Investigation of the Complex Relationship Between Plasma Rotation and Resistive Wall Mode Stabilization in NSTX

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US-Japan Workshop on
MHD Control, Magnetic Islands and Rotation
November 24, 2008
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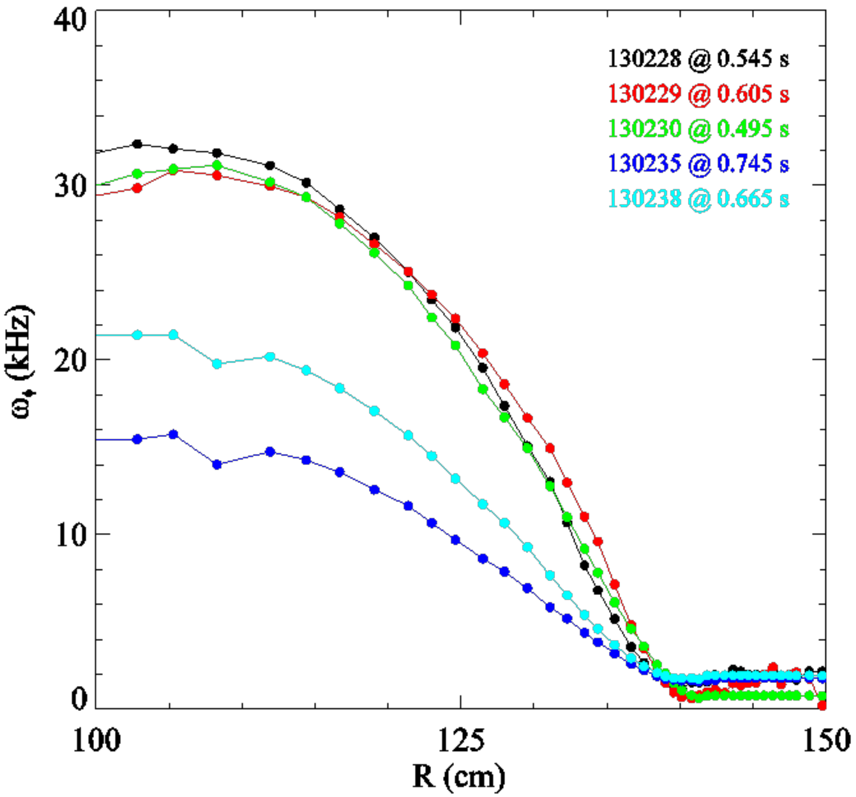
Resistive Wall Mode stabilization in NSTX may be explained by kinetic theory

- Motivation
  - The relationship between plasma rotation and RWM stability in NSTX is more complex than simple models suggest, and kinetic theory has the potential to explain it.
- Outline
  - The MISK code is used to calculate RWM stability.
  - Kinetic theory matches NSTX experimental results of instability at moderate rotation.
  - Kinetic theory can match the evolution of a shot.
  - DIII-D results suggest the importance of hot ions.

The relationship between rotation and RWM stability is not straightforward in NSTX



Examples of plasma rotation profiles at the time of RWM instability.

- RWM observed in NSTX with a variety of plasma rotation profiles.
  - Including with ω = 0 at the q = 2 surface.
  - This does not agree with "simple" theories.
  - What can kinetic theory tell us about the relationship between rotation and stability?

The Modifications to Ideal Stability by Kinetic Effects (MISK) code

- Written by Bo Hu, University of Rochester
  - Hu, Betti, PRL, 2004 and Hu, Betti, and Manickam, POP, 2005
- Uses a perturbative calculation, with marginal stability eigenfunction from the PEST code.
- Calculation of δW_K includes the effects of:
  - Trapped Ions
  - Trapped Electrons
  - Circulating Ions
  - Alfven Layers
  - Hot Ions

$$\gamma_K \tau_w = -\frac{\delta W_\infty + \delta W_K}{\delta W_b + \delta W_K}$$

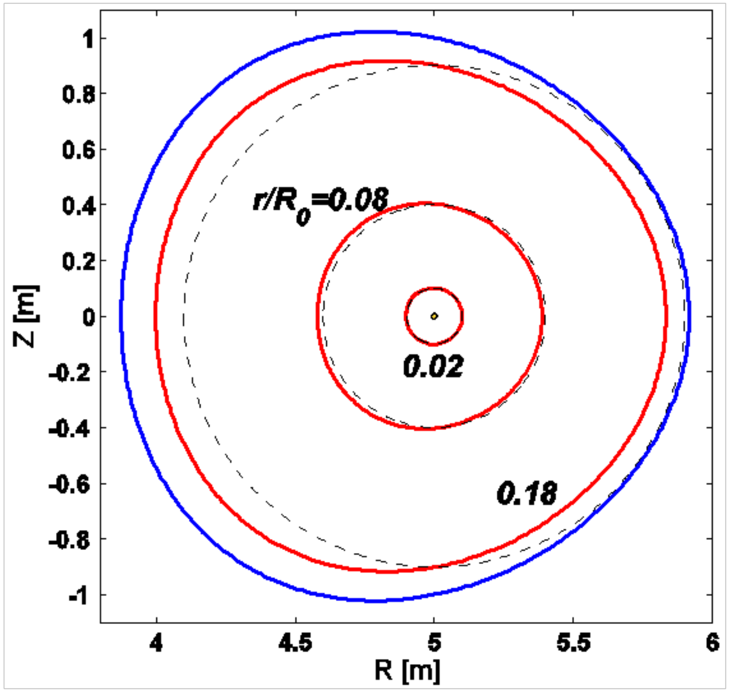
PEST (δW_b) MISK (δW_K)

(Hu, Betti, and Manickam, PoP, 2005)

MISK and MARS-K were benchmarked using a Solov'ev equilibrium

$$\mu_0 P(\psi) = -\frac{1+\kappa^2}{\kappa R_0^3 q_0}\psi, \quad F(\psi) = 1$$

$$\psi = \frac{\kappa}{2R_0^3 q_0}\left[\frac{R^2 Z^2}{\kappa^2} + \frac{1}{4}\left(R^2 - R_0^2\right)^2 - a^2 R_0^2\right]$$



(Liu, ITPA MHD TG Meeting, Feb. 25-29, 2008)

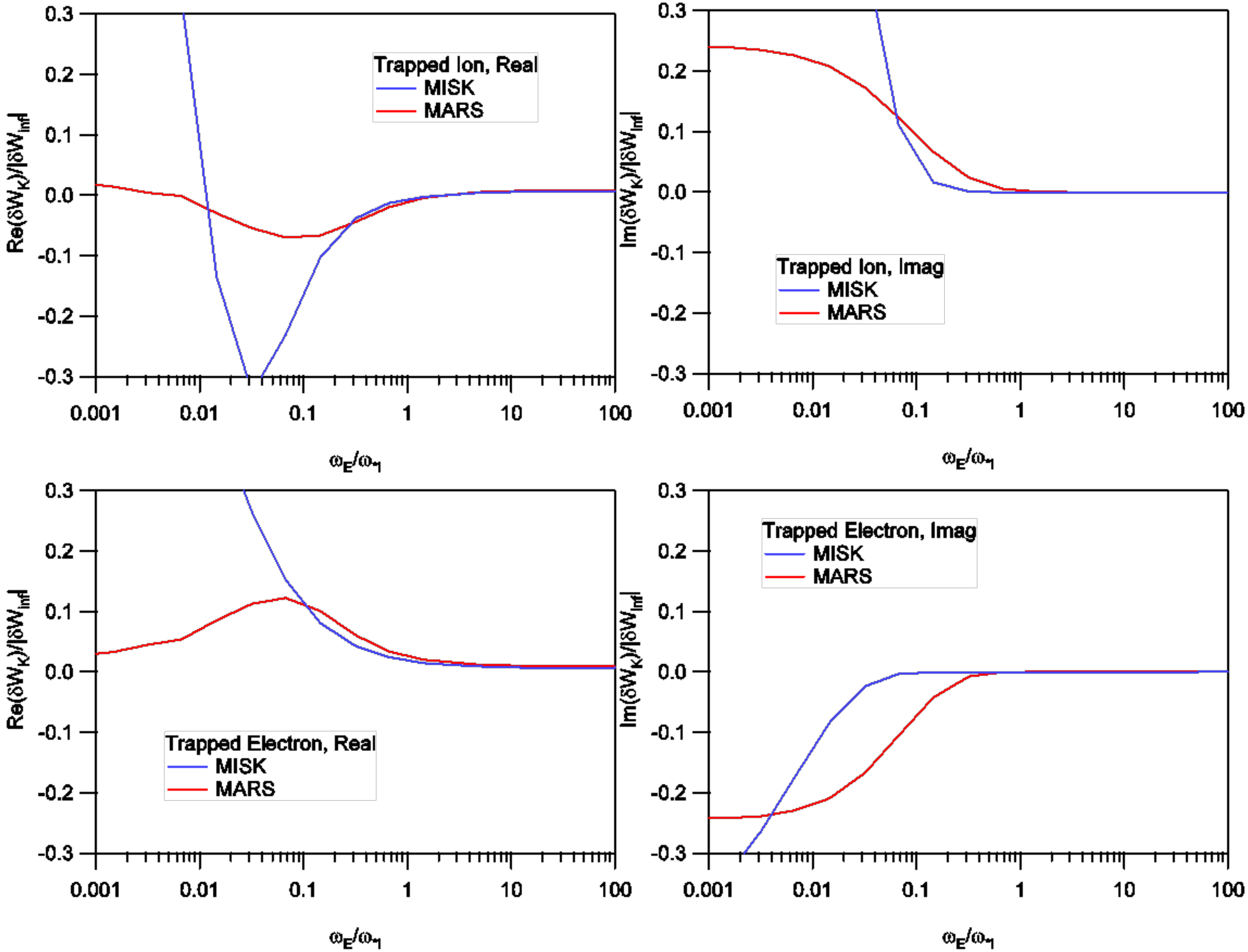
- Simple, analytical solution to the Grad-Shafranov equation.
- Flat density profile means $\omega_{*N} = 0$.
- Also, ω, γ, and ν_{eff} are taken to be zero for this comparison, so the frequency resonance term is simply:

$$\delta W_K \propto \int \left[\frac{\left(\hat{\varepsilon} - \frac{3}{2}\right)\omega_{*T} + \omega_E}{\langle\omega_D\rangle + l\omega_b + \omega_E}\right]\hat{\varepsilon}^{\frac{5}{2}} e^{-\hat{\varepsilon}} d\hat{\varepsilon}$$

MISK and MARS-K match well at reasonable rotation

Good match for trapped ions and electrons at high rotation, but poor at low rotation.

The simple frequency resonance term denominator causes numerical integration problems with MISK that don't happen with realistic equilibria.



The dispersion relation can be rewritten in a convenient form for making stability diagrams

The kinetic contribution has a real and imaginary part, so:

$$Re(\gamma_K \tau_w) = -\frac{\delta W_\infty \delta W_b + (Im(\delta W_K))^2 + Re(\delta W_K)(\delta W_\infty + \delta W_b + Re(\delta W_K))}{(\delta W_b + Re(\delta W_K))^2 + (Im(\delta W_K))^2}$$

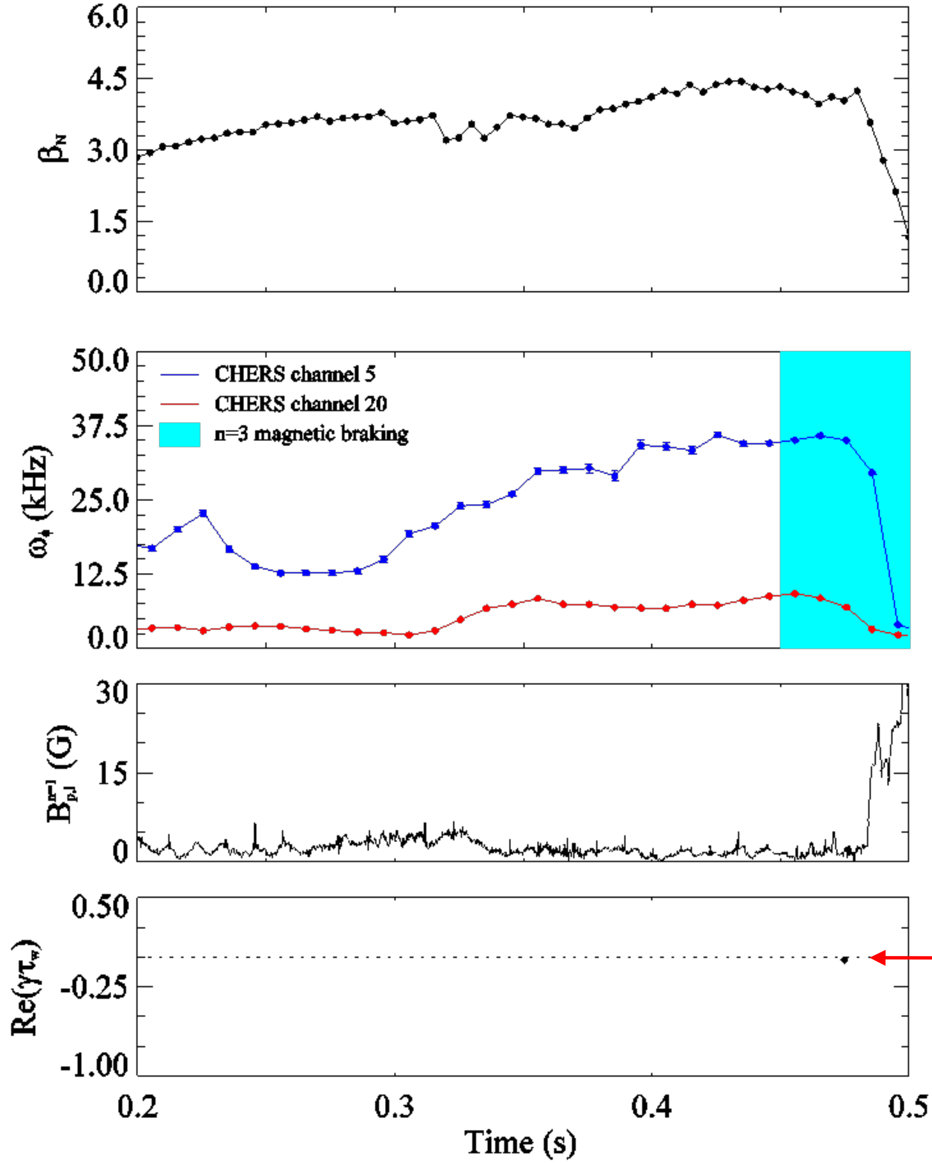
$$\left(Re\left(\delta W_K\right) - a\right)^2 + \left(Im\left(\delta W_K\right)\right)^2 = r^2$$

On a plot of $Im(\delta W_K)$ vs. $Re(\delta W_K)$, contours of constant $Re(\gamma\tau_w)$ form circles with offset a and radius r.

$$a = \frac{1}{2}\left(\delta W_b + \delta W_\infty\right) + \frac{1}{2}\left(\delta W_b - \delta W_\infty\right)\frac{Re(\gamma\tau_w)}{1 + Re(\gamma\tau_w)}$$

$$r = \frac{1}{2}\left(\delta W_b - \delta W_\infty\right)\frac{1}{1 + Re(\gamma\tau_w)}$$

Example calculation: NSTX shot 121083 @ 0.475s

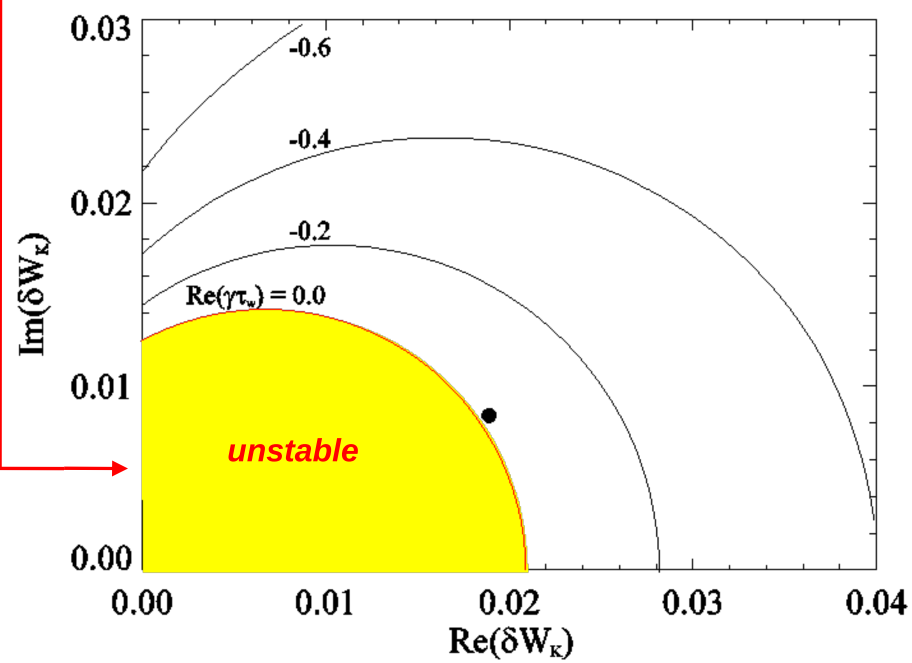


$$\delta W_\infty = -2.09 \times 10^{-2} \quad \textit{(PEST)}$$

$$\delta W_b = 7.42 \times 10^{-3} \quad \textit{(PEST)}$$

$$Re(\delta W_K) = 1.89 \times 10^{-2} \quad \textit{(MISK)}$$

$$Im(\delta W_K) = 8.41 \times 10^{-3} \quad \textit{(MISK)}$$



The Kinetic Approach to δW

$$mn\frac{d\mathbf{u}}{dt} = \mathbf{j}\times\mathbf{B} - \nabla\cdot\mathbf{P}$$ starting with a momentum equation…

$$\omega^2\frac{1}{2}\int mn|\boldsymbol{\xi}|^2 d\mathbf{V} = \frac{1}{2}\int \boldsymbol{\xi}^*\cdot\left(\tilde{\mathbf{j}}\times\mathbf{B}+\mathbf{j}\times\tilde{\mathbf{B}} - \nabla\tilde{p}_F - \nabla\tilde{p}_K\right)d\mathbf{V} \qquad \omega^2 K = \delta W_F + \delta W_K$$

…splitting into fluid and kinetic pressures $$\delta W_K = -\frac{1}{2}\int \boldsymbol{\xi}^*\cdot\nabla\tilde{p}_K d\mathbf{V}$$

For trapped ions:

Volume — Pitch Angle — Energy

$$\delta W_K^{ti} = \int_0^{\Psi_a}\frac{d\Psi}{B_0}\left(\frac{p}{1+\frac{T_e}{T_i}}\right)\frac{\sqrt{\pi}}{2}\sum_{l=-\infty}^{\infty}\int_{B_0/B_{max}}^{B_0/B_{min}} d\Lambda\hat{\tau}_b$$
$$\times\int_0^\infty\left[\frac{\omega_{*N}+\left(\hat{\varepsilon}-\frac{3}{2}\right)\omega_{*T}+\omega_E-\omega-i\gamma}{\langle\omega_D\rangle+l\omega_b-i\nu_{\rm eff}+\omega_E-\omega-i\gamma}\right]\hat{\varepsilon}^{5/2}e^{-\hat{\varepsilon}}d\hat{\varepsilon}$$
$$\times\left\langle\left(2-3\frac{\Lambda}{B_0/B}\right)(\boldsymbol{\kappa}\cdot\boldsymbol{\xi}_\perp)-\left(\frac{\Lambda}{B_0/B}\right)(\nabla\cdot\boldsymbol{\xi}_\perp)\right\rangle^2$$

(Hu, Betti, and Manickam, PoP, 2006)

The "perturbative" approach may be sufficiently accurate

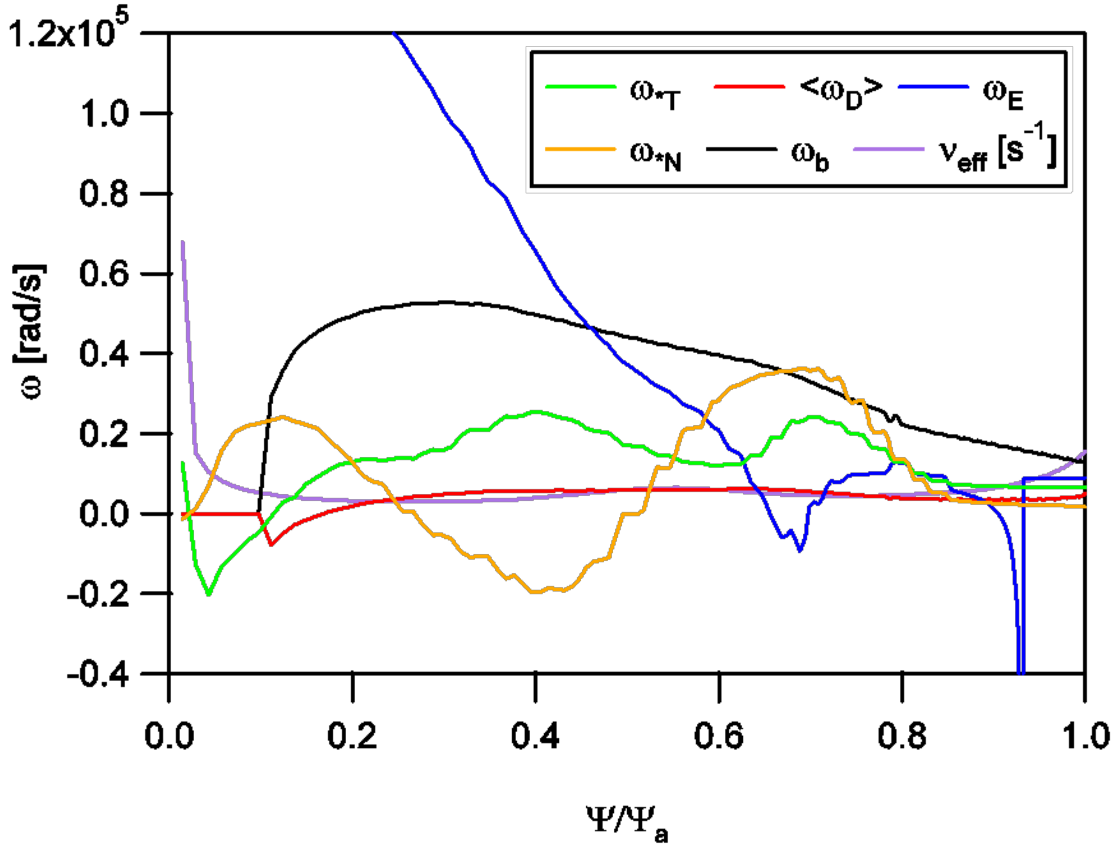
- There are two major differences between the two approaches
  - First, the self-consistent approach includes the effects of kinetics on the eigenmode.
  - Second, the self-consistent approach finds multiple roots.
  - The first effect may not be important. The second can, in principle, be found with the perturbative approach as well.

$$\delta W_K^{ti} = \int_0^{\Psi_a} \frac{d\Psi}{B_0} \left(\frac{p}{1 + \frac{T_e}{T_i}} \right) \frac{\sqrt{\pi}}{2} \sum_{l=-\infty}^{\infty} \int_{B_0/B_{max}}^{B_0/B_{min}} d\Lambda \hat{\tau}_b$$
$$\times \int_0^{\infty} \left[\frac{\omega_{*N} + \left(\hat{\varepsilon} - \frac{3}{2}\right) \omega_{*T} + \omega_E - \omega - i\gamma}{\langle \omega_D \rangle + l\omega_b - i\nu_{\text{eff}} + \omega_E - \omega - i\gamma} \right] \hat{\varepsilon}^{5/2} e^{-\hat{\varepsilon}} d\hat{\varepsilon}$$
$$\times \left\langle \left(2 - 3\frac{\Lambda}{B_0/B}\right) (\boldsymbol{\kappa} \cdot \boldsymbol{\xi}_\perp) - \left(\frac{\Lambda}{B_0/B}\right) (\boldsymbol{\nabla} \cdot \boldsymbol{\xi}_\perp) \right\rangle^2$$

(Hu, Betti, and Manickam, PoP, 2006)

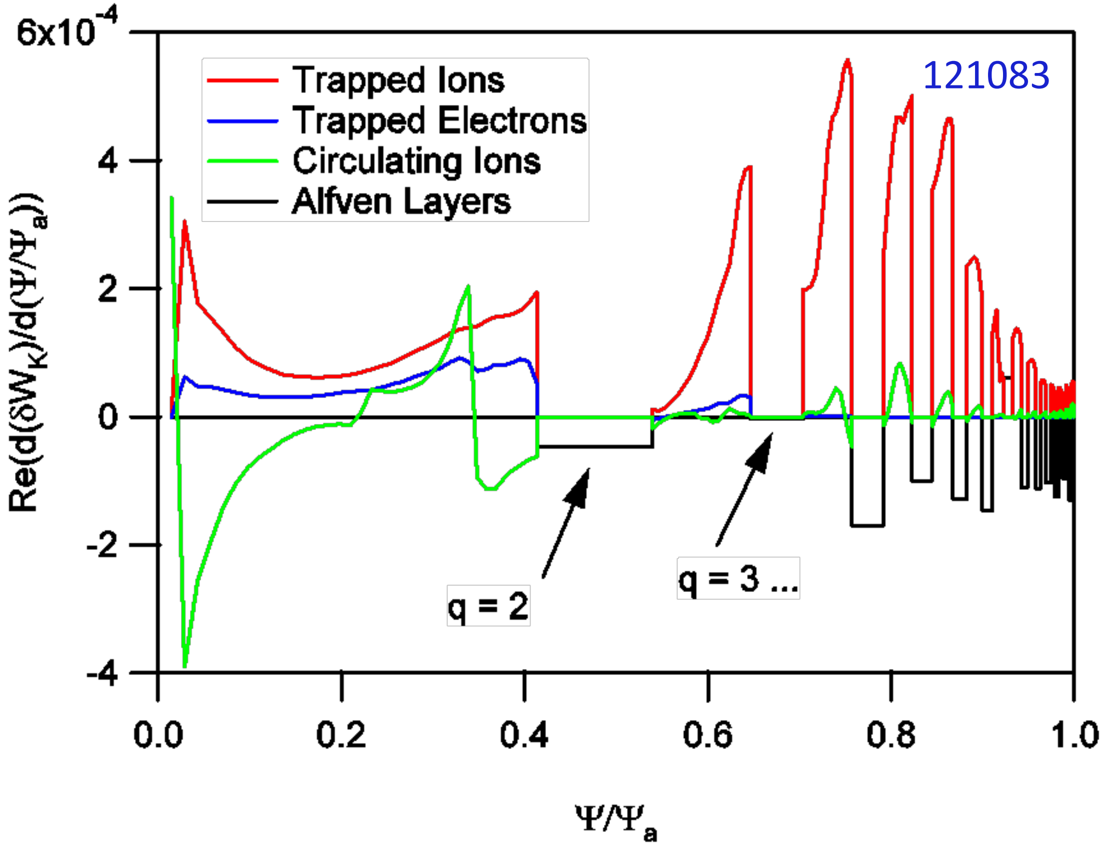
Stabilization arises from the resonance of various plasma frequencies

Frequency resonance term: $$\delta W_K \propto \int \left[\frac{\omega_{*N} + \left(\hat{\varepsilon} - \frac{3}{2}\right)\omega_{*T} + \omega_E - \omega - i\gamma}{\langle \omega_D \rangle + l\omega_b - i\nu_{\text{eff}} + \omega_E - \omega - i\gamma} \right] \hat{\varepsilon}^{\frac{5}{2}} e^{-\hat{\varepsilon}} d\hat{\varepsilon}$$



The collision frequency shown is for thermal ions, the bounce and precession drift frequencies are for thermal ions and zero pitch angle.

Strong trapped ion stabilization comes from the outer surfaces



Contributions to Re(δW_K)

- A large portion of the kinetic stabilization comes from q>2, where ω_{*N} and ω_{*T} $\geq\omega_E$.

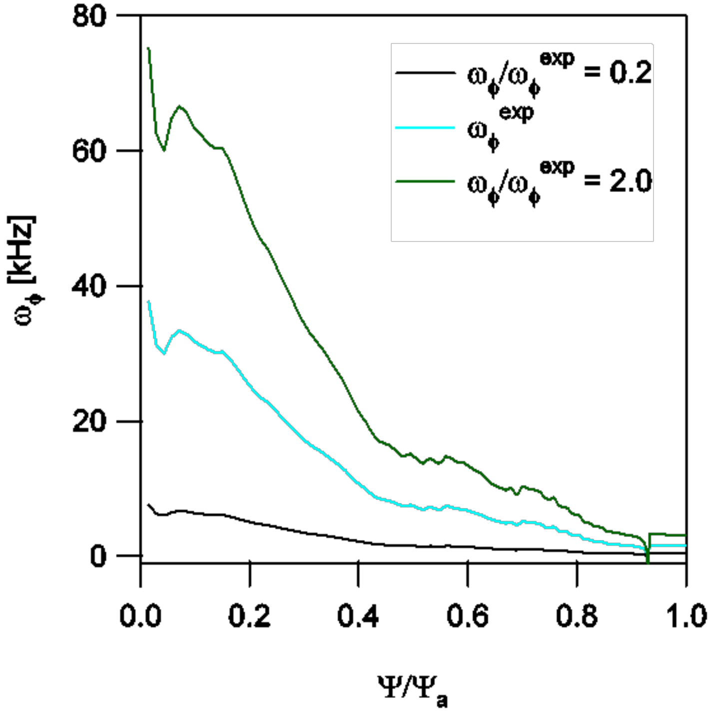
The flat areas are rational surfaces (integer q ± 0.2) where the contributions have been zeroed out and dealt with analytically through calculation of inertial enhancement by shear Alfven damping.

(Zheng et al., PRL, 2005)

The effect of rotation on stability is more complex with kinetic theory

$$\delta W_K \propto \int \left[\frac{\omega_{*N} + \left(\hat{\varepsilon} - \frac{3}{2}\right)\omega_{*T} + \omega_E - \omega - i\gamma}{\langle \omega_D \rangle + l\omega_b - i\nu_{\text{eff}} + \omega_E - \omega - i\gamma} \right] \hat{\varepsilon}^{\frac{5}{2}} e^{-\hat{\varepsilon}} d\hat{\varepsilon} \qquad \omega_E = \omega_\phi - \omega_{*N} - \omega_{*T}$$

Using self-similarly scaled rotation profiles, we can isolate and test the effect of rotation on stability.



121083

The effect of rotation on stability is more complex with kinetic theory

$$\delta W_K \propto \int \left[\frac{\omega_{*N} + \left(\hat{\varepsilon} - \frac{3}{2}\right)\omega_{*T} + \omega_E - \omega - i\gamma}{\langle \omega_D \rangle + l\omega_b - i\nu_{\text{eff}} + \omega_E - \omega - i\gamma} \right] \hat{\varepsilon}^{\frac{5}{2}} e^{-\hat{\varepsilon}} d\hat{\varepsilon} \qquad \omega_E = \omega_\phi - \omega_{*N} - \omega_{*T}$$

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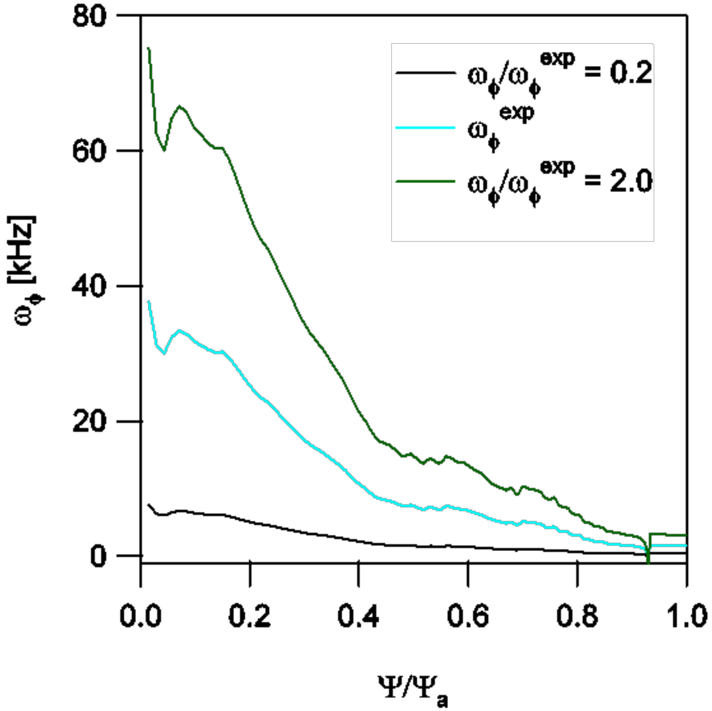
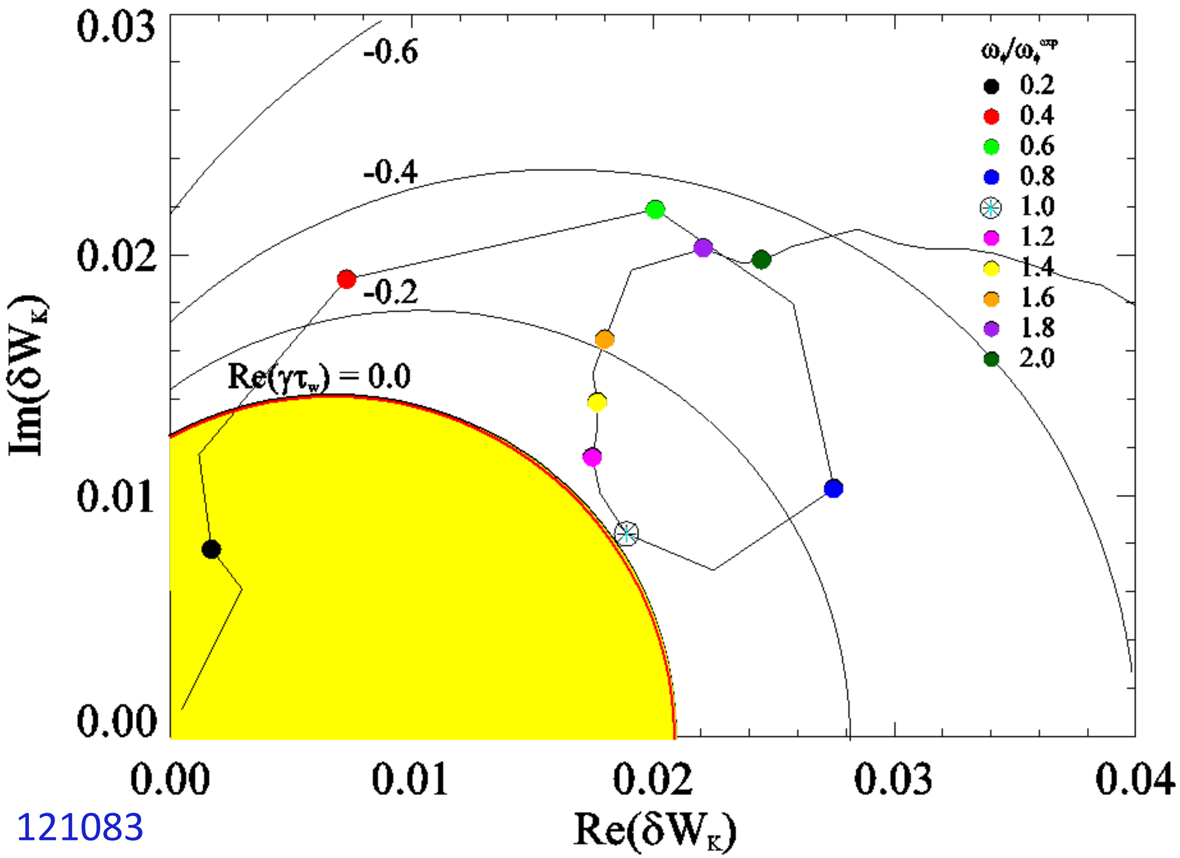
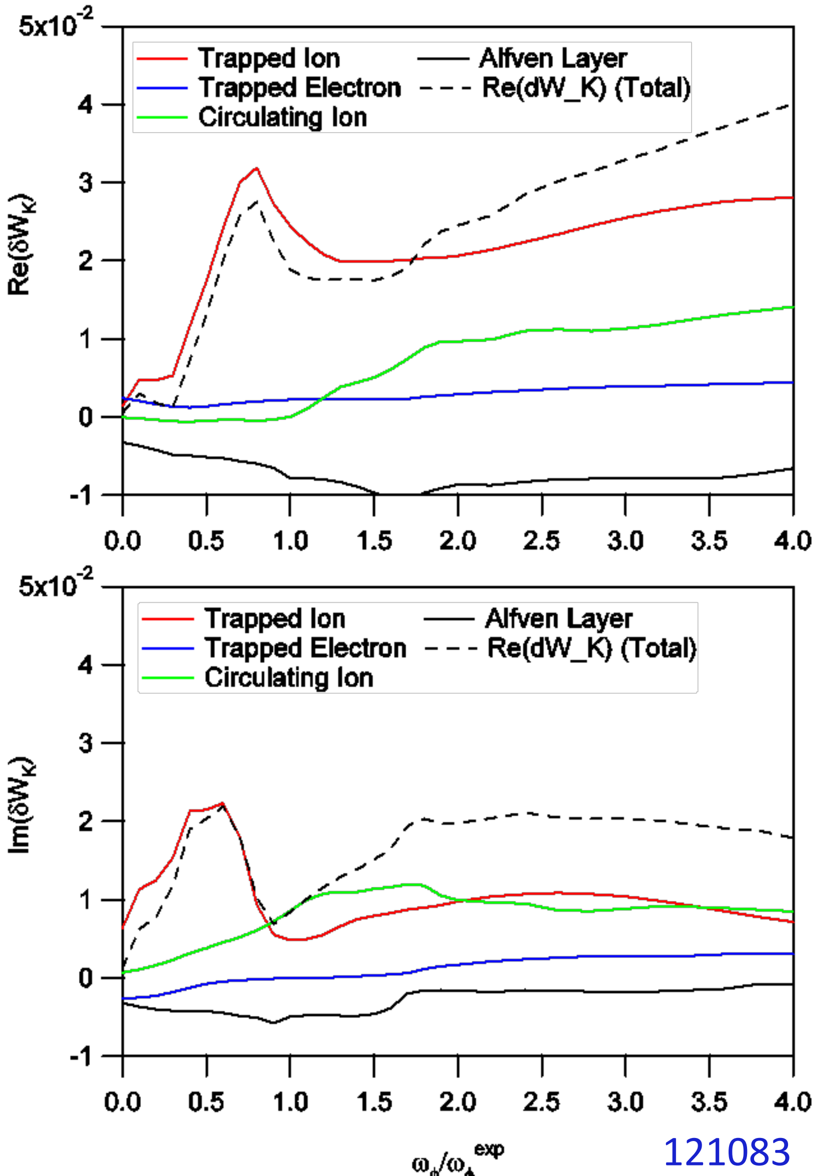




121083

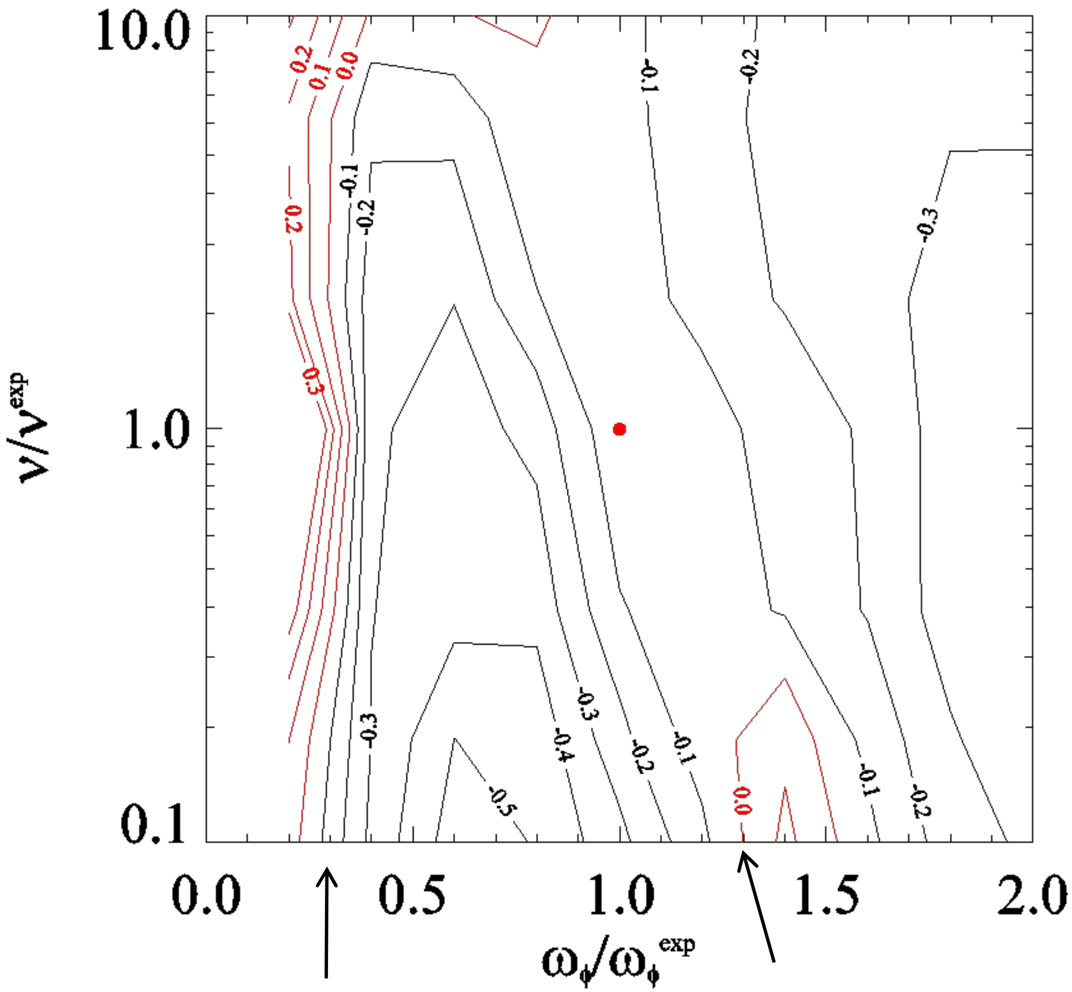
Resonances with multiple particle types contribute to the trends



- For $\omega_\phi/\omega_\phi^{exp}$ from 0 to 0.6 stability increases as the real and imaginary trapped ion components increase.
- From 0.6 to 0.8 the real part increases while the imaginary part decreases, leading to the turn back towards instability.
- For rotation levels above the experimental value, trapped ion and circulating ion components rise, leading to strong stability.

Isolating and testing the effect of collisionality reveals band of marginal stability

contours of $Re(\gamma\tau_w)$



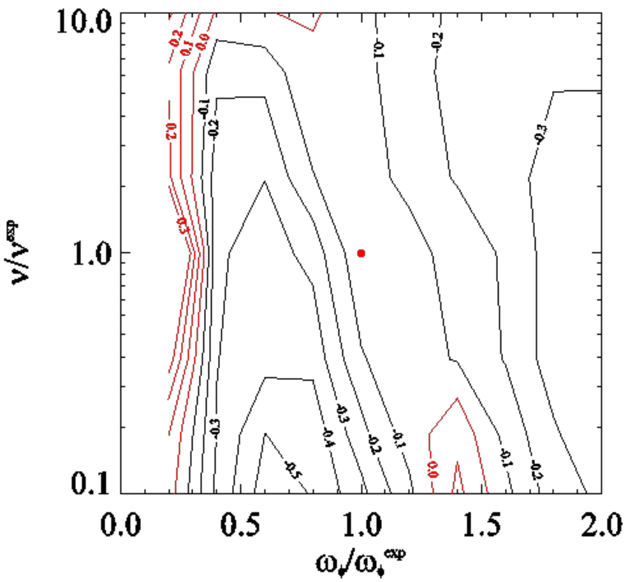
Similar to a "critical" rotation threshold.

Band of marginally stable moderate rotation.

- Density and temperature profiles are self-similarly scaled while keeping β constant (ie, n*2 and T/2 or n/2 and T*2).
- We find a low rotation "critical" threshold.
- However, there is also a band of marginally stable moderate rotation, and it is here that the experiment goes unstable.

Analysis of other NSTX shots shows the same characteristic behavior

121083



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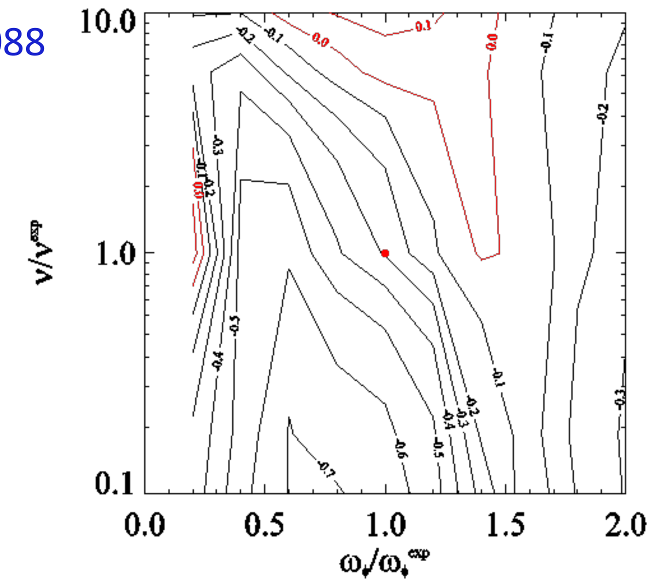

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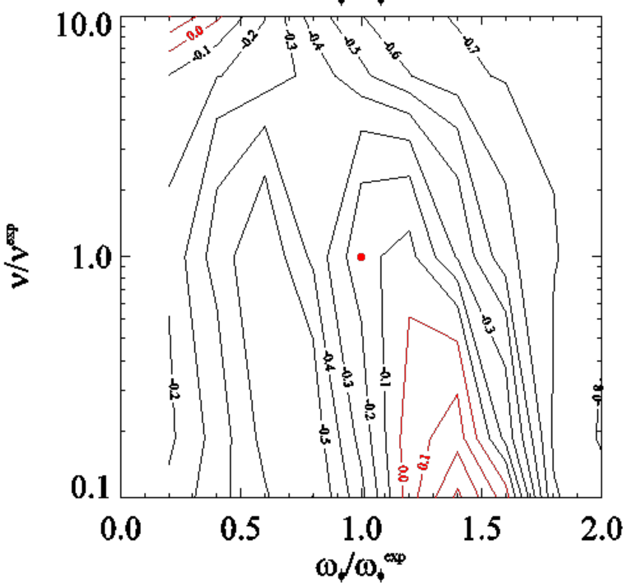
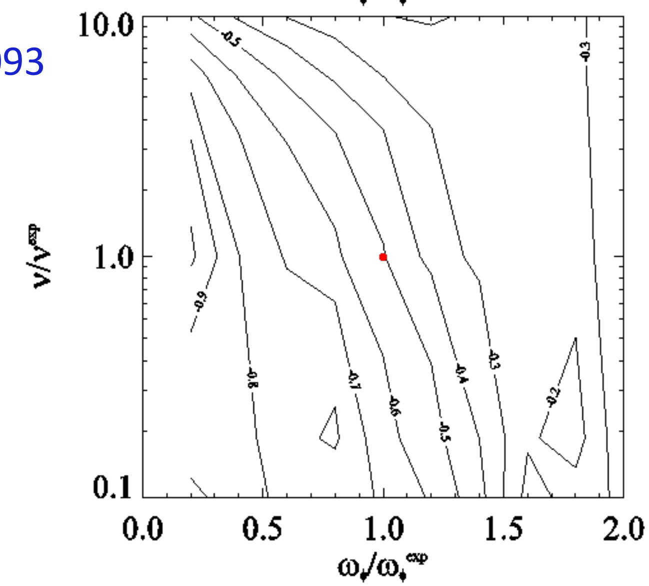
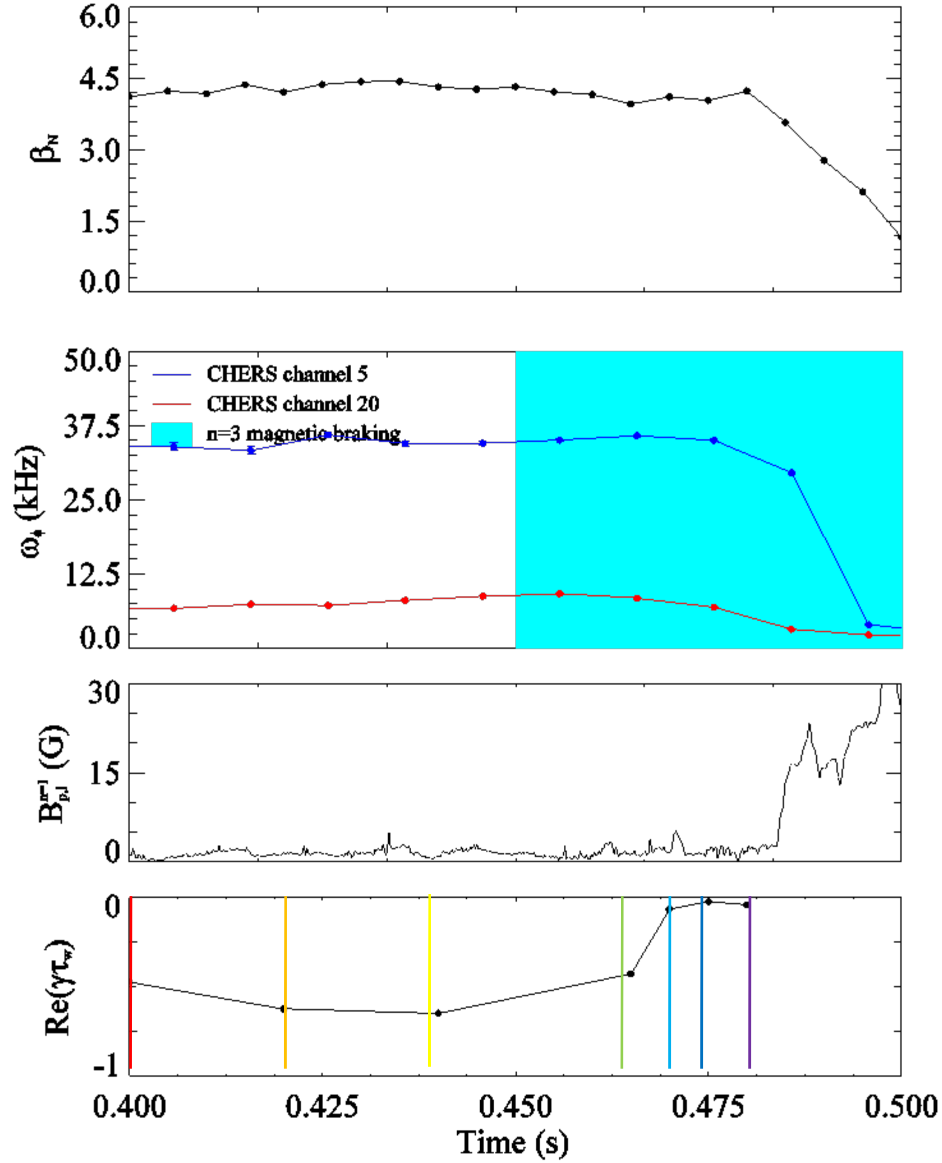


121093

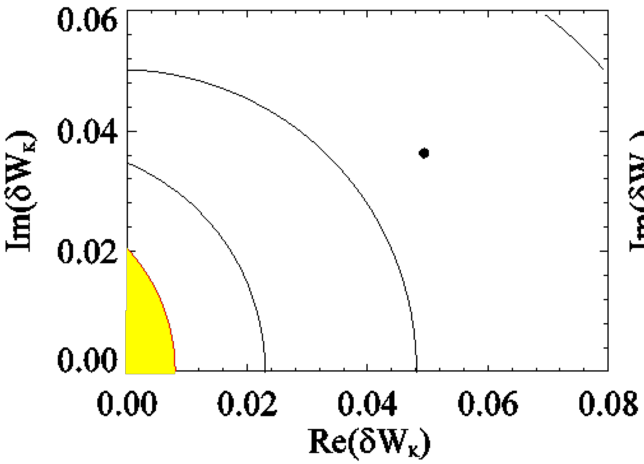


Kinetic prediction of instability matches experimental result

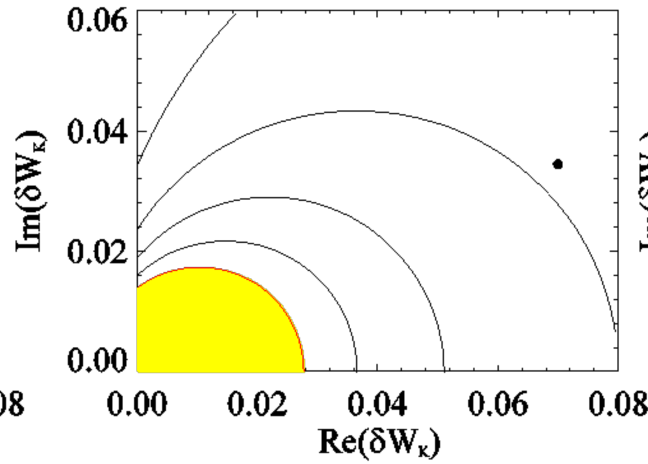
- Examining the evolution of a shot
  - For NSTX shot 121083, β and ω_ϕ are relatively constant leading up the the RWM collapse.
  - Calculation of the RWM kinetic growth rate for multiple equilibria shows a turn towards instability just before the RWM.



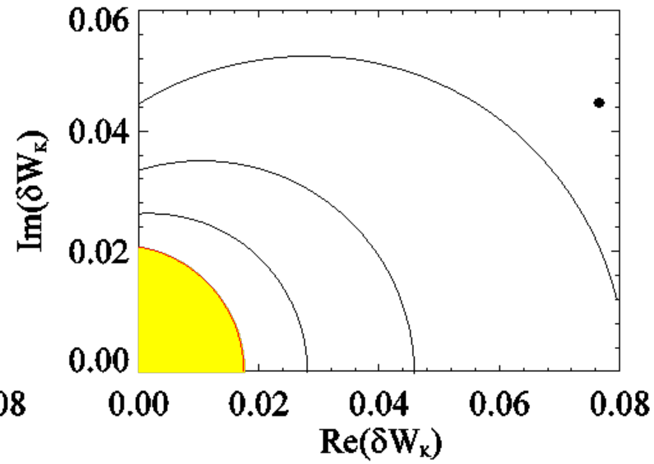
As time progresses the stabilizing δW_K decreases



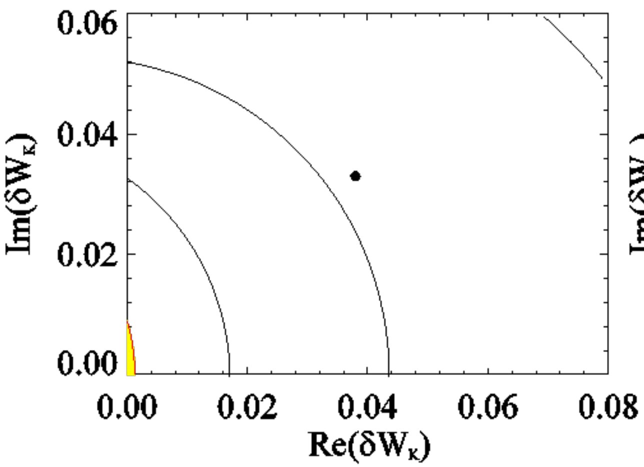
0.400



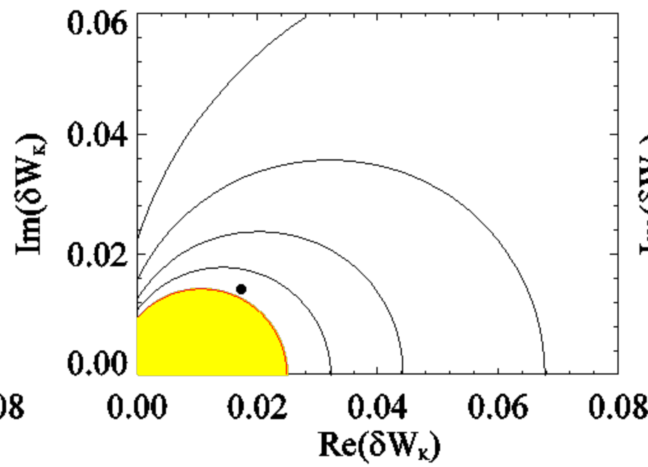
0.420



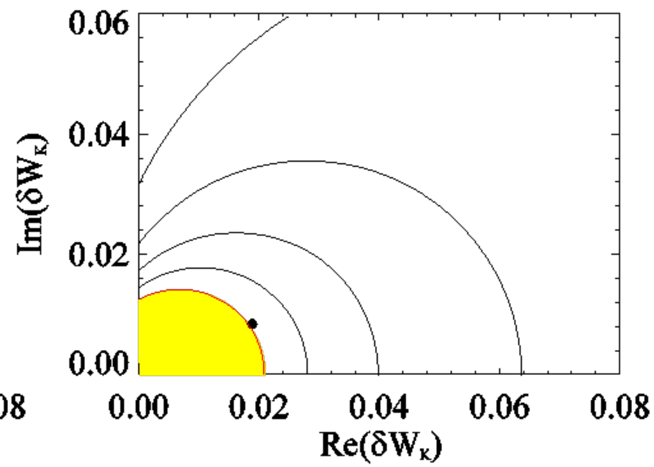
0.440



0.465

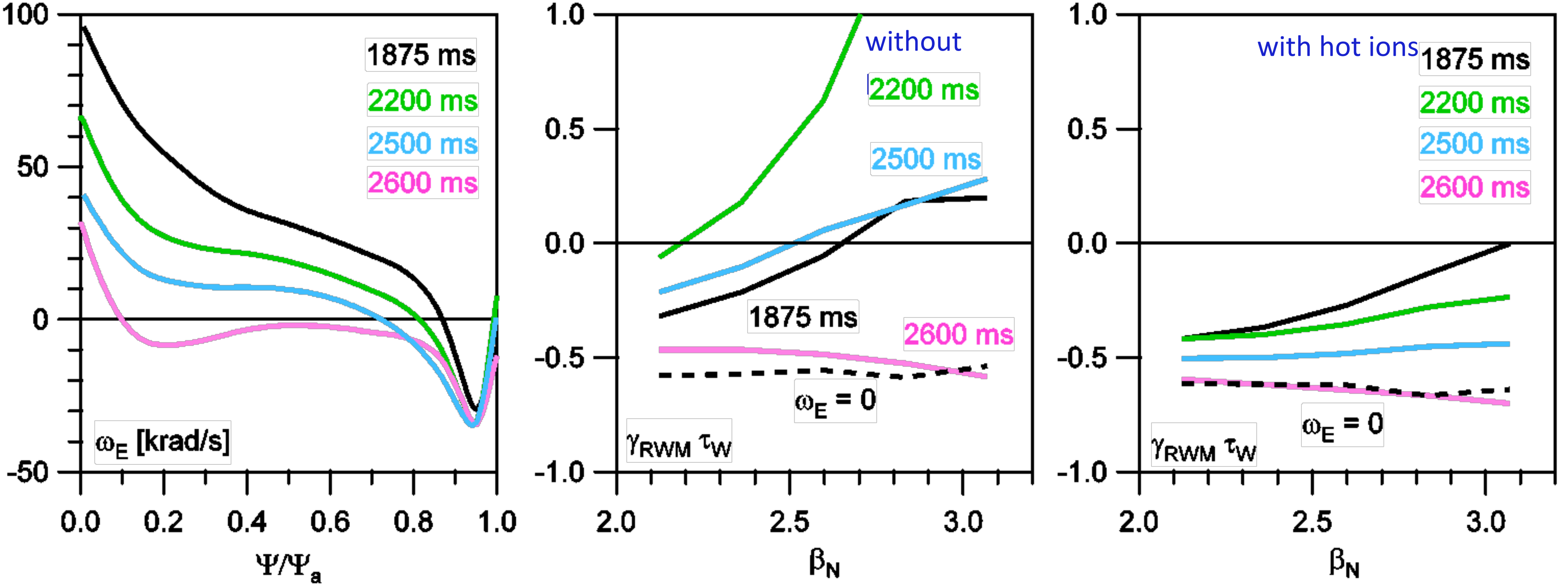


0.470



0.475

Hot ions have a strongly stabilizing effect for DIII-D



- Using the equilibrium from DIII-D shot 125701 @ 2500ms and rotation from 1875-2600ms, MISK predicts a band of instability at moderate rotation without hot ions, but complete stability with hot ions.
- This could help to explain why DIII-D is inherently more stable to the RWM than NSTX, and possibly why energetic particle modes can “trigger” the RWM.

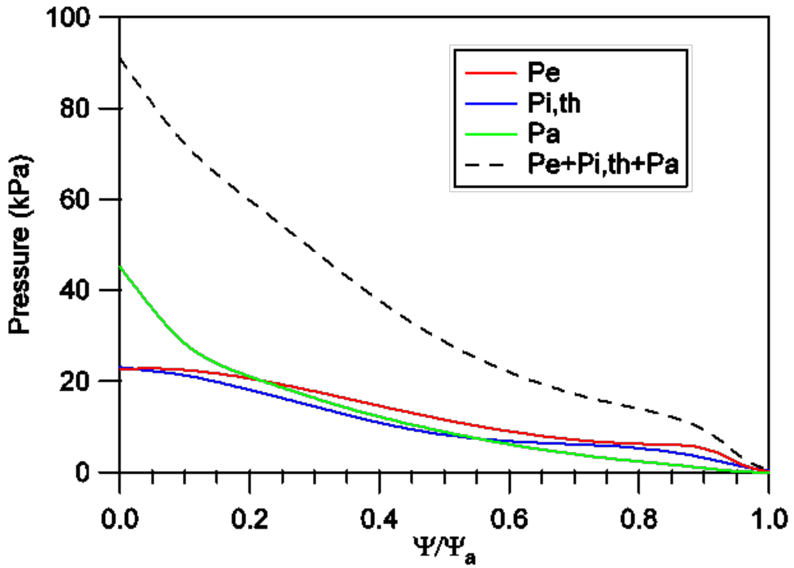
(Matsunaga et al., IAEA, 2008)

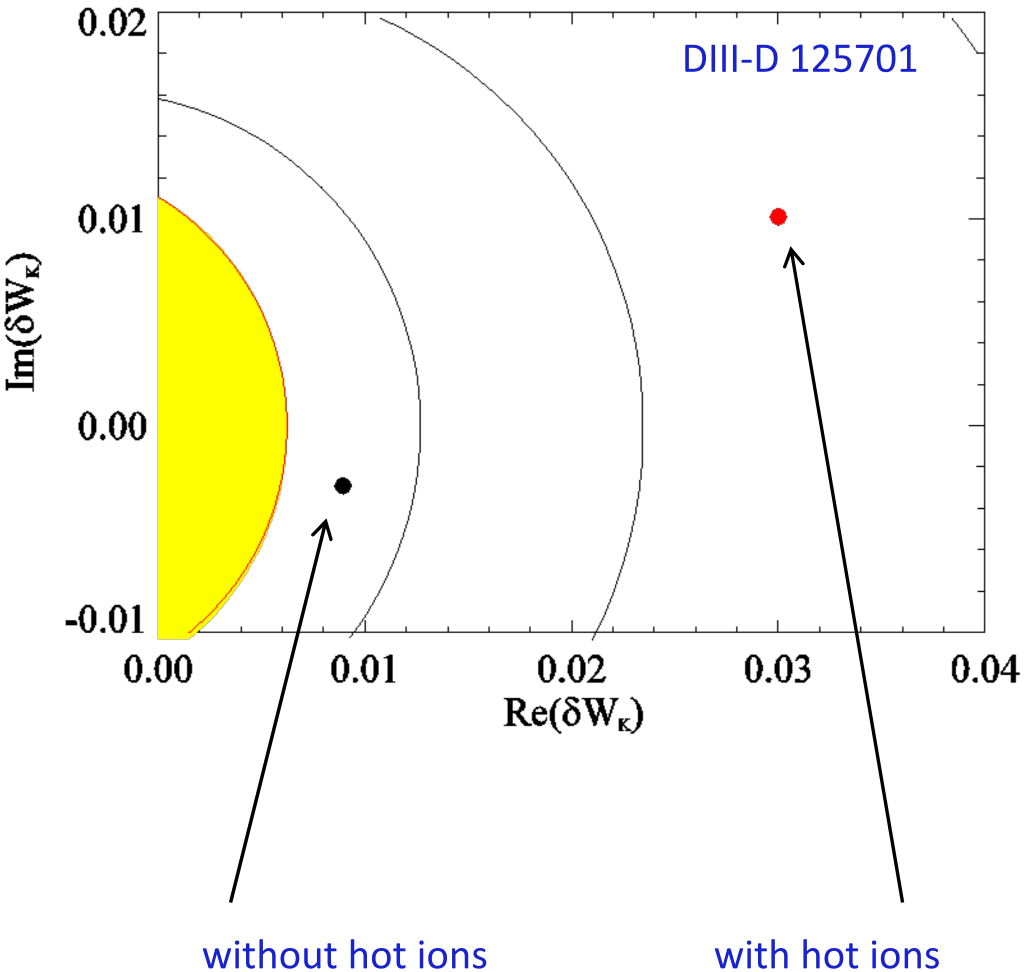
Summary

- Kinetic effects contribute to stabilization of the RWM, and may explain the complex relationship between plasma rotation and stability in NSTX.
- The MISK code is used to calculate the RWM growth rate with kinetic effects.
- Calculations match the experimental observation of instability at moderate rotation and the evolution of a discharge from stable to unstable.
- DIII-D results suggest the importance of hot ions.

Extra Slides

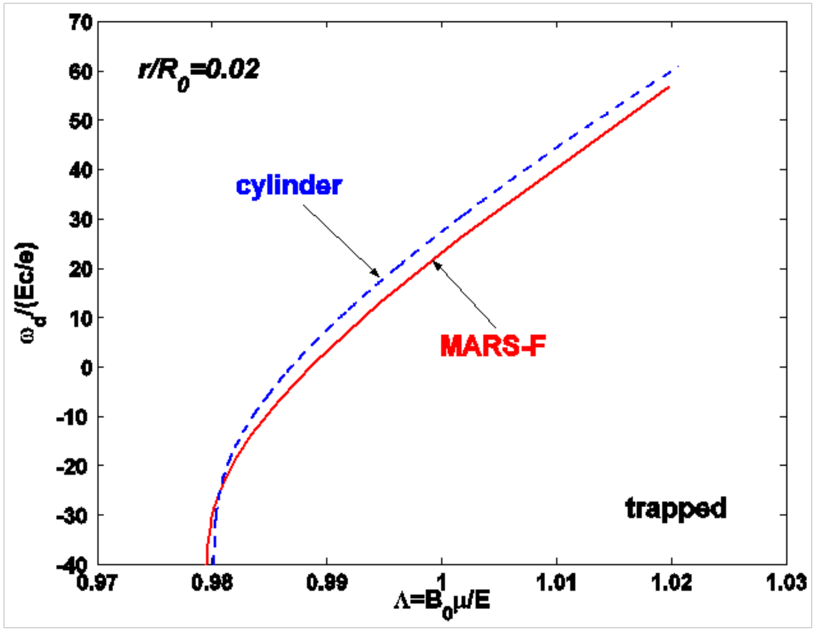
Hot ions contribute a large fraction of β to DIII-D, and have a strong stabilizing effect





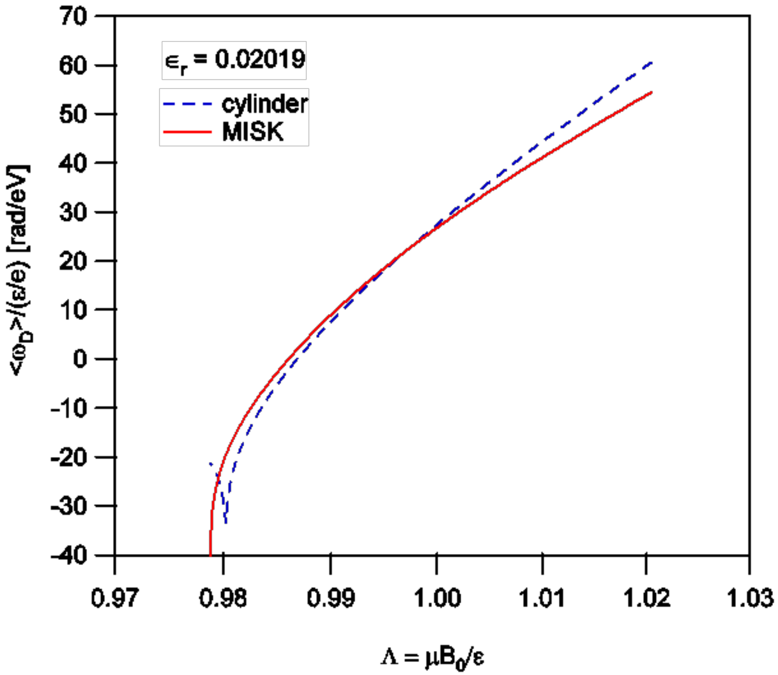
– Hot ions not yet implemented for NSTX.

Drift frequency calculations match for MISK and MARS-K



MARS

(Liu, ITPA MHD TG Meeting, Feb. 25-29, 2008)



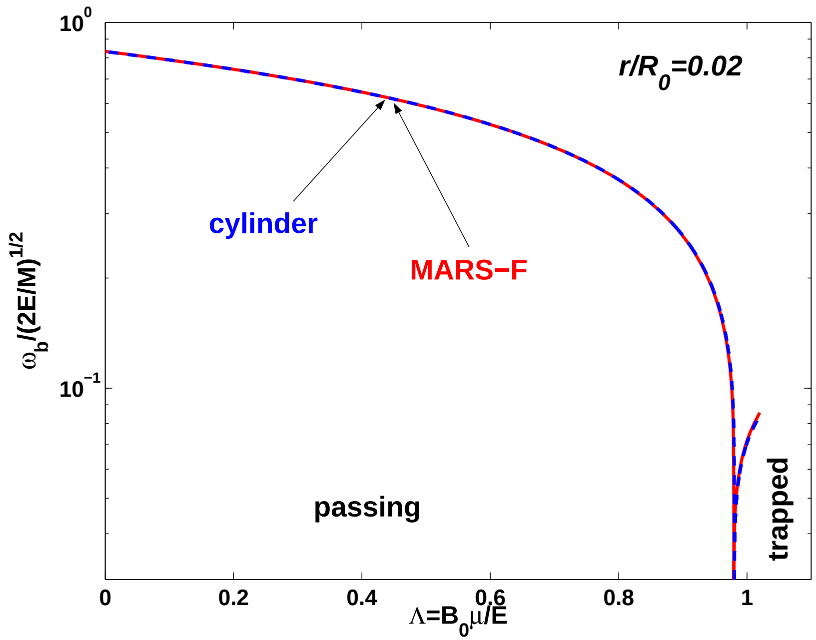
MISK

large aspect ratio approximation *(Jucker et al., PPCF, 2008)*

$$\frac{\langle \omega_D \rangle}{\varepsilon / e} = \frac{2q\Lambda}{R_0^2 B_0 \epsilon_r}\left[(2s+1)\frac{E\left(k^2\right)}{K\left(k^2\right)} + 2s\left(k^2-1\right) - \frac{1}{2}\right] \qquad k = \left[\frac{1-\Lambda+\epsilon_r\Lambda}{2\epsilon_r\Lambda}\right]^{\frac{1}{2}}$$

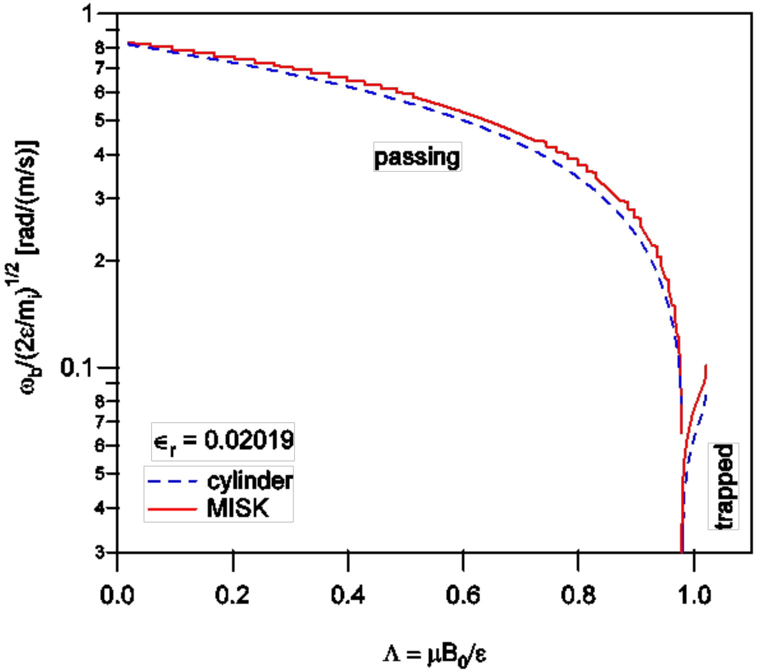
here, ϵ_r is the inverse aspect ratio, s is the magnetic shear, K and E are the complete elliptic integrals of the first and second kind, and $\Lambda = \mu B_0/\varepsilon$, where μ is the magnetic moment and ε is the kinetic energy.

Bounce frequency calculations match for MISK and MARS-K



MARS

(Liu, ITPA MHD TG Meeting, Feb. 25-29, 2008)



MISK

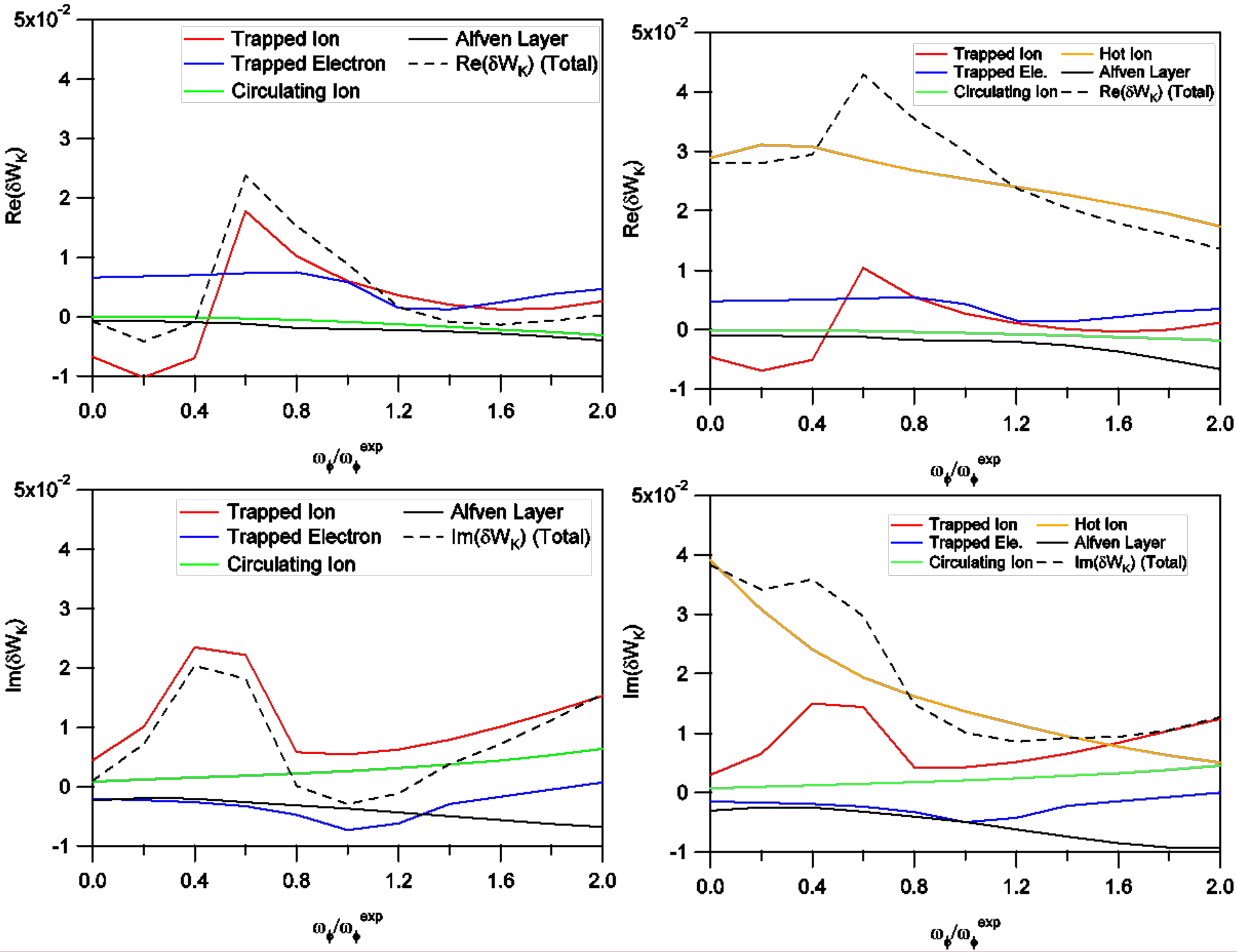
large aspect ratio approximation *(Bondeson and Chu, PoP, 1996)*

$$\frac{\omega_b}{\sqrt{2\varepsilon/m_i}} = \frac{\sqrt{2\epsilon_r\Lambda}}{4qR_0}\frac{\pi}{K(k)} \quad \text{(trapped)}$$

$$\frac{\omega_b}{\sqrt{2\varepsilon/m_i}} = \frac{\sqrt{1-\Lambda+\epsilon_r\Lambda}}{2qR_0}\frac{\pi}{K(1/k)} \quad \text{(circulating)}$$

$$k = \left[\frac{1-\Lambda+\epsilon_r\Lambda}{2\epsilon_r\Lambda}\right]^{\frac{1}{2}}$$

Components of δW_K without and with hot ions



hot ion contribution is large.

Hot ions are included in MISK with a slowing-down distribution function

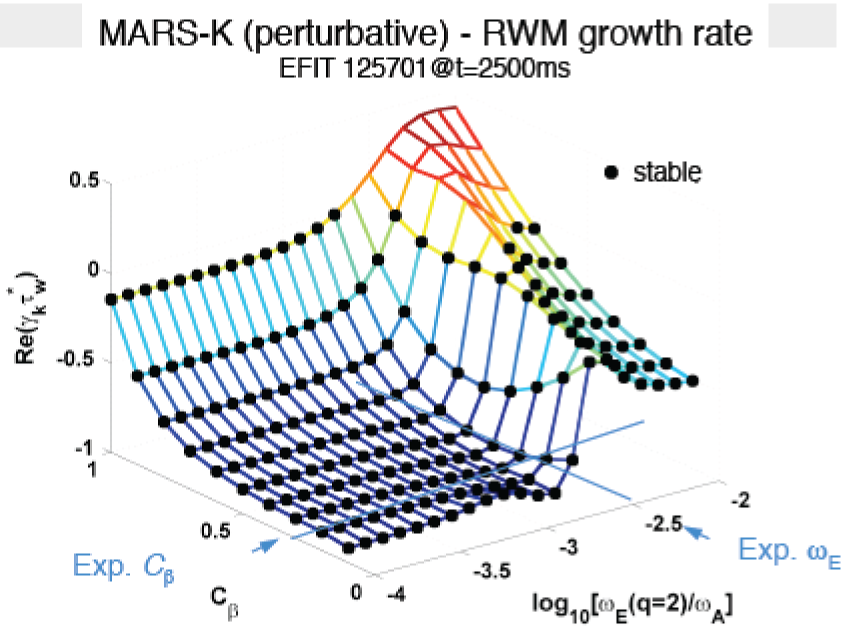
$$\delta W_K^a \propto \left(\int \frac{\hat{\varepsilon}_a^{\frac{1}{2}}}{\hat{\varepsilon}_a^{\frac{3}{2}} + \hat{\varepsilon}_c^{\frac{3}{2}}} d\hat{\varepsilon}_a \right)^{-1} \int \left[\frac{\omega_{*N}^a + \omega_f^a + \frac{1}{\hat{\varepsilon}_a^{\frac{3}{2}} + \hat{\varepsilon}_c^{\frac{3}{2}}} \frac{d}{d\Psi} \hat{\varepsilon}_c^{\frac{3}{2}} + \frac{\frac{3}{2}\hat{\varepsilon}_a^{\frac{1}{2}}}{\hat{\varepsilon}_a^{\frac{3}{2}} + \hat{\varepsilon}_c^{\frac{3}{2}}} Z_a \left(\omega_E^a - \omega - i\gamma \right)}{\langle \omega_D^a \rangle + l\omega_b^a - i\nu_{\text{eff}}^a + Z_a \left(\omega_E^a - \omega - i\gamma \right)} \right] \frac{\hat{\varepsilon}_a^{\frac{5}{2}}}{\hat{\varepsilon}_a^{\frac{3}{2}} + \hat{\varepsilon}_c^{\frac{3}{2}}} d\hat{\varepsilon}_a$$

(Hu, Betti, and Manickam, PoP, 2006)

$$\hat{\varepsilon}_c = \left(\frac{3\sqrt{\pi}}{4} \right)^{\frac{2}{3}} \left(\frac{m_a}{m_i} \right) \left(\frac{m_i}{m_e} \right)^{\frac{1}{3}} \left(\frac{T_e}{\varepsilon_a} \right) \qquad \omega_f^a = \int \frac{\hat{\varepsilon}_a^{\frac{1}{2}}}{\hat{\varepsilon}_a^{\frac{3}{2}} + \hat{\varepsilon}_c^{\frac{3}{2}}} d\hat{\varepsilon}_a \frac{d \left(\left(\int \frac{\hat{\varepsilon}_a^{\frac{1}{2}}}{\hat{\varepsilon}_a^{\frac{3}{2}} + \hat{\varepsilon}_c^{\frac{3}{2}}} d\hat{\varepsilon}_a \right)^{-1} \right)}{d\Psi}$$

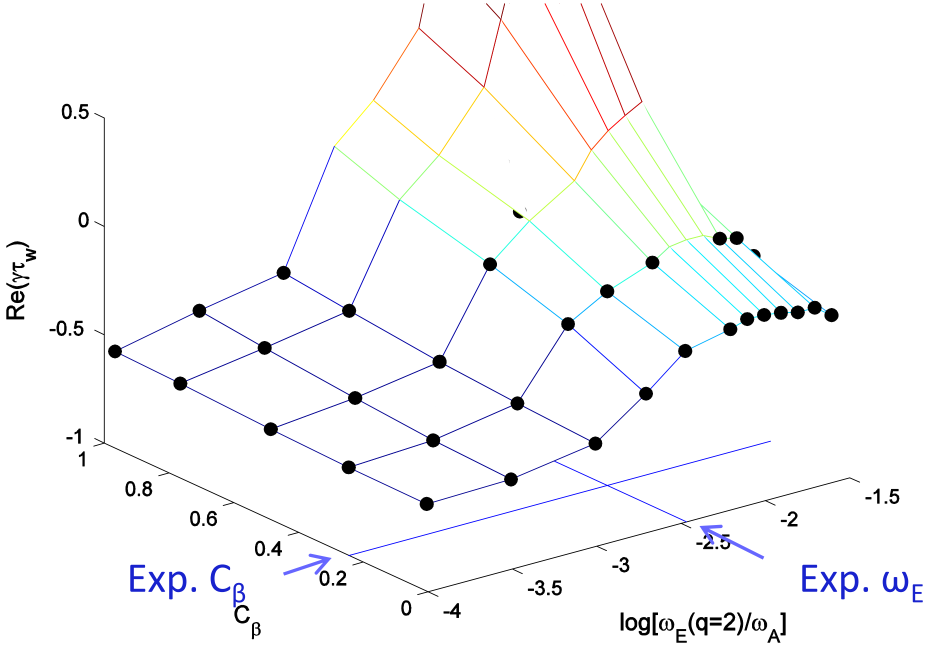
- Profiles of p_a and n_a, calculated by onetwo are used both directly and to find ε_a.
- The hot ion pressure is subtracted from the total pressure for the other parts of the calculation.

MARS-K (perturbative) and MISK produce similar results for the same DIII-D case



MARS

(Reimerdes, PO3.00011, Wed. PM)



MISK

- Results are qualitatively similar. Main differences are unstable magnitude and behavior at low ω_E.
- This case does not include hot ions.

Future work: inclusion of pitch angle dependent collisionality

$$\left[\frac{\partial}{\partial t} + \mathbf{v_g} \cdot \boldsymbol{\nabla}\right] \tilde{f} + \left[\frac{\tilde{\mathbf{F}}}{m}\right] \frac{\partial f}{\partial \mathbf{v}} = C(\tilde{f})$$

Collisionality enters through drift kinetic equation

$$C(\tilde{f}) = \nu_{\mathrm{eff}} \tilde{f}$$

Effective collisionality can be included in various forms:

$$\nu_0 = 0 \text{ (collisionless)}$$

$$\nu_1(\Psi) = \frac{n_i e^4 \ln \Lambda_e}{12\pi^{\frac{3}{2}} \epsilon_0^2 m_i^{\frac{1}{2}} T_i^{\frac{3}{2}} \epsilon_r} \text{ (no energy dependence)}$$

(MARS)

$$\nu_2(\Psi, \varepsilon) = \nu_1 \hat{\varepsilon}^{-\frac{3}{2}} \text{ (simple energy dependence)}$$

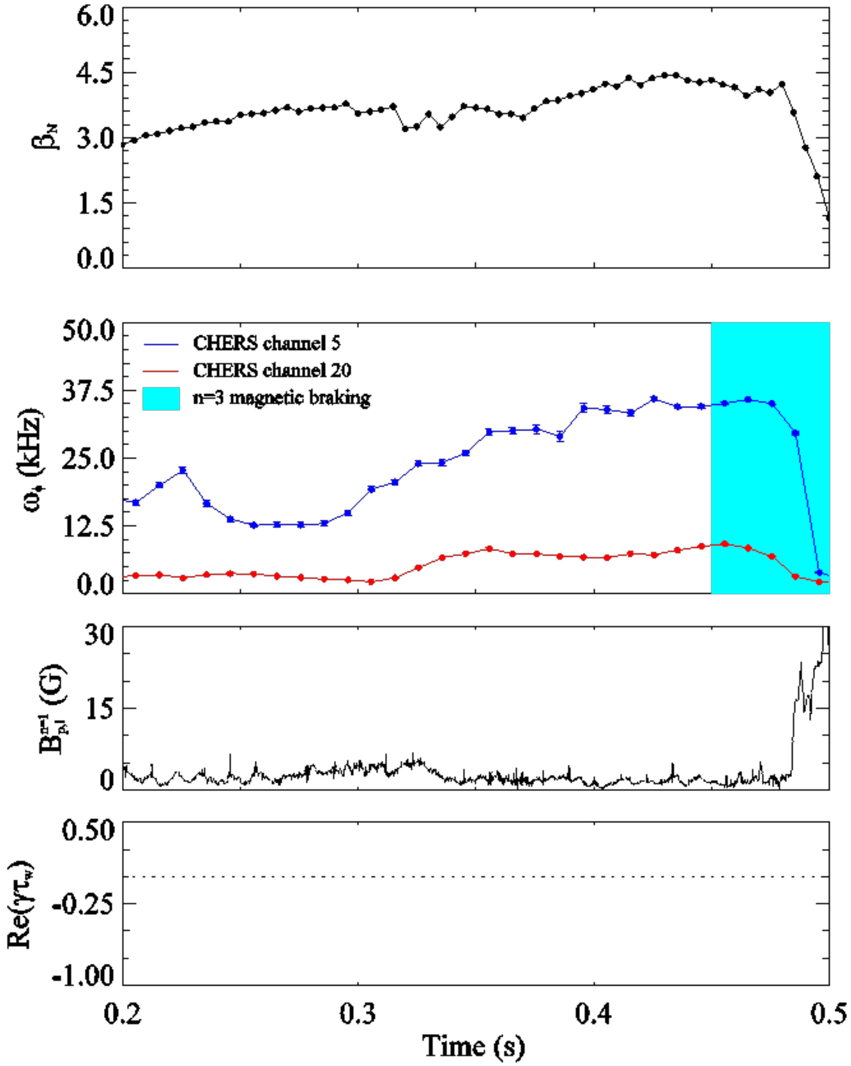
(MISK)

$$\nu_3(\Psi, \varepsilon, \Lambda) = 2\nu_1 \left[Z_{\mathrm{eff}} + \frac{1}{\sqrt{\pi \hat{\varepsilon}}} e^{-\hat{\varepsilon}} + \frac{1}{\sqrt{\pi}} \left(2 - \hat{\varepsilon}^{-1}\right) \int_0^{\sqrt{\hat{\varepsilon}}} e^{-t^2} dt\right] \times \frac{\partial}{\partial \Lambda} \left(\Lambda \sqrt{\frac{B_0}{B} - \Lambda}\right) \frac{\partial}{\partial \Lambda} \left(\sqrt{\frac{B_0}{B} - \Lambda}\right) \text{ (Lorentz operator, pitch angle dependence)}$$

(Future?)

(Fu et al., POFB, 1993)

RWM can be experimentally unstable in NSTX



- RWM observed in NSTX on magnetic diagnostics at the time of β and ω_ϕ collapse.
  - Does kinetic theory predict that the RWM growth rate becomes positive at this time?
  - We will calculate $\gamma\tau_w$, the normalized kinetic growth rate, with the MISK code.

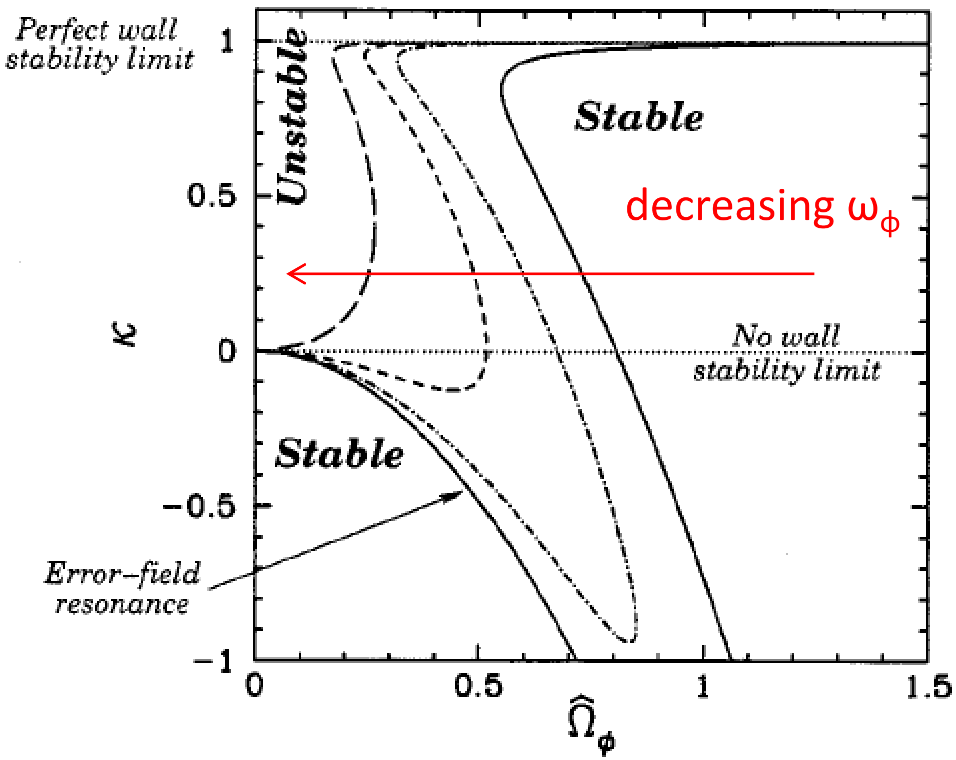
Previous “simple” models led to a “critical” rotation

$$\left[(\hat{\gamma} - i\hat{\omega}_\phi)^2 + \nu_* (\hat{\gamma} - i\hat{\omega}_\phi) + (1-\kappa)(1-md)\right](\hat{\gamma} S_* + 1 + md) = 1 - (md)^2$$

(Fitzpatrick, PoP, 2002)

toroidal plasma rotation

- Fitzpatrick simple model
  - Plasma rotation increases stability, and for a given equilibrium there is a “critical” rotation, above which the plasma is stable.



The RWM Energy Principle is modified by kinetic effects

$$\gamma_F \tau_w = -\frac{\delta W_\infty}{\delta W_b}$$

(Haney and Freidberg, PoF-B, 1989)

$$\gamma_K \tau_w = -\frac{\delta W_\infty + \delta W_K}{\delta W_b + \delta W_K}$$

(Hu, Betti, and Manickam, PoP, 2005)

PEST → δW_∞ ; MISK → δW_K

The kinetic contribution has a real and imaginary part, so:

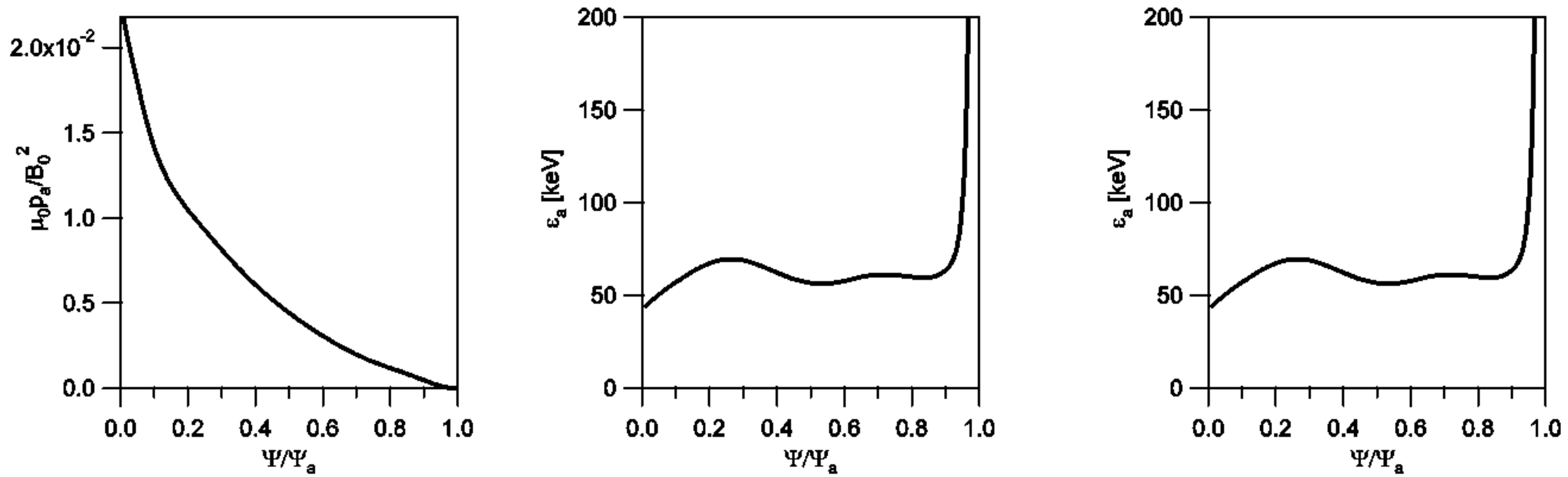
$$Re(\gamma_K \tau_w) = -\frac{\delta W_\infty \delta W_b + (Im(\delta W_K))^2 + Re(\delta W_K)(\delta W_\infty + \delta W_b + Re(\delta W_K))}{(\delta W_b + Re(\delta W_K))^2 + (Im(\delta W_K))^2}$$

Inputs to the hot ion contribution

Given p_a and n_a, find ε_a, by iteration:

$$\varepsilon_a = \frac{p_a}{p_i}\frac{n_i}{n_a}\frac{\overline{T_i}}{\overline{T_a}}$$

$$\overline{T_a} = \frac{2}{3}\left(\int \frac{\hat{\varepsilon}_a^{\frac{3}{2}}}{\hat{\varepsilon}_a^{\frac{3}{2}} + \hat{\varepsilon}_c^{\frac{3}{2}}} d\hat{\varepsilon}_a\right)\left(\int \frac{\hat{\varepsilon}_a^{\frac{1}{2}}}{\hat{\varepsilon}_a^{\frac{3}{2}} + \hat{\varepsilon}_c^{\frac{3}{2}}} d\hat{\varepsilon}_a\right)^{-1} \qquad \hat{\varepsilon}_c = \left(\frac{3\sqrt{\pi}}{4}\right)^{\frac{2}{3}}\left(\frac{m_a}{m_i}\right)\left(\frac{m_i}{m_e}\right)^{\frac{1}{3}}\left(\frac{T_e}{\varepsilon_a}\right)$$



input input result*

* note that the rise at the edge is due to n_a going to zero faster than p_a. This didn't affect the results in this case (see Ψ > 0.9 on plot on next page), but it is something to pay attention to.

Difference between hot and thermal ion calculations

Let us examine the difference between the trapped ion and hot ion contributions:

$$\delta W_K^{hi} = \sum_l \sqrt{2}\pi^2 \int \frac{d\Psi}{B_0} n_a \varepsilon_a \int d\Lambda \hat{\tau}_b^{hi} |\langle H^{hi} \rangle|^2 I^{hi}$$

$$\delta W_K^{ti} = \sum_l \frac{\sqrt{\pi}}{2} \int \frac{d\Psi}{B_0} n_i T_i \int d\Lambda \hat{\tau}_b^{ti} |\langle H^{ti} \rangle|^2 I^{ti}$$

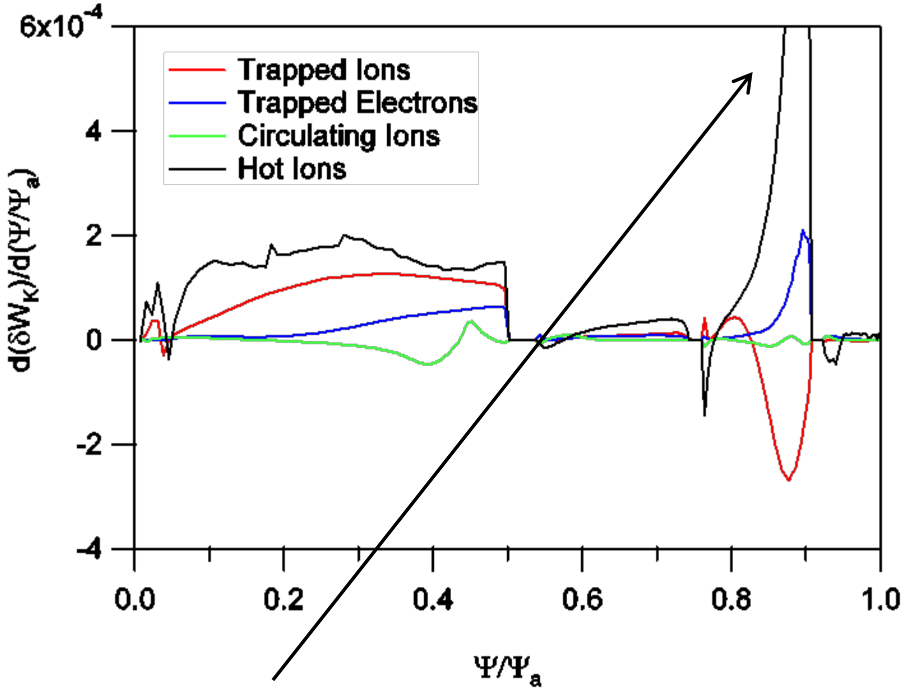
$$\frac{\delta W_K^{hi}}{\delta W_K^{ti}} \approx (2\pi)^{\frac{3}{2}} \frac{n_a}{n_i} \frac{\varepsilon_a}{T_i} \frac{I^{hi}}{I^{ti}}$$

large

small

large to very large, near edge

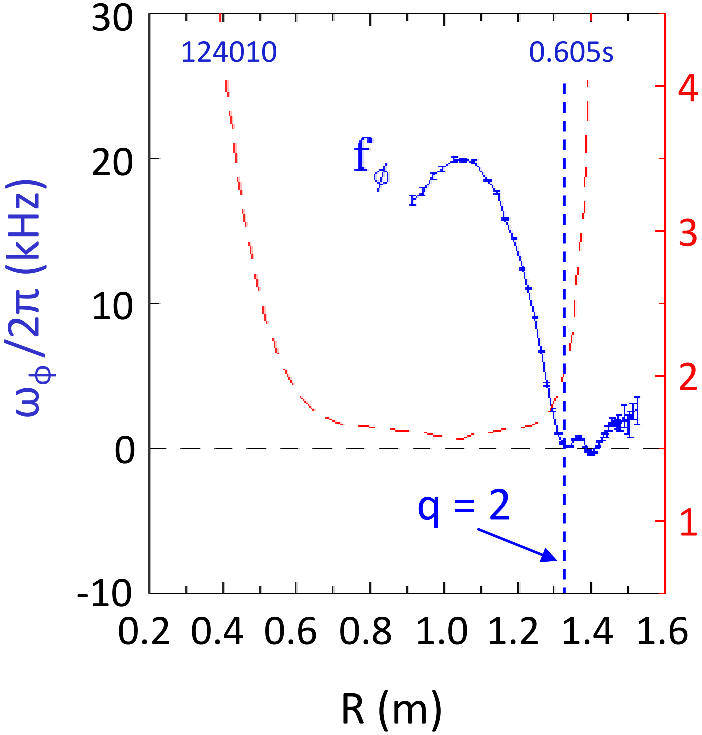
comparable magnitudes, but different Ψ distribution.

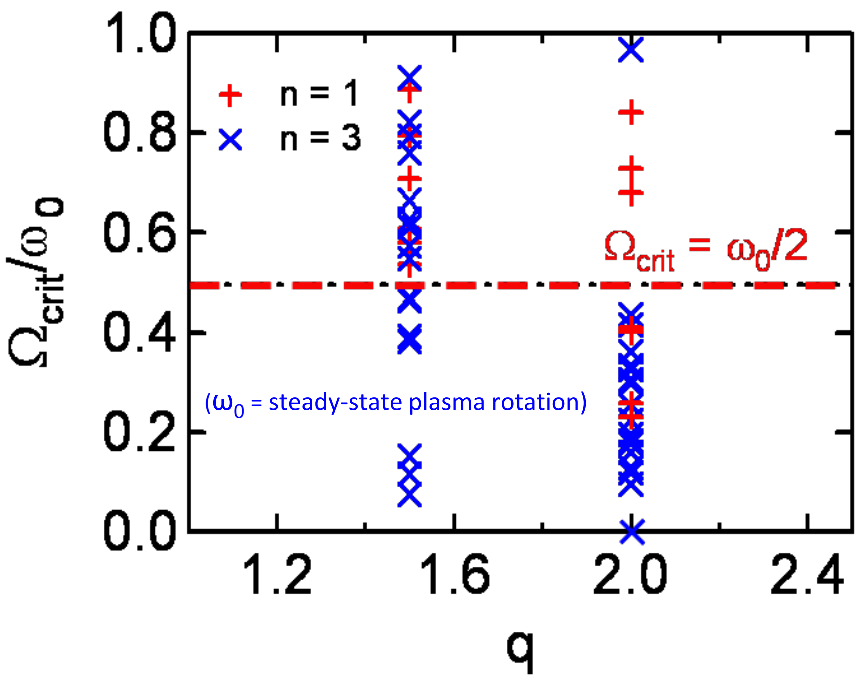


Large ε_a near the edge amplifies any hot ion energy integral contribution. For this not to happen, hot ion density near the edge would have to be very low.

Non-resonant magnetic braking is used to probe RWM stabilization physics

– Scalar plasma rotation at q = 2 inadequate to describe stability.
  - Marginal stability, $\beta_N > \beta_N^{no\text{-}wall}$, with $\omega_\phi^{q=2} = 0$
– Ω_{crit} doesn't follow simple $\omega_0/2$ rotation bifurcation relation.





(Sontag, NF, 2007)