

# Error field scalings and sideband effects - an update

Richard Buttery<sup>\*</sup>, Y Gribov<sup>3</sup>, D. F. Howell<sup>\*</sup>,  
T. C. Hender<sup>\*</sup>, R. J. La Haye<sup>1</sup>, J. T. Scoville<sup>1</sup>,  
COMPASS-D<sup>\*</sup> and DIII-D<sup>1</sup> teams,  
and JET-EFDA contributors<sup>2</sup>.

<sup>\*</sup>EURATOM/UKAEA Fusion Association, Culham Science Centre, United Kingdom.

<sup>1</sup> General Atomics, San Diego, USA.

<sup>2</sup> Annex 1 of Pamela, J., 2003 Proc. 19th Int. Conf. on Fusion Energy (Lyon, 2002, IAEA).

<sup>3</sup> ITER-Naka Joint Working Site, JAPAN.

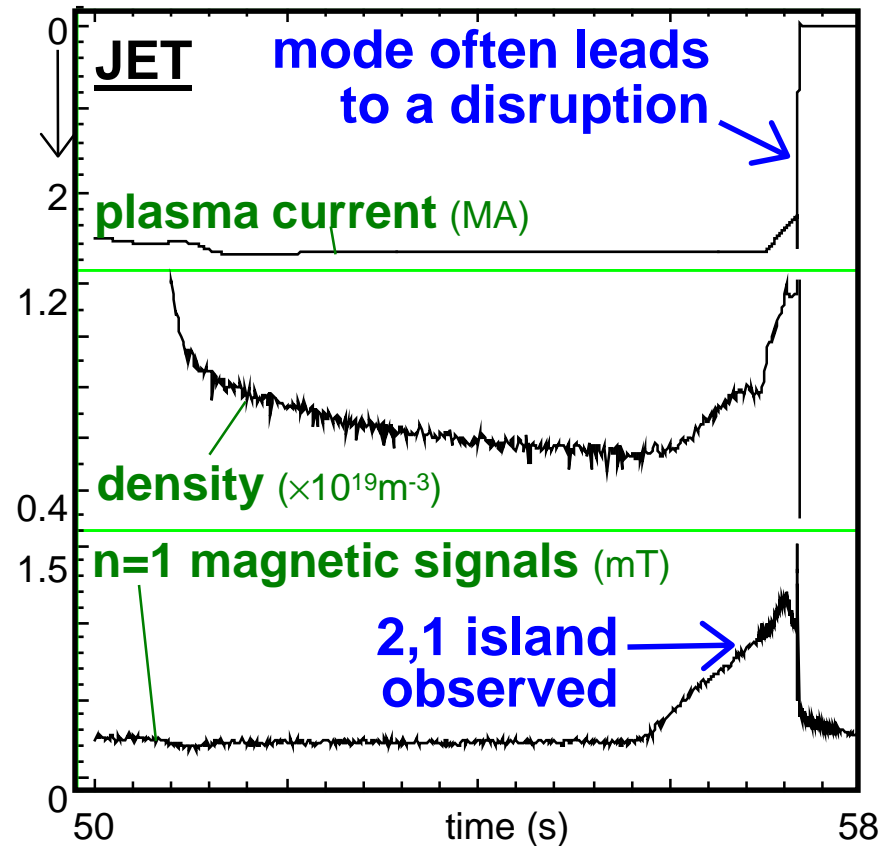
# Outline

---

- Review and update error field studies
  - motivation - the error field problem
  - sideband effects → update from DIII-D
  - cross machine scalings → how do they do against C-Mod?
  - Underlying theory → cf new JET studies of rotation dependence
  - heating and mode avoidance → increased sensitivity at high  $\beta$
  - ITER
  - → plus interesting observation of EF role in disruptions
- Updates to conclusions

# Error field modes a problem for ITER

- Error fields naturally arise from asymmetries in coil design and position
  - overcomes plasma inertia and viscosity to drive mode growth
  - usually on  $q=2$  surface
  
- Modes limit low density **Ohmic** operation:
  - required for H-mode access for ITER or other next step device
  
- Issues:
  - level of sensitivity for ITER?
  - which harmonics trigger modes?
  - what is likely ITER error?
  - how do we measure error fields and prevent modes?



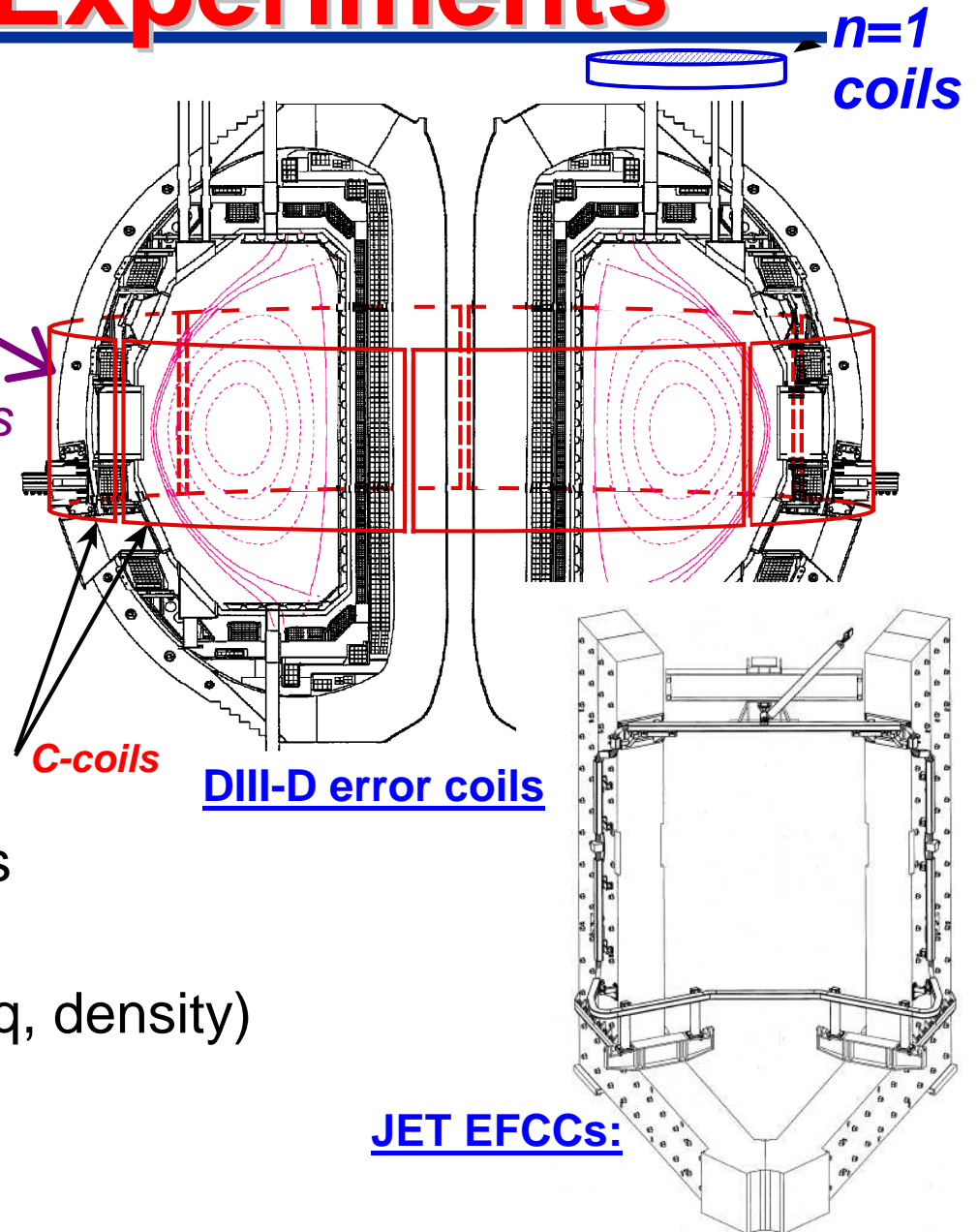
# Error Field Experiments

Error fields simulated by using dedicated asymmetric coils:

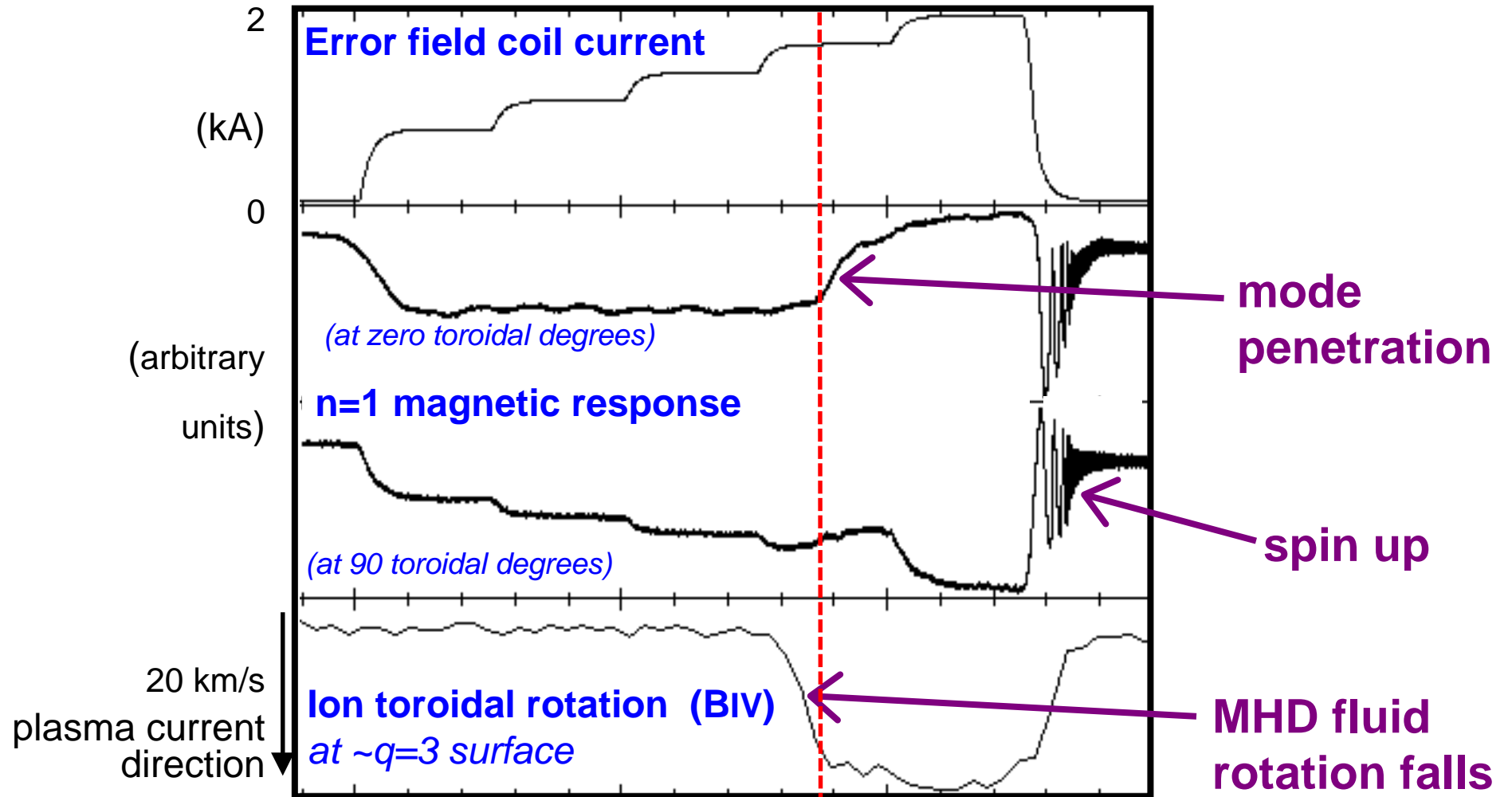
- DIII-D had 'C' coils (now also 'I')
- COMPASS-D has 8/10 bars in each quadrant - *ideal for sideband studies*
- JET has lower saddles + EFCCs - *vital point for ITER scaling*

**Current in these ramped until mode formed:**

- changing coil/bar combinations changes harmonic mix
- scan various parameters (TF, q, density)
- in Ohmic L-mode



# Error field mode penetration on COMPASS-D





# Sidebands have strong effect

## Scan 3,1 field for 2,1 penetration:

- 3,1 field penetrates at similar levels to 2,1
- mode always forms on  $q=2$
- two coupling mechanisms
- (1,1) harmonics also play a role

**Model must include these effects . . .**

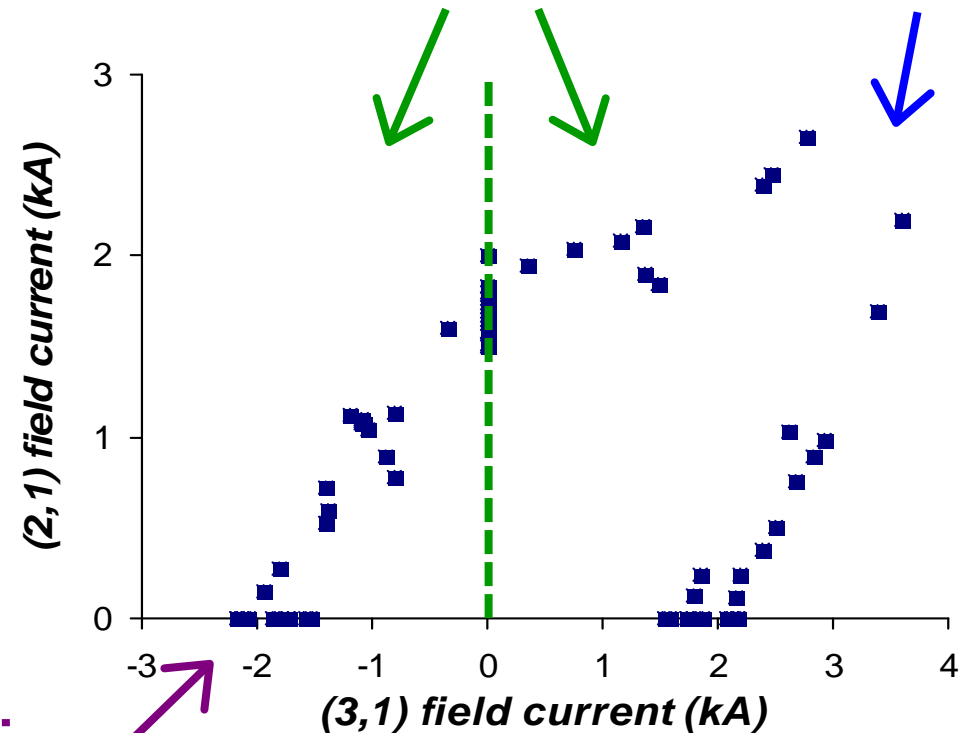
**coupling is strong:**

- effect of 3,1 is comparable to 2,1

## COMPASS-D: Ohmic L-mode

**asymmetry indicates toroidal coupling**

**closure indicates viscous coupling**



# Good fit obtained for DIII-D DND plasmas

Fit for critical torque on DIII-D in DND gave:

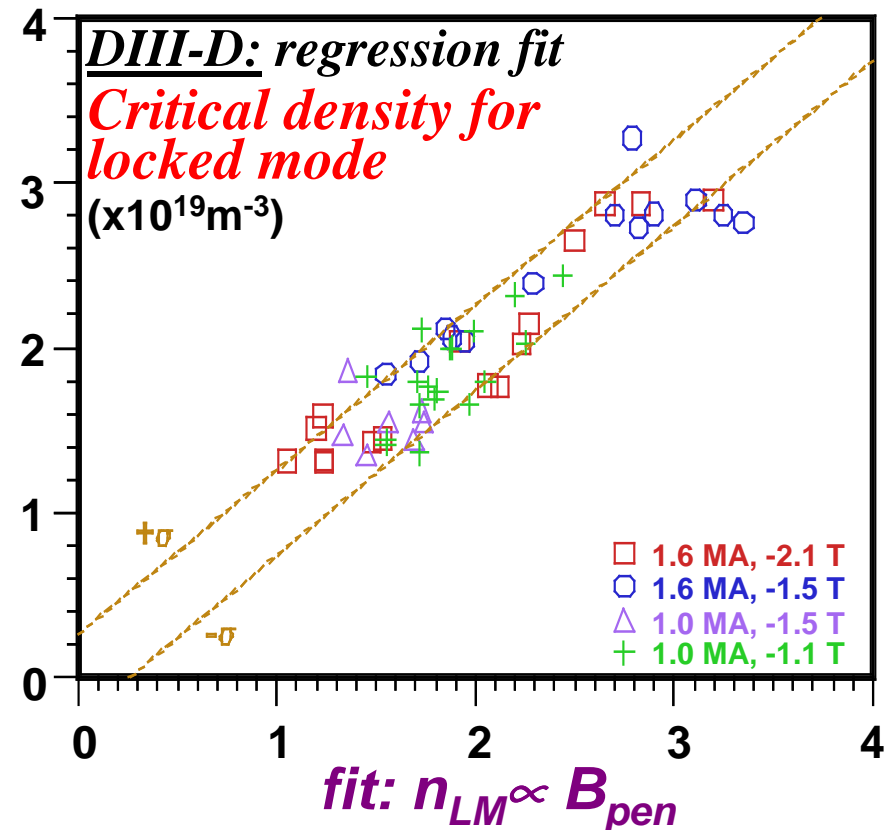
$$T \propto B_{\text{pen}}^2 \underbrace{|\mathbf{B}_{2,1} - 0.05\mathbf{B}_{3,1} + 0.25\mathbf{B}_{1,1}|^2}_{\text{Toroidal coupling}} + \underbrace{0.28|\mathbf{B}_{1,1}|^2 + 0.51|\mathbf{B}_{3,1}|^2}_{\text{Viscous drag}}$$

equivalent threshold of pure 2,1 field

- Equivalent sideband effects are seen in COMPASS-D SND plasmas

⇒ Separate sideband correction desirable in any next step device

$n_{LM}$  data



# DIII-D I coils can test sideband fits

- COMPASS-D SND sideband weights differ from DIII-D DND:

$$\text{C-D: } T_{tot}(q = 2) \propto B_{pen}^2 = 0.13|B_{1,1}|^2 + 1.69|B_{3,1}|^2 + |(0.21 + 0.02i)B_{1,1} + B_{2,1} + (0.85 - 0.17i)B_{3,1}|^2$$

← much more 3/1

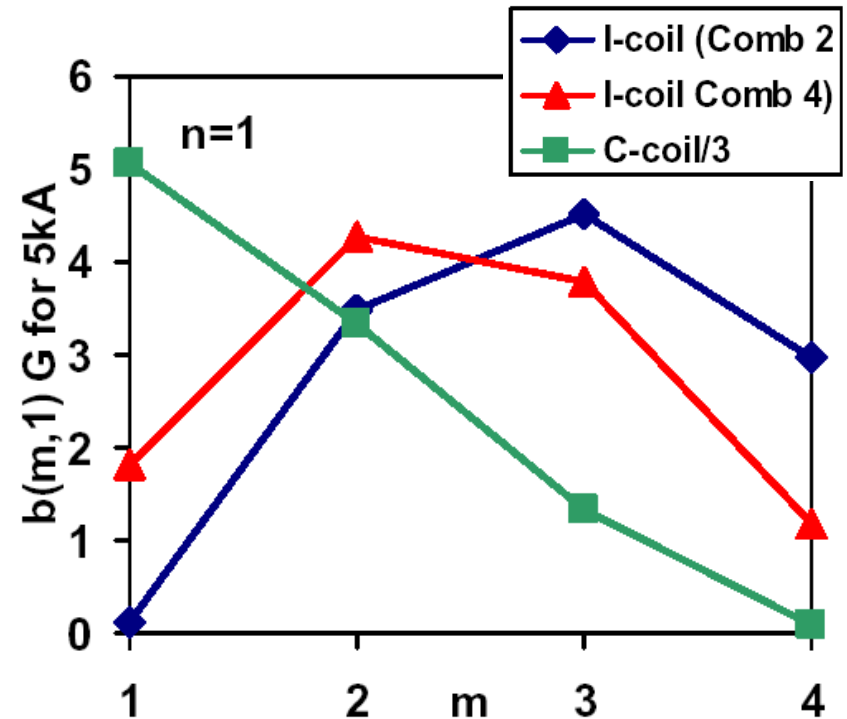
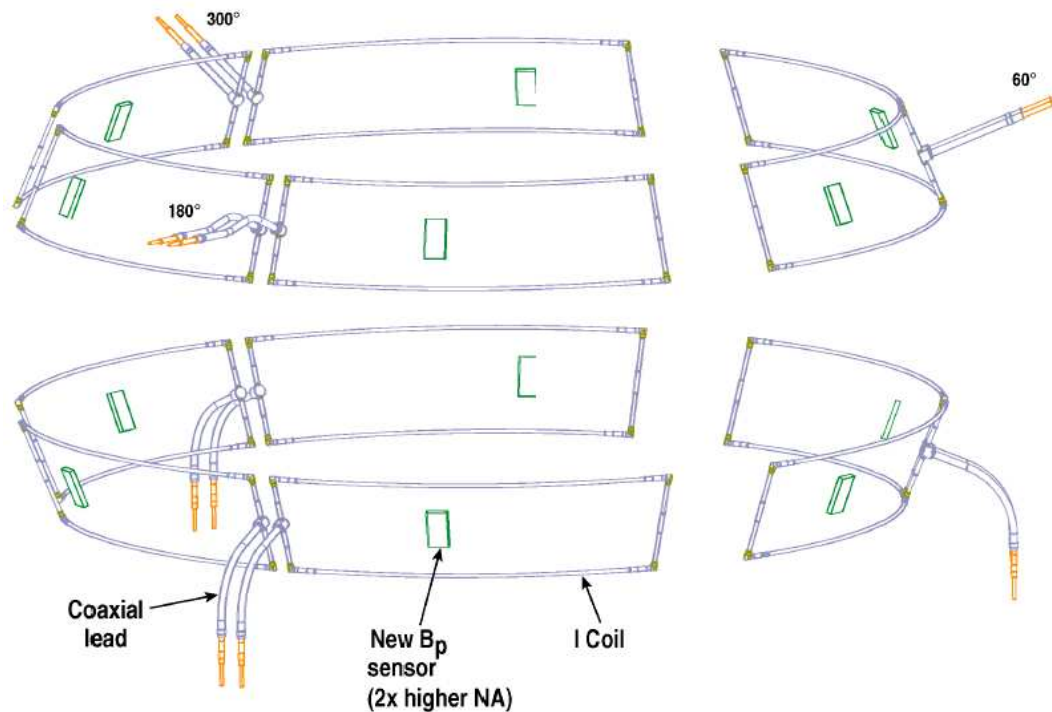
$$\text{DIII-D: } T \propto B_{pen}^2 \propto |B_{2,1} - 0.05B_{3,1} + 0.25B_{1,1}|^2 + 0.28|B_{1,1}|^2 + 0.51|B_{3,1}|^2$$

- *Note also discrepancies between COMPASS-D scalings and JET/DIII-D - different physics regime?*
- New I coils provide opportunity to test the fits in 'ITER' conditions and SND



# DIII-D I coils can test sideband fits

- I coils very flexible wiring with more sidebands possible:



# New I coils sideband fits

- New I coil SND results cf old DIII-D DND C coil results:

with complex mode-coupling coefficients:

$$B_{\text{pen}}^2 = 0.0 B_{11\_q1}^2 + |(-.55 - 0.1i) B_{11} + B_{21} + (-.68 - .32i) B_{31}|^2 + 0.1 B_{31\_q3}^2$$

more 1/1

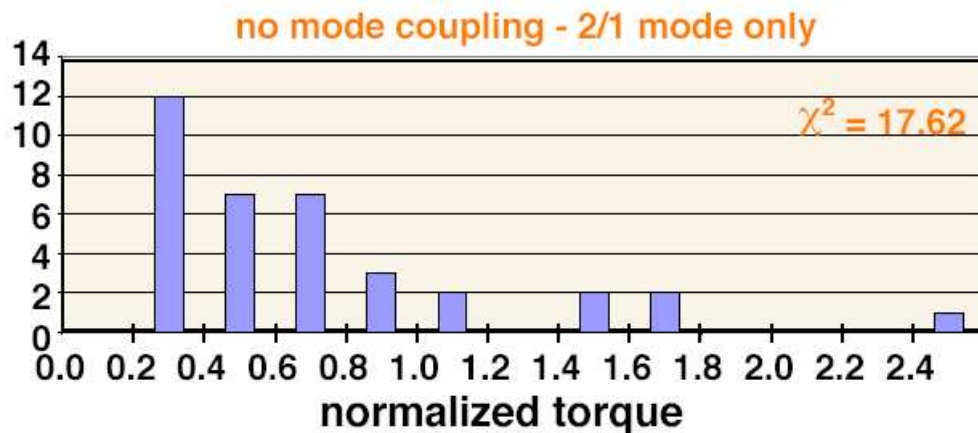
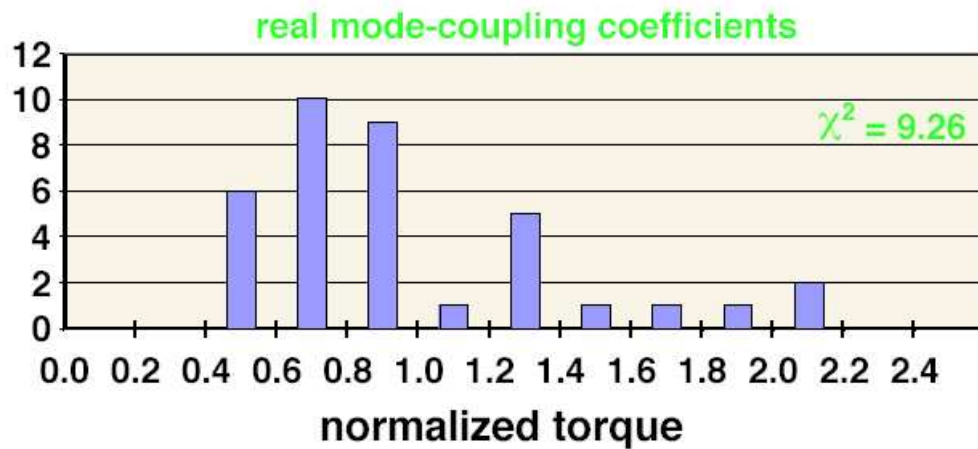
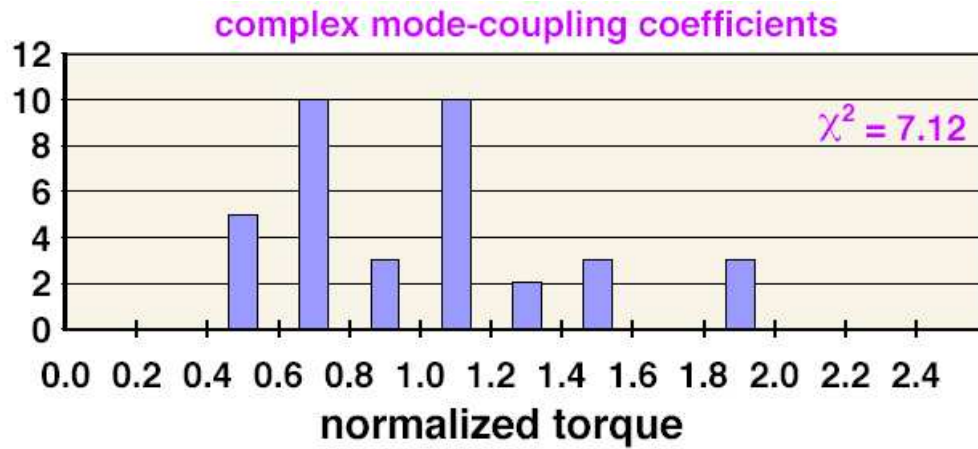
less viscous coupling

force real mode-coupling coefficients:

$$B_{\text{pen}}^2 = 0.0 B_{11\_q1}^2 + |-.60 B_{11} + B_{21} - .74 B_{31}|^2 + 0.1 B_{31\_q3}^2$$

- Considerable differences between old and new and C-D
  - but profiles changed (Ohmic H mode)
  - calculated intrinsic error has changed:

	<u>m=1</u>	<u>m=2</u>	<u>m=3</u>	<u>m=4</u>
new	.086 @ 209°	.811 @ 126°	.326 @ 333°	1.01 @ 74°
old	2.03 @ 75°	2.31 @ 73°	1.60 @ 323°	0.85 @ 28°



---

**Fits show sidebands are important**

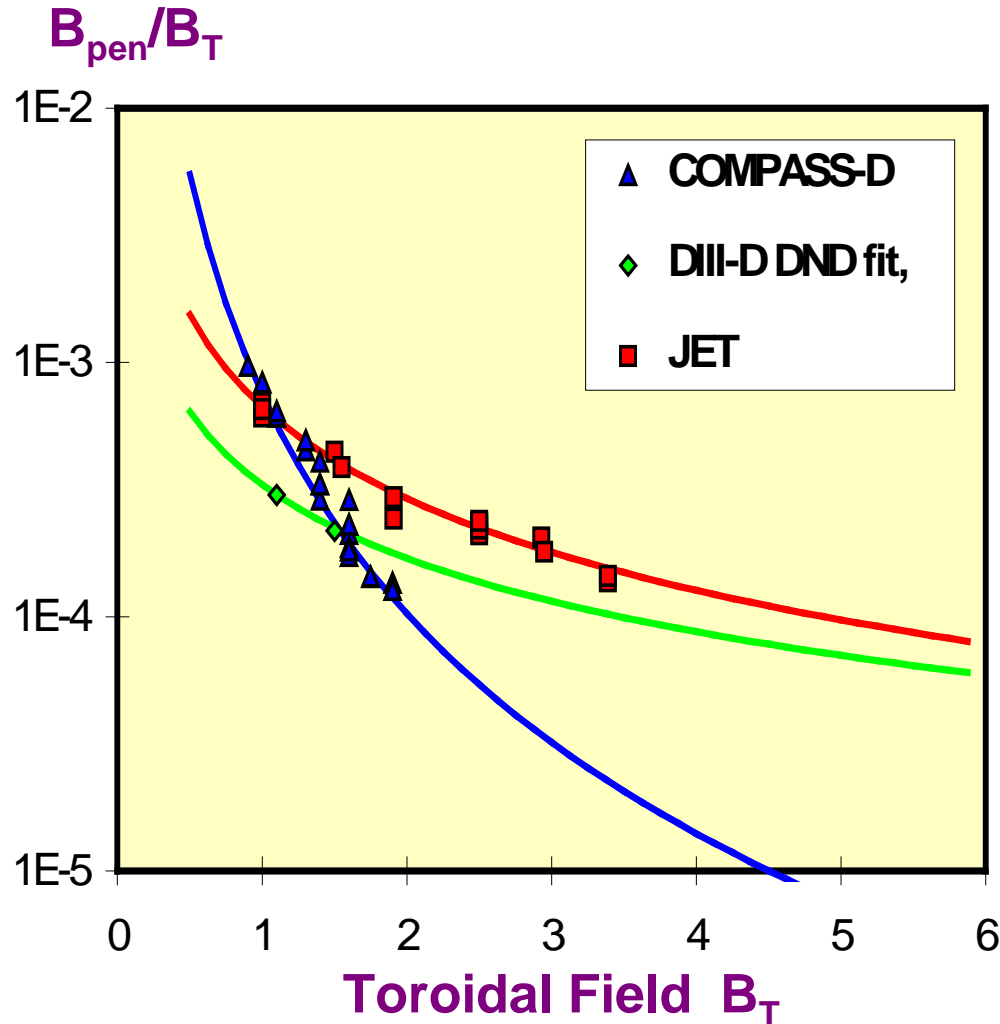
**Compare fits using sidebands with results only including 2/1 harmonics**

## **Summary on sideband effects**

---

- New DIII-D I coil results confirm sidebands are important
  - but fits differ from C-coil and COMPASS-D results
  - ...and give different intrinsic error measurements
    - *...although agree on phase*
- Further analysis ongoing to resolve discrepancies
  - is it result of more data or different shapes, etc.
  - compare each fit from COMPASS-D & DIII-D against other machine's data
  - more data required?
- More JET experiments in January 2004 will help by using new EFCC coils set (hopefully with saddles simultaneously)

# Toroidal field scalings



- JET shows weak TF scaling
- DIII-D's points support this
- Clear difference in gradient with COMPASS-D
  - may be due to different rotation behaviour (- ion losses?)
- **Highlights uncertainty in scaling**
  - *intrinsic problem of using global scaling for local behaviour*

# Comparison of scalings

- Scale multimode threshold, using power law form:

- $B_{\text{pen}} / B_T \propto n^{\alpha_n} R^{\alpha_R} B^{\alpha_B} q^{\alpha_q}$

- deduce  $\alpha_R = 2\alpha_n + 1.25\alpha_B$  from dimensional considerations, in line with approach for confinement scaling

*(Connor and Taylor NF 17 1047)*

<u>Exp scans:</u>	$\alpha_n$	$\alpha_B$	$\alpha_q$
JET	0.94	-1.2	0.05
DIII-D	0.99	-0.96	0.83
COMPASS-D	~1	-2.9	1.6

- density scaling consistent across machines
- q scalings expected to differ due to different mode spectra

# Extrapolates to moderate ITER sensitivity

- *Use data from JET to minimise extrapolation:*
  - ⇒ Gives  $\alpha_R = 0.4 \pm 0.2$
  - ⇒ Predicts ITER threshold of  $B_{\text{pen}}/B_T \sim \underline{1.25 \times 10^{-4}}$   
 at ITER reference density of  $2 \times 10^{19} \text{m}^{-3}$  and  $q_{95} = 3.3$   
 (using COMPASS-D SND sideband formula)
- Comparable with DIII-D DND extrapolations ( $\sim 2 \times 10^{-4}$ )
- COMPASS-D shows uncertainties in scaling - changes in regime can cause different rotation behaviour and TF scaling
  - ⇒ *Important to examine rotation and  $\beta$  effects*
    - additional heating studies on **JET**, **COMPASS-D** and **DIII-D** . . .



# Extrapolation to CMOD

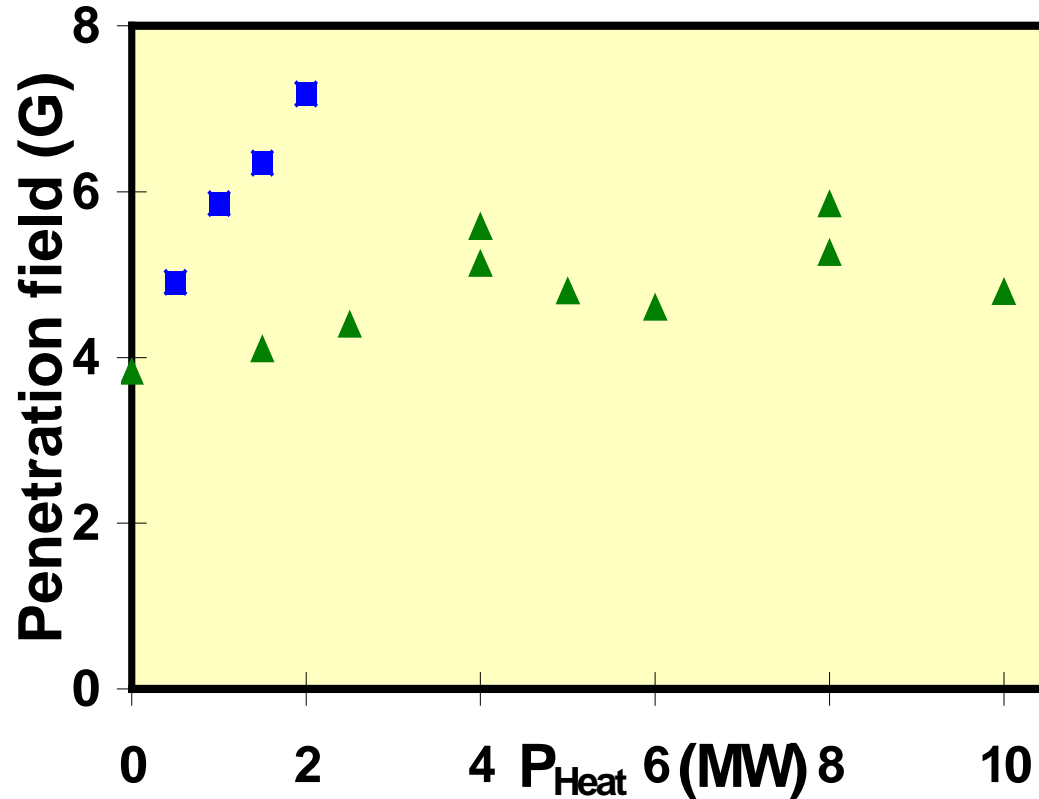
- Take TF & density dependence fits and scale to CMOD parameters:  $B_T=4T$   $n_e=5e19m^{-3}$

Step	JET fit	DIII-D fit
Fit formula	$B_{pen}=5.23B^{-0.2}(n_e/1.6e19)$	$B_{pen}=(n_e/0.6e19)(q/3.3)^{0.79}(B/2.1)^{0.01}$
$B_{pen}$ at CMOD TF and $n_e$	12.5 Gauss for R=2.9m	8.4 Gauss for R=1.67m
Size scale exponent	0.4	0.76
$B_{pen}$ CMOD	6.9 Gauss	4.2 Gauss

- Combined 1/1, 2/1 and 3/1 field using fits from C-D or DIII-D
  - on JET 2/1 is about ~70% of  $B_{pen}$
  - on DIII-D 2/1 is about ~80% of  $B_{pen}$
- ➔ These are in right ball park for CMOD 2/1 thresholds
  - uncertainties:  $\alpha_R$ , sideband form, CMOD spectrum, shape,  $q_{95}$
- JET-CMOD identity experiments desirable - and planned!



## Heating: NBI rotation is key effect



NBI and RF to separate  $\omega$  and  $\beta$  dependence:

- modest rise in  $\beta$  has weak effect
- H-mode transition has little effect
- rotation from NBI raises threshold

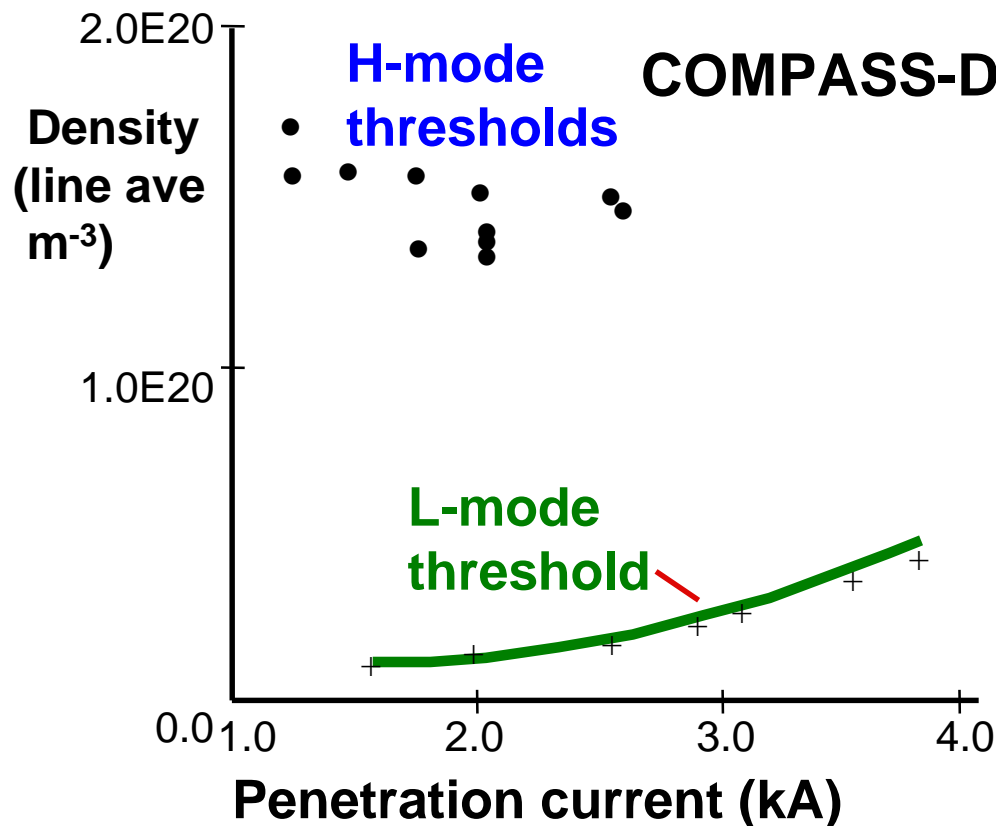
⇒ *NBI rotation possible technique for avoidance on ITER*

# Role of rotation

---

- Modes form when resonant surface is braked by resonant response to EF to half it's natural frequency
  - tiny static island induced by EF
  - viscous forces try to keep bulk plasma rotating slipping about island
    - *this opposes island growth*
  - torque exerted through island and viscosity by EF brakes plasma
  - if rotation slows enough, island can grow, increasing torque and bifurcating to a locked mode state
  - threshold scales as  $B_{\text{pen}} \sim B_T \omega_0 \tau_A (\tau_{\text{rec}} / \tau_v)^{1/2}$ 
    - $\omega_0$  often taken to be electron diamagnetic rotation
- Expect cancelling natural q=2 rotation should lower thresholds
  - *didn't see in normal B operation on JET so look in reverse B...*

# Error fields penetrated easily when rotation stopped in COMPASS-D Ohmic H-mode



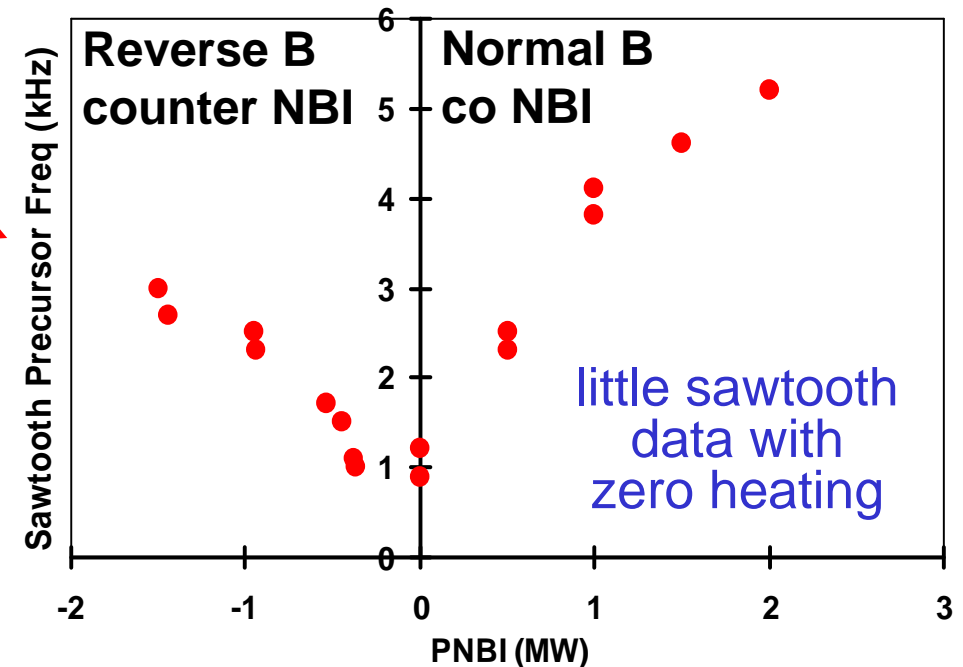
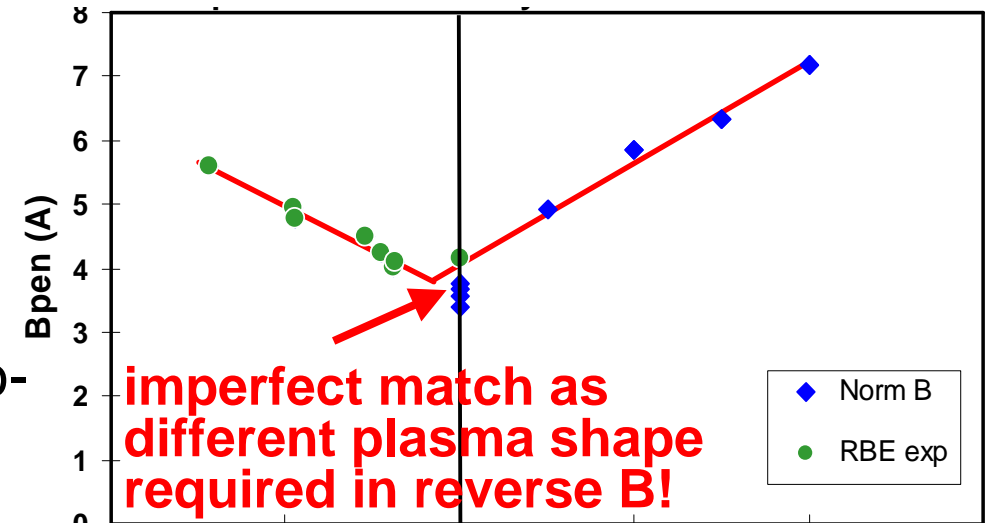
## Ohmic ELM-free H-mode

- Error-field penetrates at 10x density for L-mode penetration
- MHD fluid slows leading to a very low penetration threshold
- Confirms key role of rotation

*- Leaves open question of what will happen to rotation in ITER H-mode?*

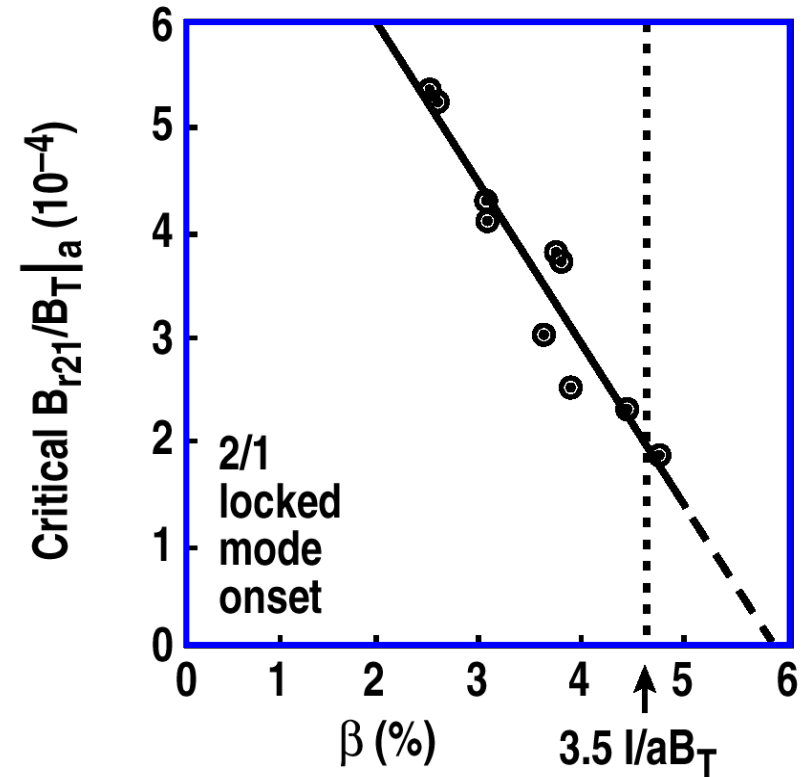
# JET rotation scan

- Threshold does not plunge towards 0 at any power!
- Possible minima for counter injection power:
  - 0 → 0.5 MW: strong rise with co-, no rise with counter
  - natural plasma rotation in ion (co-) direction
    - get less core rotation for given  $P_{\text{NBI}}$  in reverse B
    - also observed with charge exchange measurements
    - Ohmic modes spin up in ion direction
  - all consistent with counter NBI opposing a naturally co-dir'n  $\omega$



# What about at high $\beta$ ?

- Old result from DIII-D shows thresholds fall as ideal limit approached
  - generally locked modes
  - error field amplification effect?
- *What about proximity to NTM?*



1992 La Haye experiments



## Error field thresholds fall at high $\beta$

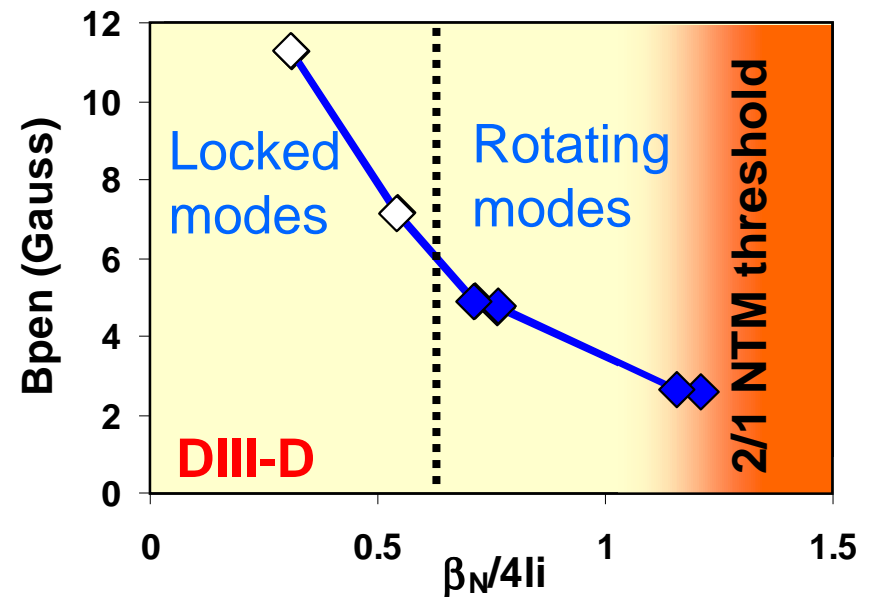
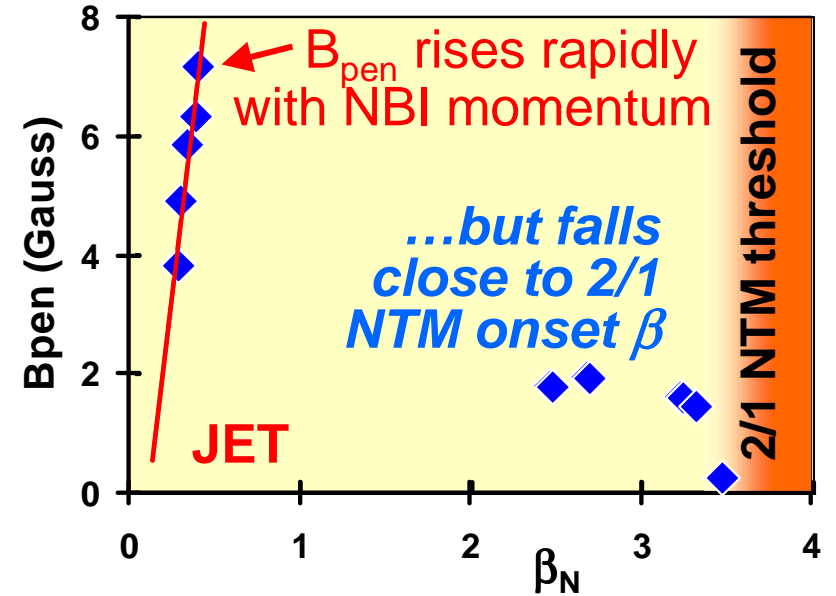
- Error field thresholds fall close to the 2/1 NTM onset  $\beta$  on JET

⇒ increased error field sensitivity

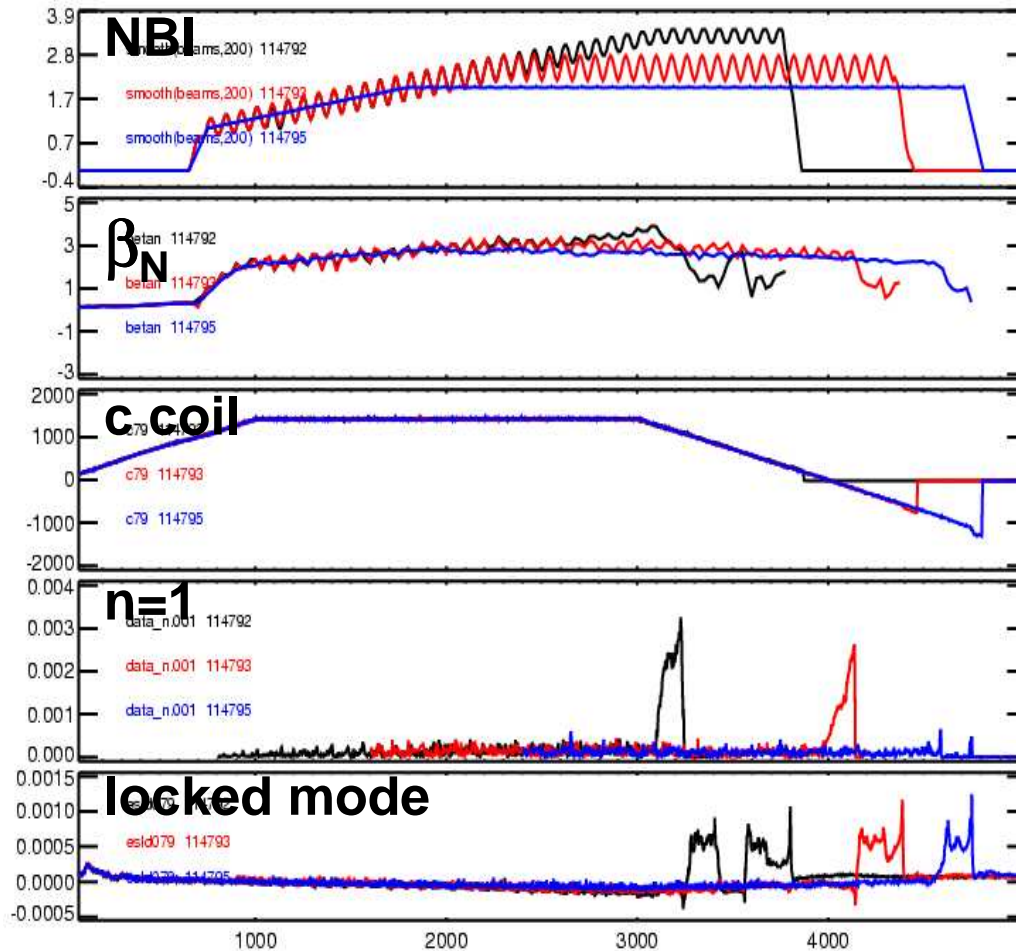
- Also observed on D3D...

— here modes formed rotating at intermediate  $\beta_N$

⇒ EF is directly assisting NTM onset mechanism



# DIII-D: error field - NTM interplay



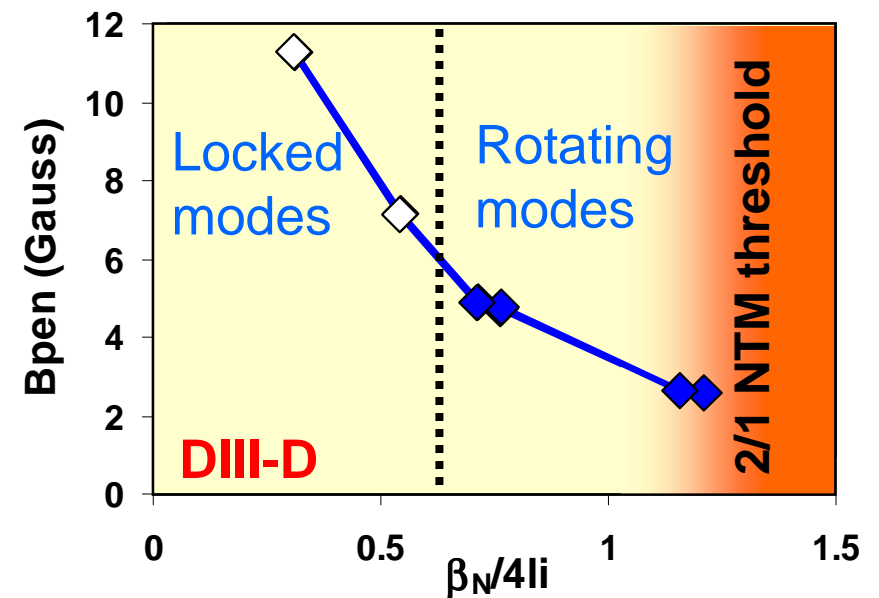
- Fix at subcritical  $\beta$

- apply EF ramp

⇒ EFs assist NTM onset

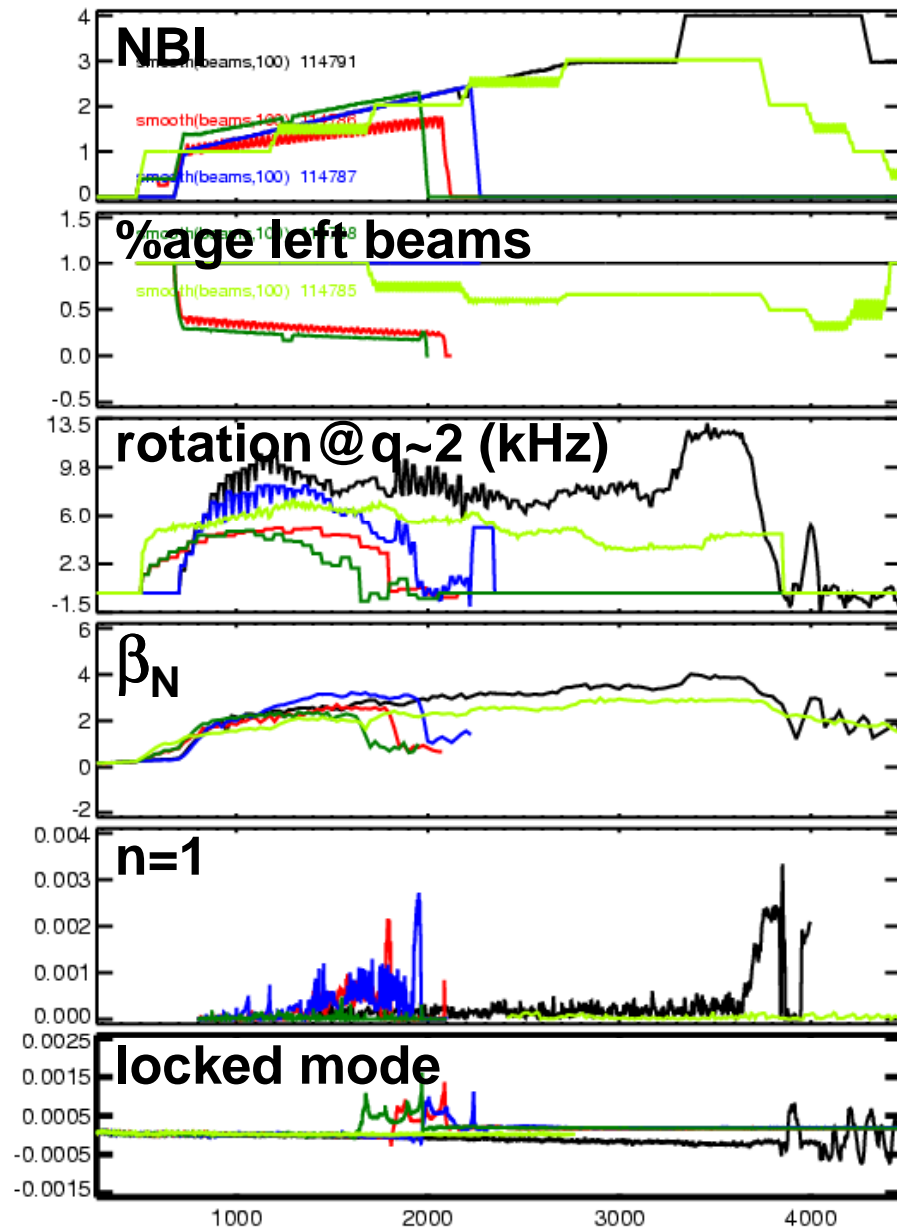
- at medium  $\beta$  modes born rotating - NTMs

- low  $\beta$ : standard locked modes





## NTMs at lower $\beta$ with pre-existing errors



- Pre-existing error field in coloured cases (**not black**)
- Leads to error field modes during beam ramp up
  - plasma slows first
  - goes rapidly to locked mode
- Using ‘left’ beams helps avoid modes

*Why no mode for light green?*

# Original conclusions

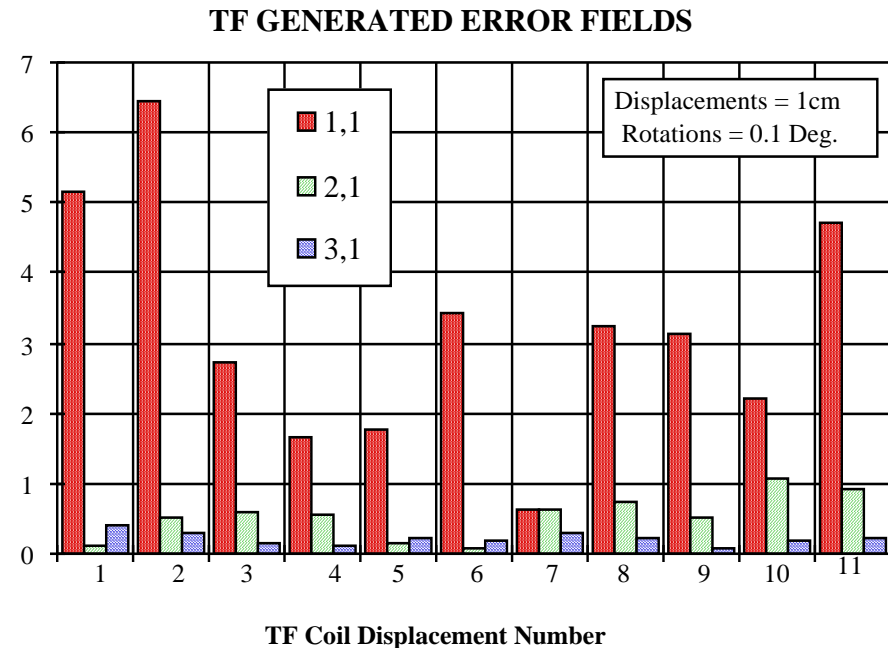
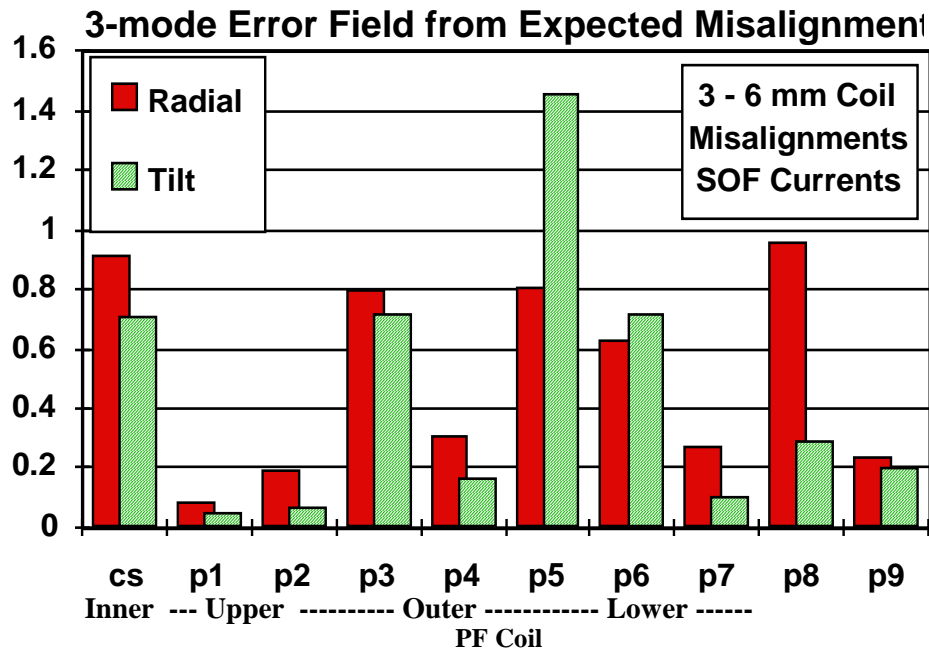
---

- Error fields limit low density operation in tokamaks
  - likely to be a problem for H-mode access on ITER or other next step devices
- (1,1) (2,1) and (3,1) harmonics all contribute to mode formation
- Predicted ITER sensitivity comparable to likely intrinsic error  
( $B_{3mode} \approx 1 \cdot 10^{-4}$ )
  - some level of correction important
    - *3 harmonic design was a prudent approach*
- Neutral beams may also assist mode avoidance

# Likely ITER error field $B_{3mode} \approx 1 \cdot 10^{-4}$

Design studies have identified key sources of error on ITER:

- calculate 3-mode error according to:  $B_{3mode} = (0.2B_{1,1}^2 + B_{2,1}^2 + 0.8B_{3,1}^2)^{1/2}$
- combine to give probability distribution:  $P(B_{m,n}) = \int f(B_{m,n}) dB_{m,n}$

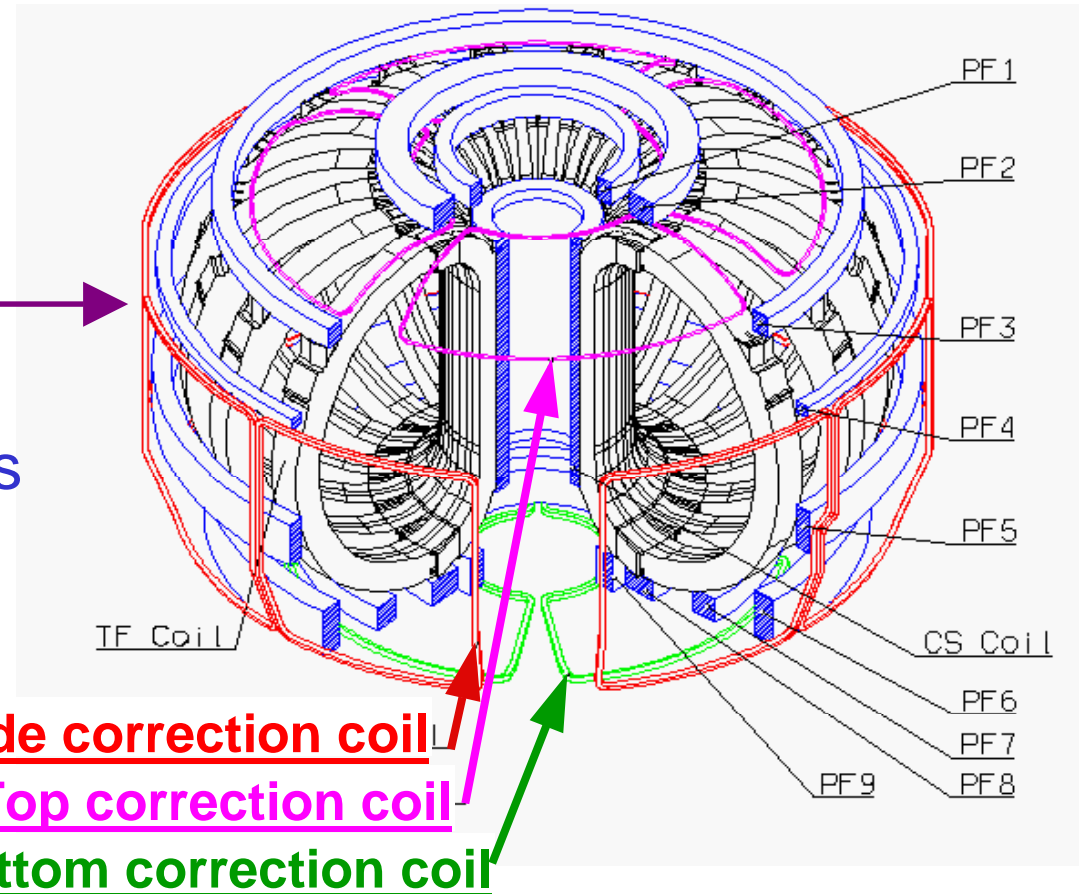


- *plus additional sources from test blanket and NBI shields*

# ITER Correction System

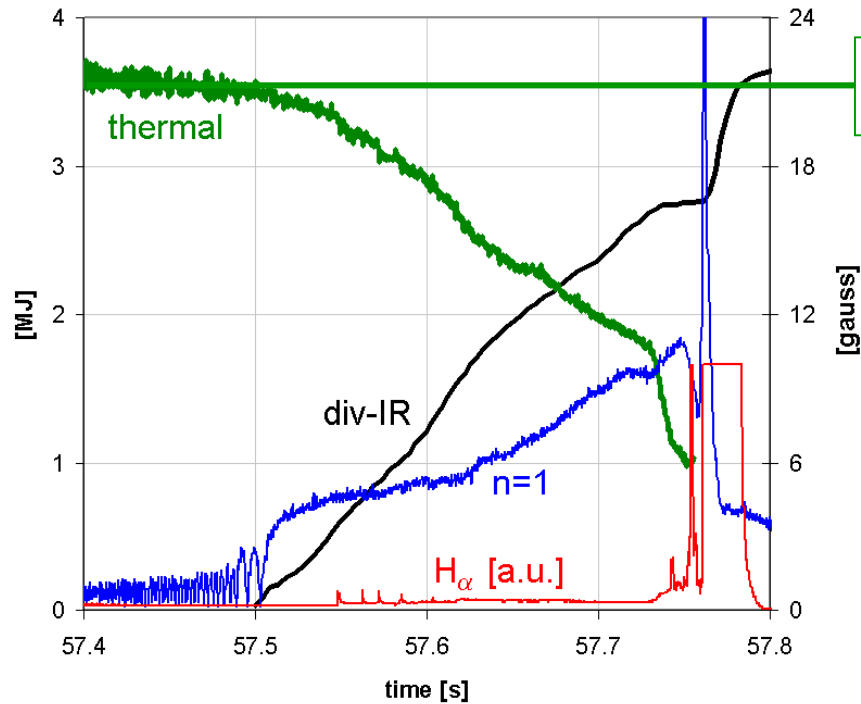
ITER intrinsic error comparable to predicted tolerance

- sideband harmonics also play a role
- ⇒ 3-harmonic correction designed as prudent approach
- capable of reducing errors of  $\delta B/B_T \sim 12 \times 10^{-5}$  down to  $2 \times 10^{-5}$



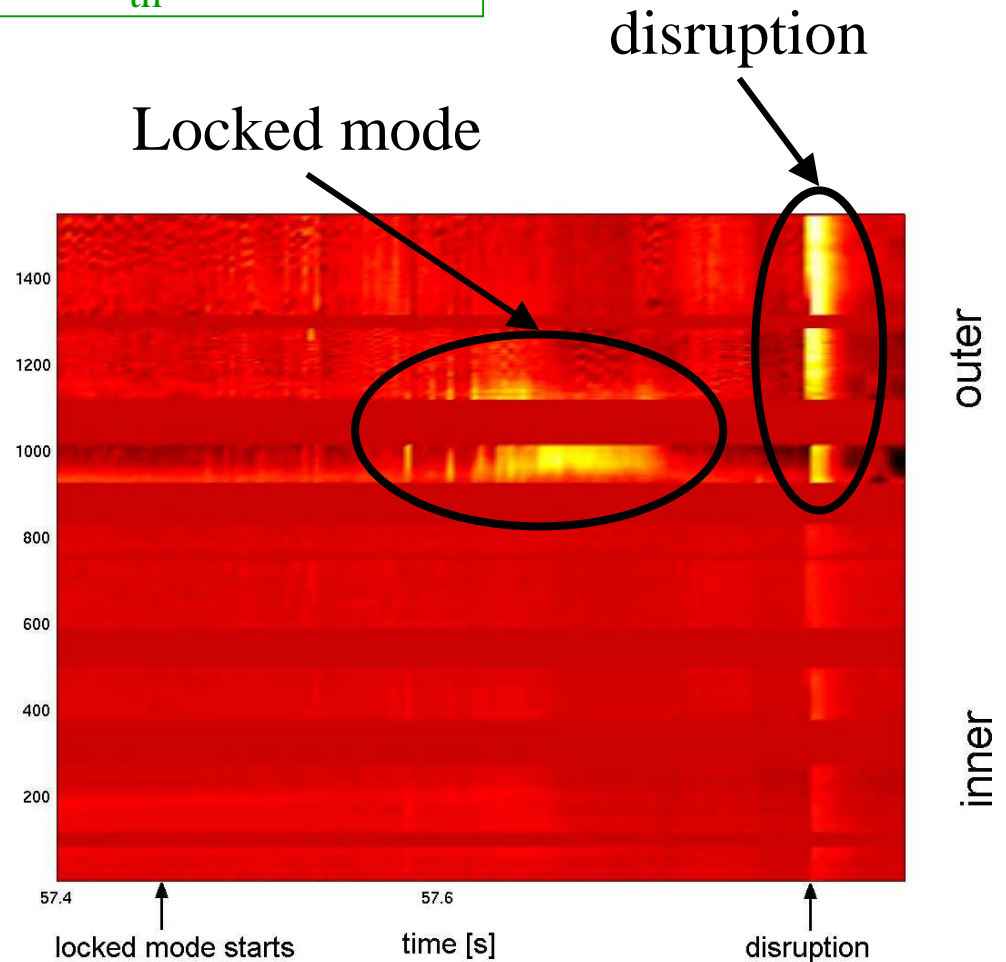


# Disruption heat loads - divertor



All  $E_{th}$  to divertor

Large area of plasma-target interaction



# Updated conclusions

- ITER still looks conservatively designed for H-mode access
  - scaling does well for CMOD - *further 'identity' tests a high priority*
- Some uncertainties in sideband physics
  - discrepancies with old/new DIII-D and COMPASS-D scalings
  - but sidebands clearly still important
- But increased sensitivity close to NTM thresholds
  - can lead to locked modes with less error field
  - + rotating NTMs more easily triggered if high error field
- Rotation appears to play a role
  - though not as strong as in theory (could not get  $B_{pen} \rightarrow 0$ )
  - natural rotation possibly in ion diamagnetic direction on JET
- Locked modes during disruptions spread heat loading and send all the energy to divertor

- Reserve slides...

# Origin of dimensional scaling constraint

[Connor and Taylor (NF 17 1047):]

Take given basic equations for a regime (eg Collisional Vlasov at high  $\beta$ )

- there are only certain (typically 1-3) transformations on basic parameters that leave them invariant
- these must also leave quantities such as  $Q(\text{basic params})$  invariant
- thus each transform places a constraint on  $Q$ 's dependencies
- if power law form for  $Q$  get *linear relations between coefs*

$\Rightarrow$  showed  $B\tau$  obeyed relation ( $B \equiv$  cyclotron freq  $\Rightarrow B\tau$  dimensionless),

Generally true for all dimensionless quantities, *provided regime or physics does not change over scaling region.*

Relevant regime is Ohmic collisional high  $\beta$  (*this works for confinement scaling, so use here*) and dimensionless quantity is:

$$B_{\text{pen}}/B_T \sim n^{\alpha_n} B^{\alpha_B} R^{\alpha_R} \quad \text{and constraint is} \quad 8\alpha_n + 5\alpha_B - 4\alpha_R = 0$$

(Collisional high  $\beta$  encompasses collisionless and low beta regimes - most general approach.

Ohmic heating forces  $T(n,R,B)$  and so all regimes have same constraint.

This approach has been seen to work experimentally for confinement scalings.)

# Mechanism for additional ion losses on C-D?

Toroidal rotation from neoclassical theory:

$$v_{\phi} = \frac{-T_i}{m_i \Omega_{i\theta}} \left( \frac{d \ln P_i}{dr} + \frac{e}{T_i} \frac{d\phi}{dr} + k \frac{d \ln T_i}{dr} \right) \quad (\text{Hinton \& Hazeltine, 1976})$$

simplifies to:  $v_{\phi} \sim E_r / B_{\theta} \propto E_r / I_p$ , if radial electric field dominates  
 (otherwise  $P$  and  $T$  terms dominate which may scale differently with TF,  $I_p$ )

Ion banana orbit:  $width = \frac{2mv_{par}}{eB_{\theta}} \approx \frac{1.6\sqrt{2mE}}{eB_{\theta}}$  for trapped particles...

calculate for high energy tail ions ( $E \sim 3 \times \text{Temperature}$ ) on  $q=2$ :

C-D: 4.5cm

JET: 3.4cm

ITER: 3.1cm

↑ comparable with distance from  $q=2$  to edge/vessel

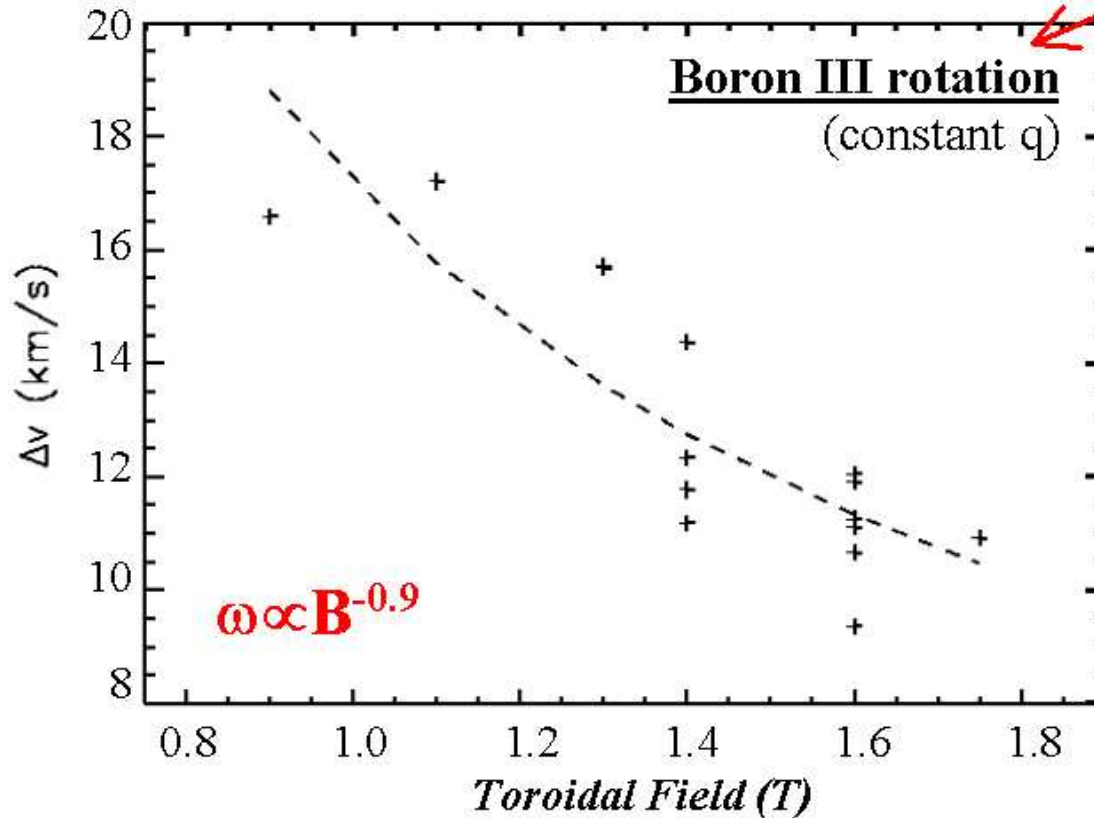
⇒ Ion loss may set up a strong electric field on C-D, not others

⇒ Gives rise to increased flow, occurs most at low TF and  $I_p$ , possibly explaining strong scaling of flows and thresholds seen on COMPASS-D



## COMPASS-D different scalings - rotation?

Different TF and q scaling may be due to changes in rotation behaviour on COMPASS-D:

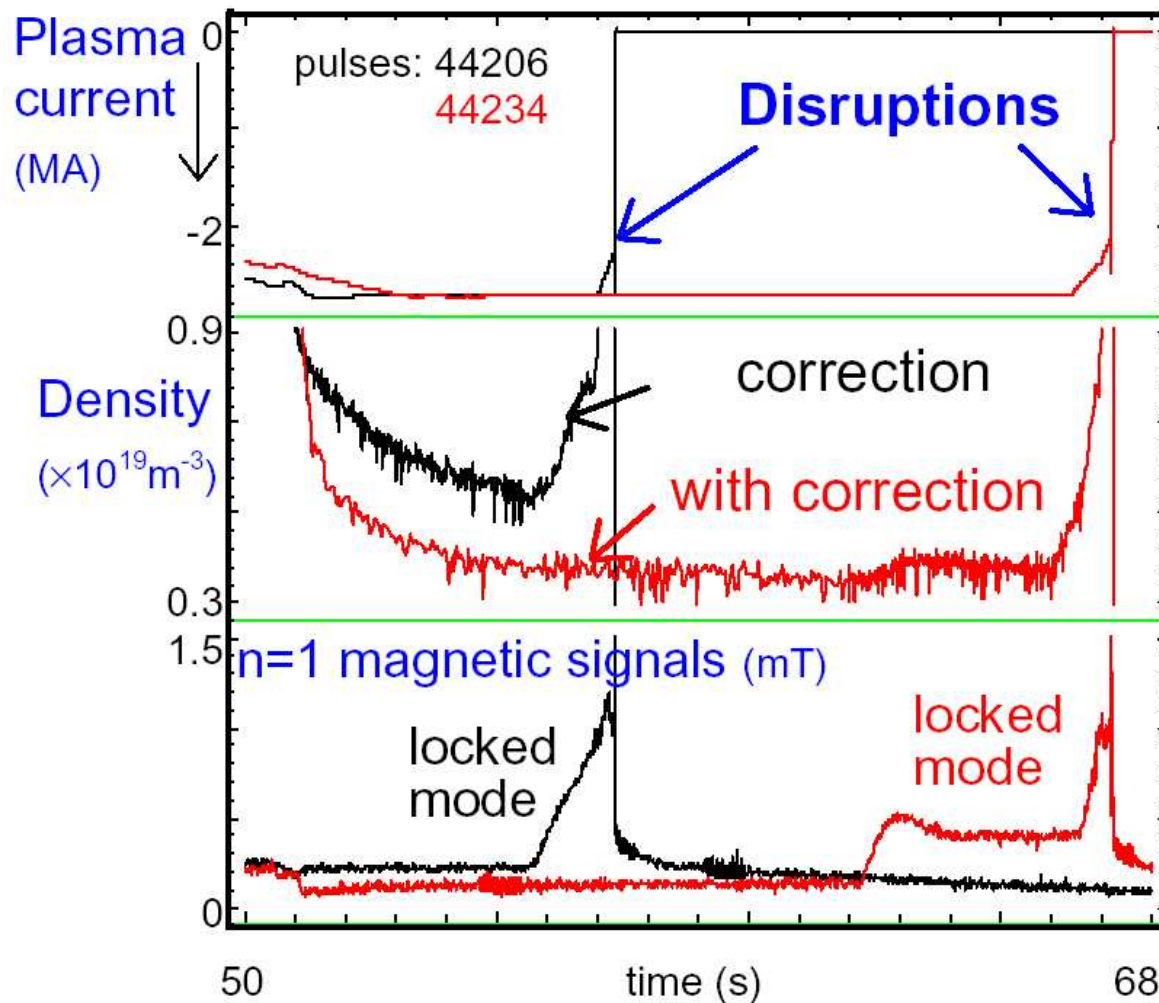


- C-D rotates much slower at higher TFs (**constant q**)
- DIII-D: C-VI shows  $\omega \propto B^{+1.0}$  (**constant I<sub>p</sub>**) cf  $B^0$  for C-D

May be related to ion losses?

- JET ripple expts showed ion losses raised threshold (and rot'n) in Ohmic plasmas
- C-D ion banana orbits may take ions out of plasma, esp at low TF. . . (EXTRA SLIDE)

# Intrinsic error correction accesses lower $n_e$



- 35% lower density achieved with correction on JET
- on DIII-D 50% lower achieved with C-coil
- non-perfect correction indicates significance of sideband effects
  - (also significantly reduced disruptions following locked modes on JET)



## Benefits of error correction

- 84% of 25 uncorrected locked modes disrupted, cf 20% of 10 corrected
- In density ramp-down pulses correction accessed much lower density:
  - 14 fringes ( $5.7 \times 10^{19} \text{m}^{-3}$ ) no correction
  - 9 fringes ( $3.7 \times 10^{19} \text{m}^{-3}$ ) with correction
- *Routine correction may reduce disruptions and aid low density access*

