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Active Resistive Wall Mode Stabilization in Low Rotation, High Beta NSTX Plasmas

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11th Workshop on Active Control of MHD Stability
Active MHD Control in ITER

November 7, 2006

PPPL - Princeton, NJ

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Physics understanding and control of pressure-amplified error fields, unstable RWMs reduce performance risks for ITER

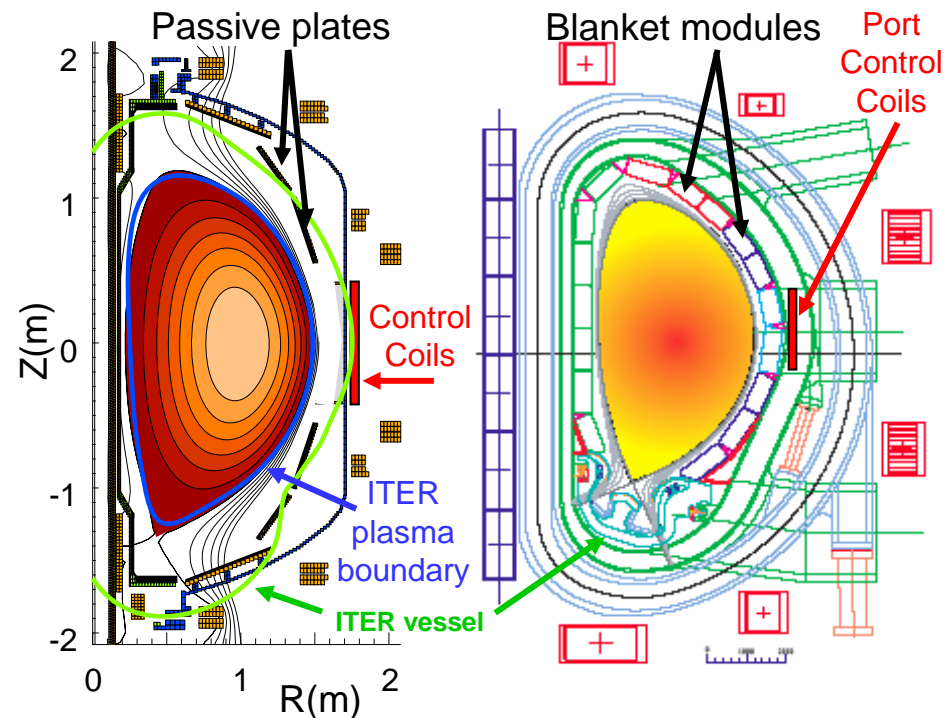
❑ RWM active stabilization

- ❑ RWM control demonstrated
- ❑ RWM actively stabilized in slowly rotating plasmas

❑ Plasma rotation control

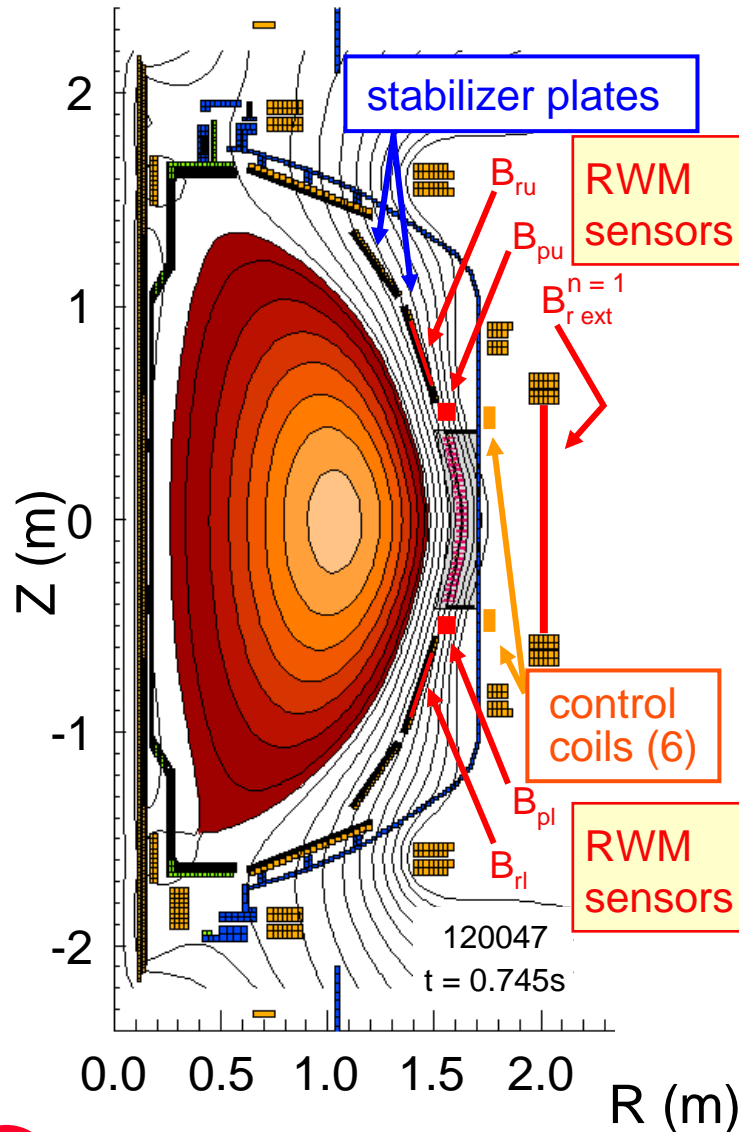
- ❑ Sustained rotation by real-time reduction of amplified error field
- ❑ Reduced rotation by non-resonant magnetic braking
- ❑ Quantitative understanding of momentum dissipation

NSTX / ITER RWM control



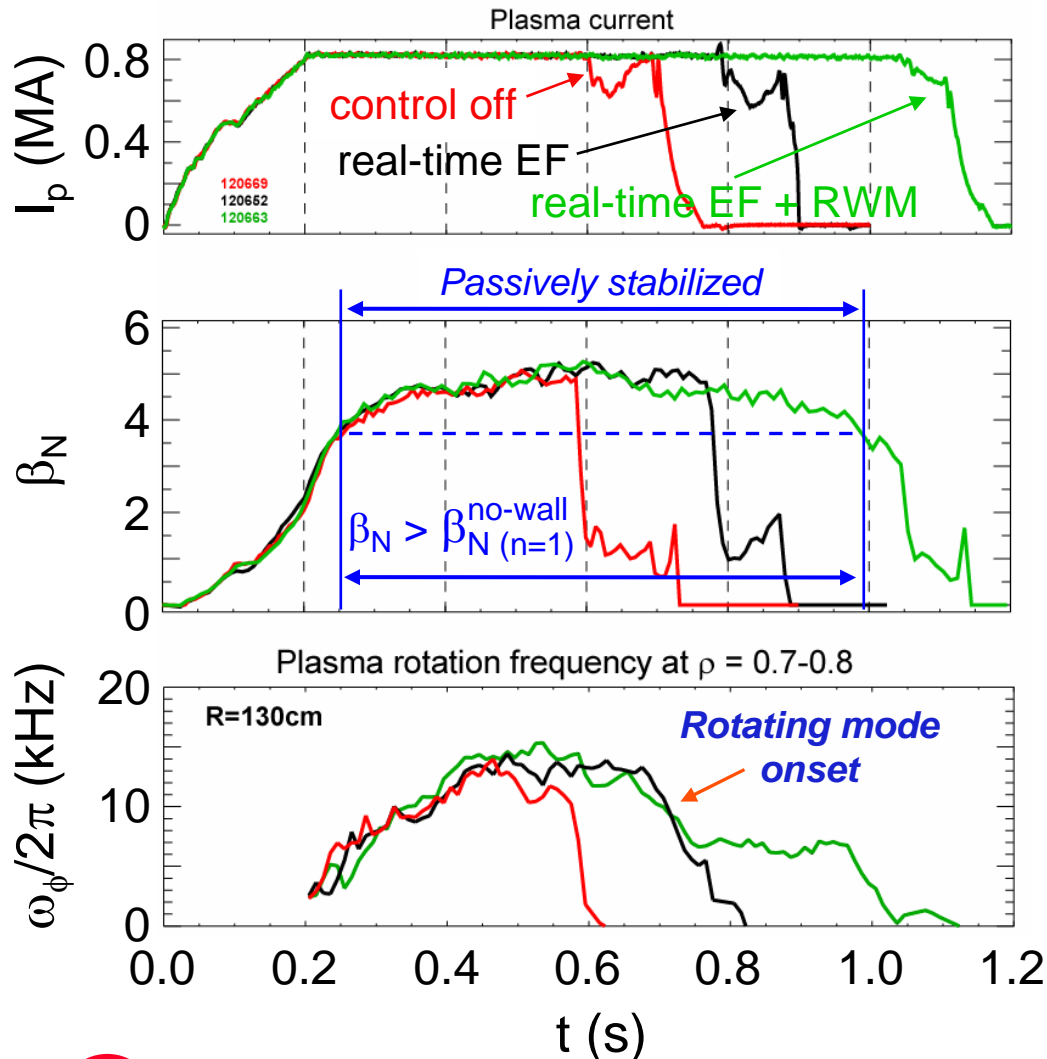
Advantage: low aspect ratio, high β plasmas provide high leverage on theory to uncover key tokamak physics

RWM Active Feedback System Installed on NSTX



- ❑ Stabilizer plates for kink mode stabilization
 - ❑ Operation at $\beta_N / \beta_N^{no-wall} > 1.5$ for pulse $\gg \tau_{wall}$ at highest $\beta_N > 6$
- ❑ External midplane control coils closely coupled to vacuum vessel
- ❑ Internal sensors can detect $n = 1 - 3$ RWM
 - ❑ Unstable $n = 1 - 3$ RWMs already observed in NSTX (Sabbagh, et al., NF 46 (2006) 635.)
 - ❑ $n > 1$ RWM studied during $n = 1$ active stabilization

Dynamic error field correction increases pulse length in strongly rotating plasmas

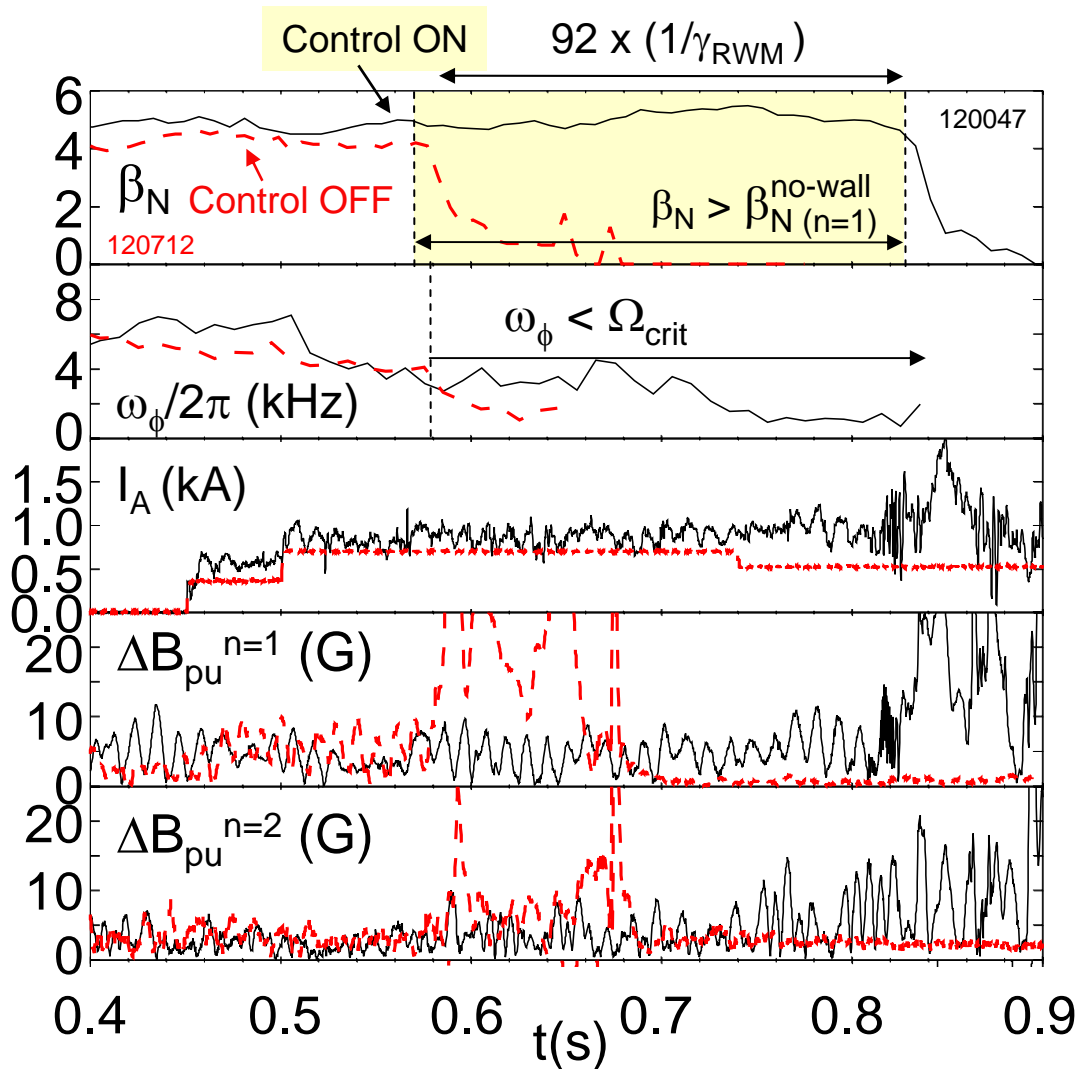


- Real-time correction of known error field (EF)
 - yields higher rotation
 - yields longer pulse

- Combined real-time EF correction + $n = 1$ RWM feedback yielded best result
 - Toroidal rotation, ω_ϕ , increase or saturation at long pulse lengths - first time for NSTX

J.E. Menard, et al., APS 2006

RWM actively stabilized at low, ITER-relevant rotation

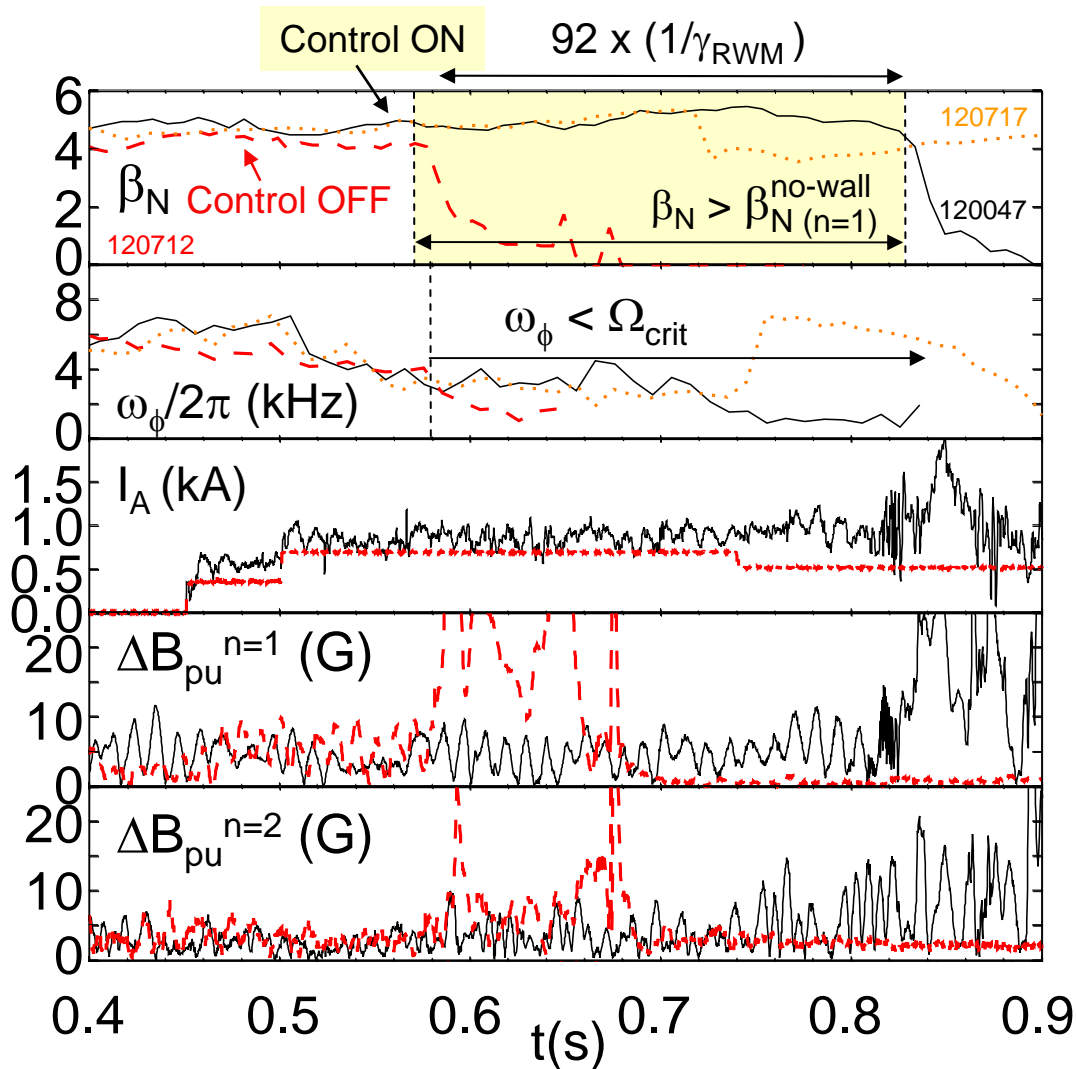


- First such demonstration in low A tokamak
 - Long duration $> 90/\gamma_{RWM}$
 - Exceeds DCON $\beta_N^{no-wall}$ for $n = 1$ and $n = 2$
 - $n = 2$ RWM amplitude increases, mode remains stable while $n = 1$ stabilized
 - $n = 2$ internal plasma mode seen in some cases

- Plasma rotation ω_ϕ reduced by non-resonant $n = 3$ magnetic braking
 - Non-resonant braking to accurately determine RWM critical rotation

(Sabbagh, et al., PRL **97** (2006) 045004.)

RWM actively stabilized at low, ITER-relevant rotation



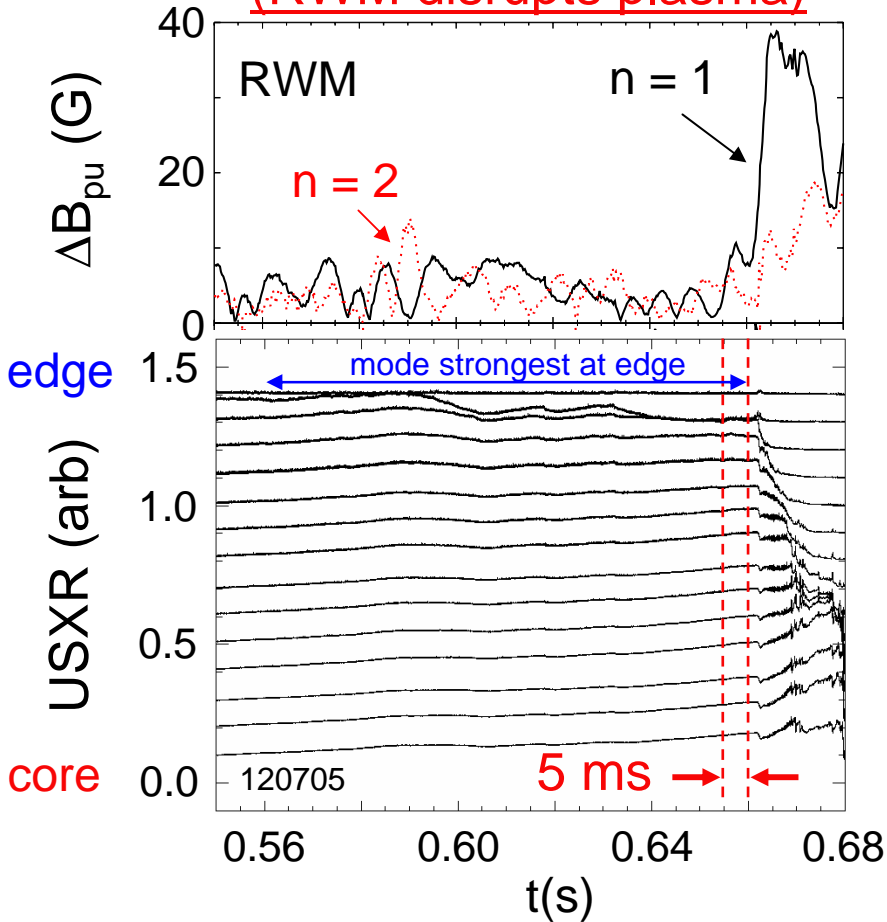
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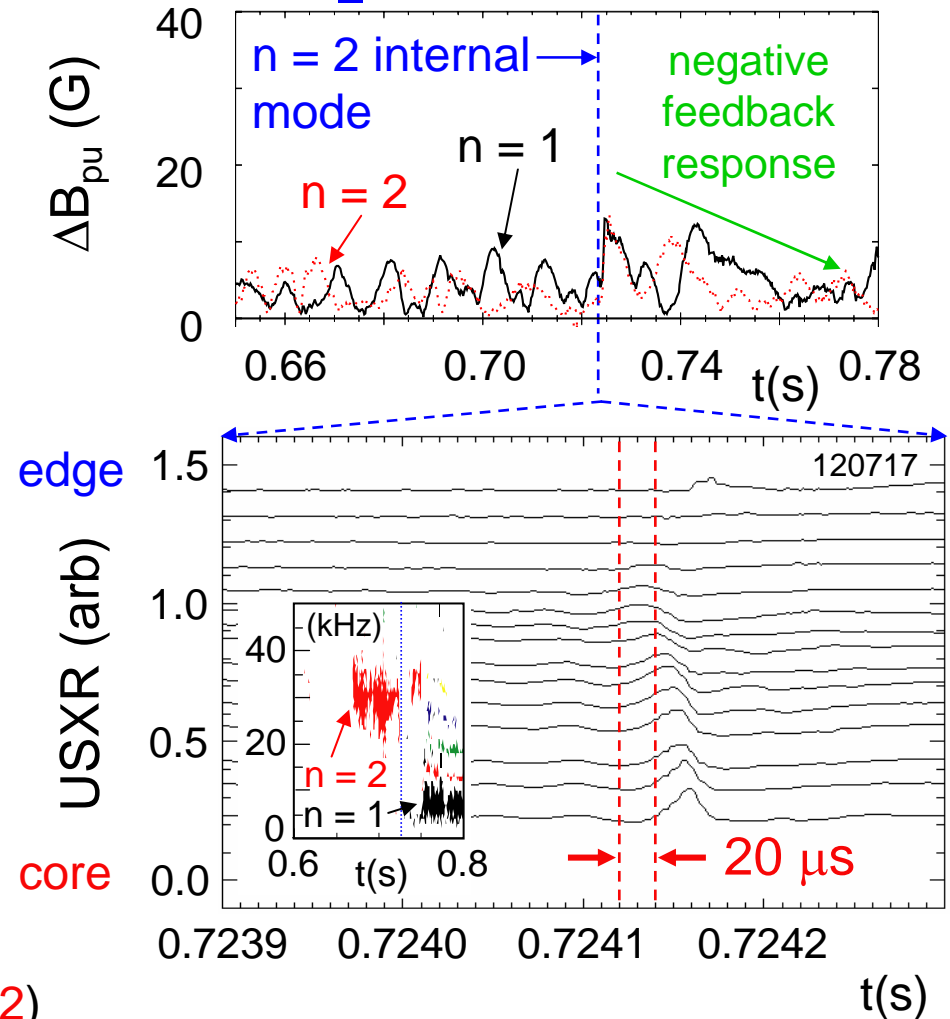
(Sabbagh, et al., PRL **97** (2006) 045004.)

$n = 2$ RWM does not become unstable during $n = 1$ stabilization

Control OFF
(RWM disrupts plasma)



Control ON
(fast β_N drop, plasma recovers)

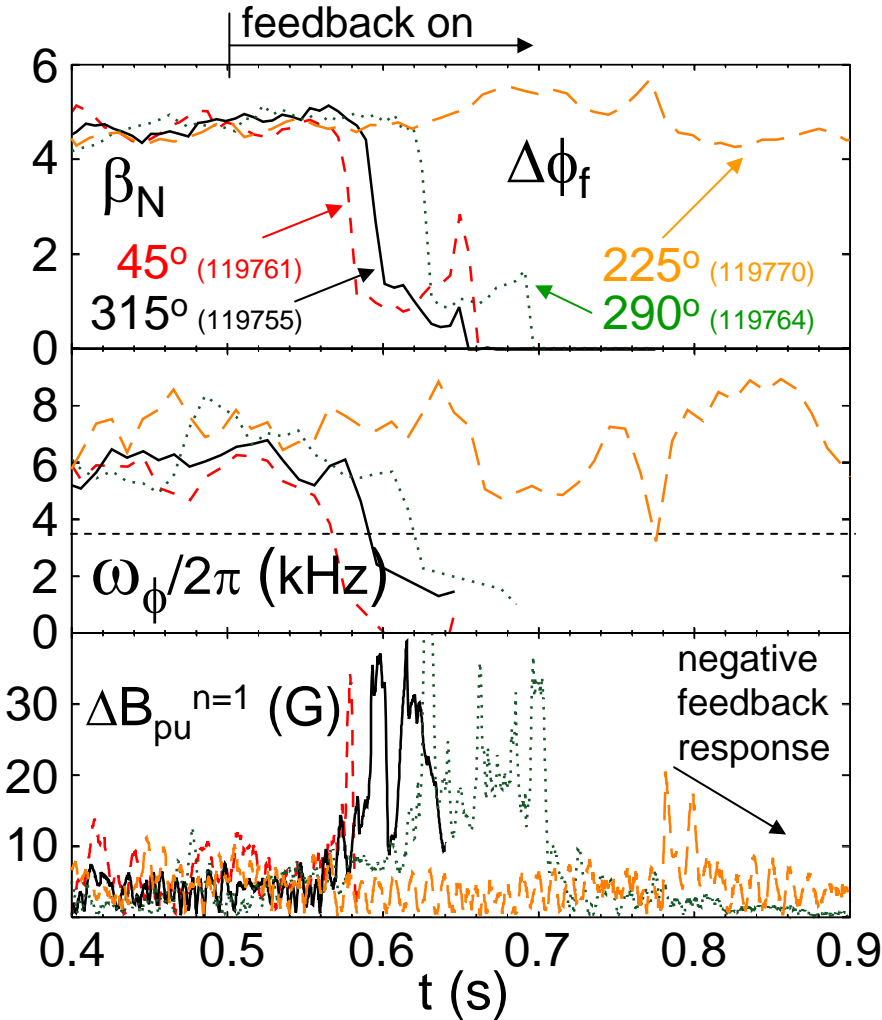


Internal mode ~ 25 kHz

Plasma rotation ~ 12 kHz ($n = 2$)



Varying relative phase shows positive/negative feedback

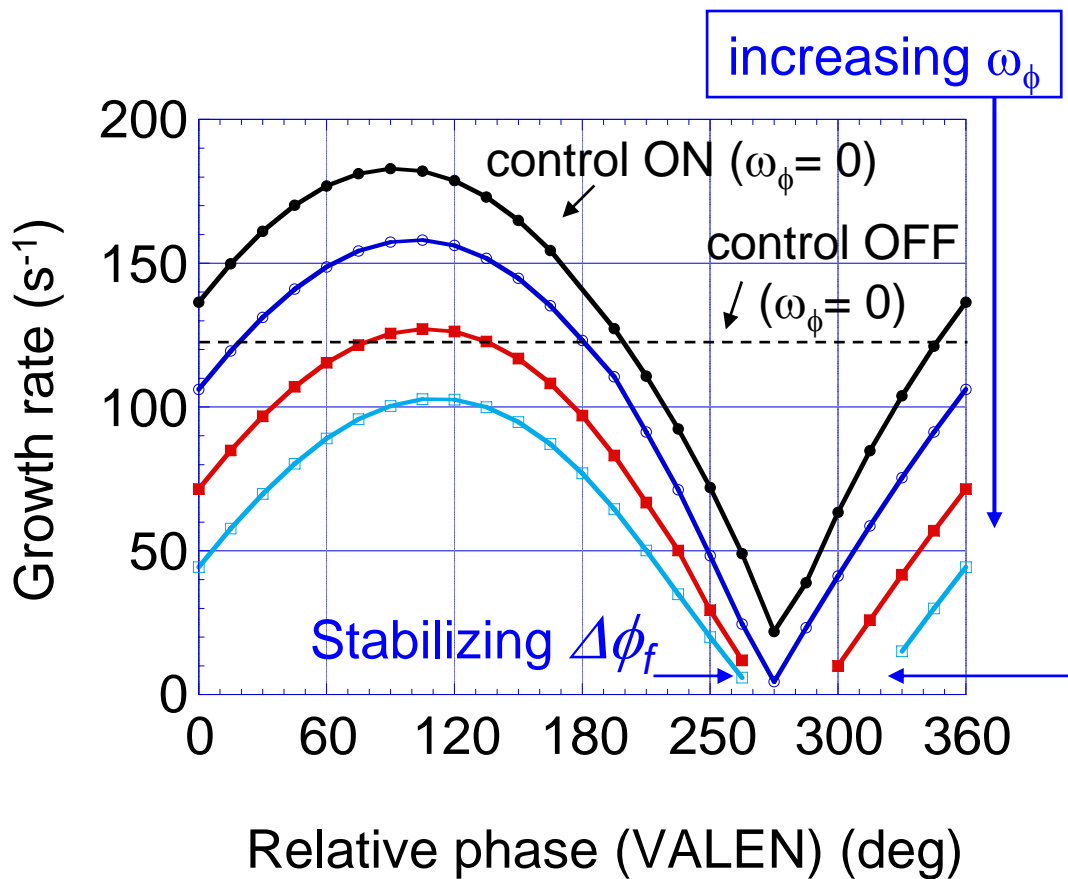


- ❑ Feedback on $n = 1$ RWM
 - ❑ Control current has relative phase $\Delta\phi_f$ to measured ΔB_p

- ❑ Phase scan shows superior settings for negative feedback
 - ❑ Pulse length increases
 - ❑ Internal plasma mode seen at $\Delta\phi_f = 225^\circ$, damped feedback system response

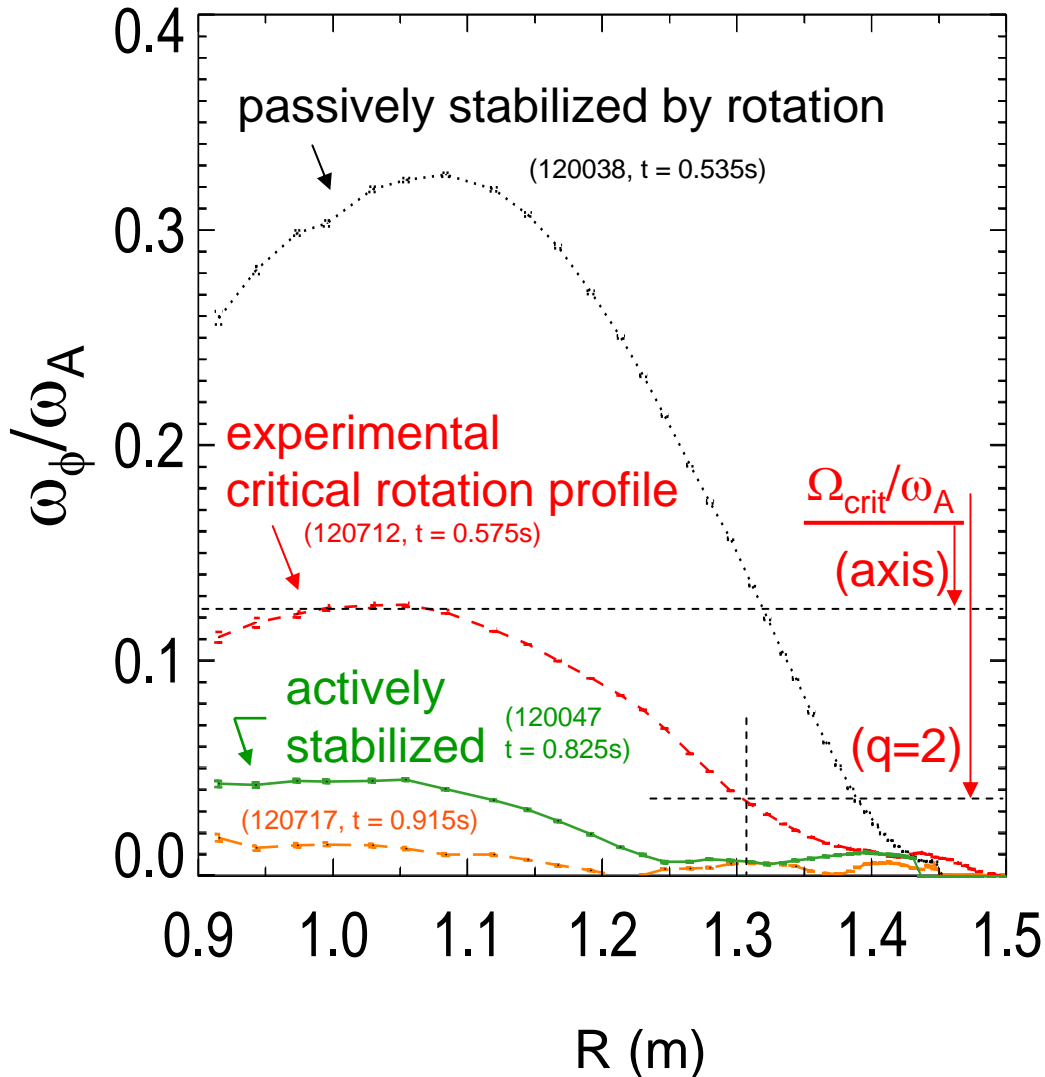
- ❑ Gain scan also performed
 - ❑ Sufficiently high gain showed feedback loop instability

VALEN-3D analysis demonstrates optimal relative phase $\Delta\phi_f$ for RWM active control



- ❑ First VALEN-3D analysis with both active and passive stabilization ($\omega_\phi > 0$)
- ❑ Unfavorable $\Delta\phi_f$ drives mode growth
- ❑ Stable range of $\Delta\phi_f$ increases with increasing ω_ϕ
- ❑ Optimal $\Delta\phi_f$ for active stabilization at $\omega_\phi = 0$ bracketed by results with $\omega_\phi > 0$.

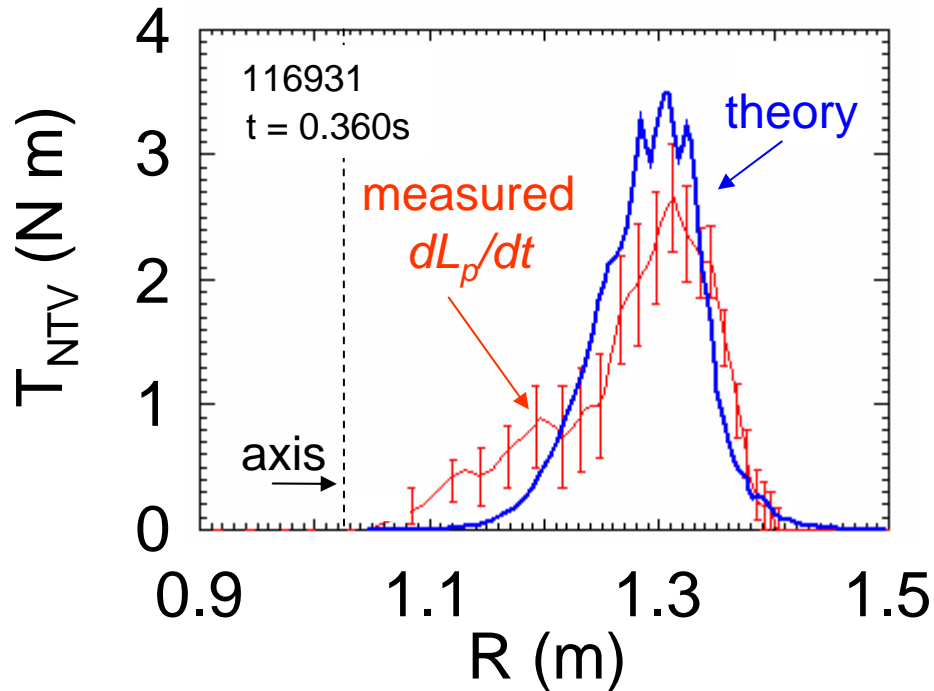
Rotation reduced far below RWM critical rotation profile



- Rotation typically fast and sufficient for RWM passive stabilization
 - Reached $\omega_\phi/\omega_A = 0.48|_{axis}$
- Non-resonant $n = 3$ magnetic braking used to slow entire profile
 - The $\omega_\phi/\omega_A < 0.01|_{q=2}$
 - The $\omega_\phi/\Omega_{crit} = 0.2|_{q=2}$
 - The $\omega_\phi/\Omega_{crit} = 0.3|_{axis}$
 - Less than $\frac{1}{2}$ of ITER Advanced Scenario 4
 $\omega_\phi/\Omega_{crit}$ (Liu, et al., NF 45 (2005) 1131.)
- Rotation profile responsible for passive stabilization, not just single radial location

Observed rotation decrease follows NTV theory

$n = 3$ applied field configuration



(Zhu, et al., PRL **96** (2006) 225002.)

- First quantitative agreement using full neoclassical toroidal viscosity theory (NTV)
 - Due to plasma flow through non-axisymmetric field
 - Computed using experimental equilibria
 - Trapped particle effects, 3-D field spectrum important

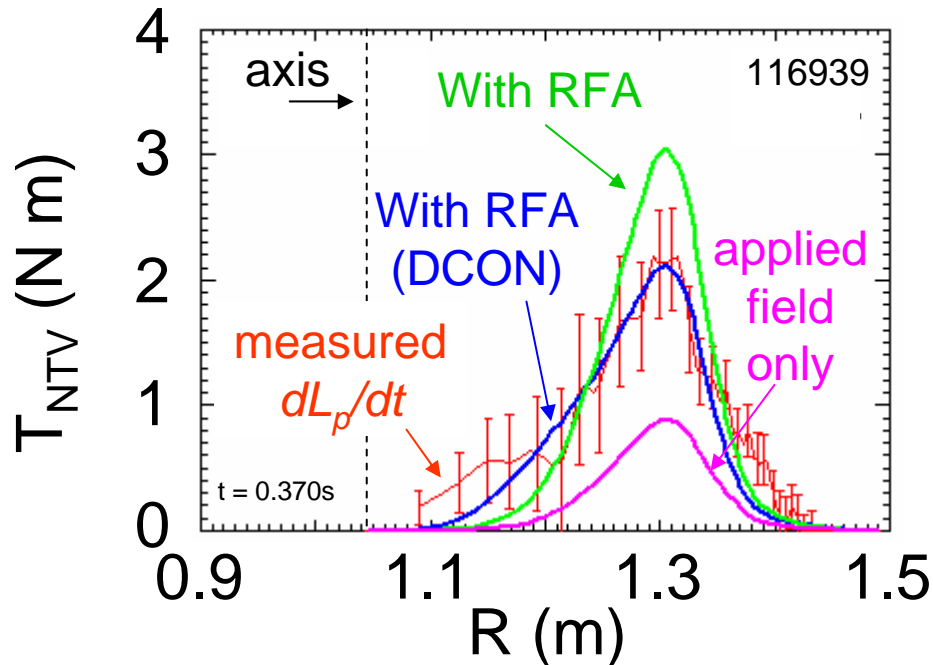
- Viable physics for simulations of plasma rotation in future devices (ITER, CTF)
 - Scales as $\delta B^2 (p_i/v_i) (1/A)^{1.5}$
 - Low collisionality, v_i , ITER plasmas expected to have higher rotation damping



Pressure-driven RFA increases NTV at high β_N

$n = 1$ applied field configuration

$$(\beta_N > \beta_N^{no-wall} = 4.4)$$



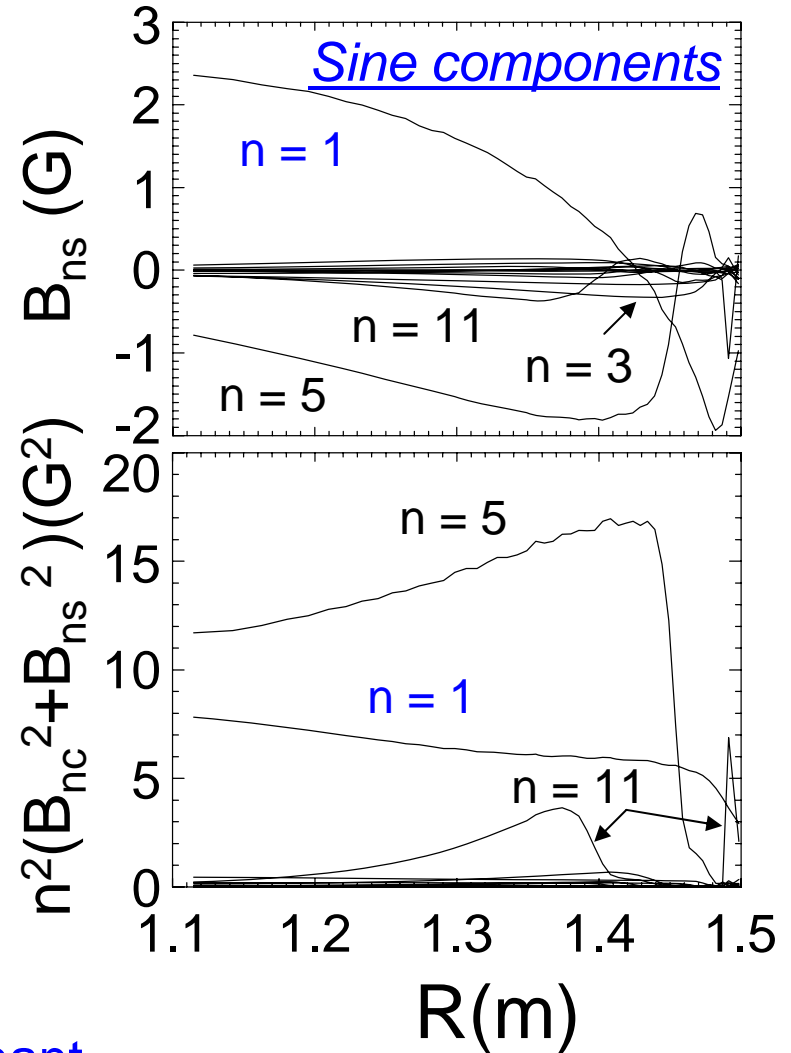
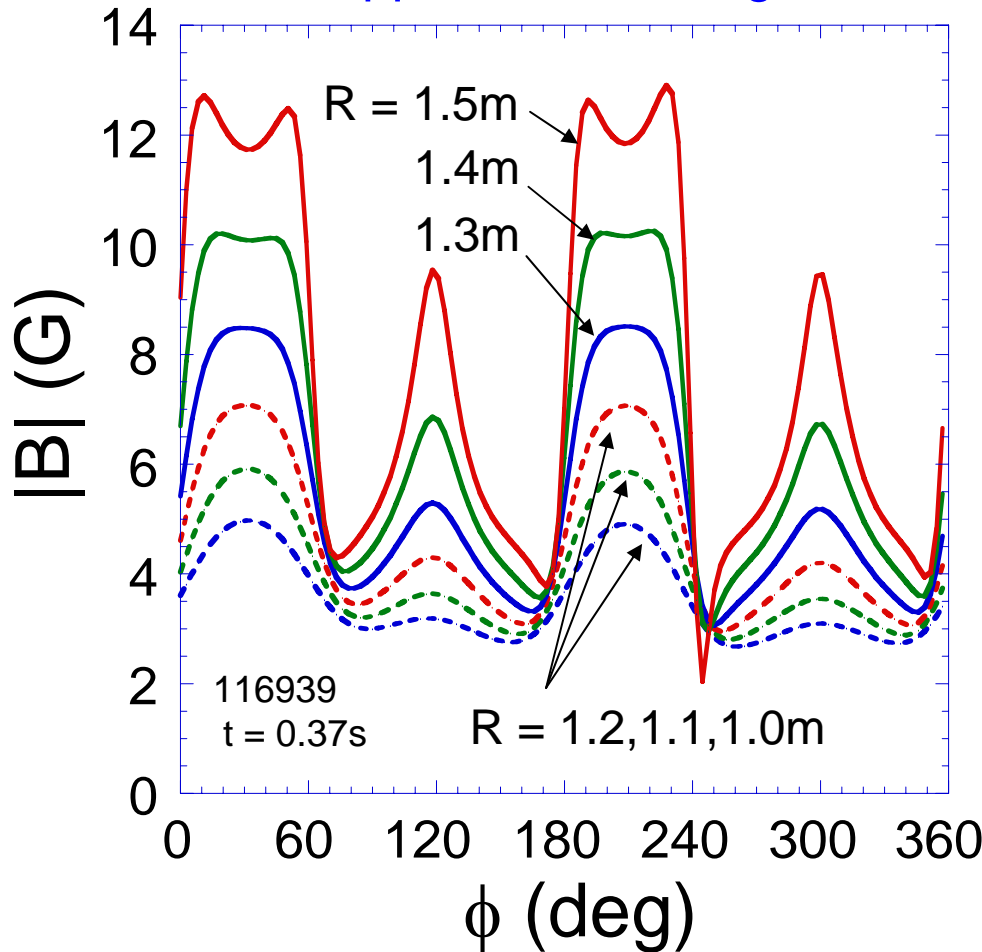
(Zhu, et al., PRL **96** (2006) 225002.)

- ❑ Resonant field amplification (RFA) increases applied field when $\beta_N > \beta_N^{no-wall}$
 - ❑ $n = 1$ has strongest RFA
- ❑ Measured RFA, or unstable RWM field yields increase in NTV rotation damping
 - ❑ Torque due to applied field alone does not match experiment as well
- ❑ Computation based on applied field, or DCON computed mode spectrum
 - ❑ DCON eigenfunction yields slightly broader profile, better experimental agreement



Accurate NTV calculation requires broad applied $|B|$ spectrum

$n = 1$ applied field configuration



□ $n = 3$ configuration has $n = 3, 9$ dominant



NSTX

Ω_{crit} not correlated with Electromagnetic Torque Model

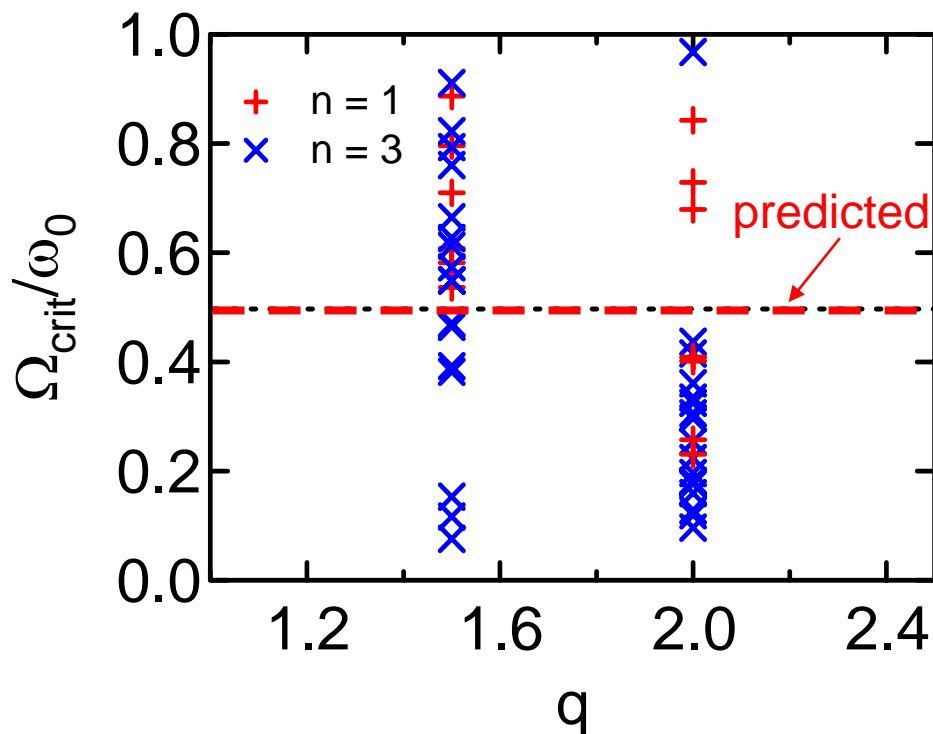
- ❑ Rapid drop in ω_ϕ when RWM unstable may seem similar to 'forbidden bands' model

- ❑ theory: drag from electromagnetic torque on tearing mode*
- ❑ Rotation bifurcation at $\omega_d/2$ predicted

- ❑ No bifurcation at $\omega_d/2$ observed

- ❑ no correlation at $q = 2$ or further into core at $q = 1.5$
- ❑ Same result for $n = 1$ and 3 applied field configuration

NSTX Ω_{crit} Database



($\omega_0 \equiv$ steady-state plasma rotation)

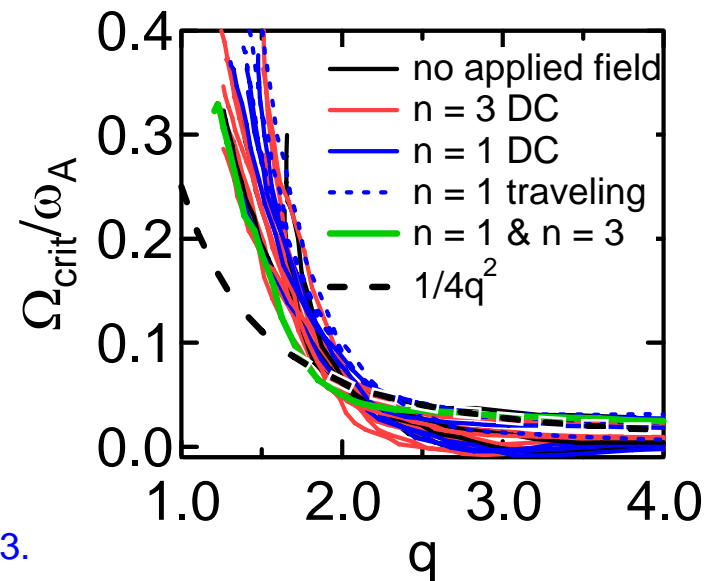
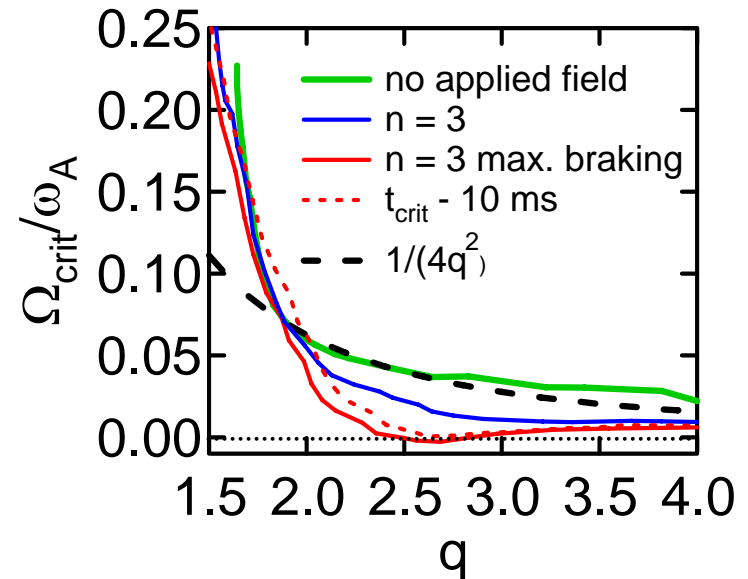
A.C. Sontag, et al., IAEA 2006

*R. Fitzpatrick, Nucl. Fusion **33** (1993) 1061.

Rotation profile shape important for RWM stability

- ❑ Benchmark profile for stabilization is

$$\omega_c = \omega_A / 4q^2^*$$
 - ❑ predicted by Bondeson-Chu semi-kinetic theory**
- ❑ $n = 3$ braking field reduces Ω_{crit} profile
- ❑ High rotation outside $q = 2.5$ not required for stability
 - ❑ Zero rotation at single q can be stable
- ❑ Scalar Ω_{crit} / ω_A at $q = 2, > 2$ not a reliable criterion for stability
 - ❑ consistent with distributed dissipation mechanism

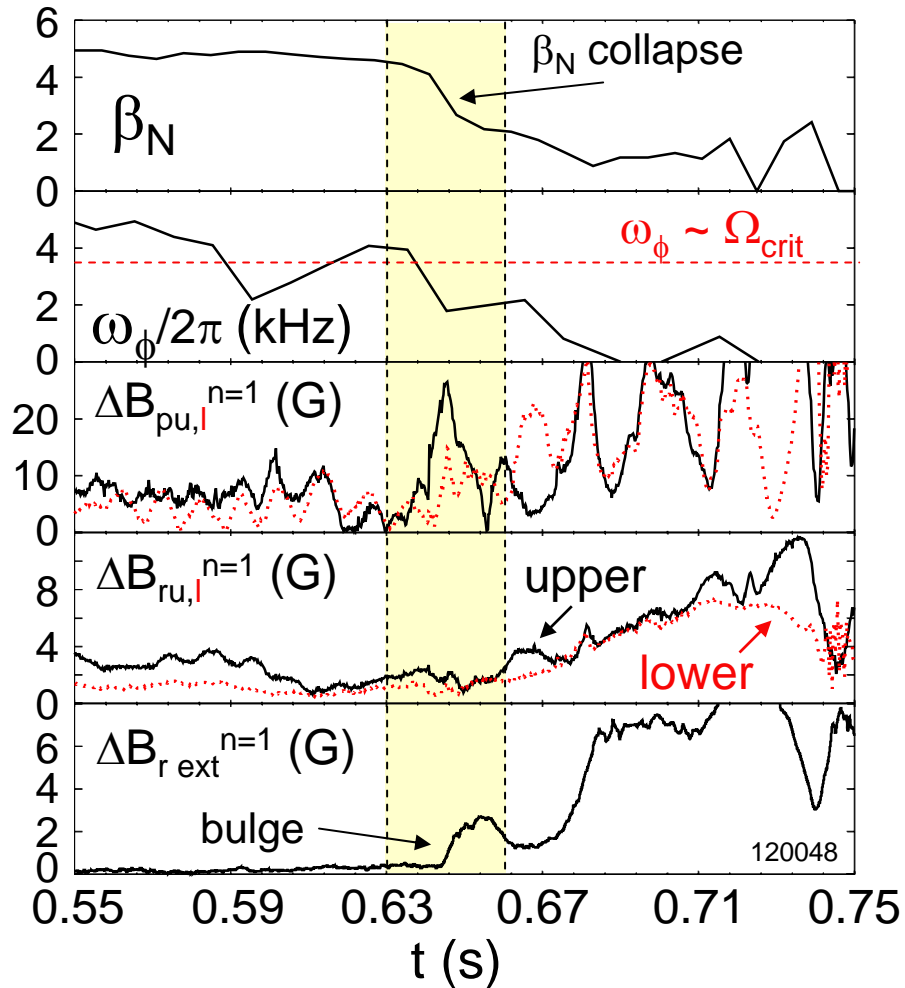


*A.C. Sontag, et al., Phys. Plasmas **12** (2005) 056112.

A. Bondeson, M.S. Chu, Phys. Plasmas **3 (1996) 3013.



RWM may change form and grow during active control



- ❑ Poloidal $n = 1$ RWM field decreases to near zero
 - ❑ Radial field increasing
- ❑ Subsequent growth of poloidal RWM field
 - ❑ Asymmetric above/below midplane
- ❑ Radial sensors show RWM bulging at midplane
 - ❑ midplane signal increases, upper/lower signals decrease
 - ❑ Theory: may be due to other stable ideal $n = 1$ modes becoming less stable

Future research will assess using combined sensors for optimization

NSTX begins RWM active stabilization research relevant to ITER, CTF and beyond

- ❑ First demonstration of RWM active stabilization in high β , low A tokamak plasmas with ω_ϕ significantly less than Ω_{crit}
 - ❑ In the predicted range of ITER
 - ❑ Positive and negative RWM feedback demonstrated by varying feedback gain and relative phase
- ❑ Stability of $n = 2$ RWM observed during $n = 1$ RWM stabilization
 - ❑ $n = 1, 2$ plasma mode sometimes observed; fast β collapse, recovery
- ❑ Plasma rotation reduction by non-resonant applied field; follows neoclassical toroidal viscosity theory
 - ❑ Full NTV calculation yielding quantitative agreement to experiment ; general momentum transport relevance
- ❑ Results continue to support Ω_{crit} as profile; scalar insufficient

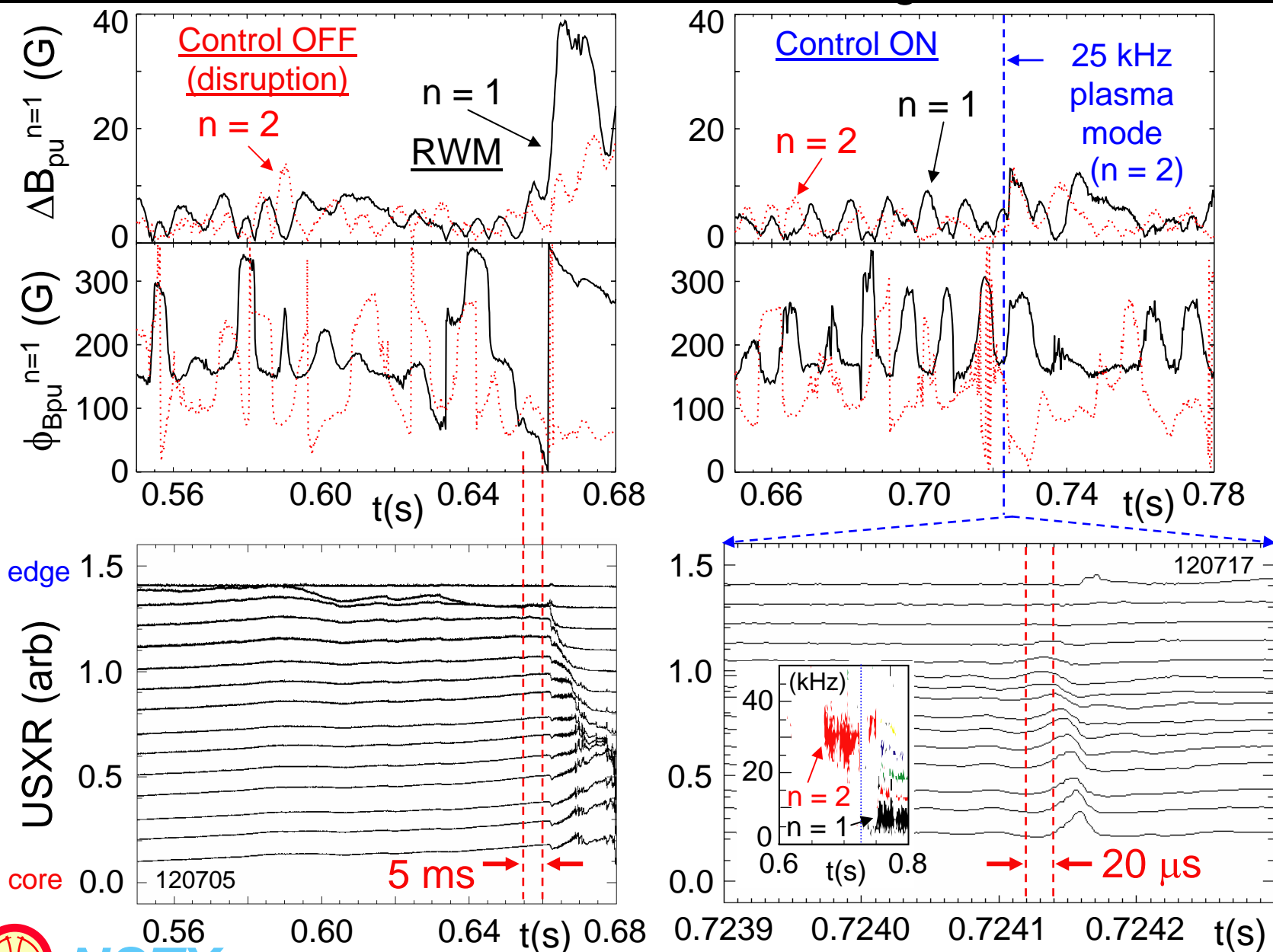
Active RWM control remains vitally important to minimize performance risk in ITER

- ❑ Agreement does not yet exist regarding Ω_{crit} , RWM stabilization physics in low-rotation plasmas
- ❑ Observed inverse dependence of Ω_{crit} on v_i indicates lower ITER collisionality may require a higher degree of RWM active stabilization (see talk by A. Sontag, this meeting)
- ❑ Similar inverse dependence of plasma momentum dissipation on v_i in NTV theory indicates ITER plasmas will be subject to higher viscosity, greater ω_ϕ reduction
- ❑ Strong δB^2 dependence of quantitatively verified NTV theory shows that error fields, resonant field amplification, ELMs need be minimized to maximize stabilizing ω_ϕ
- ❑ Pressure, q , and ω_ϕ profiles unknown for burning plasma. RWM (and ELM, error field) control reduces performance risk

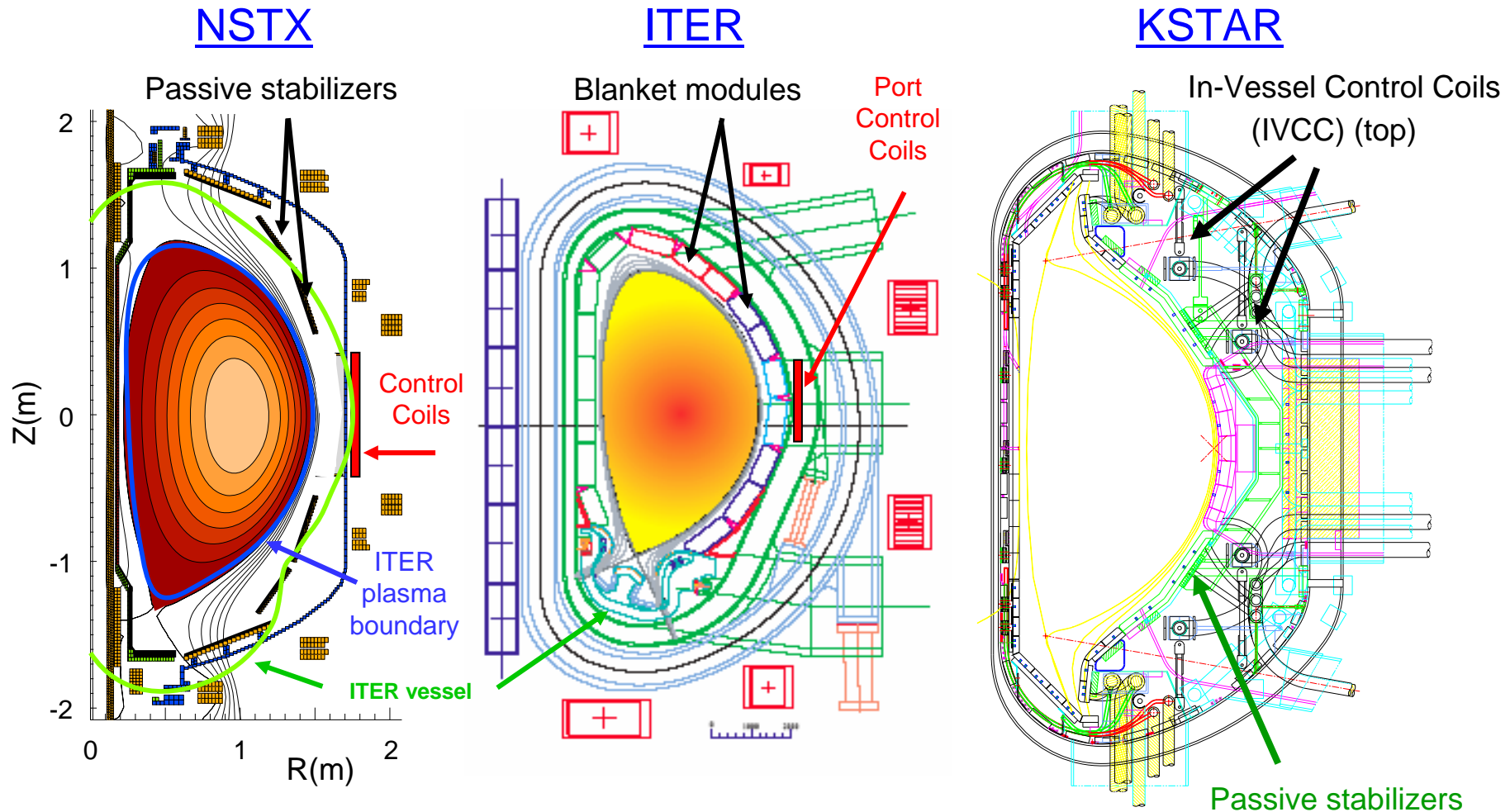
Extra slides



$n = 2$ RWM does not become unstable during $n = 1$ stabilization



NSTX global MHD mode stabilization hardware similar to designs for next-step tokamaks



□ ITER: Midplane port control coils, blanket

□ KSTAR: Midplane control coil, passive plate geometry with midplane gap



NSTX