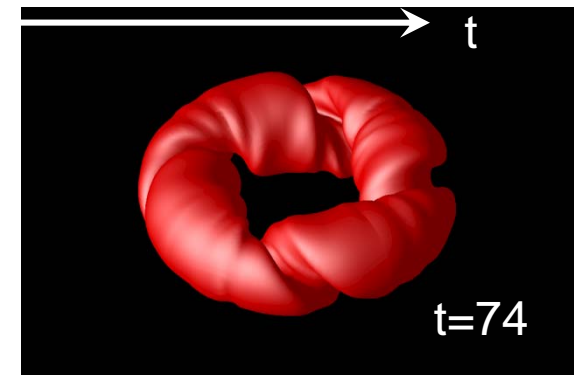
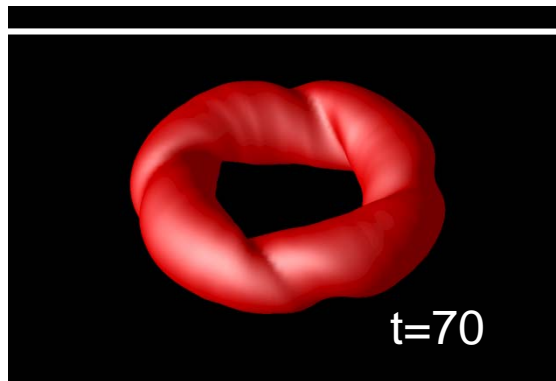
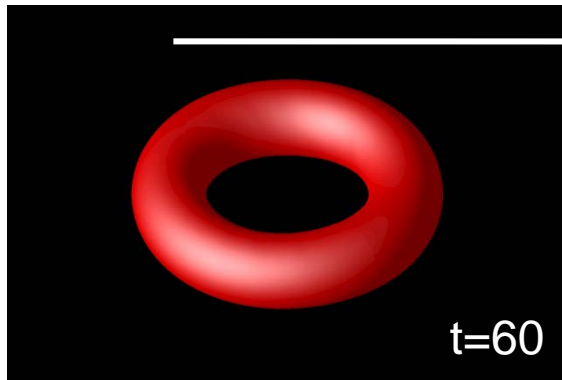
Mode spectrum at  $t=150 \tau_A$

**QSH for shallow  $F$**

=> in agreement with experiments

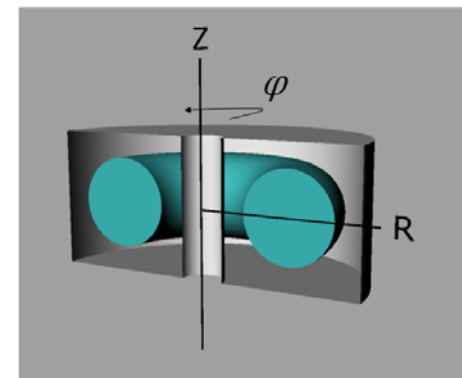
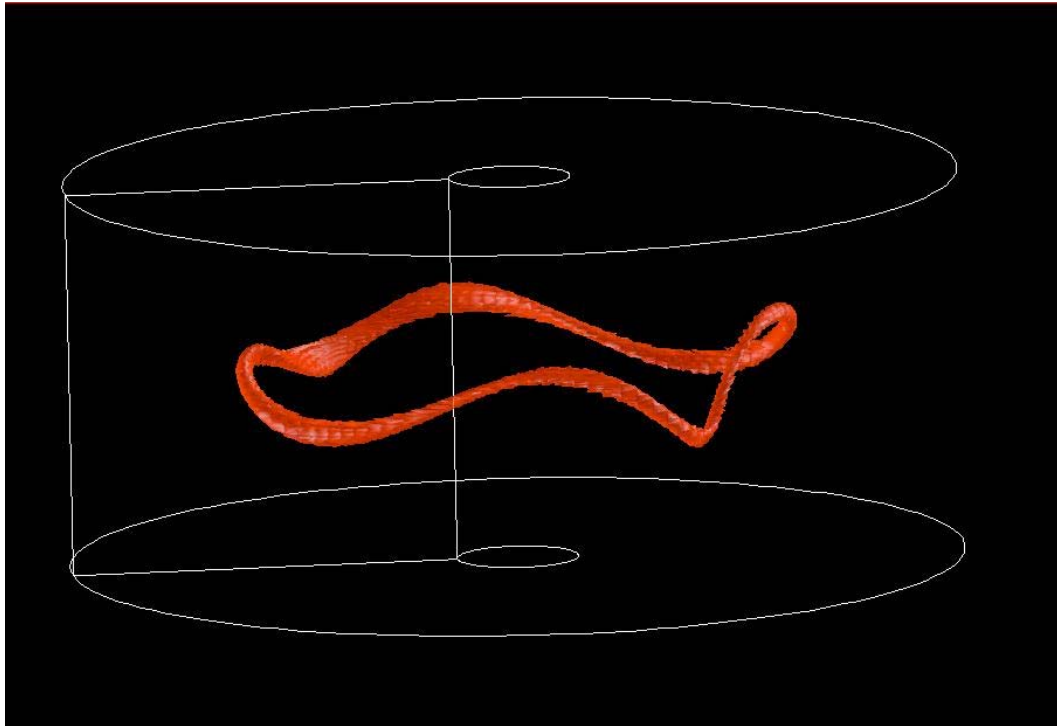
# Time evolution of equi-pressure surface to QSH

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# Possible formation of helical magnetic axis for this particular equilibrium

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# Connections to stellarator MHD

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**Self-organized RFP has a simple structure in real coordinate space**

- Is it still simple in magnetic coordinate?
- How does the field spectrum look like in magnetic coordinate space?
- Is it favorable to particle orbits?

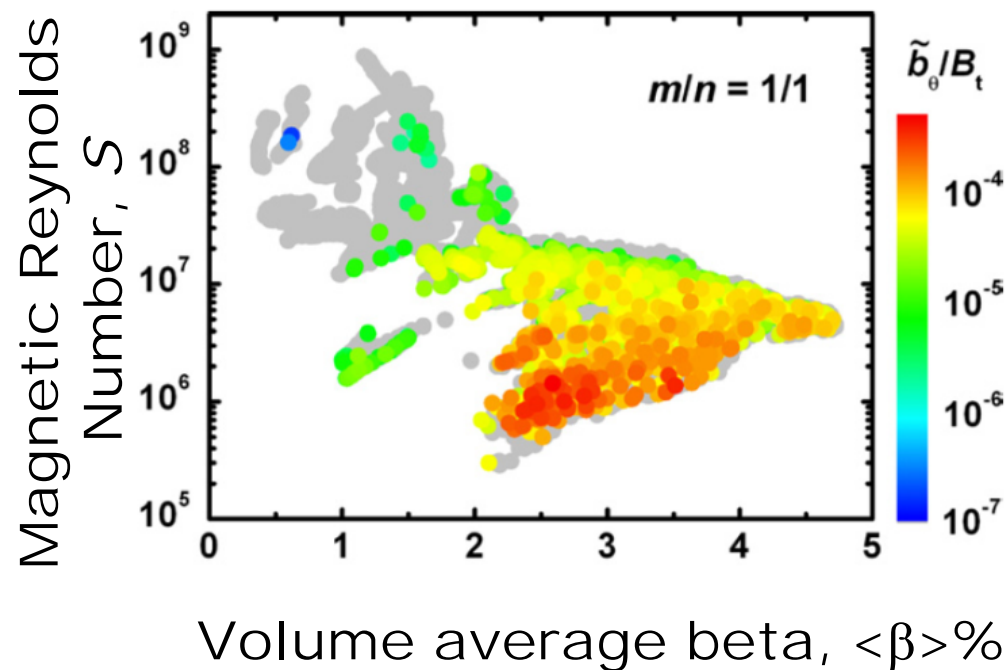
**The plasma boundary may be influential on 3-D equilibrium**

- RFP edge region may be stochastic due to the externally resonant modes
- Can we define clear outermost flux surface or shape of the plasma boundary?
- Is there possible means to determine or define the outermost flux surface in stochastic region?

**Is the self-organization concept applicable to stellarator with strong external helical field? The role of pressure-driven modes?**

# Background for MHD studies in LHD

- ✓ Rotating low-order MHD instabilities are often observed in high-beta LHD discharges.
- ✓ Fine structures or **local flattening appear in density or temperature profiles** in such high- $\beta$  plasmas with possible degradation of confinement performance.



**$m/n=1/1$  magnetic fluctuation level in  $(\beta, S)$  space**

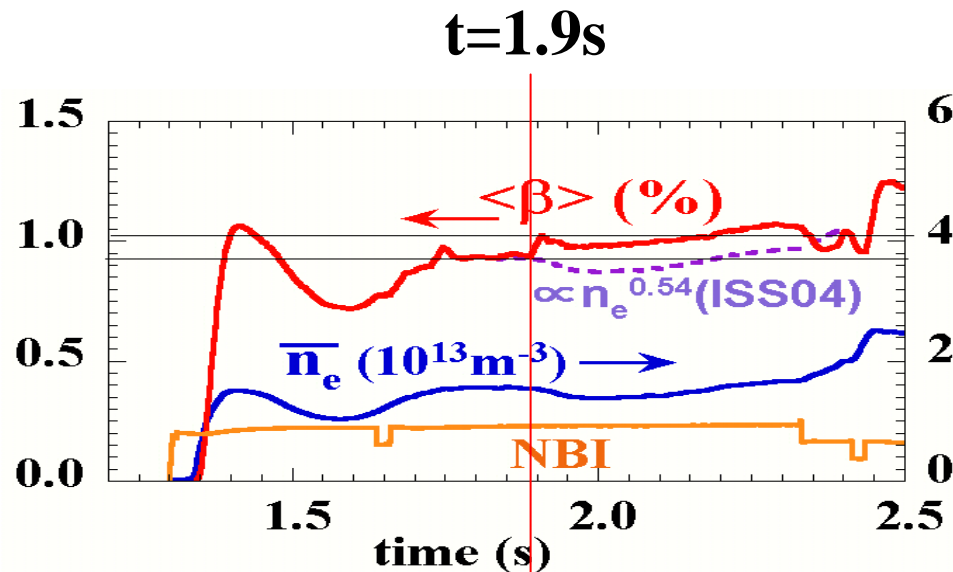
## **Objectives of internal structural studies of low-order MHD mode in LHD**

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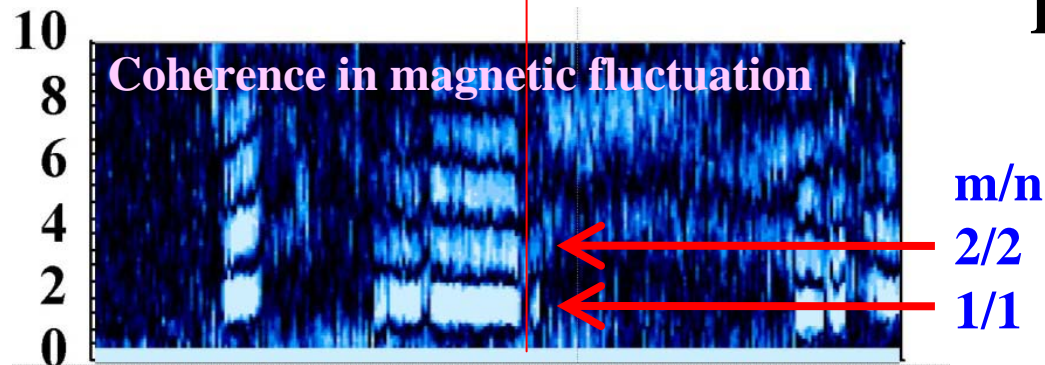
- ✓ to make quantitative evaluation of the effect of the low-order global MHD instabilities on the confinement performance**
- ✓ to find out appropriate parameter(s) characterizing the mode in expressing the relationship between the mode and its effect on confinement performance**

**In identifying the direct effect of the  $m/n=1/1$  mode, we have selected discharges which are marginally stable or unstable to the mode.**

# Example of a discharge marginally stable to m/n=1/1 mode

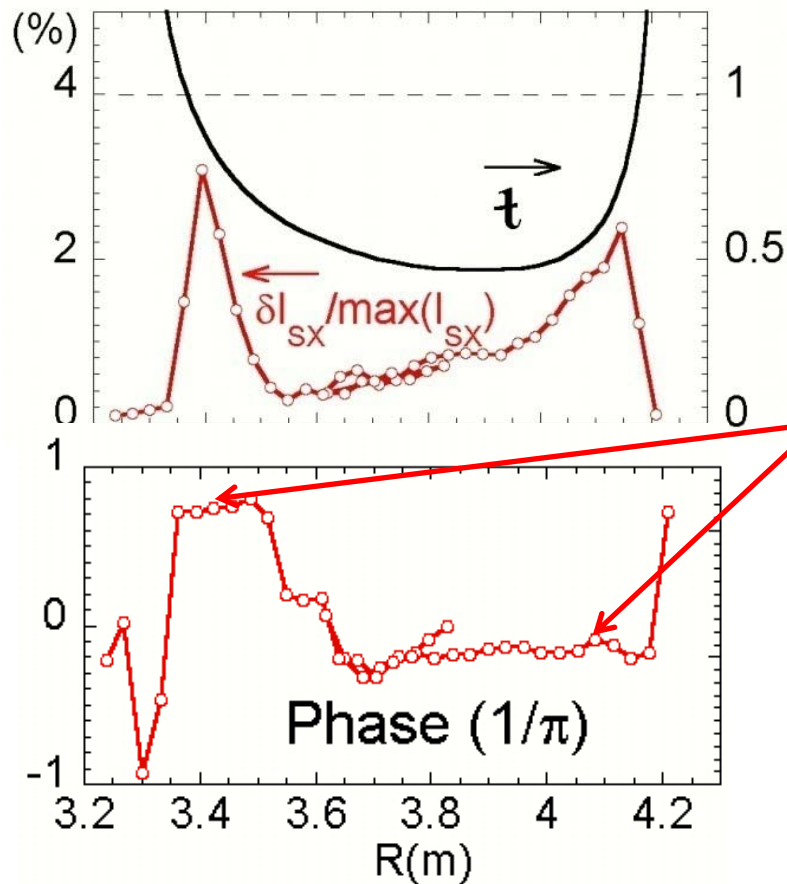


Coherent m/n=1/1 magnetic fluctuation disappears at  $t=1.9\text{s}$ , with improved  $\langle\beta\rangle$  from ISS04 scaling prediction.



Edge magnetic fluctuation level :  $\sim 0.03\%$

# Internal mode structure at $t=1.8\text{s}$ estimated from multi-chord SXR diagnostic



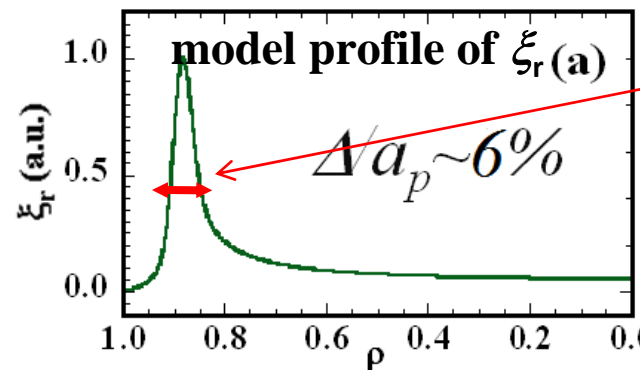
**SXR fluctuation peaks at the  $m/n=1/1$  resonant surface**

**No phase inversion across the resonant surface  
(No magnetic island?)**

**=>**

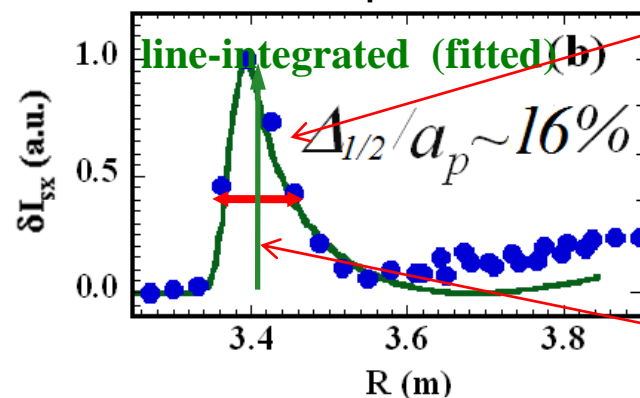
**Similarity to linear theory prediction of resistive interchange mode**

# Parameters characterizing the m/n=1/1 mode and confinement performance



**local mode width  $\Delta$ :**

FWHM from  $\xi_r$  profile



**line-integrated mode width  $\Delta_{1/2}$  :**

FWHM from normalized  $\delta I_{sx}$

**(relative) mode amplitude :**

normalized maximum  $\delta I_{sx}$

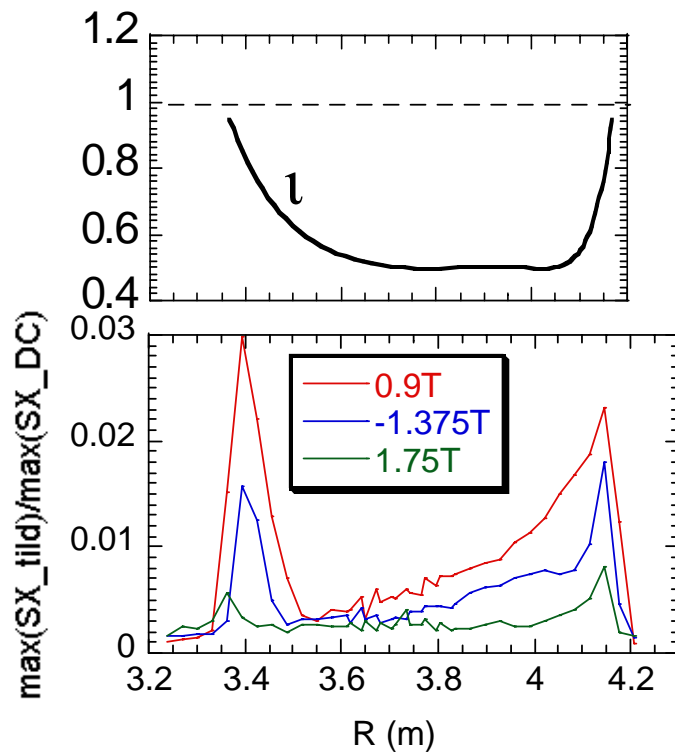
# Parameters characterizing the m/n=1/1 mode and confinement performance

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$b_{11}$ : RMS amplitude at the resonant surface based on the sheet current model where a sheet current is assumed on the resonant surface having the amplitude to reproduce experimental  $b$  on the inner wall surface

The H factor based on the ISS04 scaling as indication of confinement performance

$H_{\text{ISS04}}/H_{\text{ISS04at}^{**s}}$ : the ratio of the H factor estimated using instantaneous parameters to that in the absence of the m/n=1/1 mode at  $t^{**}$  using the ISS04 scaling.



## SXR amplitude lower than 0.3%

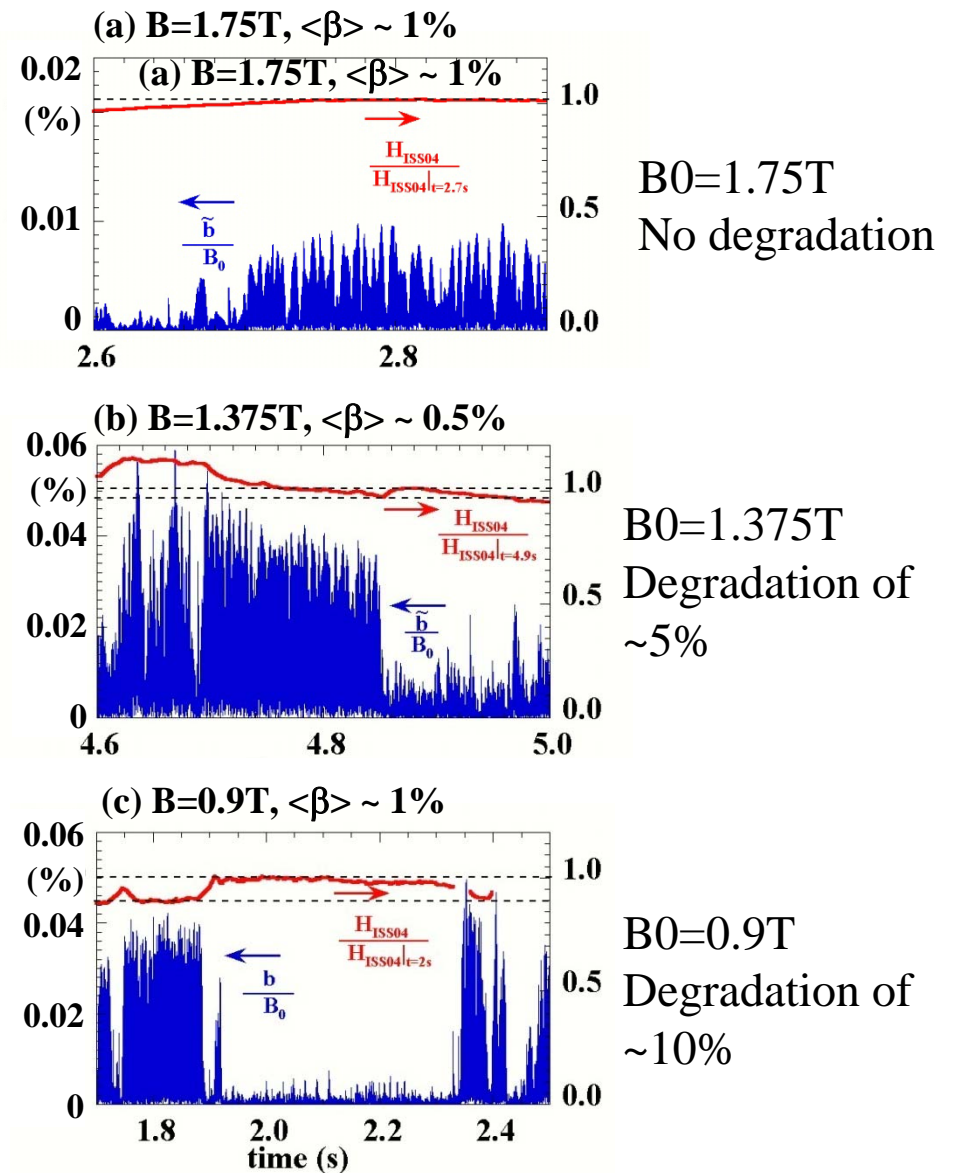
**=> no degradation**

### 1.5% mode amplitude

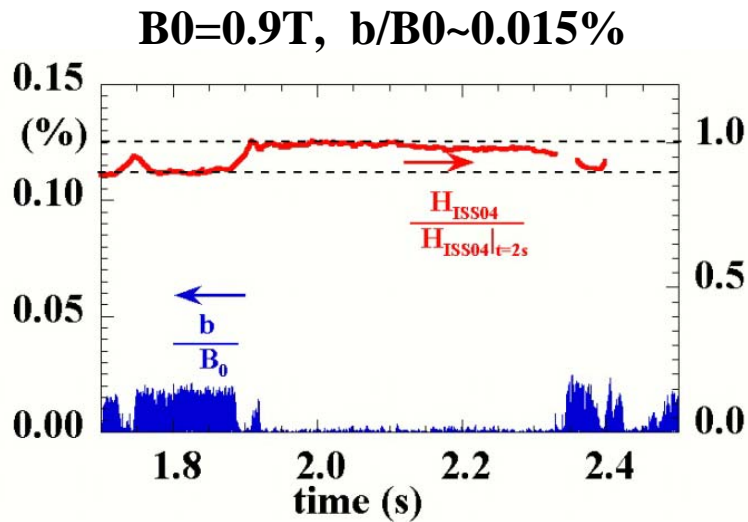
**=> ~5% degradation**

### 3% mode amplitude

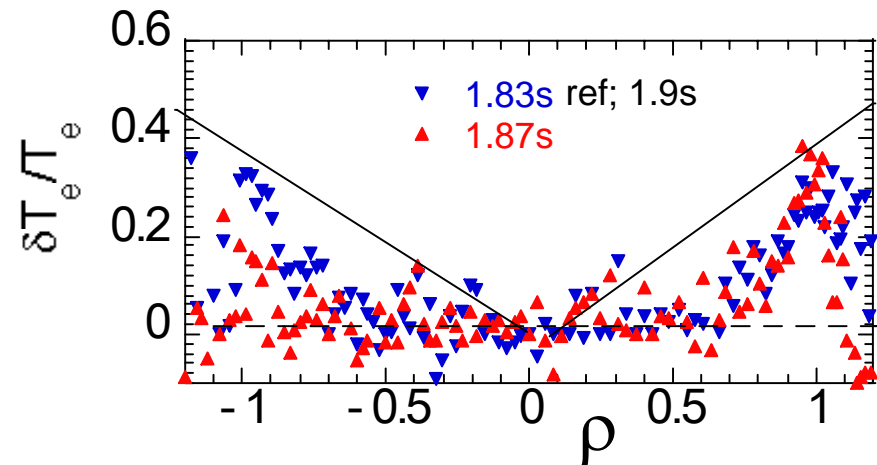
**=> ~10% degradation**



## Decrease in $T_e$ due to the $m/n=1/1$ mode is restricted only in the peripheral region



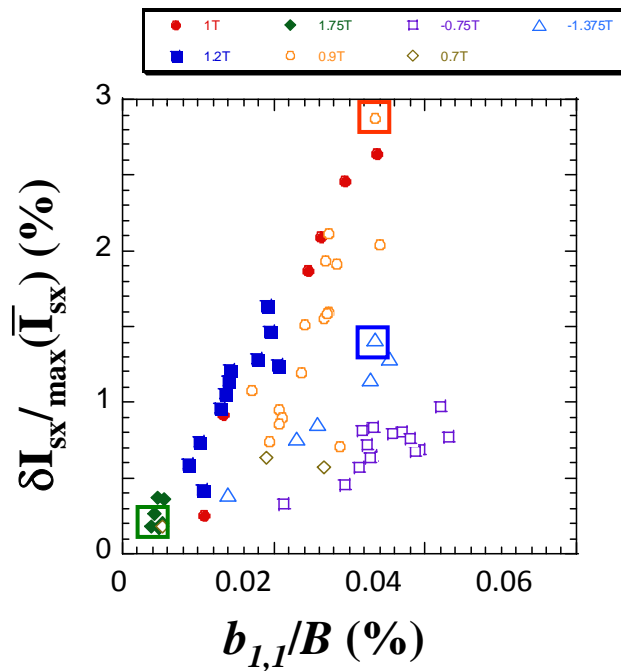
~10% degradation in H factor in the presence of the  $m/n=1/1$  mode



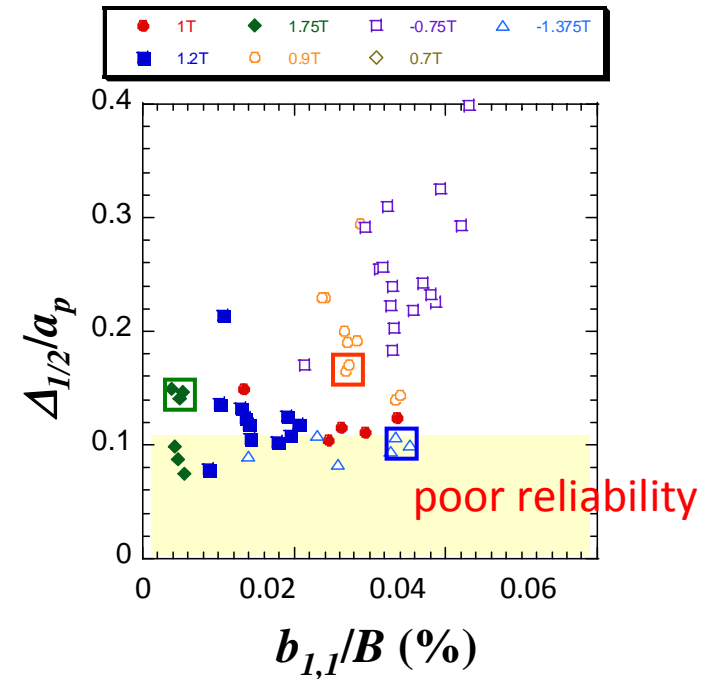
Decrease in  $T_e$  due to the presence of the  $m/n=1/1$  mode is restricted to the peripheral region  
=> decreased region corresponds to the mode width  
=> limited influence on the core region with high fusion power

# Internal structure of the m/n=1/1 mode, magnetic fluctuation level and their effect on confinement performance

Mode amplitude vs magnetic fluctuation level



Mode amplitude vs magnetic fluctuation level



- strong correlation between mode amplitude and confinement performance
- Effect on confinement performance becomes evident when the line-integrated mode width normalized to minor radius exceeds 10%

# Summary and future work

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## Summary:

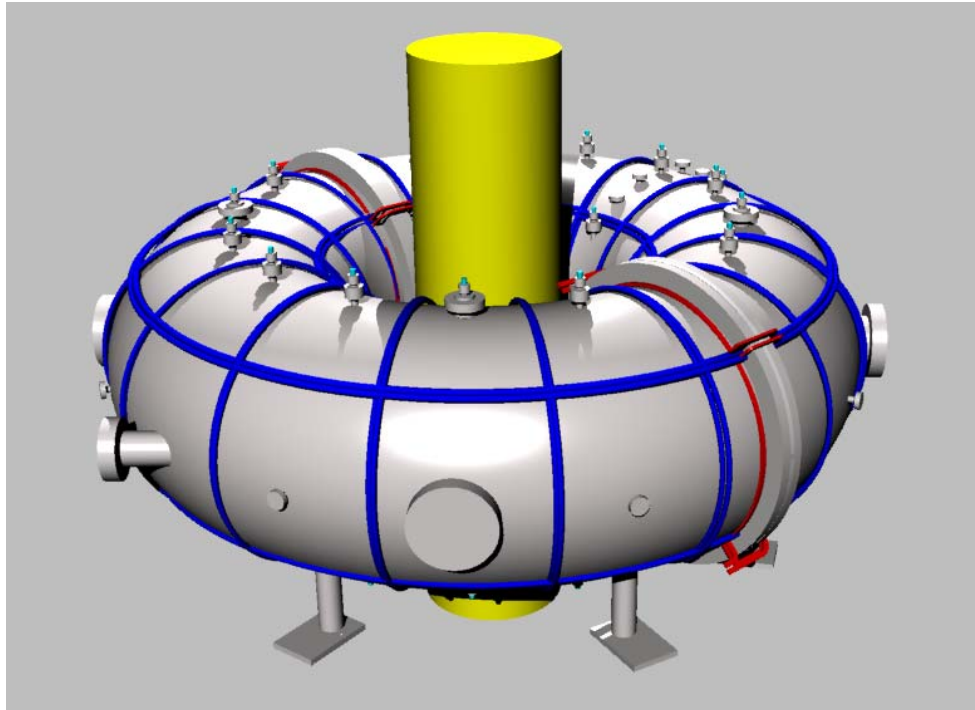
- **3-D self-organization to helical RFP is a general trend.**
- **In shallow-reversal region in RELAX, QSH with helical hot core has been attained without external control.**
- **The self-organization process in low-A RFP have been reproduced in 3-D MHD computations.**
- **In LHD, quantitative estimate of the effect of pressure driven mode on confinement has been performed. The mode is responsible for local change in temperature profile.**

## Future work in RELAX:

- **Feedback control of magnetic boundary – RWM stabilization**
- **Diagnostics – Thomson scattering for  $T_e$  measurement**
- **Supersonic gas injection for fuelling**

# Magnetic boundary control system in RELAX

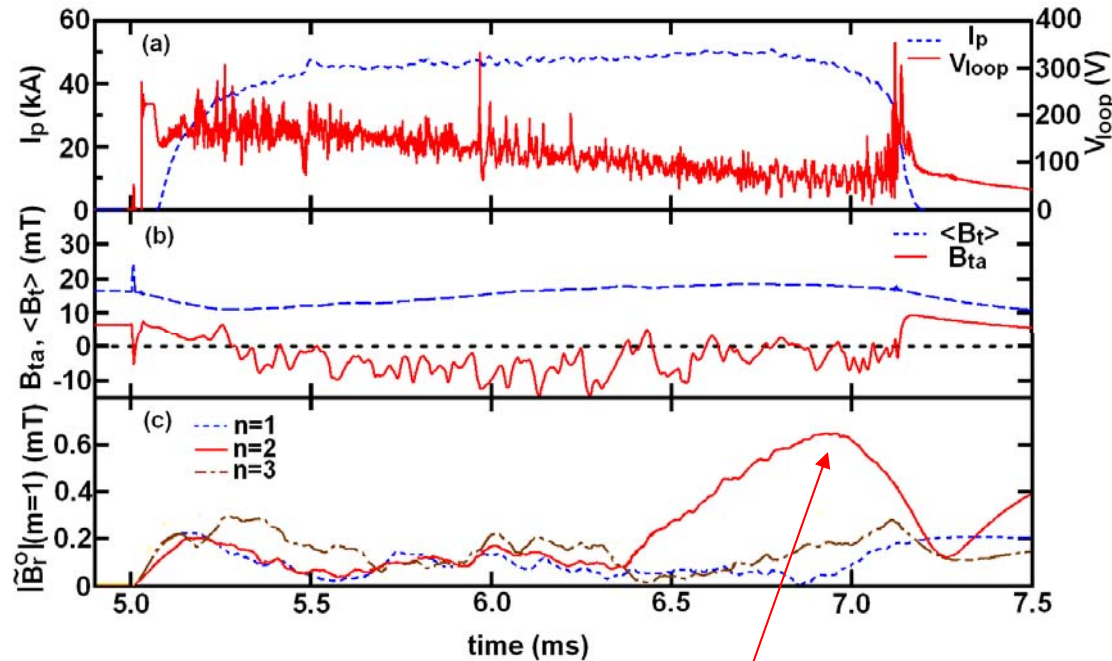
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- toroidal $\times 16$ , poloidal  $\times 4$   
saddle coil array installation  
almost completed for feedback  
control of RWM
- Construction of the power supplies  
in progress
- Introduction of digital feedback  
controllers in progress

backup

# RWM is problematic for longer pulse operation in RELAX



**Br (m=1/n=2) measured  
on the outer surface of  
the vacuum vessel**

$$\tilde{B}_r(a) / B_p(a) \approx 1 - 1.5\%$$

$$\Rightarrow I_p \text{ starts decreasing}$$

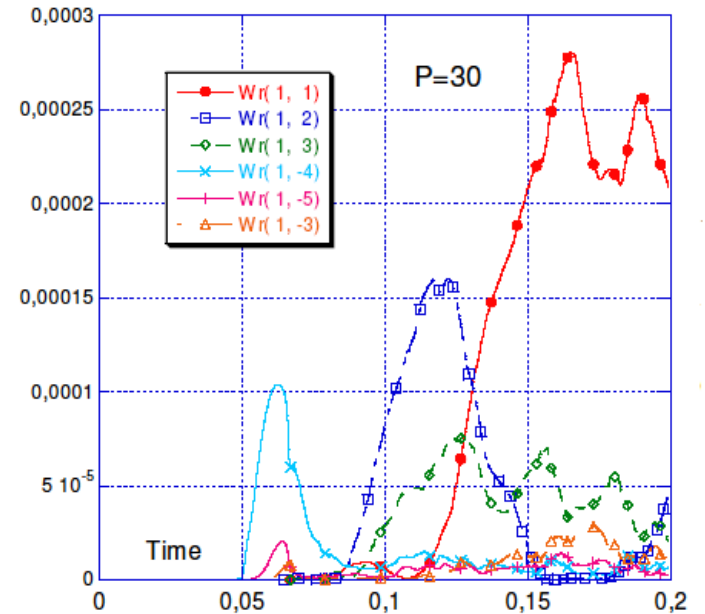


Fig.6: Radial mode energy vs. time for the RW simulation with  $\tau_w = 0.02$  ( $P=30$ ,  $S=3 \cdot 10^4$ ).

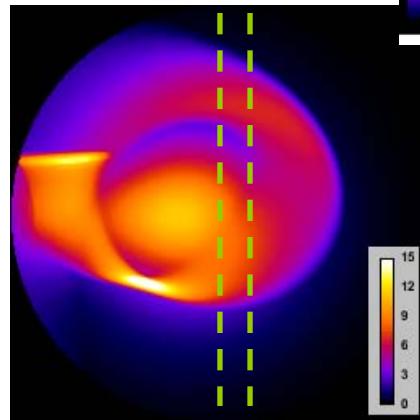
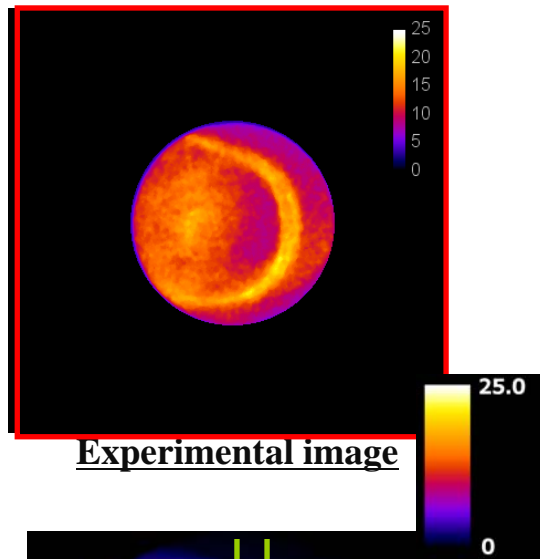
**3-D nonlinear MHD simulation predicts  
RWM in RELAX (Paccagnella, 2008)**

Initial growth of  $m=1/n=-4$  resonant mode,  
followed by growth of the non-resonant  
 $m=1/n=2$  external kink mode with growth  
time of resistive wall time constant

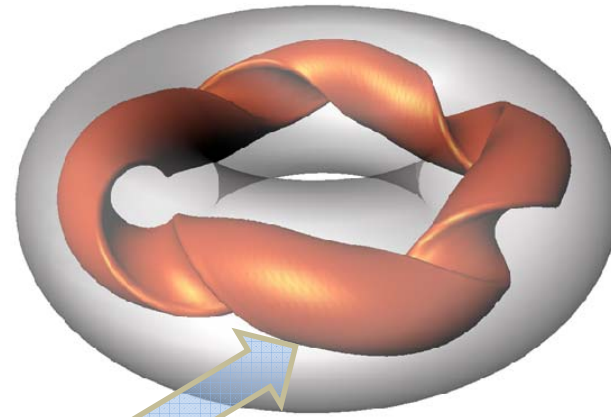
# Characteristic phenomena in shallow-reversal regions

## - SXR tangential images suggest helical hot core in QSH -

<  $m=1/n=4$  mode dominant case >

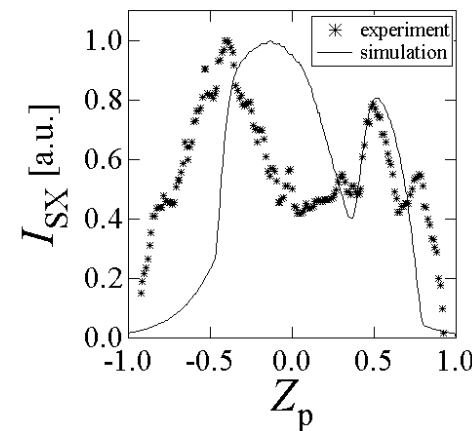
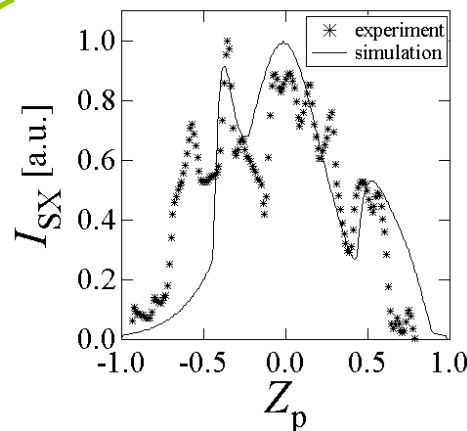


( $\alpha=3$ ,  $w=15\text{cm}$ ,  $r_{n=4}=14\text{cm}$ )



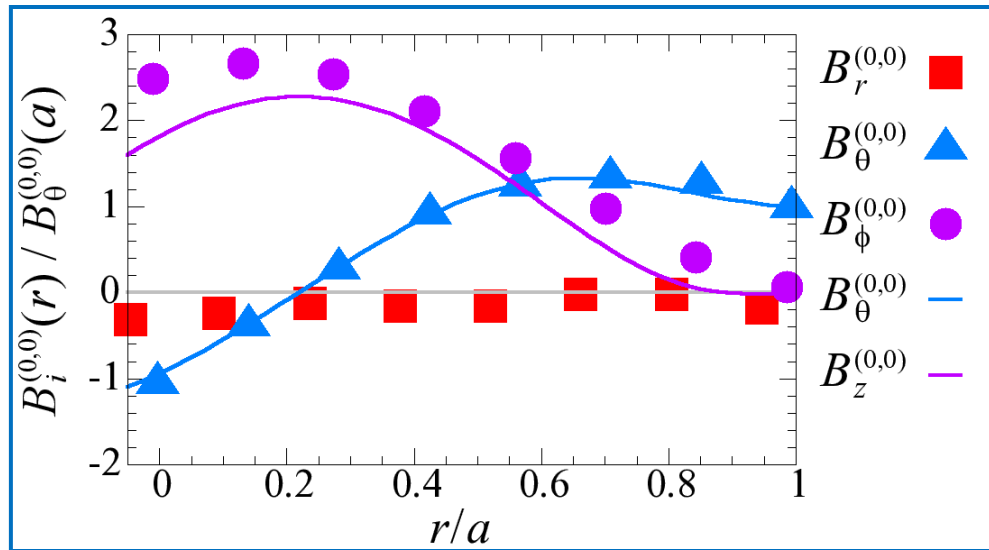
line-of-sight

helical deformation and line-of-sight of SXR pin-hole camera

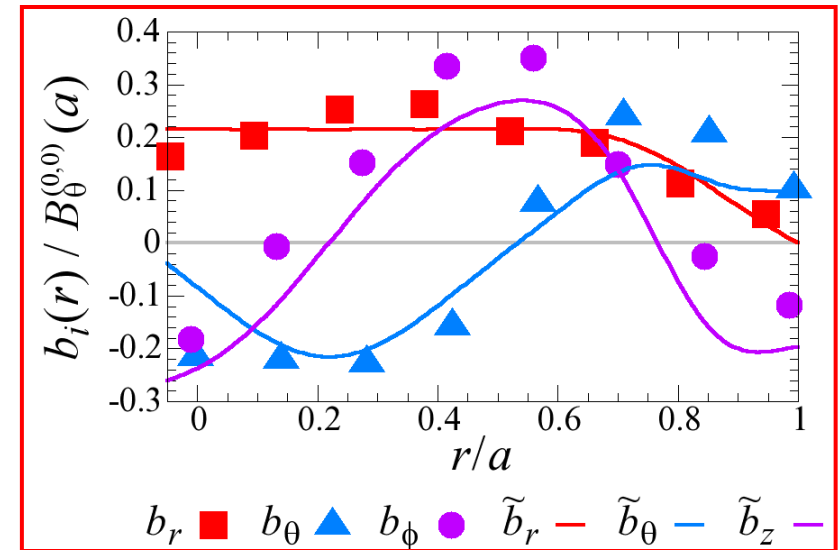


# Magnetic field profile in shallow-reversal regions is characterized by Helical Ohmic Equilibrium state

- **Symbols:** measured profiles with radial array of magnetic probes.
- **Solid lines:** Ohmic helical equilibrium solution (Paccagnella (2000)), details are in the next slide.
- Shafranov shift  $\Delta/a \sim 0.2$ .
- The helical structure rotates at a frequency of  $\sim 10$  kHz.



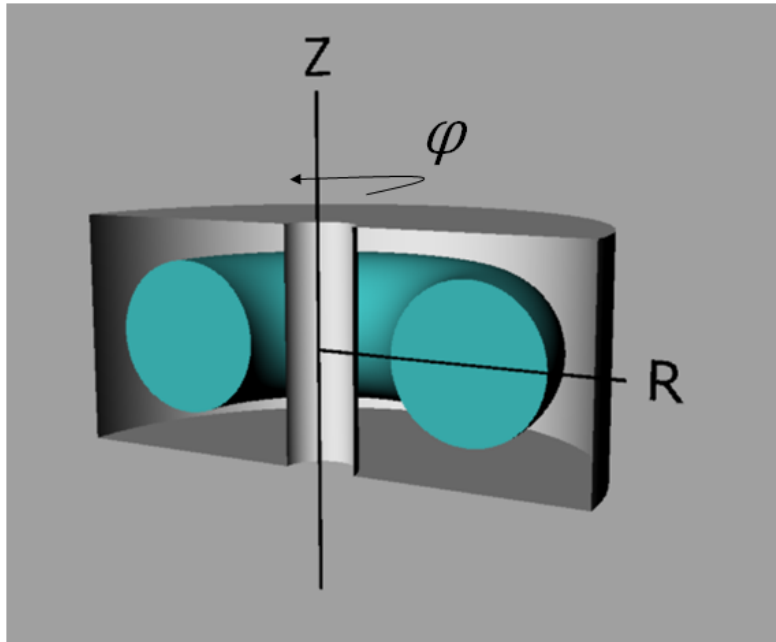
Axisymmetric components



$m=1$  helical components

# 3-D MHD computation for A=2 RFP

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## Cylindrical coordinate (R, $\phi$ , Z)

$$(N_R \times N_\phi \times N_Z) = (153 \times 128 \times 153)$$

## Boundary conditions

- Perfectly conducting and no-slip wall at all boundaries of the computation region
- $\mathbf{B}_\perp = \text{const.}$   $\mathbf{V} = 0$ ,  $\mathbf{j} = 0$  on the boundary

## Initial equilibrium

Based on reconstructed equilibrium from experiments using RELAXFit code (axisymmetric equilibrium)

## Initial perturbations

- Random white noise in velocity field is given at  $t=0$