Plans for MHD control in JT-60SA based on JT-60U results

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Outline

- 1. Concepts of JT60-SA
- 2. MHD control of JT60-SA
 - 1. Plasma shape control
 - 2. Internal profile control
 - 3. Control by ECCD: ST, NTM
 - 4. Control by coils: RWM
- 3. Summary

Mission of JT-60SA Program

JT-60SA (JT-60 Super Advanced) is a combined program of

ITER Satellite Tokamak Program of JA-EU
Japanese National Program



ITER : India, China, Korea, US, Russia, EU, Japan



Mission of JT-60SA is to support and supplement ITER toward DEMO.

Contributors to JT-60SA Program

JA-EU satellite Tokamak Working Group: S. Matsuda¹, M. Kikuchi¹, Y. Takase², Y. Miura¹, M. Matsukawa¹, S. Sakurai¹, F. Romanelli¹⁵, J. Pamela¹⁵, D. Campbell¹⁵, C. Sborchia¹⁶, J.J. Cordier¹⁷, S.Clemento¹⁸, 1) Japan Atomic Energy Agency, 2) U. Tokyo, 3) EFDA, 4) IPP-MGP, 5) CEA Cadarache, 6) CEC Brussels

JT-60SA Design Team: M. Kikuchi¹, M. Matsukawa¹, H. Tamai¹, S. Sakurai¹, K. Kizu¹, K. Tsuchiya¹, A. Sukegawa¹, Y. Kudo¹, T. Ando¹⁾, H. Matsumura¹⁾, F. Sato¹⁾, Y. Miura¹⁾, S. Ishida¹⁾, T. Fujita¹⁾, G. Kurita¹⁾, Y. Sakamoto¹⁾, A. Oikawa¹⁾, T. Yamazaki¹⁾, H. Kimura¹⁾, K. Shinohara¹, S. Konoshima¹, T. Ozeki¹, O. Naito¹, T.Takizuka¹, K.Hamamatsu¹, K.Shimizu¹, K, Ohasa¹, N.Hayashi¹, N. Aiba¹, K.Kiyono¹, T.Oshima¹, S.Sakata¹, M.Sato¹, Y. Kamada¹, H. Kubo¹, Y. Koide¹, S. Ide¹, N. Asakura¹, H. Takenaga¹, K. Hoshino¹, H. Kawashima¹, T. Hatae¹, A. Isayama¹, M. Takechi¹, T. Suzuki¹, T. Nakano¹, N. Oyama¹, K. Kamiya¹, H. Urano¹, Y. Kawano¹, T. Kondo¹, G. Matsunaga¹, M. Yoshida¹, K. Fujimoto¹, Y. Kojima¹, Y. Tsukahara¹, H. Sunaoshi¹, S. Kitamura¹, Y. Kashiwa¹, S. Chiba¹, Y. Ishii¹, T. Matsumoto¹), Y. Idomura¹), N. Miyato¹), S. Tokuda¹, K. Tobita¹), Y. Nakamura¹), M. Sato¹), S. Nishio¹), N. Hosogane¹), N. Miya¹), T. Yamamoto¹), K. Kurihara¹), K. Shimada¹), T. Terakado¹), T. Shibata¹), H. Oomori¹), J. Okano¹), H. Furukawa¹), Y. Terakado¹), K. Shibata¹), T. Matsukawa¹, A. Sakasai¹, K. Masaki¹, N. Hayashi¹, T. Sasajima¹, J. Yagyu¹, Y. Miyo¹, N. Ichige¹, Y. Suzuki¹, H. Takahashi¹, T. Fujii¹, S. Moriyama¹⁾, M. Seki¹⁾, Y. Ikeda¹⁾, M. Kawai¹⁾, N. Akino¹⁾, M. Hanada¹⁾, N. Ebisawa¹⁾, M. Kazawa¹⁾, F. Okano¹⁾, M. Kamada¹⁾, K. Usui¹⁾, A. Honda¹⁾, M. Komata¹⁾, K. Mogaki¹⁾, M. Kuriyama¹⁾, H. Ninomiya¹⁾, M. Inutake⁶⁾, N. Yoshida³⁾, Y. Takase²⁾, K. Nakamura³⁾, M. Sakamoto³⁾, M. Ichimura⁵, T. Imai⁵, Y. M. Miura⁷, H. Horiike⁷, A. Kimura⁸, H. R. Shimada⁹, H. Tsutsui⁹, M. Matsuoka¹⁰, Y. Uesugi¹¹, K. Ida¹², A. Sagara¹²⁾, A. Nishimura¹²⁾, A. Shimizu³⁾, K. Sato³⁾, Hashizume⁶⁾, K. Okano¹³⁾, Y. Kishimoto^{1,8)}, H. Azechi⁷⁾, S. Tanaka²⁾, K. Yatsu⁵⁾, S. Itoh³⁾, M. Fujiwara^{1,12}, M. Akiba¹, K. Okuno¹, R. Andreani¹⁴, J. Bialek⁴, G. Navratil⁴ 1) JAEA, 2) U. Tokyo, 3) Kyushu U., 4) Columbia University, 5) U. Tsukuba, 6) Tohoku U., 7) Osaka U., 8) Kyoto U., 9) Tokyo Institute of Technology, 10) Mie U., 11) Kanazawa U., 12) NIFS, 13) CRIEPI, 14) IPP-MGP

1st satellite tokamak WG at NAKA



National Review Committee



World families of fully superconducting tokamaks



Basic Machine Parameter of JT-60SA



JT-60SA is optimized for contribution to ITER

EF-7

Lower divertor

- 1. ITER-shape plasma is produced by EF-7 coil .
- 2. Almost same Greenwald density of ITER
- 3. Lower divertor to match ITER tri-angularity.
- 4. ITER type mono-block divertor for outer target.
- 5. Electron heating with 10MW N-NB+7MW ECRF
- 6. Co-CTR rotation control by T-NBI

7. NTM control at two frequency (110GHz&140GHz)

Parameter	ITER	JT-60SA
Plasma Current Ip	15 MA	3.5MA
Toroidal Field Bt	5.3T	2.59T
Major Radius Rp	6.2 m	3.16 m
Minor Radius a	2.0 m	1.02 m
Aspect Ratio A	3.1	3.1
Elongation κ_{95}	1.70	1.7
Triangularity δ_{95}	0.33	0.33
Safety Factor q ₉₅	3.0	3.0
Greenwald density n _g	1.2x10 ²⁰ m ⁻³	1.1x10 ²⁰ m ⁻³

 $P=41MW >> P_{th}^{L-H} \sim 20MW$

Supplemental role of JT-60SA for DEMO

High beta steady-state operation



- 2. Upper divertor to match high tri-angularity.
- 3. Down-shifted N-NBCD to form reversed shear.
- 4. Stabilizing plates (SP) for RWM control
- 5. 6 set of 3 poloidal n=1,2 RWM control coils
- 6. Ferritic steel on SP to simulate DEMO



Latest Design of Torus Hall Layout



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Plasma shape control

New equilibrium coil



EF1 and EF6 have been moved away from the mid-plane for off-axis N-NB and machine maintenance

⇒decrease in 'squareness' and triangularity controllability



Plasma shape control (2) Divertor geometry



For ITER-like shape



Stability near the edge: ELM (high-S case)



- In JT-60SA, the shaping parameter S [8] will reach to about 6.
- By changing the pressure and current profiles near the pedestal, the stability of tokamak edge plasma is investigated.
- The range of *n* of the MHD modes to be analyzed is from 1 to 60, and infinity (ballooning mode stability).
- In this strong shaping and small aspect ratio equilibrium, the stability of ideal ballooning modes will be more stable than that in ITER-like Eq.

Stability Analysis Result (high-S)



- In high-S Eq., the maximum pressure gradient (=5.61) becomes larger than that in ITER-like Eq.(=4.36), as expected.
- The n number of the peeling-ballooning mode determining the α limit is about 20, which is larger than that in ITER-like Eq.

Equilibrium of JT-60SA (ITER-like Eq.)



- The profiles of dp/d ψ and (j_{//}) are same as those in high S Eq.
- The effect of the sharpness is also investigated.
- The range of *n* of the MHD modes to be analyzed is from 1 to 60, and infinity (ballooning mode stability).
- By changing the shape at the top, we investigate the stabilizing effect of the sharpness σ . ($\sigma = (1/r_c)/(1/a)$, r_c : curvature radius, *a*: minor radius)

Stability Analysis Results (ITER-like)



- In this ITER-like equilibrium, the maximum pressure gradient (α₉₄=4.36) is determined by the stability of peelingballooning modes whose *n* are from 8 to 13.
- The marginally unstable eigenfunction, whose n=8, localizes near 0.9<s<1.0 (0.86<ρ_{vol}<1.0).

Effect of the Sharpness in ITER-like Eq.



- As mentioned in [5], the sharpness σ has an impact on the stability of the ballooning mode and the peeling ballooning mode.
- In this case, the stable limit of α increases from 4.36 to 4.76 as σ_{up} changes from 2.04 to 3.94.
- The *n* number and the width of the eigenfunction determining the α limit is similar to those of the eigenfunction in ITER-like Eq.

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Profile control by H&CD systems



Profile control by N-NB system

Since q_{min}>1 or suppression of sawtooth instability is believed necessary to achieve high beta and high confinement, it has been decided to shift the N-NBI beam lines downward, to obtain a broad current profile with weak/reversed magnetic shear



Stability improvement by profile control

The q_{min} increases with a larger shift of N-NB; q_{min} ~1.6 for $Z_{NNB} = -0.6$ m and ~2 for $Z_{NNB} = -0.9$ m (2.4 MA full CD, A=2.8, $Z_{axis}=0.07$ m, β_N ~4, f_{BS}~70%).

The ideal wall limit increases with q_{min} .

The power deposition of N-NB gets hollows as the N-NB position is moved downward.

Considering the MHD stability and central heating, the shift of 0.6 m is taken as a standard point for the present design.



N-NB system(3)

Far off-axis current drive for achieving qmin ~2 is possible by using a USN configuration with its magnetic axis elevated

Near on-axis heating is also possible by using a LSN configuration with the axis lowered.





Extended region for full-CD operation

A = 2.65 and a shift of 60 cm of N-NB (Z_{NNB} = -0.6 m). A full CD operation with a high normalized density of ne-bar/n_{GW} = 0.86 (= 4.9x10¹⁹ m⁻³) is possible at I_p = 2.4 MA for HH_{y2} =1.33 and total power of 41 MW. The fraction of the bootstrap current ('BS') is f_{BS} = 0.70. The resultant q profile has a broad weak shear region with q_{min} ~ 1.5.

A, q95, κ , β_N , β , HHy2 and fBS are very close (within 10%) to those in a DEMO design with a slim center solenoid proposed by JAEA



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ECRF System in JT-60U



Antenna for EC #1-3

Antenna for EC #1-3



The real-time NTM stabilization system: the detection of the mode location and the optimization of the injection angle of electron cyclotron (EC) wave are performed in real time



NTM has been completely stabilized in JT-60U m/n=2/1 by ECCD at q=2 (ρ~0.6)



Further optimization of ECCD location enabled complete stabilization at J_{EC}/J_{BS}~0.5

different EC location & current



TOPICS simulation with modified Rutherford equation well reproduces experimental results.





The consistent analysis shows: • ECCD width has stronger effect than amount of EC-driven current. • Precise ECCD control has enabled complete stabilization with smaller value of j_{EC}/j_{BS} : $J_{EC}/J_{BS} \sim 0.5$ [A. Isayama, IAEA FEC 2006]

ECRF system of JT-60SA



140 GHz, 4 MW for high-density plasmas 3rd resonance (small CD) appears below 2.2 T

110 GHz, 3 MW: ECCD for 1.7-2.5 T

A wide range of Bt covered by two frequencies.



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Installation of ferritic steel tiles in JT-60U

- Reduction of ripple loss decreased fast ion loss in the large volume plasma and increased the net heating power.
- The region of high β_N discharges was extended and integrated performance in long high β_N plasmas was improved by the increase of the net heating power, the plasma close to the wall, and the change of the rotation.



β_N reaches ideal wall limit.

- High β_p ELMy H-mode plasma : B_T =1.58 T, I_p =0.9 MA, δ_0 ~20 cm (d/a=1.2)
- Increase in net heating power due to the FSTs installation allows to access high β_N up to 4.2 with I_i =0.8-1.
- n=1 (m~3) mode at high beta region.
- Growth time of $1/\gamma \sim 1$ ms (< $\tau_w \sim 10$ ms) before collapse.
- RWM is suppressed by plasma rotation (100km/s at r/a=0.3).



Small critical rotation velocity of V_C/V_A~0.3% is found for suppressing RWM.

 Less counter rotation due to the FSTs installation enables to change the rotation in co-direction close to zero.



RWM control coils

RWM stabilisation by feedback control of sector coils (VALEN code analysis^{*})



18 coils, 20 kAT, 1 G (10-4 T) of m=3, n=1 component of radial magnetic field.

The result of VALEN analysis for n=1 modes.

A = 2.8,

 $r_w/a = 1.3$ on the outer mid-plane.

p(r) and j(r) consistent with the ACCOME analysis, where $q_{min} \sim 2.1$.

 $\beta_N = 4.3$ is expected, which is close to the ideal wall limit ($\beta_N^{ideal-wall} = 4.42$), while the no-wall limit is $\beta_N^{no-wall} = 2.56$. Hence $C_\beta = (\beta_N - \beta_N^{no-wall})/(\beta_N^{ideal-wall} - \beta_N^{no-wall}) = 0.9$ and very efficient RWM stabilization is expected.



JT-60SA is planed as a largest SC tokamak to support and supplement ITER toward DEMO.

Design on some components in JT-60SA including poloidal field coils, RWM control coils, ECRF and N-NB systems, have been optimized for plasma control toward high beta steady-state operation.

Shape control: an additional equilibrium field coil of EF7 Profile control: lowered beam line of N-NB and ECRF NTM control: two frequencies ECRF RWM control: coils along the port hole have been decided.

The construction of JT-60SA will take 7 years and 3 years of experiments are foreseen in 10 years of BA period.

RWM is suppressed by plasma rotation

JT-60U =

- β_N is kept constant and change the tangential NB from ctr-NB to co-NB.
- Rotation can be controlled by changing tangential NB combination
- Disruption or collapse occurs at Vt~10 km/s ->n=1 mode grow with $1/\gamma \sim 10$ ms .
- The mode suppressed after $\beta_N < \beta_{N no-wall limit}$
- To investigate the effect of beta on critical rotation, we change the constant β_N.





Scientific Mission (II) (long term)

Expand operation regime of high beta steady-state for DEMO Long pulse operation : ~8 hours

- 1. Particle exhaust for long pulse and development of DEMO relevant PFC
- 2. Development of reliable operation and reduction of disruption probability



Overview of the MARG2D code

- MARG2D solves the eigenvalue problem associated with the 2-D Newcomb equation with the artificial weight function[1].
- With the vector potential method for calculating the vacuum energy, MARG2D realizes to identify the stability of the wide *n* range ideal MHD modes[2,6].
- By using both the artificial weight function and the physical inertia, the marginal stability condition can be identified, and the compression-less growth rate can be estimated.

$$N\mathbf{Y} = -\lambda R\mathbf{Y}, \quad (\mathbf{Y} = r)$$
$$R_{l,m} \propto \begin{cases} (m/q - n)^2 \dots l = m \\ 0 \qquad l \neq m \end{cases}$$

m : poloidal mode number

- *n* : toroidal mode number
- q: safety factor

 $(\xi_r) \quad \begin{array}{l} N : \text{Newcomb operator} \\ \xi_r : \text{displacement in the } r \text{ direction} \end{array}$

$$\blacktriangleright R\mathbf{Y} = \rho(\mathbf{Y} + (\mathbf{V} - \beta \mathbf{Y}))$$

- **V** : displacement vector in the $\nabla \psi \times \mathbf{B}$ direction
- ρ : plasma density
- β : unorthogonality of the coordinate

[6] N.Aiba et al., TH/P8-1, 21st IAEA FEC, Chengdu, China (2006).