### **RWM Control Code Maturity**

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EFCC is the fusion research arm of the United Kingdom Atomic Energy Authority



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### **RWM Control Code Maturity**

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### **Outline**

#### Introduction

#### Formulations behind advanced RWM control code

- Passive stabilisation: drift-kinetic-MHD hybrid
- Active control: 3D conductors + damping physics

#### **Examples of recent results**

- Kinetic effects from fast ions (DIII-D)
- New CarMa implementation: 3D walls + kinetic effects
- Simulation of RWM feedback in RFX

#### **Summary**



### Introduction: RWM code maturity issues

- Significant progress has been made in recent years towards realistic modelling of the RWM
- Two major issues have been attacked in developing RWM control codes
  - Passive control: mode damping physics (plasma rotation, drift-kinetic modeparticle resonances)
  - Active control: 3D structure of walls and feedback coils, realistic (and advanced) control logics
- Eventual goal is to have a code capable of solving above issues simultanously and consistently
- CarMa code is ready to perform this task.
- There are a few other RWM codes, with one or another advanced features. (AEGIS-K, LIGKA, MISK, MISHKA+HAGIS, VALEN, STARWALL, NMA,KINX,...)



### Introduction: MARS-F/K & CarMa codes

- MARS-F is single fluid, full MHD, full toroidal eigenvalue code with feedback options, including
  - Toroidal shear flow of plasma
  - Complete thin wall model (2D wall)
  - Continous description of control coils along toroidal angle
- MARS-K adds drift-kinetic terms into MHD in a self-consistent manner, including kinetic contributions from
  - Trapped and passing thermal particles [Liu PoP08]
  - Trapped fast ions (reported here)
- CarMa couples MARS-F/K with a 3D eddy current code CARIDDI
  - Coupling strategy is a key issue



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## Self-consistent formulation couples drift kinetic effects with fluid MHD via perturbed kinetic pressure tensors

#### MHD-kinetic hybrid formulation

$$(\gamma + in\Omega)\xi = \mathbf{v} + (\xi \cdot \nabla\Omega)R^2\nabla\phi,$$

$$\rho(\gamma + in\Omega)\mathbf{v} = (\nabla \cdot \mathbf{p} + \mathbf{j} \times \mathbf{B} + \mathbf{J} \times \mathbf{Q} - \rho \left[2\Omega\hat{\mathbf{Z}} \times \mathbf{v} + (\mathbf{v} \cdot \nabla\Omega)R^2\nabla\phi\right],$$

$$(\gamma + in\Omega)\mathbf{Q} = \nabla \times (\mathbf{v} \times \mathbf{B}) + (\mathbf{Q} \cdot \nabla\Omega)R^2\nabla\phi,$$

$$(\gamma + in\Omega)p = -\mathbf{v} \cdot \nabla P,$$

$$\mathbf{j} = \nabla \times \mathbf{Q},$$

$$\mathbf{p} = p\mathbf{I} + p\|\hat{\mathbf{b}}\hat{\mathbf{b}} + p\|(\mathbf{I} - \hat{\mathbf{b}}\hat{\mathbf{b}}),$$

$$\mathbf{thermal} \qquad \mathbf{fast ion}$$

$$p\|e^{-i\omega t + in\phi} = \sum_{e,i} \int d\Gamma M v_{\parallel}^2 f_{N_t}^1(\xi_{\perp}) + \int d\Gamma M v_{\parallel}^2 f_{N_t}^1(\xi_{\perp}),$$

$$p_{\perp}e^{-i\omega t + in\phi} = \sum_{e,i} \int d\Gamma \frac{1}{2} M v_{\perp}^2 f_{N_t}^1(\xi_{\perp}) + \int d\Gamma \frac{1}{2} M v_{\perp}^2 f_{N_t}^1(\xi_{\perp})$$



### Equilibrium distribution functions of particles

- Assume Maxwellian distribution for thermal particles  $f_{th}^0 = N \left(\frac{M}{2\pi T}\right)^{3/2} e^{-\epsilon_k/T}$
- And slowing-down distribution for fast ions  $\longrightarrow f_h^0 = \begin{cases} \frac{C}{\varepsilon_k^{3/2} + \varepsilon_c^{3/2}} & 0 < \varepsilon_k < \varepsilon_0 \\ 0 & \varepsilon_k > \varepsilon_0 \end{cases}$
- Same assumptions made in MISK code
- Radial profile of hot ion's energy distribution determined given hot ion's density and pressure profiles

$$\varepsilon_0 = \begin{cases}
3.5 \text{MeV} & \alpha - \text{particles} \\
\varepsilon_h(\psi) & \text{hot ions}
\end{cases}$$



### Perturbed distribution functions of particles

- Perturbed distribution functions for thermal and fast ions derived by solving linear perturbed drift kinetic equation, following [Antonsen PF82] and [Porcelli PoP94].
- Neglected finite drift orbit width and FLR
- Consider various resonances between mode and particles
- Thermal particles:
  - Bounce resonance of ions
  - Precession drift of ion&electron
  - Trapped and passing particles
- Fast ions:
  - Precession drift of trapped ions

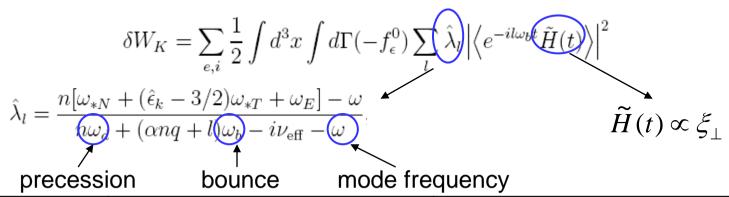
$$f_{th}^{1} \propto \frac{n[\omega_{*N}^{th} + (\hat{\varepsilon}_{k} - 3/2)\omega_{*T}^{th} + \omega_{E}] - \omega}{n\omega_{d}^{th} + (\alpha nq + l)\omega_{b}^{th} - i\nu_{\text{eff}} - \omega}$$

$$f_h^1 \propto \frac{n(f_{h,\psi}^0/Zef_{h,\varepsilon}^0) - \omega}{n\omega_d^h - \omega}$$



### MARS-K includes both perturbative & selfconsistent approach

Drift kinetic energy perturbation [Antonsen82, Porcelli94]

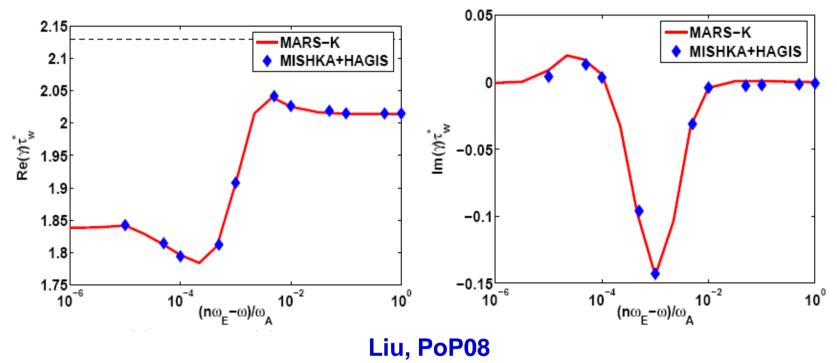


	perturbative	self-consistent
eigen-function $\xi_{\perp}$	ideal-kink or fluid RWM	modified by kinetic effects self- consistently
eigen- frequency $\omega$	$\omega = 0$ or $i\gamma_f$	$\omega = i\gamma$ monlinear eigenvalue formulation
other damping	shear Alfven damping included a-posteriori	continuum damping included via MHD terms



## MARS-K perturbative approach benchmarked against MISHKA+HAGIS code

- Choose a Soloviev equilibrium with circular-like shape
- Both codes run with perturbative approach



Validates approximation of neglecting banana width for kinetic RWM



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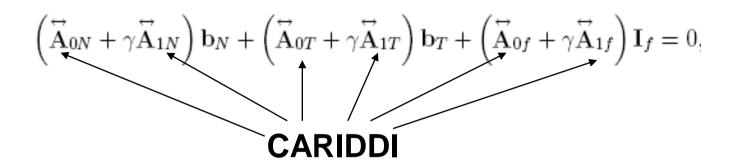
### Formulations behind CarMa code

- CarMa couples MARS-F/K to 3D eddy current code CARIDDI, via two coupling schemes
  - "Forward" coupling scheme
  - "Backward" coupling scheme
- Key ideas of forward coupling [Portone PPCF08]
  - Introduce a coupling surface just outside plasma surface
  - Replace plasma magnetic response by a virtual surface current response at coupling surface
  - Effectively "condense" MHD equations into eddy current equations
- Forward coupling resembles in some aspects VALEN multimode approach. Like VALEN and STARWALL, consider static plasma response, hence no plasma inertia, (fast) rotation, kinetic effects



### Formulations behind CarMa code

- Backward coupling scheme "condenses" conductors' (dynamic) response into MHD equations
  - as a non-trival boundary condition at coupling surface just outside plasma.
- Key idea is to transform conductors' dynamic response into a linear BC [Smith PoP08, Guazzotto PoP08, Liu PoP08]
  - Hugely beneficial for linear MHD eigenvalue formulation
  - Allows easy inclusion of all effects beyond ideal MHD in CarMa





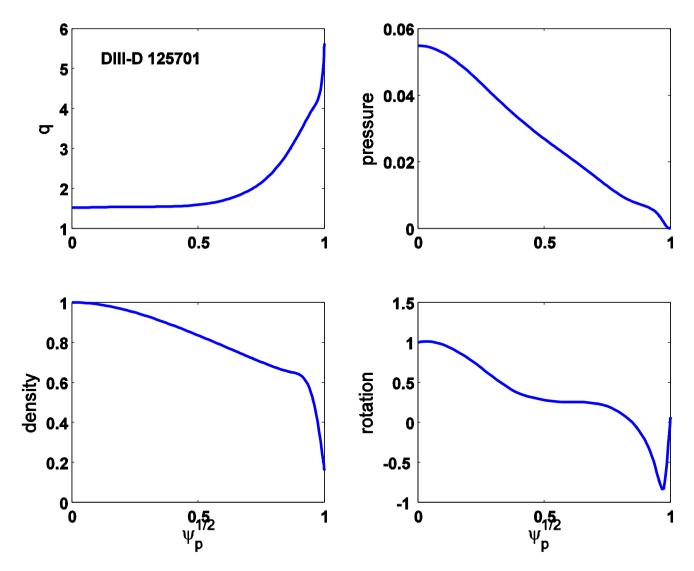
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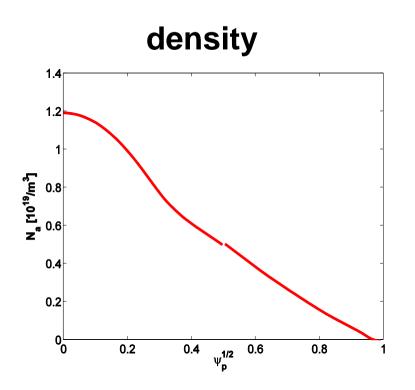


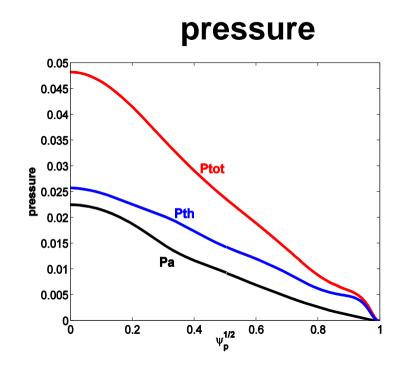
### **Equilibrium profiles from DIII-D 125701**





### Fast ion density and pressure profiles

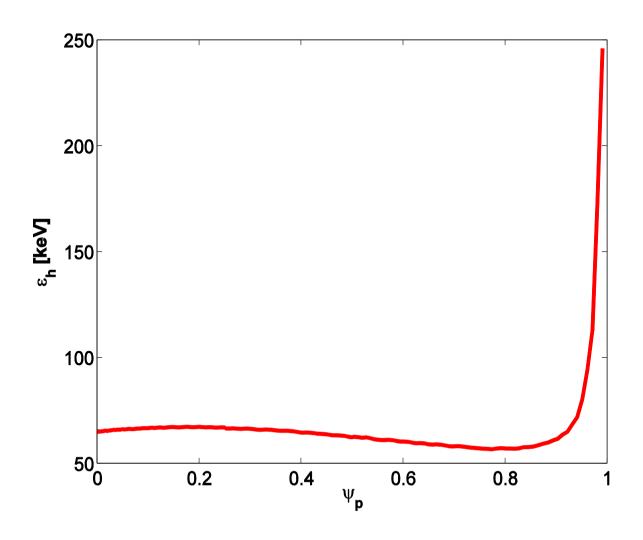




Fast ion contributes nearly half of total pressure

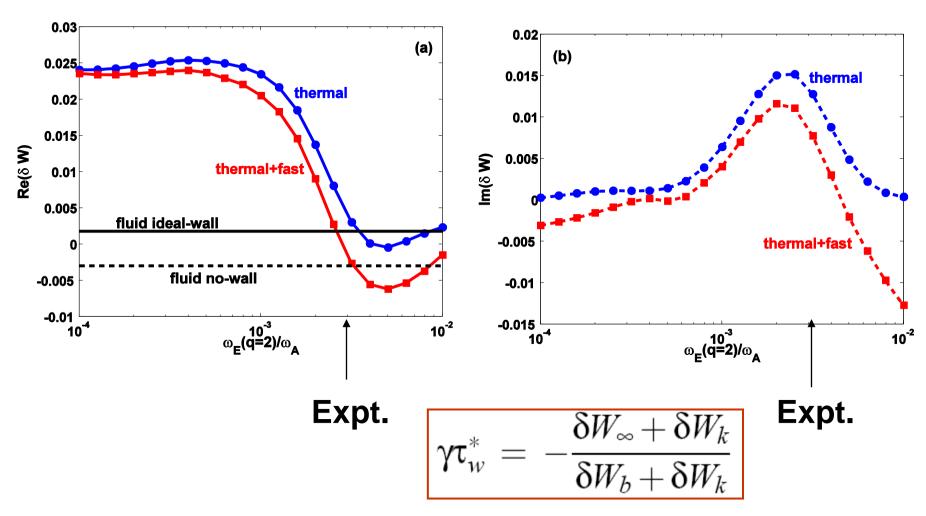


## Computed radial distribution of fast ion energy





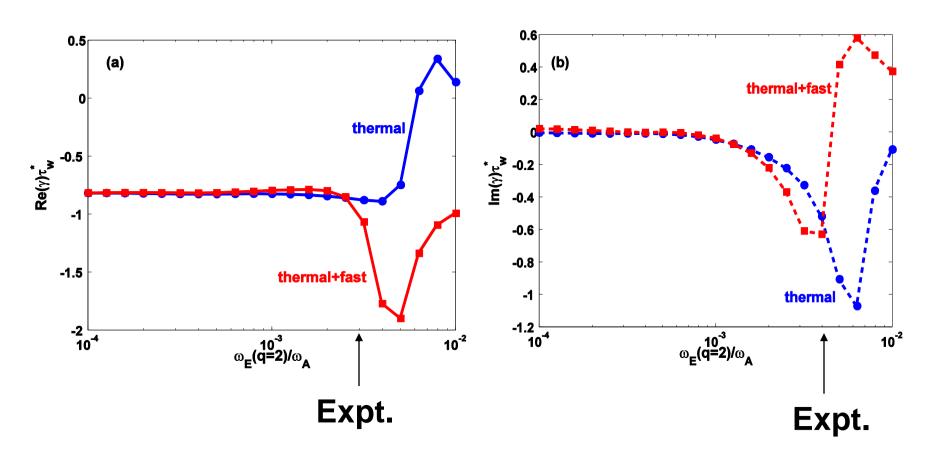
### Perturbative approach: energy



Fast ions give minor contribution at slow rotation



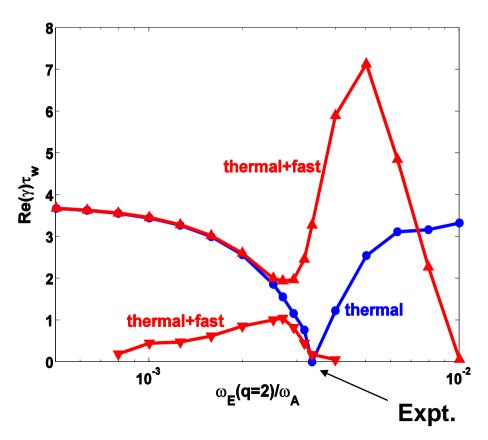
### Perturbative approach: eigenvalue



- Fast ions are slightly destabilising at slow rotation, stabilising at faster rotation
- RWM at expt. rotation is stable w/ or w/o fast ions



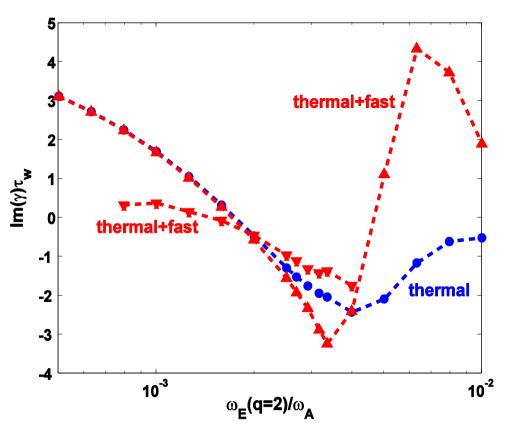
### Self-consistent approach: growth rate



- w/o fast ions, SC approach predicts a narrow window of stabilisation
- w/ fast ions, a second unstable branch appears



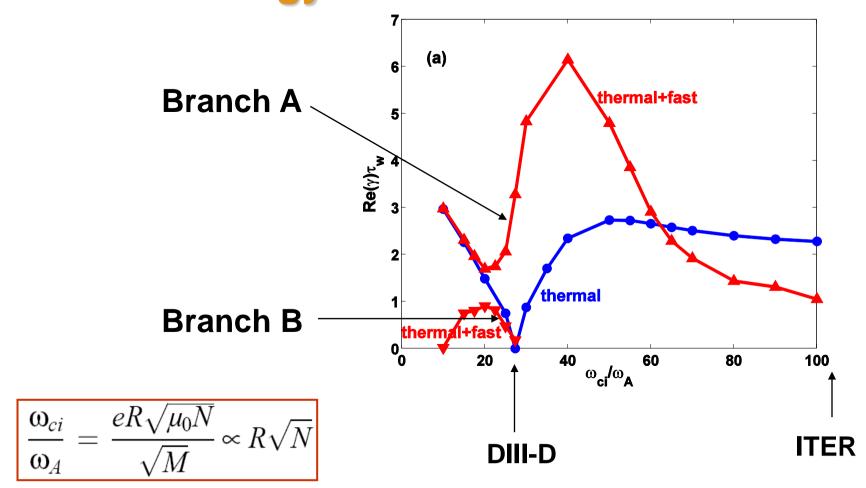
### Self-consistent approach: mode frequency



 Fast ion driven (unstable) branch can have different direction of mode rotation



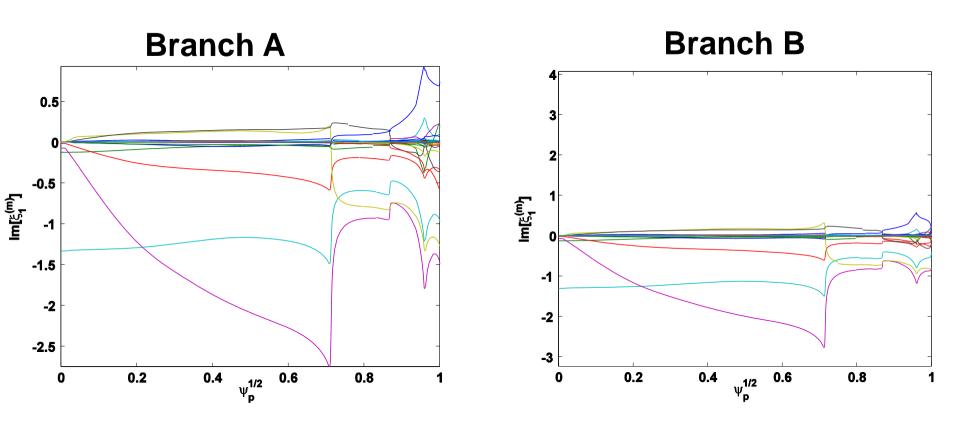
## Self-consistent approach: sensitivity on ion gyroradius



Fast ions can be stabilising in ITER



## Self-consistent approach: compare eigenmode structure

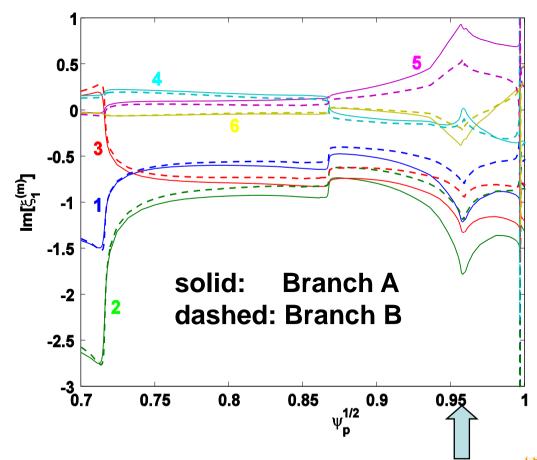


 Less unstable mode (branch B) more peaked at last rational surface



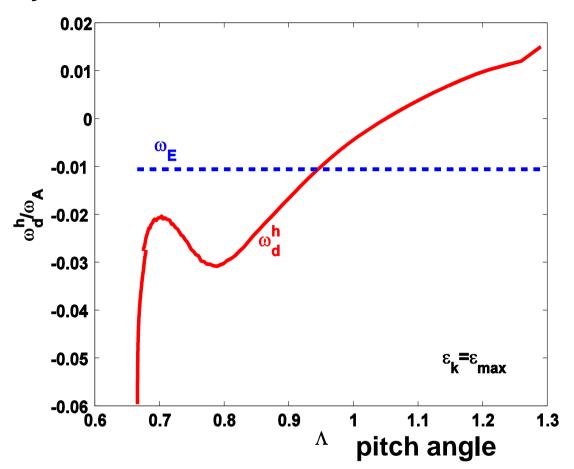
## Self-consistent approach: compare eigenmode structure

 Eigenmodes differ mostly in the outer region of the plasma, outside the q=2 surface



### Resonance condition for fast ions?

Fast ion precession drift frequency comparable to experimental rotation frequency





### **Outline**

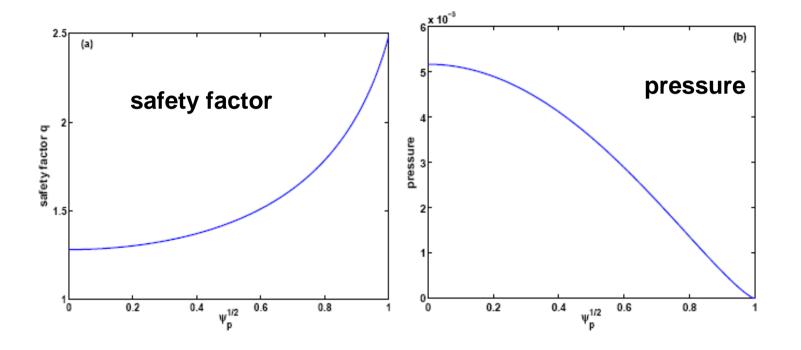
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## CarMa capable of solving RWM problems with damping + 3D wall + feedback

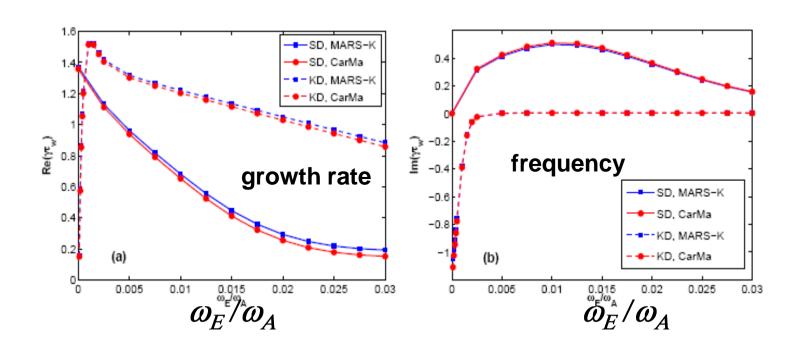
Choose a test equilibrium with circular cross section, and R/a=5





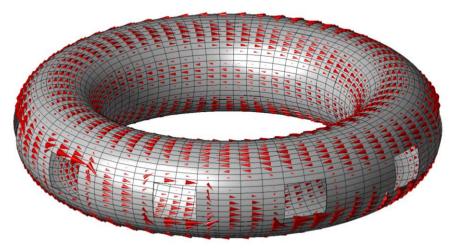
### CarMa benchmarked for damping models

- Compute passive stability versus plasma rotation
  - Parallel sound wave damping
  - kinetic damping
- With 2D wall, CarMa results almost identical to MARS-K



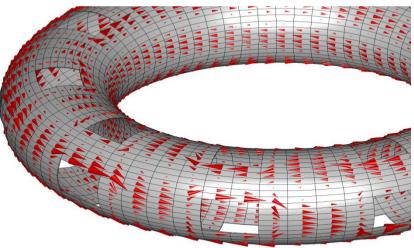


### Models of 3D wall with holes



**Model A** 

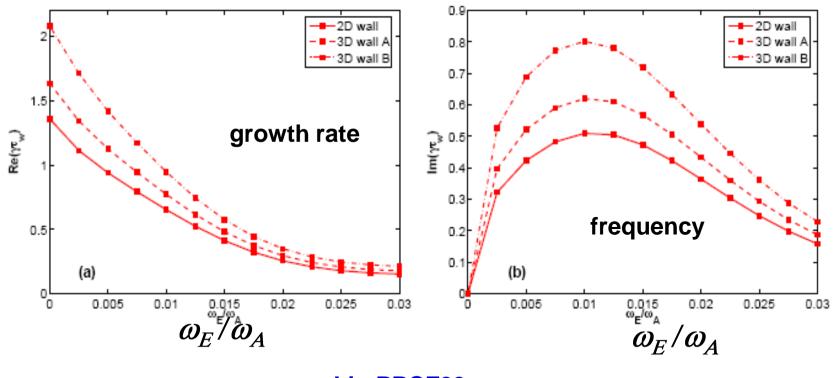
**Model B** 





## Wall holes increase both mode growth rate and frequency

 Effect of 3D wall on RWM passive stability, with plasma rotation and strong sound wave damping

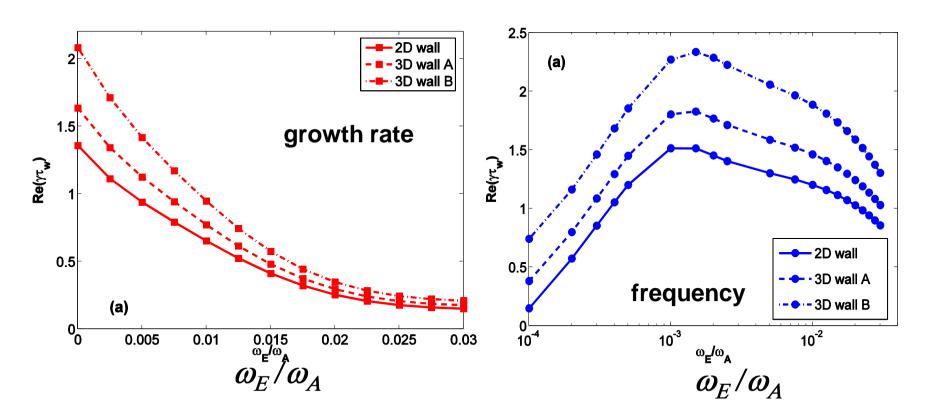




Liu PPCF09

## Wall holes increase both mode growth rate and frequency

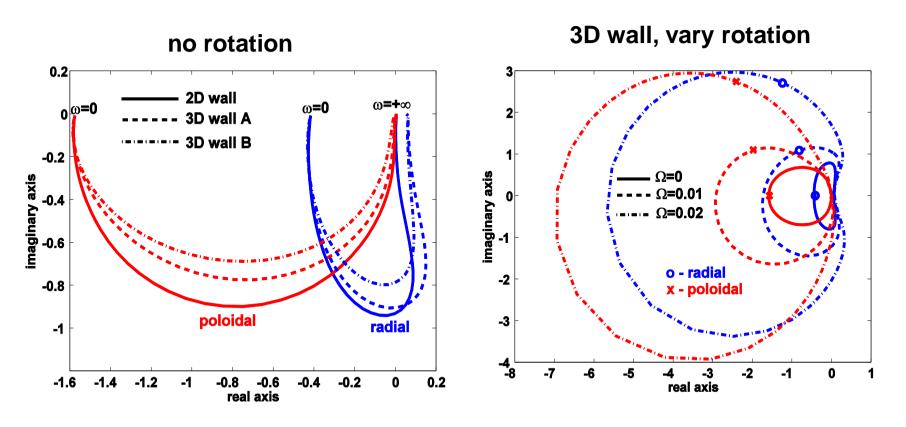
 Effect of 3D wall on the RWM passive stability with plasma rotation and drift-kinetic damping model





## Nyquist diagrams show effects of 3D wall on feedback performance

- 3D wall reduces closed loop stability margin
- Synergy between rotation and feedback, in presence of 3D wall



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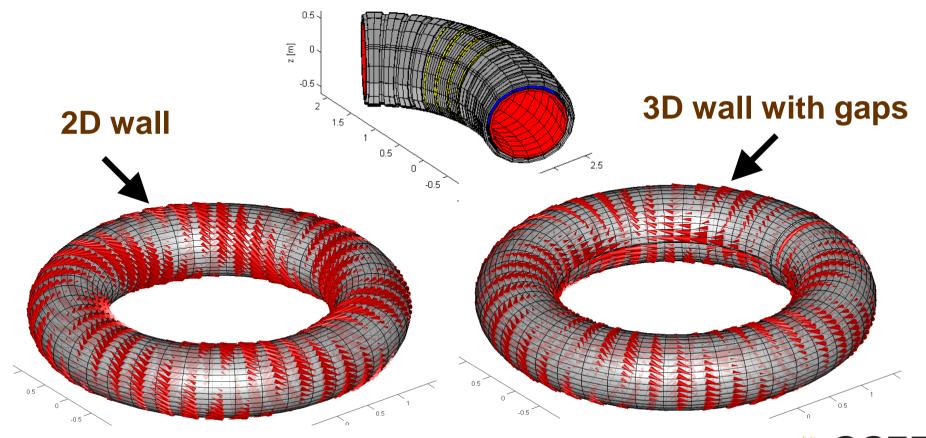
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## CarMa used for realistic modelling of RFX feedback experiments

- CarMa computed wall eddy current pattern
- For an unstable mode with n=3,m=1

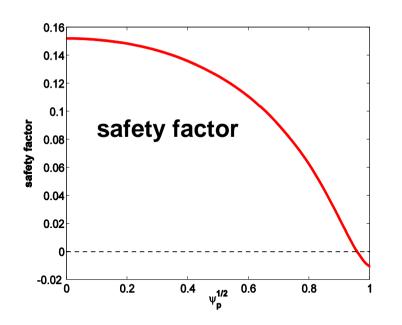
**Villone PRL08** 

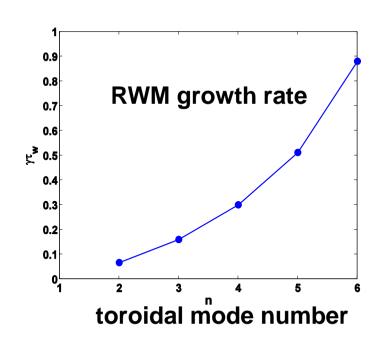




## CarMa used for realistic modelling of RFX feedback experiments

 A typical RFX plasma equilibrium with shallow field reversal: THETA=1.426, F=-0.060

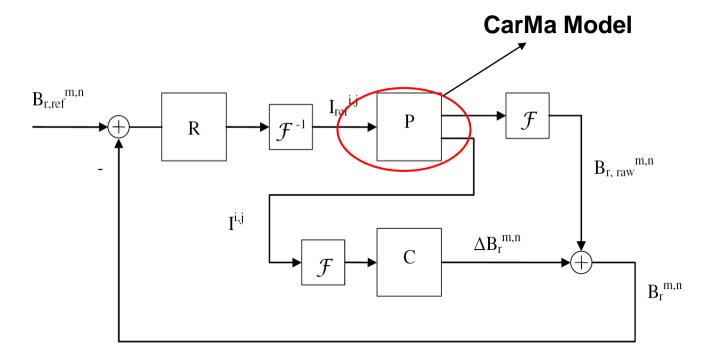






### Feedback loop in RFX experiments

 CarMa has been used to represent the plant in a full closed loop simulation of the RWM control system, yielding closed loop growth rates, that are directly compared with experiments



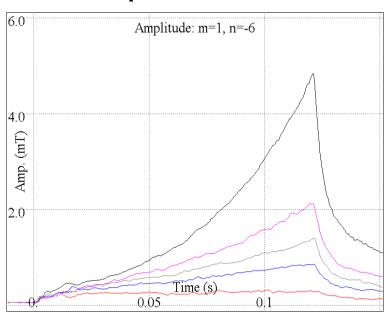
**Marchiori EPS09** 



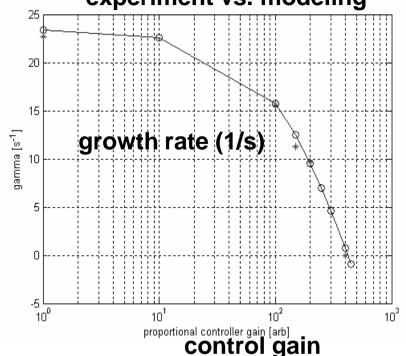
### RWM feedback modelling in RFX

 Comparing experimental vs. CarMa simulated closed-loop growth rates of RWM

#### expriment



#### experiment vs. modeling

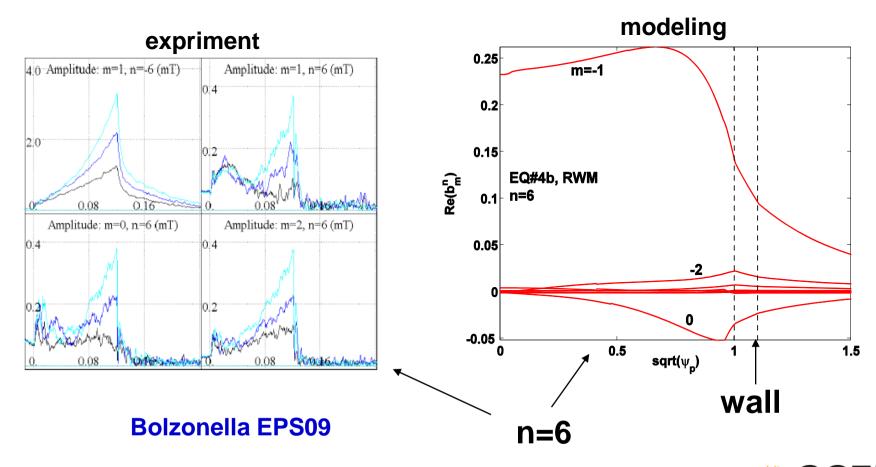


**Marchiori EPS09** 



## Eigenmode structure can also be compared between modeling and expt.

 MARS-F computed amplitude ratio between poloidal harmonics (due to toroidicity) match reasonably well experiments





### **Summary**

- Self-consistent approach generally predicts less kinetic stabilisation on RWM compared to perturbative approaches
- Fast ions contribution, whilst can give additional stabilisation in perturbative approach, can also drive an extra unstable branch in selfconsistent approach
- A new version of CarMa code allows consistent modelling of 3D conductors, kinetic effects and feedback stabilisation all together
- The RFX feedback experiments well modelled by CarMa
- Other aspects of RWM code maturity, not discussed here
  - Capability of treating thick walls and ferromagnetic materials
  - Multimodal analysis (e.g. n=0+1, n=1+2+3, etc.) ...



### **Discussions**

#### Are RWM codes mature enough? (No)

- MARS-K: finite banana orbit effects. FLR. collision?
- AEGIS-K, LIGKA more advanced in kinetic physics
- Nonlinear physics (RWM couple to other modes)?
- CarMa: iron saturation (nonlinear) ?

#### MARS-K results vs. DIII-D experiment?

- Finite banana effects? FLR?
- Collision?
- Fast particle distribution function? Profiles?
- Evidence of fast ion driven RWM in other experiments?

#### What about prediction for ITER?

- Seems no complete stabilisation with thermal particles
- MARS-K modelling with alpha particles in progress
- Waiting for new reference plasma for steady state scenario...

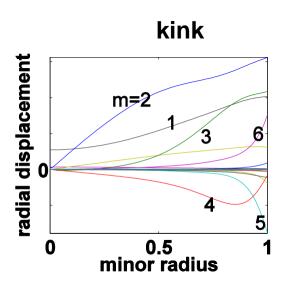


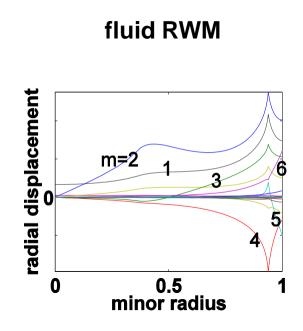
## Backup slides ...

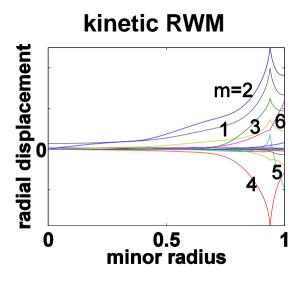


## Kinetic effects do modify mode eigenfunction for a test toroidal equilibrium

• Test toroidal Soloviev equilibrium: R/a=3, = 1.6  $\kappa$ 

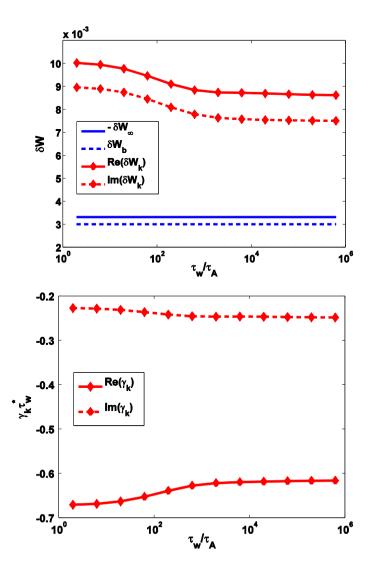




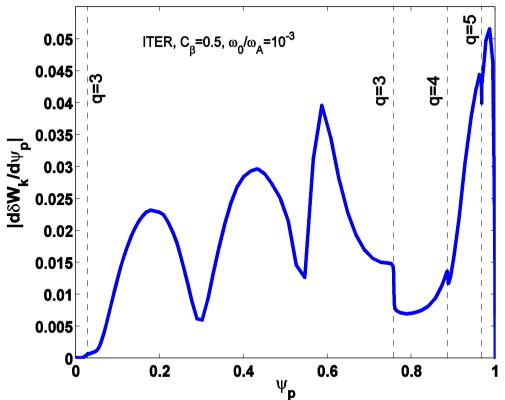




### ITER: perturbative



 Radial distribution of drift kinetic energy perturbation with precessional resonances





## CarMa allows multimodal analysis Multimodal ITER results /2 Combined n=0 (VDE) + n=1 (external kink) plasma evolution

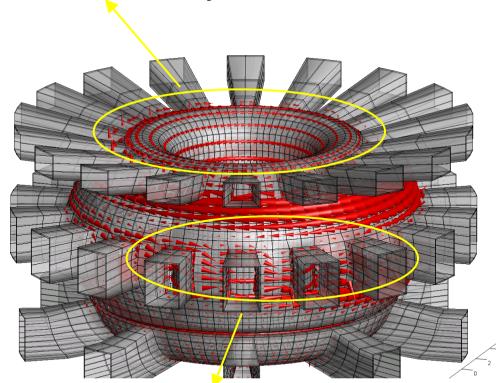
ITER Scenario 4 ( $\beta_N$ =2.7). Three unstable eigenvalues:

•  $\gamma_{1,2}$  = 10.4 s<sup>-1</sup>, eigenvectors with a n=1 spatial behaviour, shifted of  $\pi/2$  toroidally

•  $\gamma_3 = 7.66 \text{ s}^{-1}$ , eigenvector with a n=0 spatial behaviour

Almost identical to those computed with monomodal models: little (open-loop) coupling of n=0 and n=1 modes due to 3D structures.

### Predominantly n=0



**Predominantly n=1** 

