

Barrier dynamics studies in different experimental configurations and implications for advanced tokamak research

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PPPL

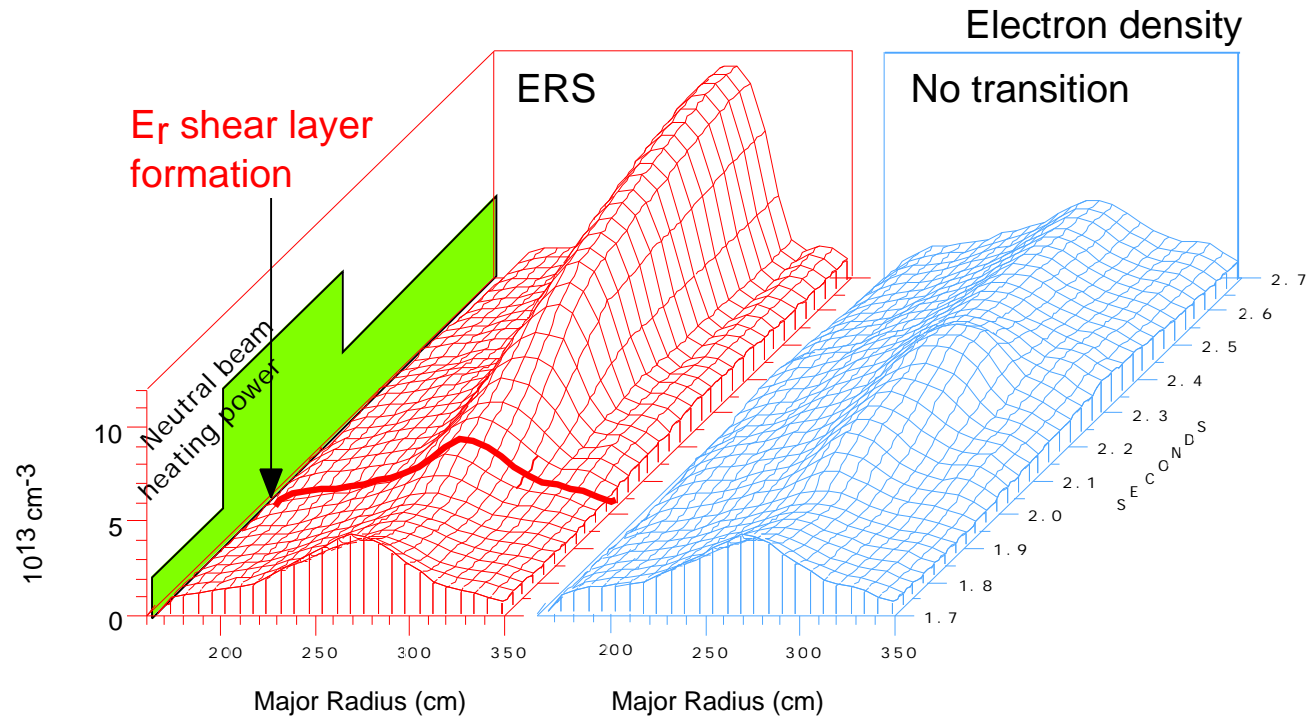
Advanced Tokamak Workshop

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General Atomics, San Diego, CA

Three points:

1. In AT modes, particle confinement is in danger of being *intolerably good*, demanding control tools
 - helium ash? other impurities?
2. How experiment generates $E \times B$ shear has a profound influence on barrier dynamics
 - barrier control is possible
 - need to ensure experiment is best aimed at AT goal
3. Models with $E \times B$ shear reproduce many of the dynamics
 - fluctuations character, transition time scales

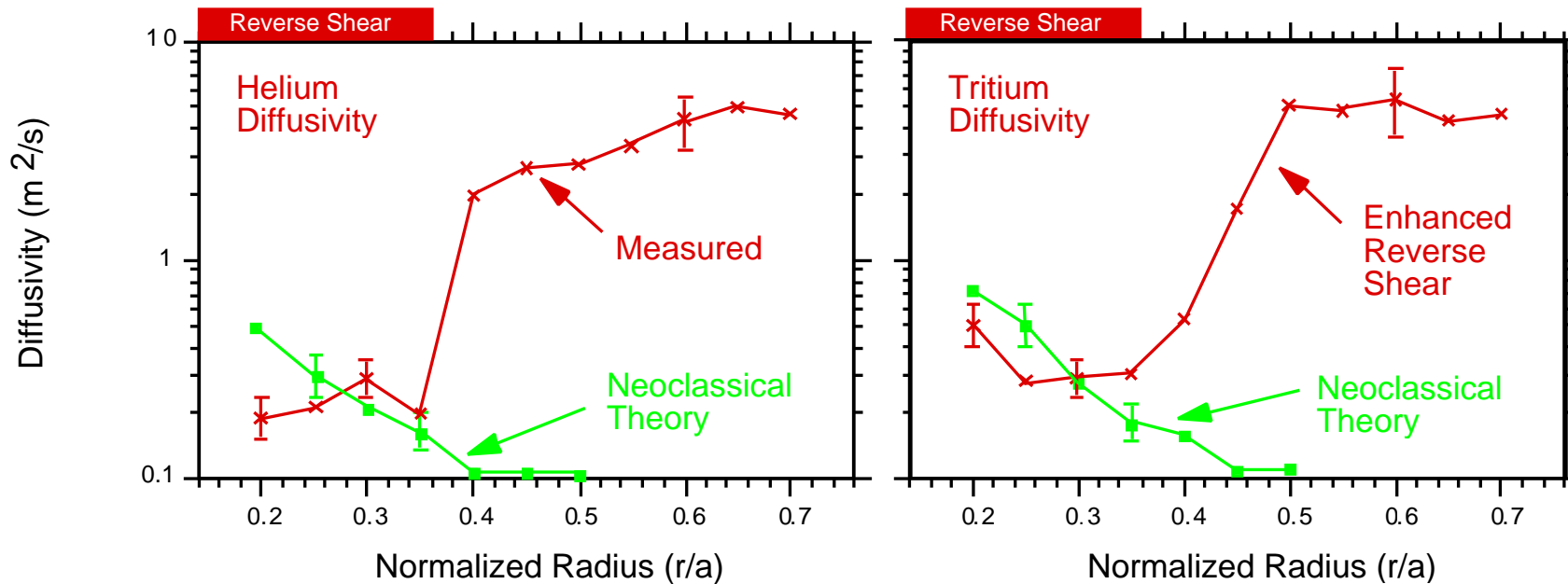
Dramatically reduced particle and energy transport in core follows narrow E_r shear layer formation



- Particle transport improves with energy transport
may need barrier control to exhaust impurities
- Trigger event is E_r shear layer formation about 2 cm wide
small scale events unrelated to device size can fundamentally alter core confinement

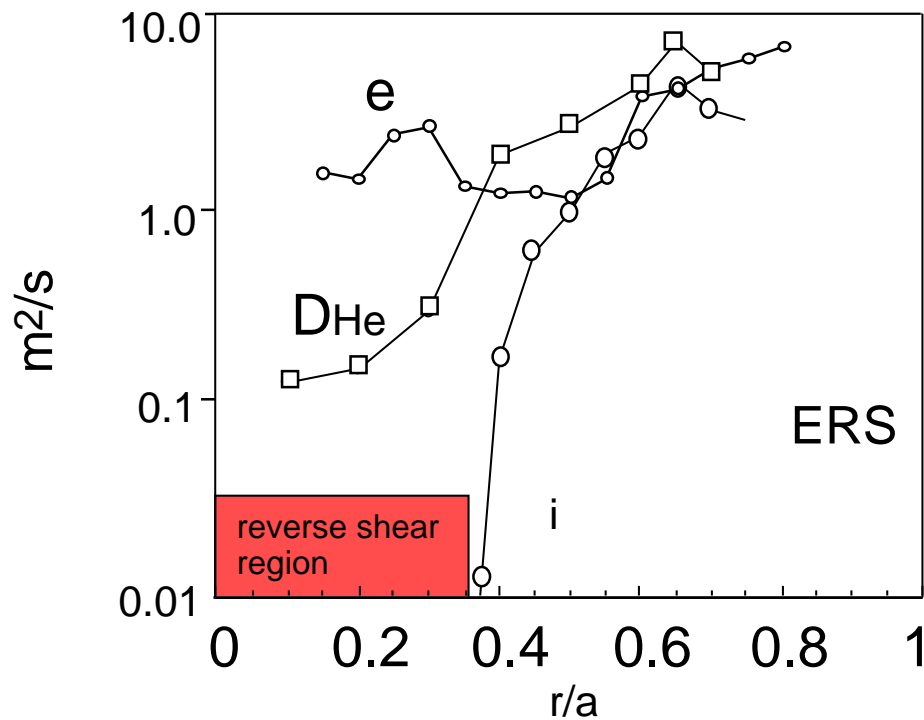
In ERS Core, Helium & Tritium Diffusivities are Similar and Agree with Neoclassical Theory

TFTR



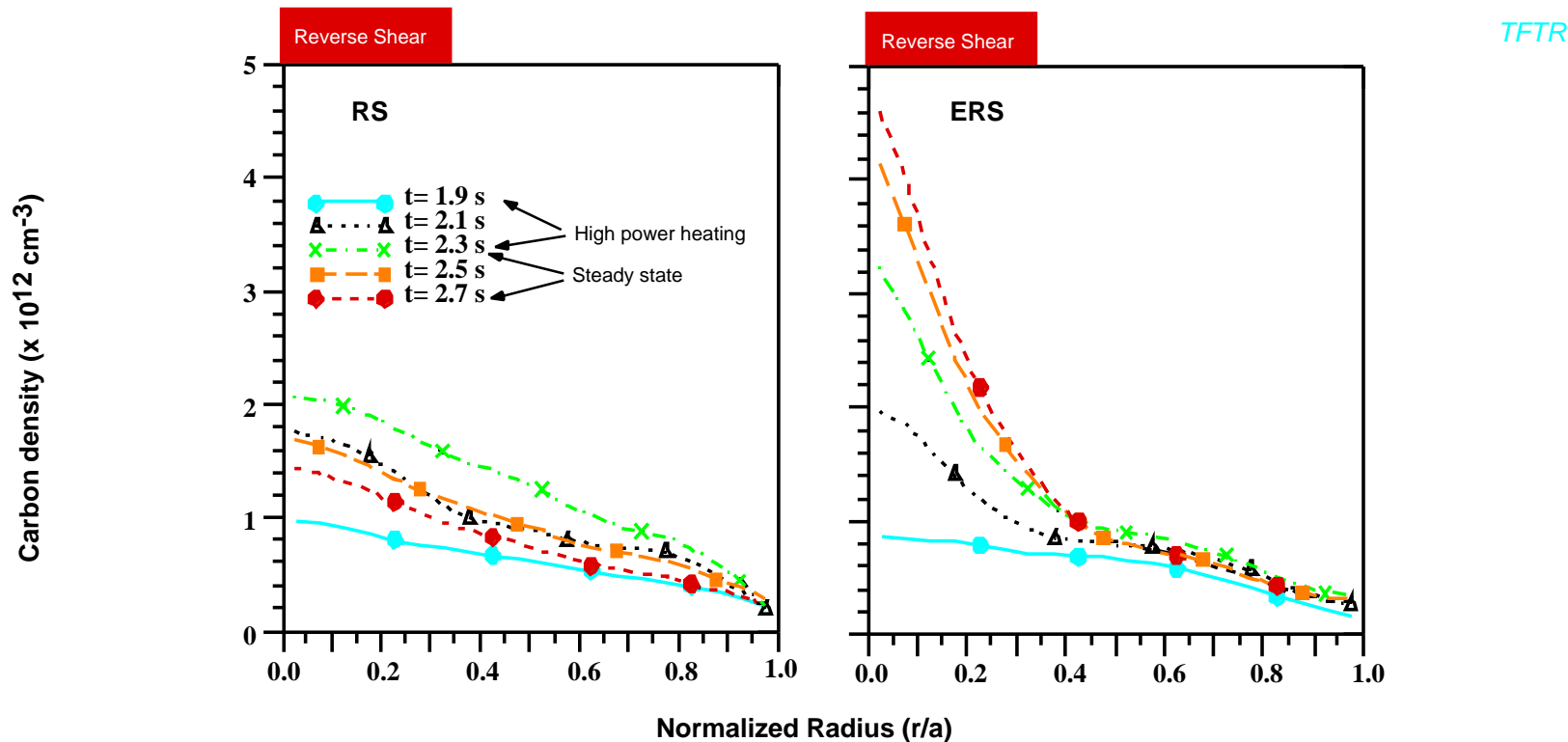
- Neoclassical helium & tritium diffusivity calculated by NCLASS using measured profiles

Assessing the helium ash problem requires understanding the observed behavior of different transport channels



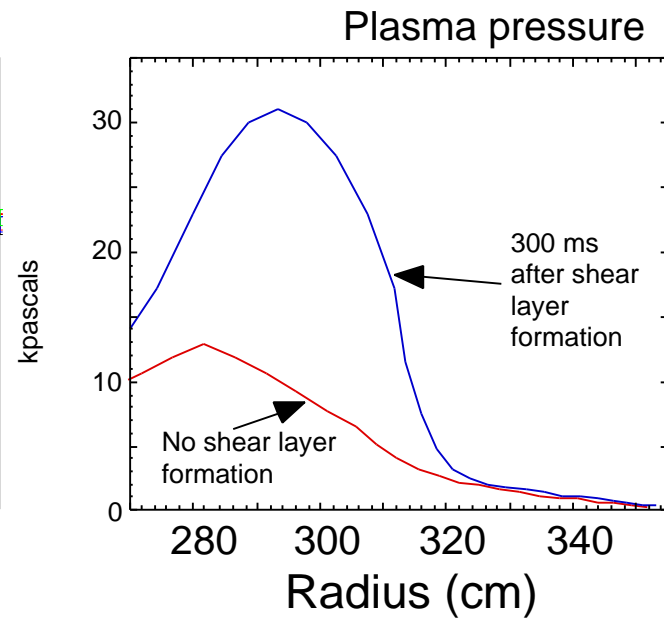
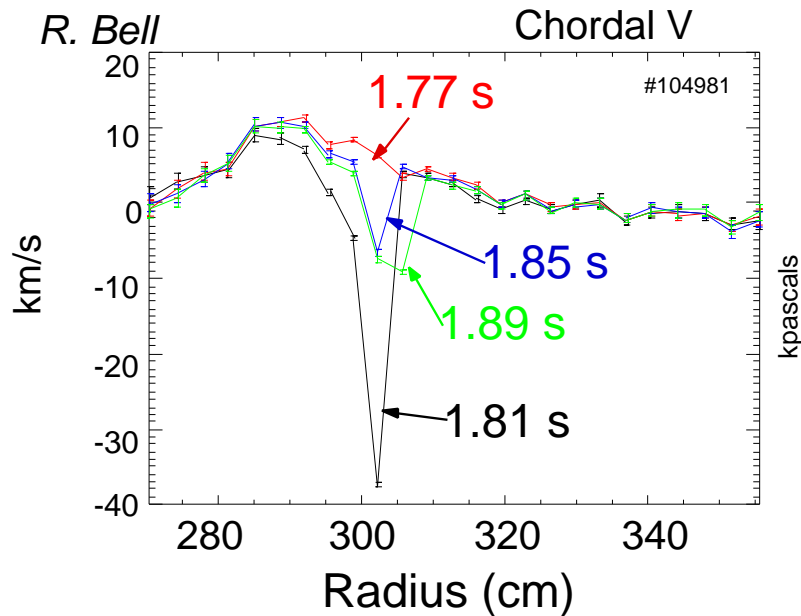
- On TFTR, $D_{He} \sim$ fluid, which would be acceptable for ash removal - BUT
- This may be an accident of having a very low i and high e

ERS Carbon Profile Peaks Due to Inward Pinch within Reverse Shear Region



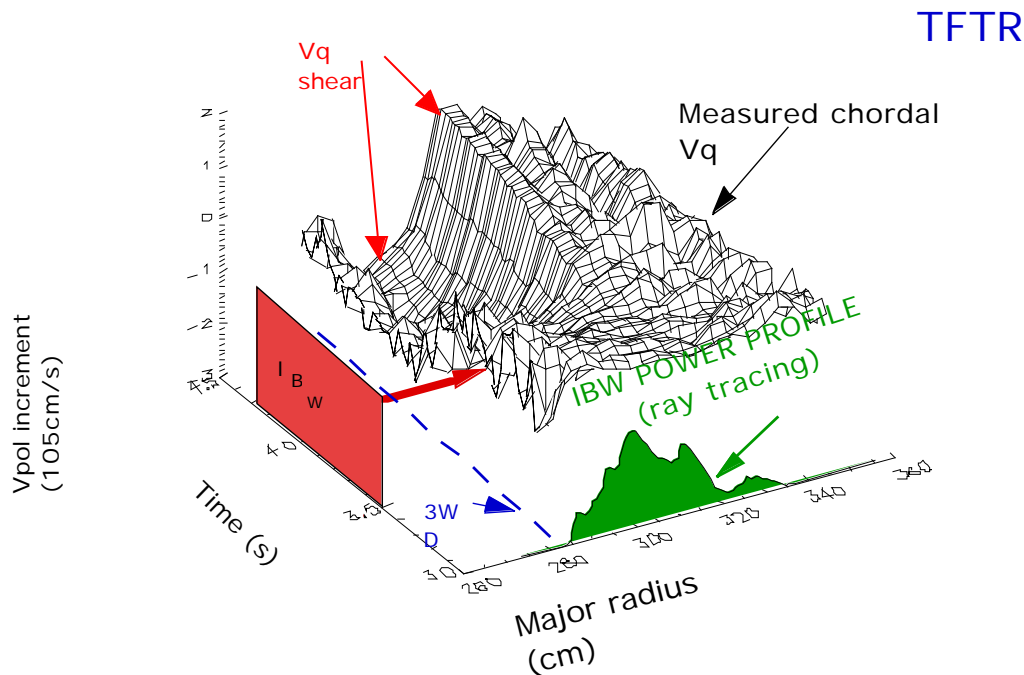
- ERS carbon profile continues to peak along with Ne after reduction in carbon source.
- RS carbon density decreases at all radii after reduction in carbon source.

E_r shear layer formation on a small scale can have large-scale consequences



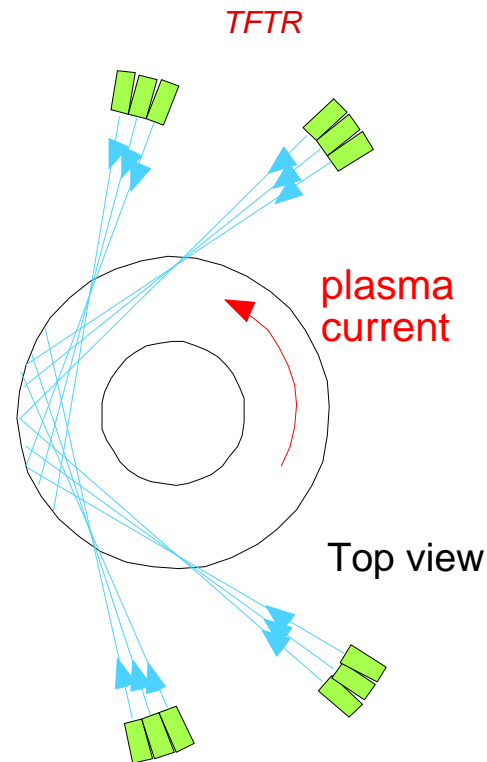
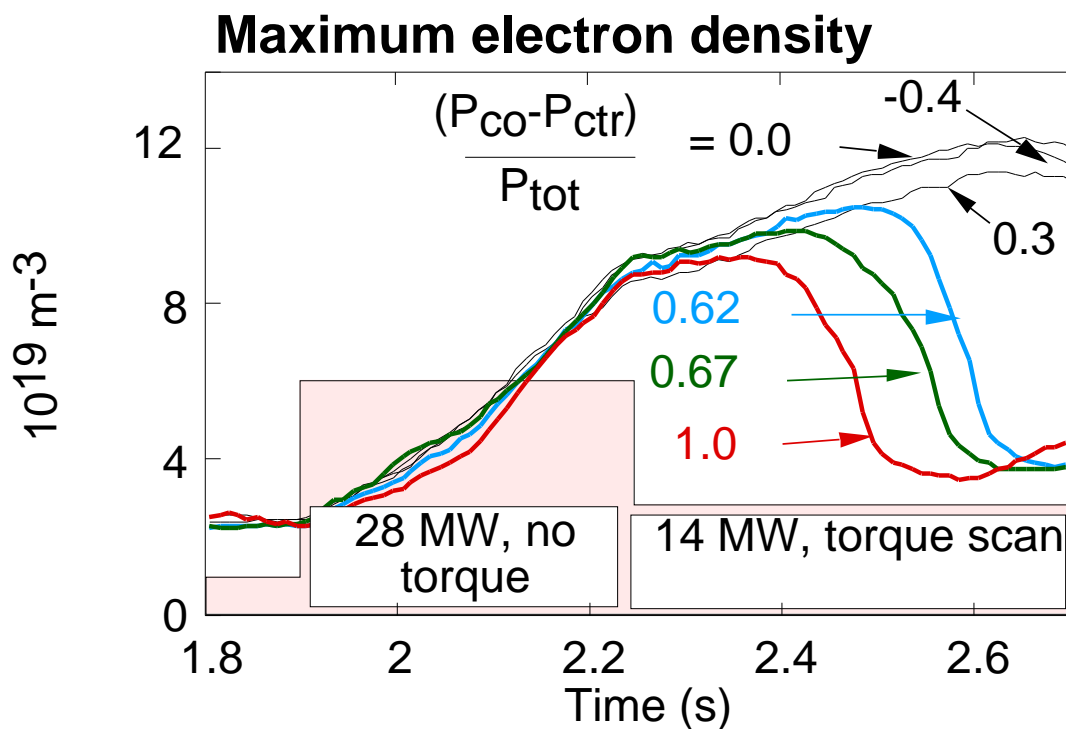
- shear layer narrower than chordal profile suggests
- plasma will respond to manipulation of E_r shear on a small scale

Poloidal Velocity Shear Produced by Ion Bernstein Waves (IBW)



- ~150 kW of IBW was coupled into TFTR out of 400 kW launched.
- Coupling may be improved with poloidal phasing or waveguide coupler. Problems likely due to surface waves.
- Calculations indicate need ~1.5 MW to form barrier

Coarse barrier control was found by varying the applied torque

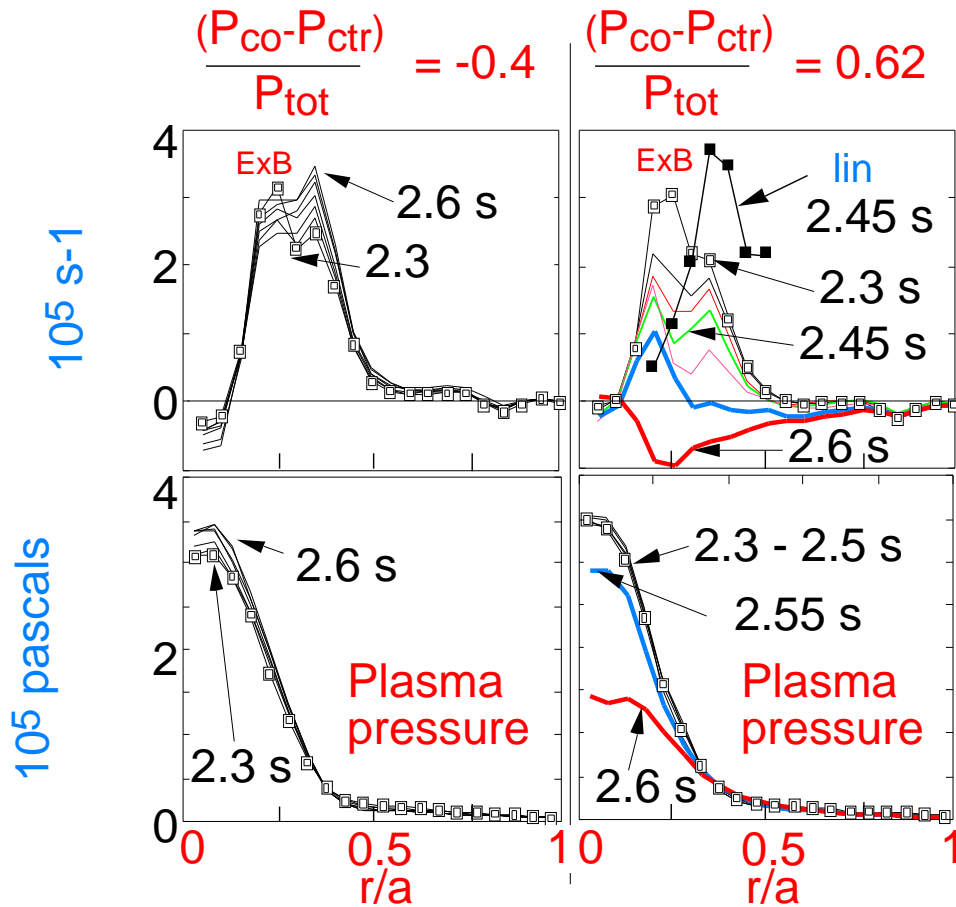


- larger fractions of injected power parallel to the current earlier losses of core confinement
- small or negative fractions antiparallel to the current no losses of core confinement

Synakowski, Phys. Rev. Lett. **78** (1997)



Pressure profile collapses only when shearing rate is forced to small values

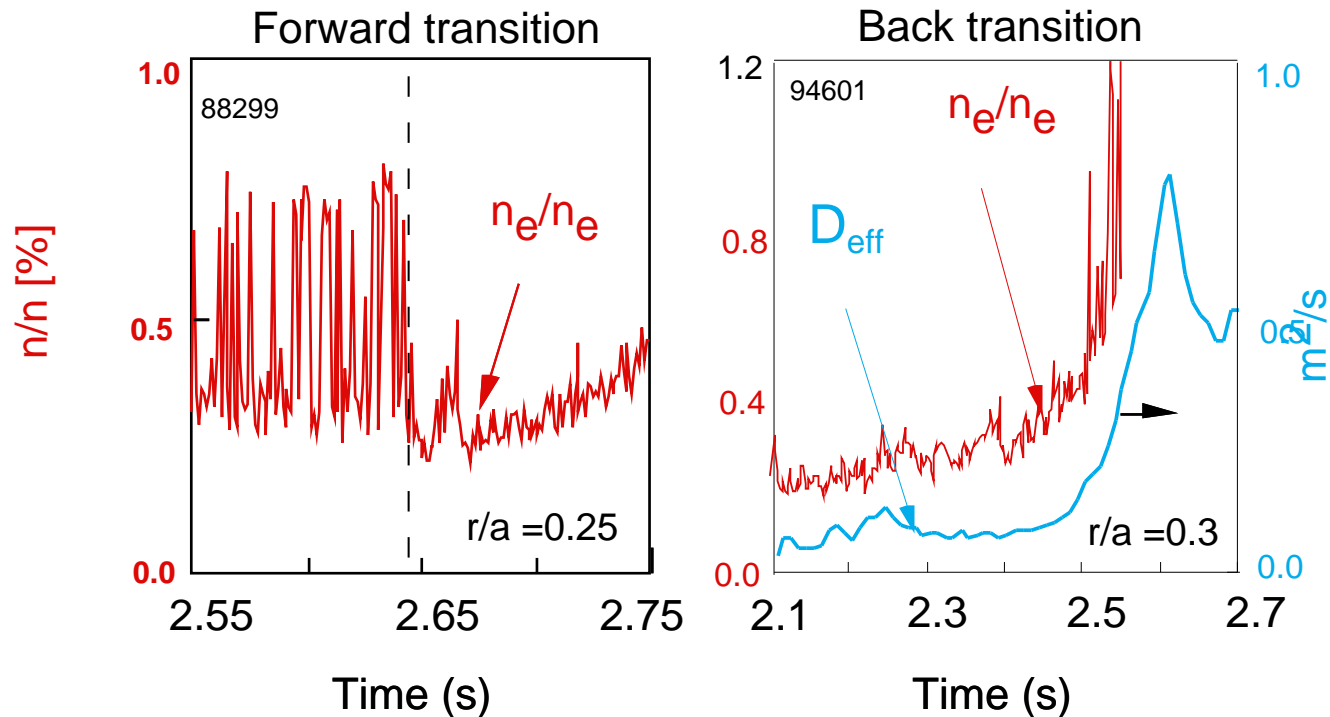


$$E_r = p / (nZe) + V B - V B$$

$$E_{\times B} = \frac{(RB)^2}{B} \frac{E_r}{RB}$$

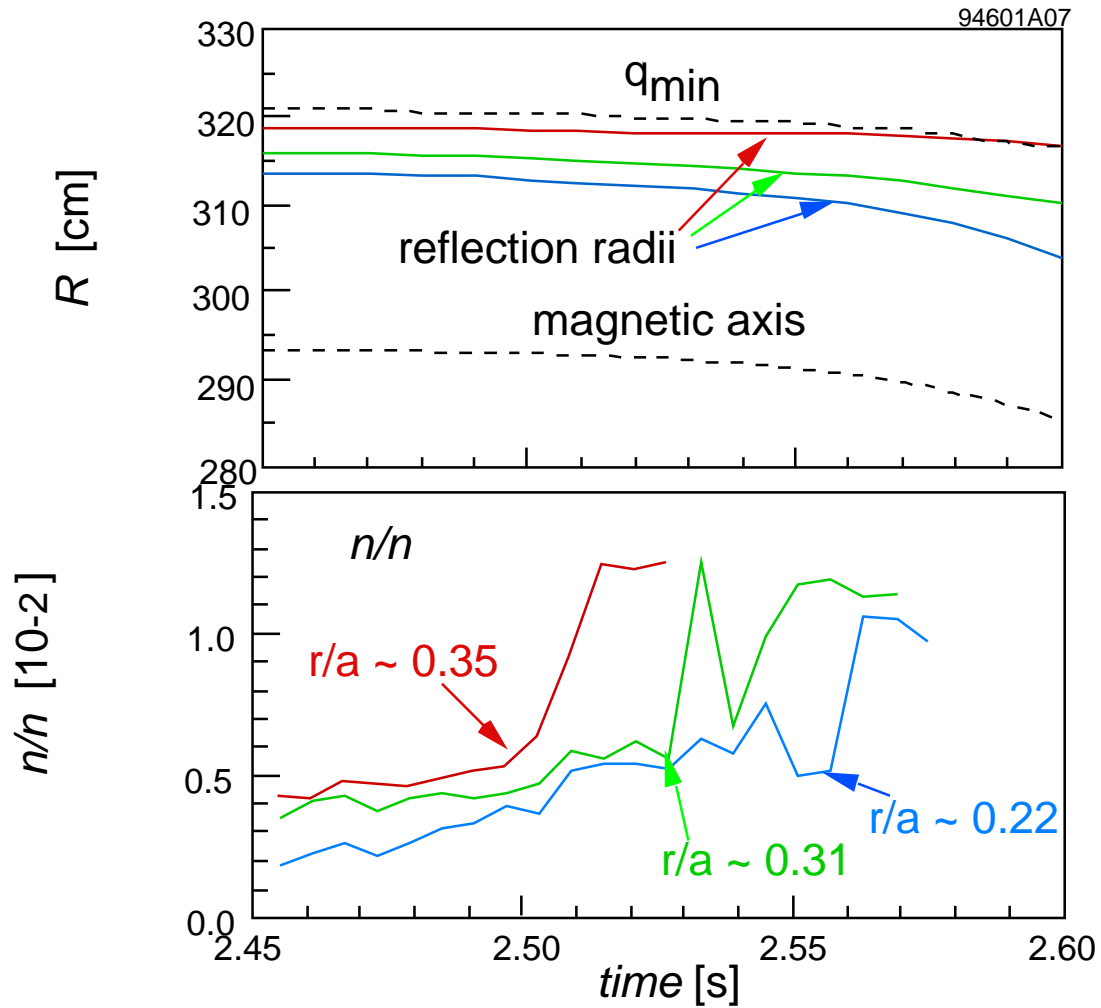
- Quantities central to i , e , and ω -induced stabilization are held constant
- Core barrier collapses when $E_{\times B} \sim 1/2 \text{ lin}$

Change in fluctuation level is fastest for forward transition



- Note difference in time interval
- *Formation*: positive feedback with ρ quick quench of fluctuations
Collapse: As barrier peels away, core is continually accessed, propping up ρ and $E \times B$ slow degradation

Measured fluctuations return from the outside in



- Fluctuations provide measure of barrier location

Calculations have been used to examine dynamics of turbulence and profiles

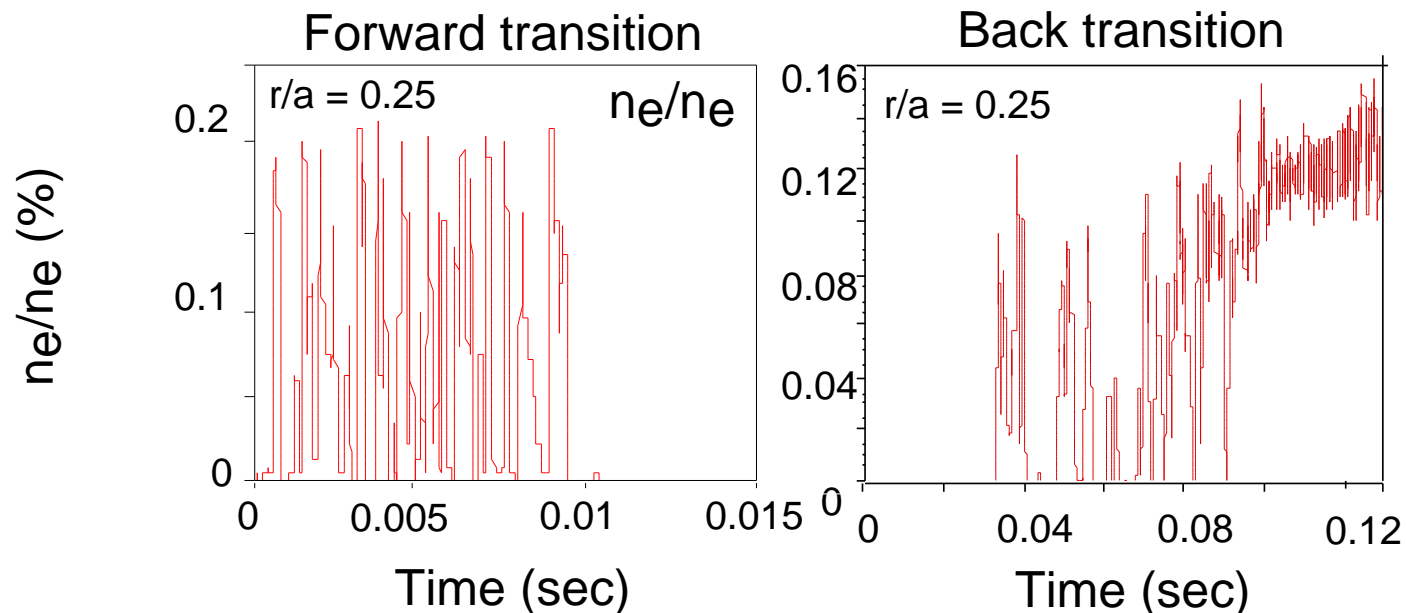
David Newman, University of Alaska at Fairbanks

- 1-D turbulence calculation: envelope model, poloidally and toroidally averaged, coupled with multimode transport model
- $E \times B$ shear explicitly included in growth and saturation of fluctuation amplitudes; phase information absent
- feedback loop between gradients, instability drive, $E \times B$ shear

Key elements for bifurcations:

- a. growth rate reduction inside q_{\min} due to reverse shear
- b. central particle and power deposition

Calculated time scale of fluctuation change is fastest for forward transition

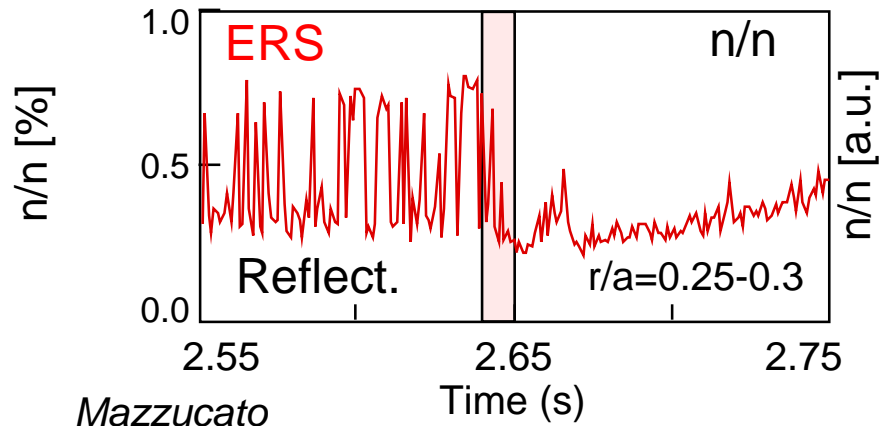


- Note difference in time intervals
- Bursting is predicted:
a consequence of interplay between instability drive and $E \times B$ shear

Core transport and fluctuation reduction can have very different time scales

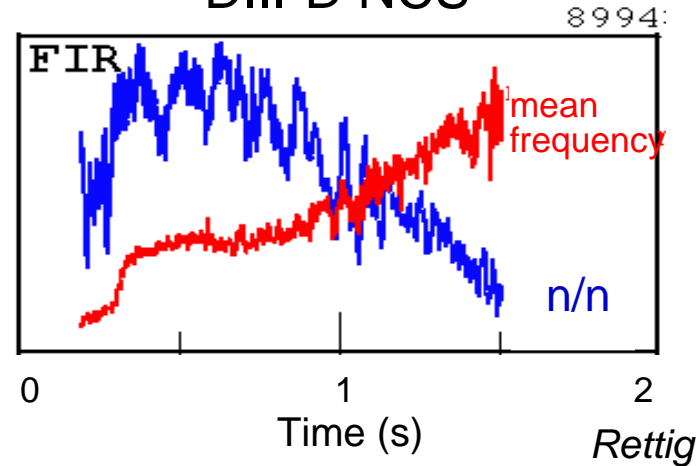
Fast

High power ERS Core (TFTR)



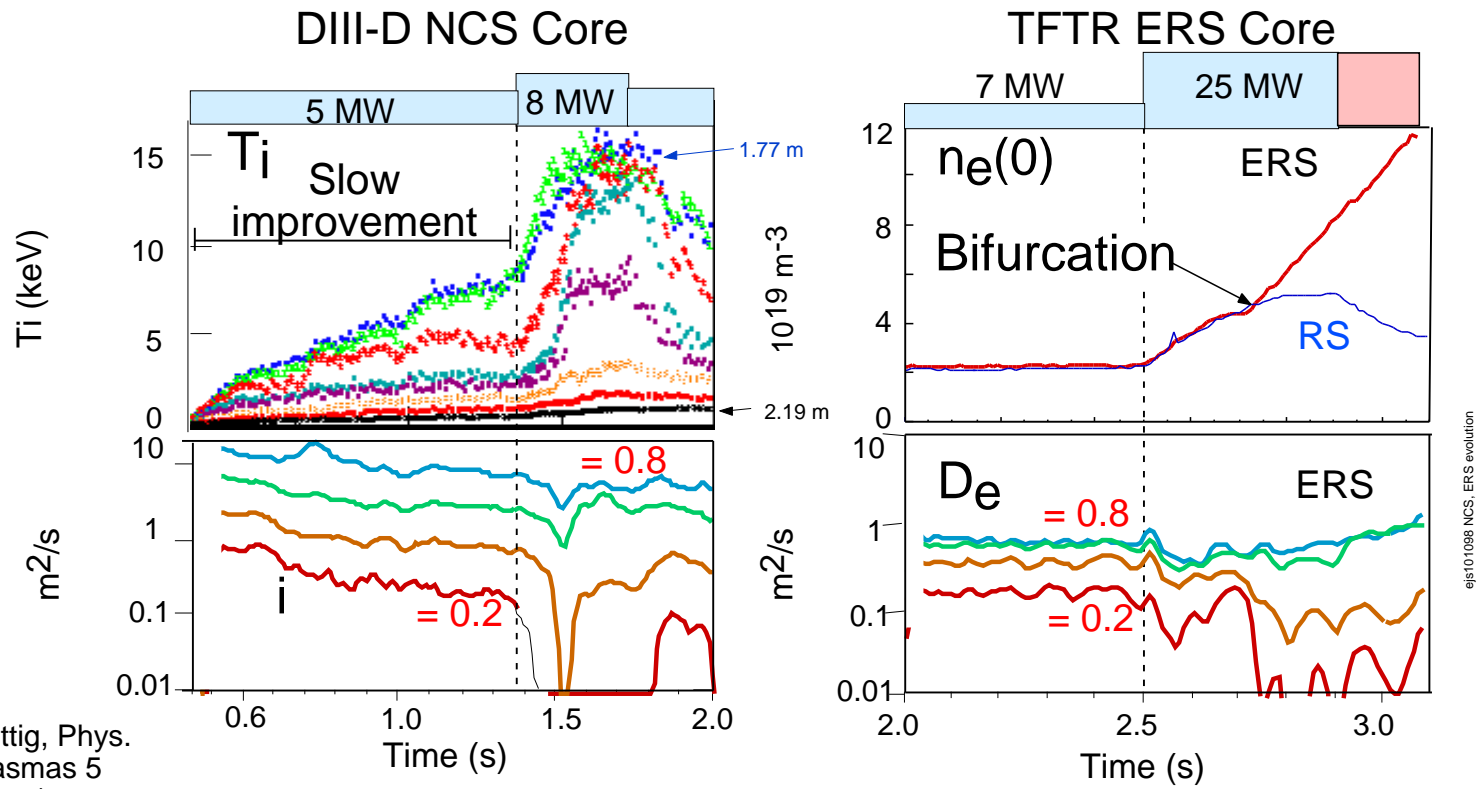
Slow

DIII-D NCS



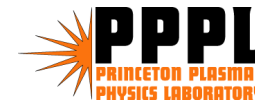
- Difference in time scale may be due to fast bootstrap with p -dominated bifurcation vs. competition between p and V in DIII-D case
- DIII-D: fluctuation, transport reduction during time when V shear is slowly increasing
- TFTR: fluctuation, transport change is “single step” in character

DIII-D NCS develops more slowly than TFTR ERS plasmas



Rettig, Phys. Plasmas 5 (1998)

- DIII-D NCS: Gradual improvement early in time
Faster improvement as V' increases with power step
- TFTR ERS: Rapid bifurcation at V' excursion



Conclusions:

1. Particle confinement provide another incentive for transport barrier control
 - neoclassical ion particle confinement in TFTR ERS
 - spotty behavior of electron thermal channel makes it even harder to predict severity of problem
2. Models with $E \times B$ shear reproduce many of the dynamics
3. Our AT programs should take advantage of flexibility in $E \times B$ shear generation if it is available
 - need to be in the configuration nearest the AT target
 - plasmas will respond to local flow shear: push RF control tools
 - flexibility of varying momentum input would enhance understanding of barrier dynamics