

Gyrokinetic simulation of tokamak edge plasma under resonant magnetic perturbations in realistic divertor geometry

Robert Hager^{a)}, C. S. Chang^{a)}, N. M. Ferraro^{a)}, A. Kleiner^{a)}, J. Lee^{b)}, R. Nazikian^{a)}

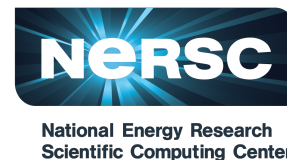
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^{a)}Princeton Plasma Physics Laboratory

^{b)}Korea Institute of Fusion Energy

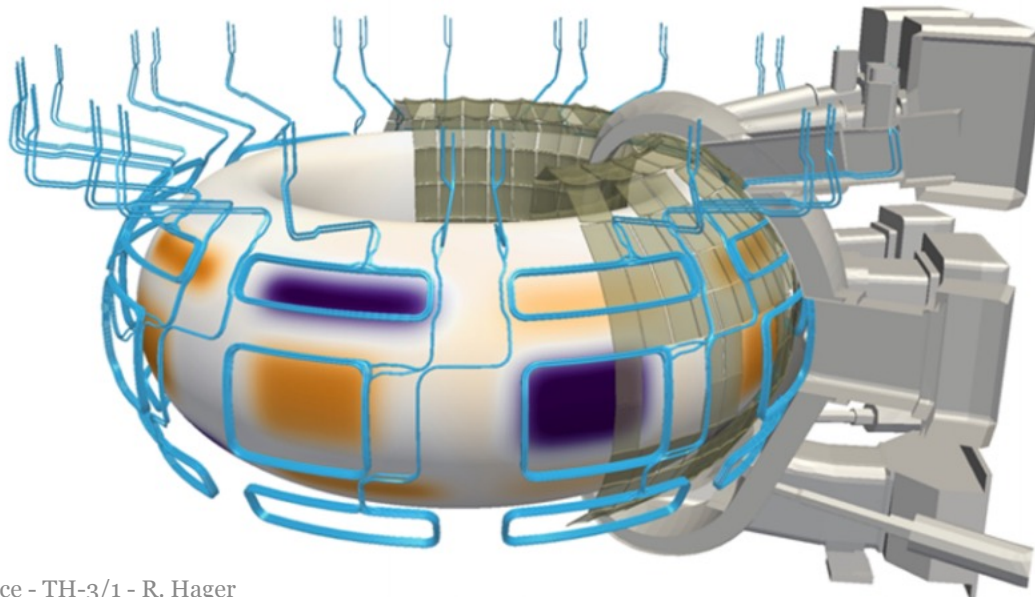
Thanks to:

Qiming Hu, Brian Grierson, Hyeon K. Park, Won-Ha Ko, the XGC team, the DIII-D Team, and the KSTAR team



- Introduction
- Numerical Approach
- Using DIII-D and KSTAR plasmas
 - Change of turbulence properties by RMPs
 - Change of transport by RMPs
- Conclusions and Discussion

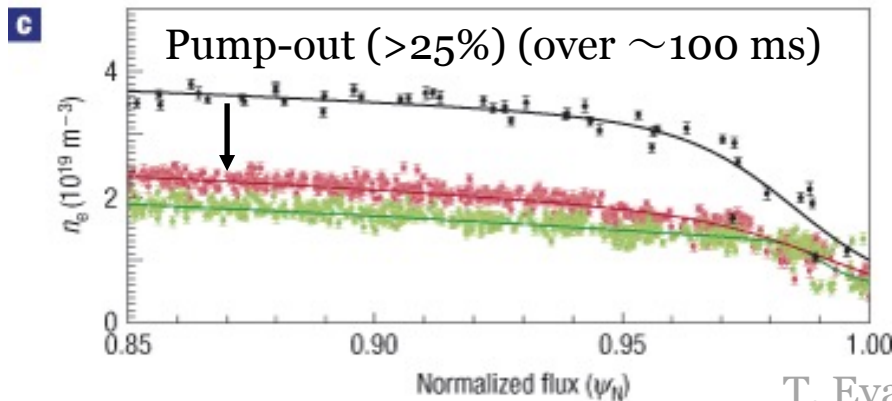
ITER 3D coils





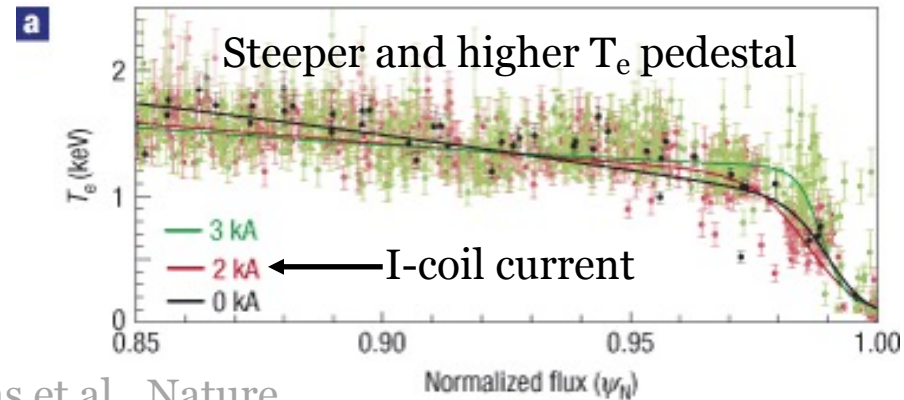
- ITER plans to use 3D fields, **R**esonant **M**agnetic **P**erturbations (RMP), for ELM suppression

What is the physics behind the density pump-out?



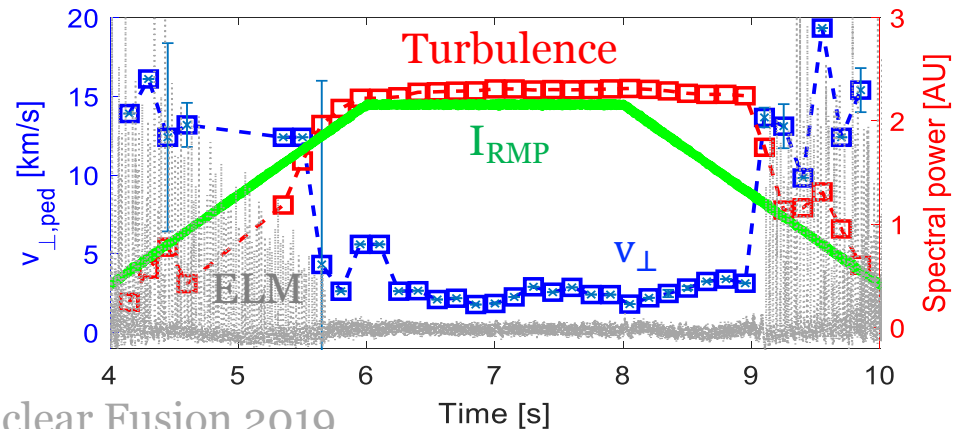
T. Evans et al., Nature 2006

Why is electron heat still confined?

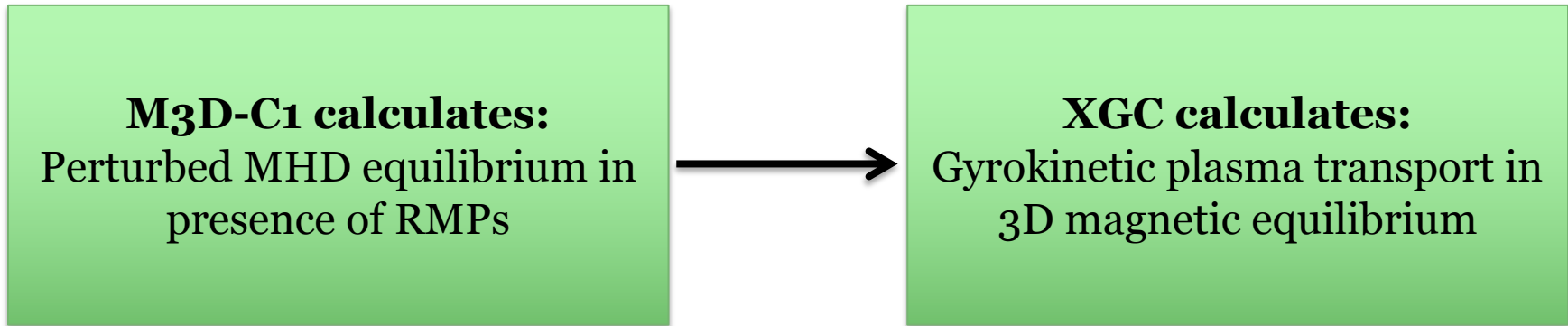


Why is turbulence intensity higher with RMPs?

Apparent correlation between I-coil current, turbulence intensity, v_{\perp} and ELM intensity



H. K. Park et al., Nuclear Fusion 2019

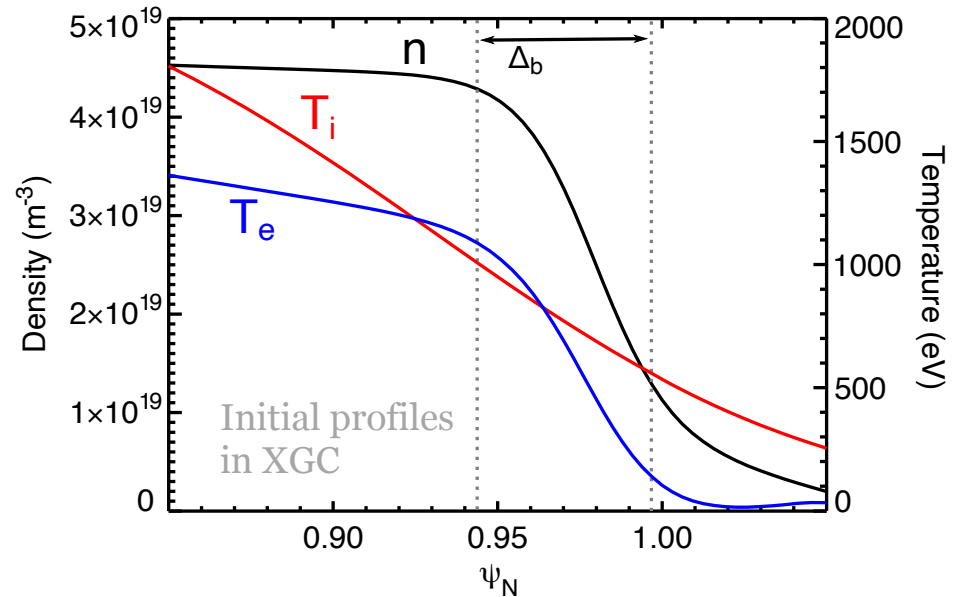
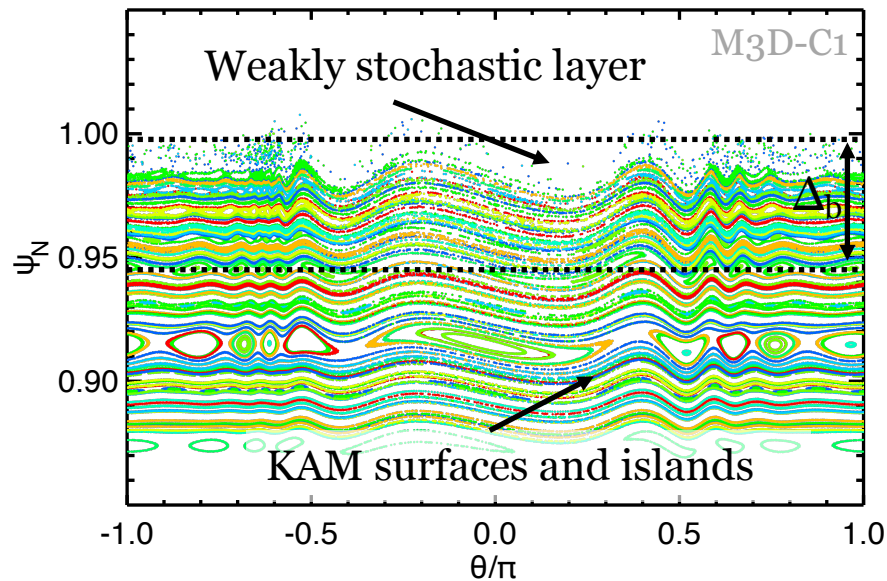


- M3D-C1 is a two-fluid extended-MHD code
 - Domain includes plasma, SOL, and surrounding vacuum region.
 - Uses Spitzer resistivity and realistic transport coefficients ($\sim 1 \frac{\text{m}^2}{\text{s}}$)
 - Perturbed equilibrium includes small magnetic islands and localized regions of stochasticity
- XGC is a **global 5D gyrokinetic**, total-f particle-in-cell code
 - Whole volume simulation including SOL, separatrix, and magnetic axis
 - Self-consistent 3D background and turbulent electric field
 - Nonlinear Fokker-Planck-Landau collision operator
 - Neutral particle recycling



- DIII-D H-mode discharge #157308 → M3D-C1 yields perturbed field with good KAM surfaces at $\psi_N \leq 0.98$
- Thin, weakly stochastic layer close to the separatrix ($\psi_N \geq 0.98$)
- δB is strong enough to affect trapped particle dynamics
- Pedestal width comparable to ion banana orbit width Δ_b

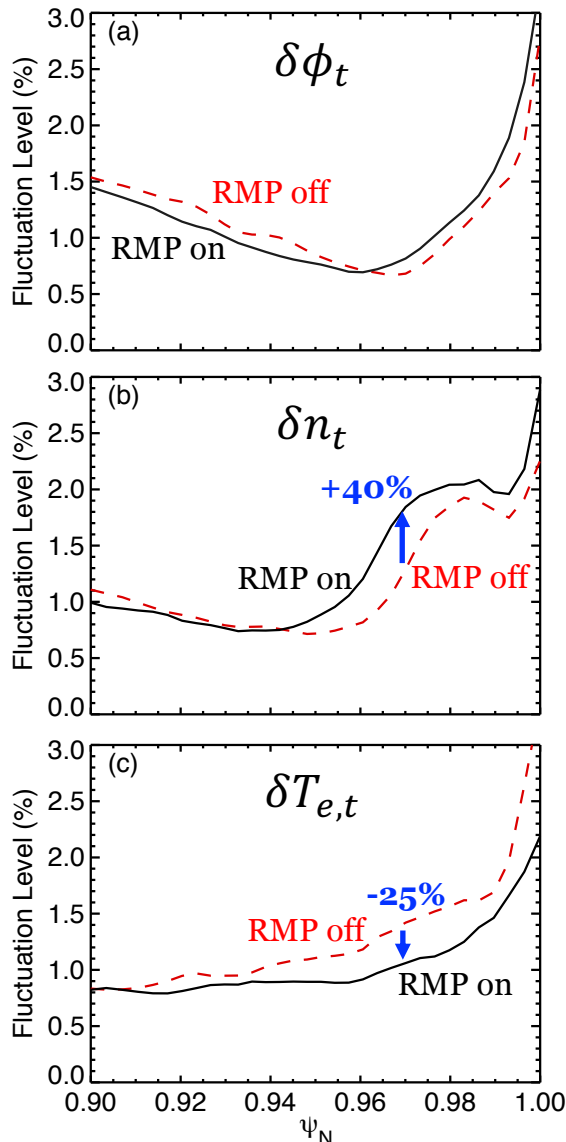
→ **Need to study kinetic physics to understand transport!**



RMPs Increase Density Fluctuations and Decrease Temperature Fluctuations in the Pedestal

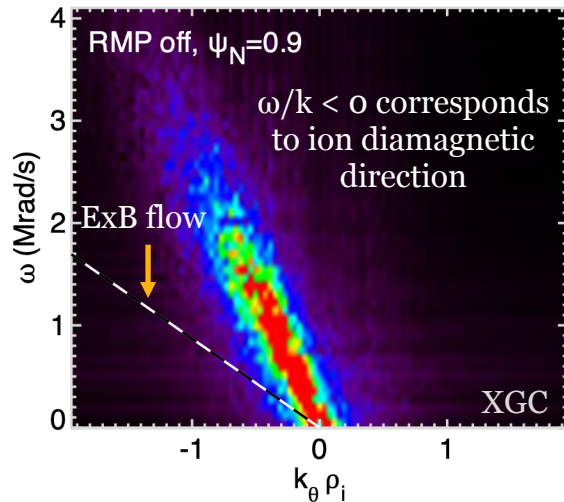


XGC

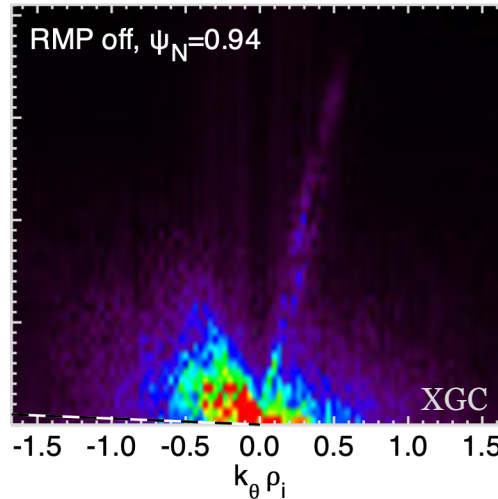


- n=3 mode (RMP) is removed to study changes in relative RMS turbulence fluctuation levels
- Potential fluctuations change only by ~10%.
- Density fluctuations increase with RMP.
- Electron temperature fluctuations decrease with RMPs.
- These changes are correlated with changes in the transport fluxes.
- δp_e is minimized by restricting δT_e .

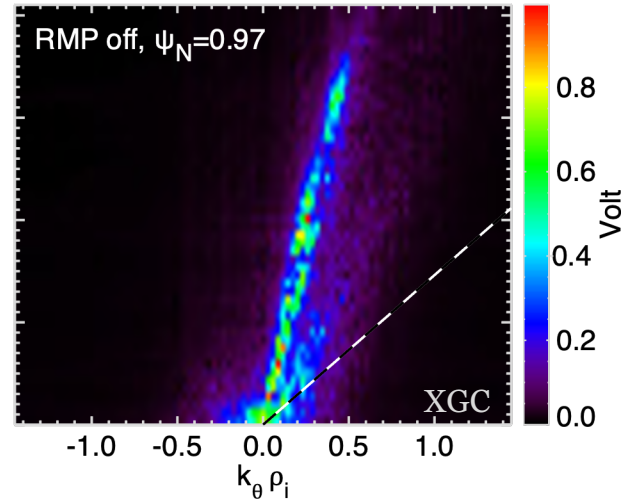
Spectra suggest enhanced TEM in pedestal slope. ITG deeper inside does not change as much (at $t \sim 0.2$ ms)



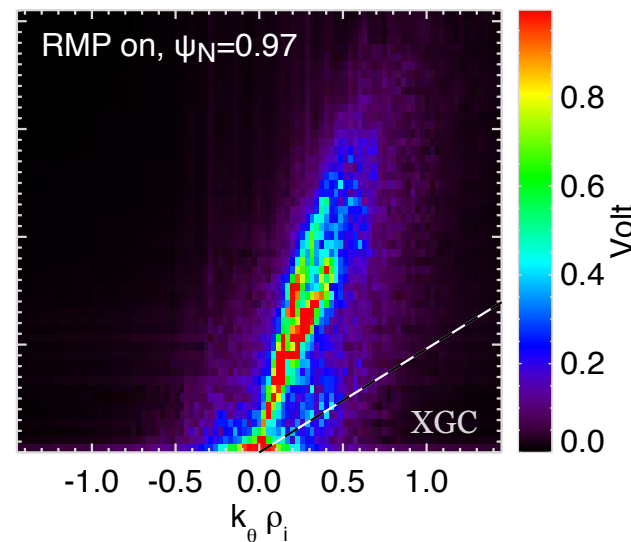
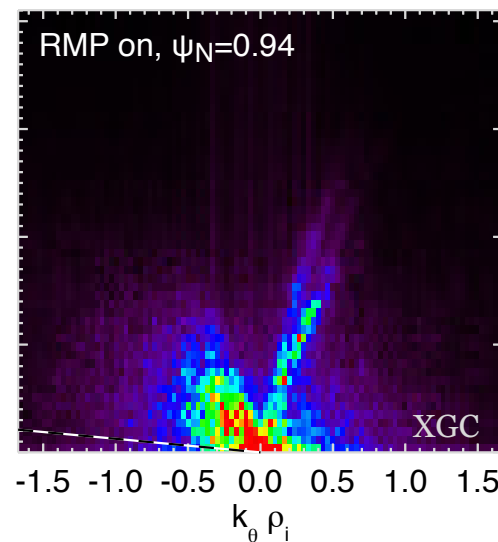
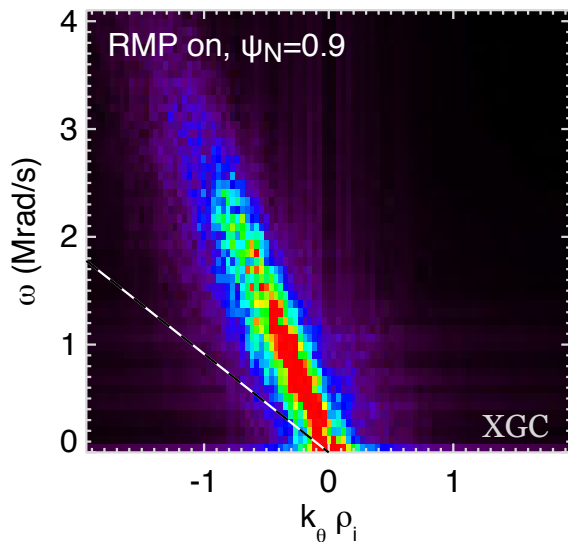
Ion mode



Transition between ion and electron mode



Electron mode





There are three main transport channels:

- **Neoclassical flux**

$$\Gamma_D = \frac{\langle \int [\nabla\psi \cdot (\mathbf{v}_D + \bar{\mathbf{v}}_{ExB}) \bar{f}] d^3v \rangle}{\langle |\nabla\psi| \rangle}$$

- **Flutter transport**

$$\Gamma_{3D} = \frac{\langle \int [\nabla\psi \cdot (\delta\mathbf{B}/|B|) v_{\parallel} \tilde{f}] d^3v \rangle}{\langle |\nabla\psi| \rangle}$$

- **Turbulent ExB flux**

$$\Gamma_{turb} = \frac{\langle \int [\nabla\psi \cdot \tilde{\mathbf{v}}_{ExB} \tilde{f}] d^3v \rangle}{\langle |\nabla\psi| \rangle}$$

$$\Gamma_{neo} \stackrel{\text{def}}{=} \Gamma_D + \Gamma_{3D}$$

$$f = \bar{f} + \tilde{f}$$

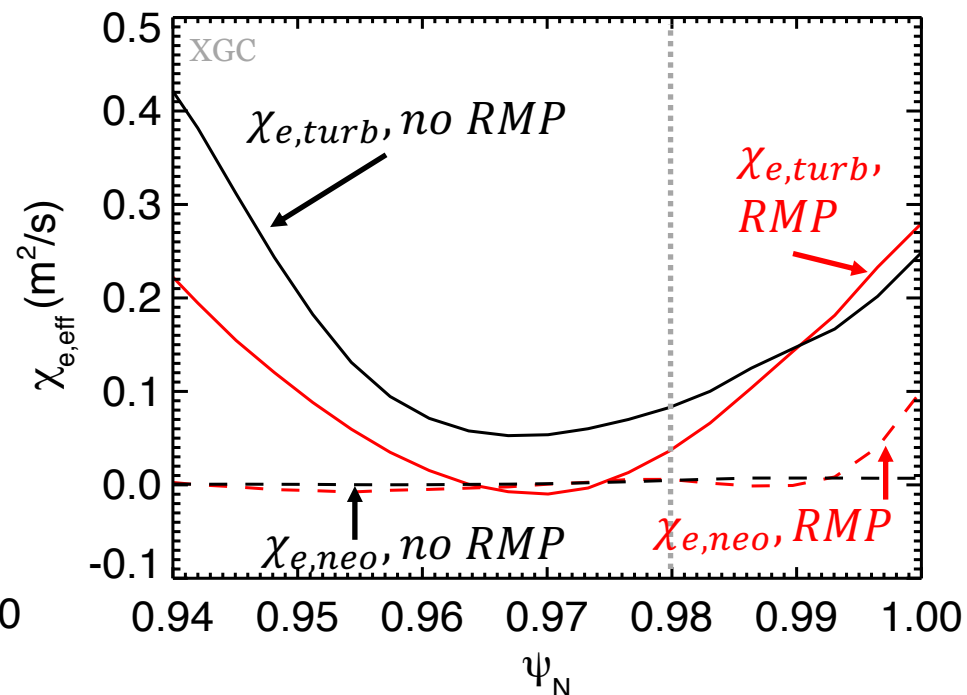
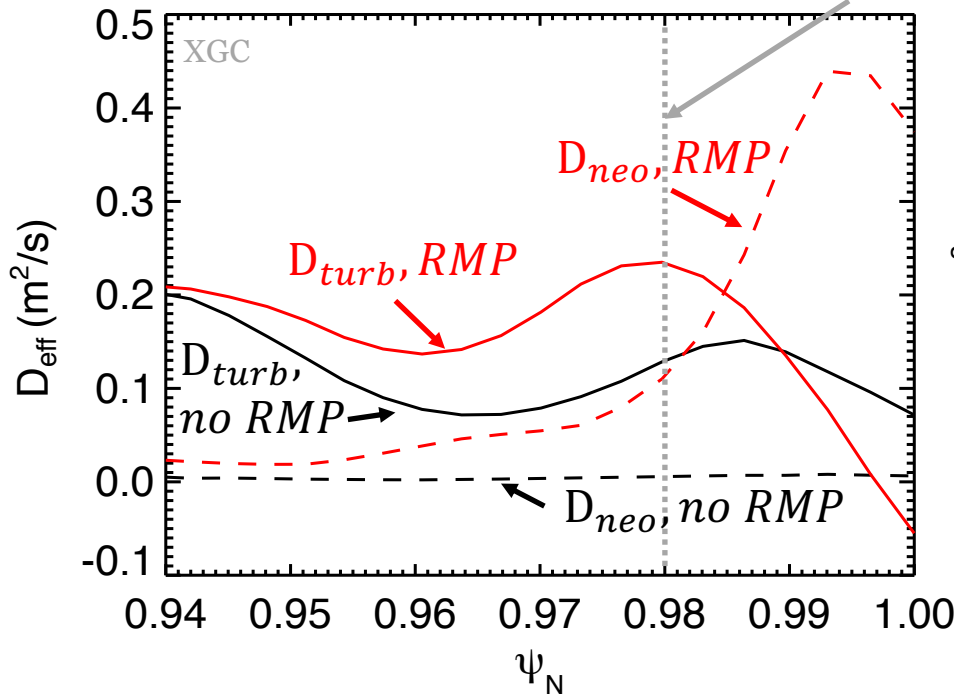
$\overline{(\dots)} \rightarrow$ toroidal average



Increased turbulent+neoclassical particle diffusivity with RMPs

Lower effective electron heat conductivity

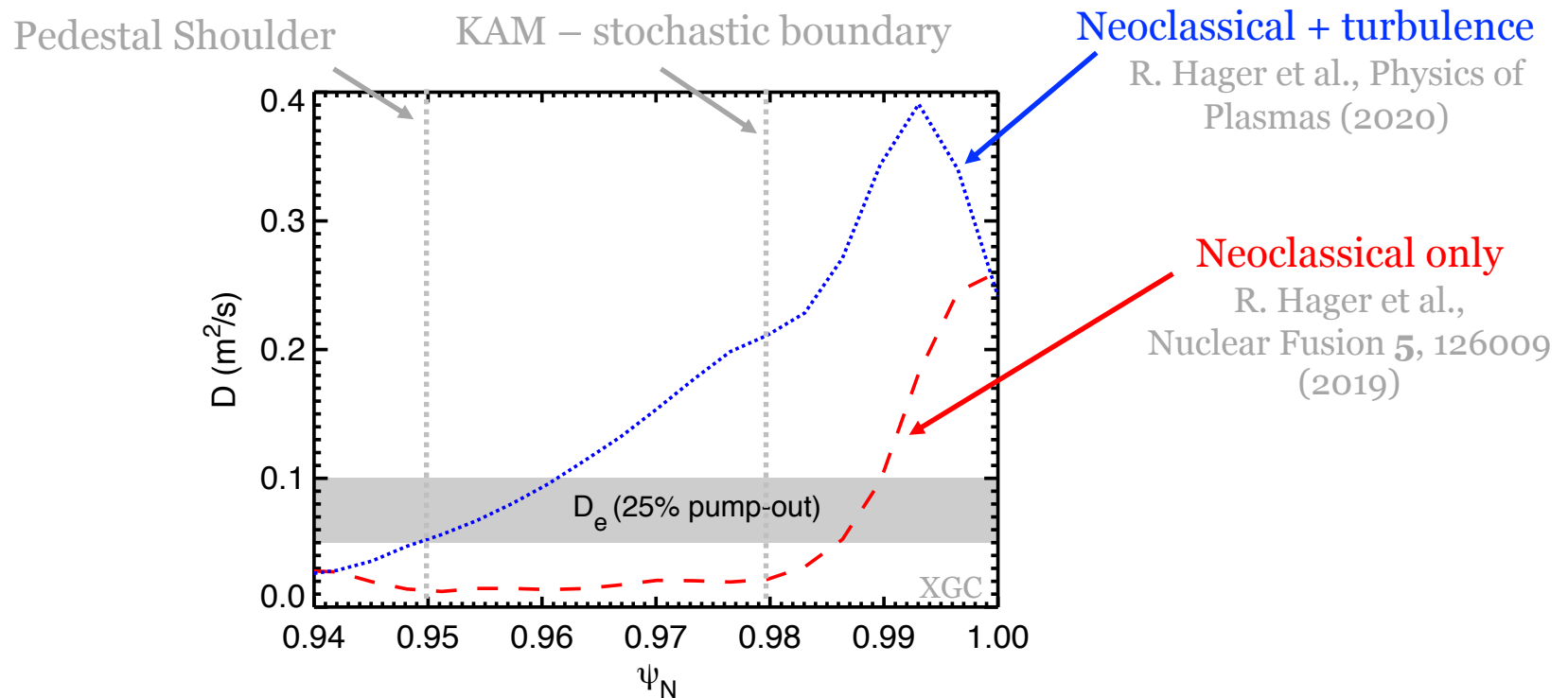
KAM – stochastic boundary



→ Electron thermal transport barrier in the steep pedestal region survives



- RMP-driven increase of **neoclassical+turbulent** particle diffusivity is largely sufficient for density pump-out in the steep pedestal region



→ Increase of turbulent transport boosts pump-out from enhanced neoclassical transport



- Find the origin of the increased turbulent particle flux density in:
 - Higher turbulence amplitude or
 - Shifted cross-phase between turbulent fluctuations
 - Which mode numbers are responsible?

Cross-spectrum: $S_{AB} = \langle \hat{A}(\psi_N, m, \varphi) \hat{B}^*(\psi_N, m, \varphi) \rangle_{tor}$

Cross-power: $P_{AB} = |S_{AB}|$

Cross-phase: $\delta\zeta_{AB} = \arctan[\Im(S_{AB})/\Re(S_{AB})]$

Turbulent transport fluxes in terms of cross-power and cross-phase:

$$\Gamma_t(\psi_N, m) = \alpha(\psi_N) P_{vn} \cos(\delta\zeta_{vn})$$

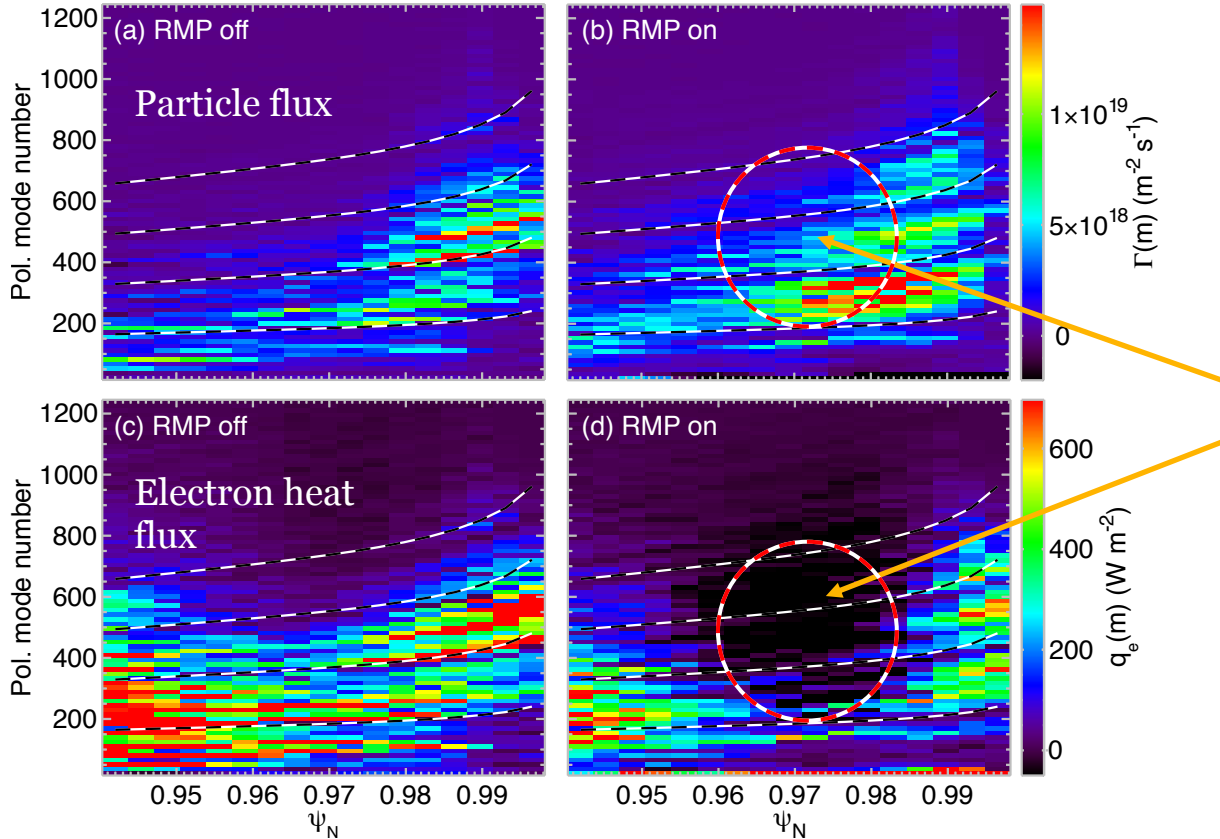
$$Q_t(\psi_N, m) = \alpha(\psi_N) \frac{3}{2} k_B [\langle n \rangle P_{vT} \cos(\delta\zeta_{vT}) + \langle T \rangle P_{vn} \cos(\delta\zeta_{nT})]$$

$$q_t(\psi_N, m) = Q_t - \frac{5}{2} k_B \langle T \rangle \Gamma_t$$

(α is a geometric factor from the flux-surface average.)



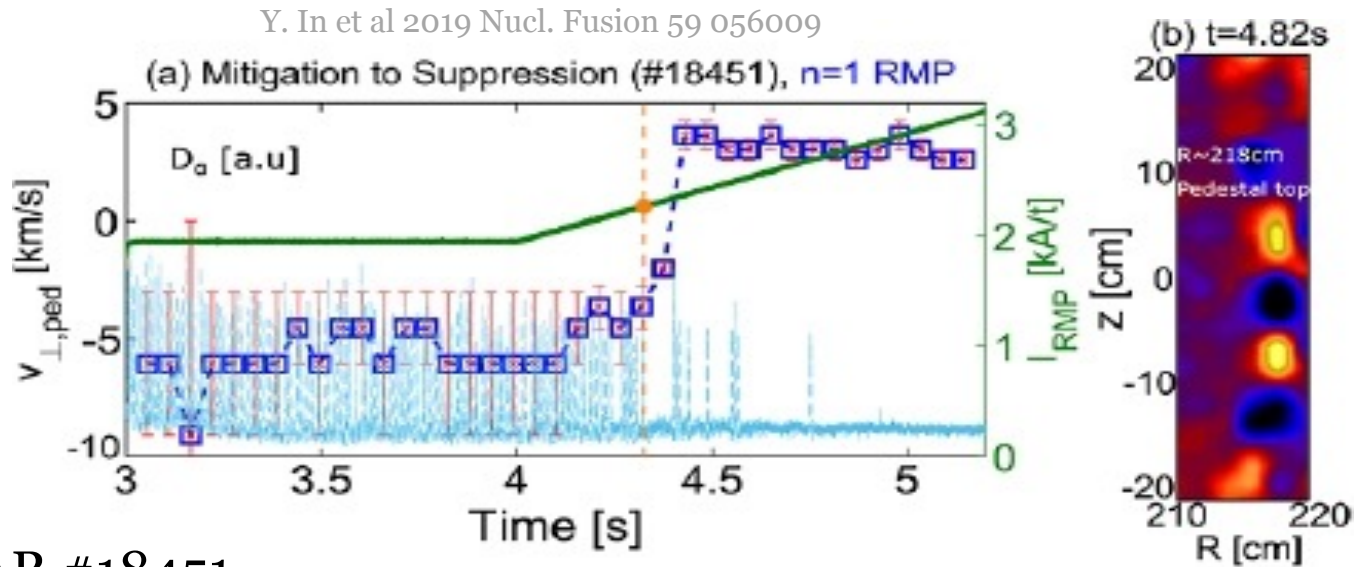
XGC cross-spectrum analysis



RMPs increase particle flux and decrease electron heat flux density around $\psi_N \sim 0.97$ at higher poloidal mode numbers $m \geq 200$ ($k_{\theta} \rho_i \gtrsim 0.2-0.3$).

→ Electron energy transport is convective, riding the particle flux, which does not alter the T_e gradient much.

R. Hager et al., Physics of Plasmas (2020)

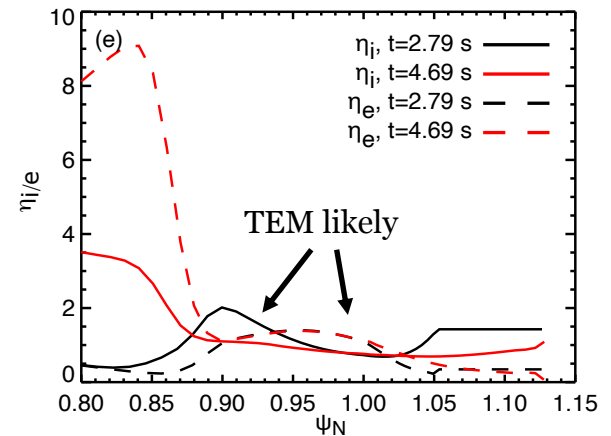
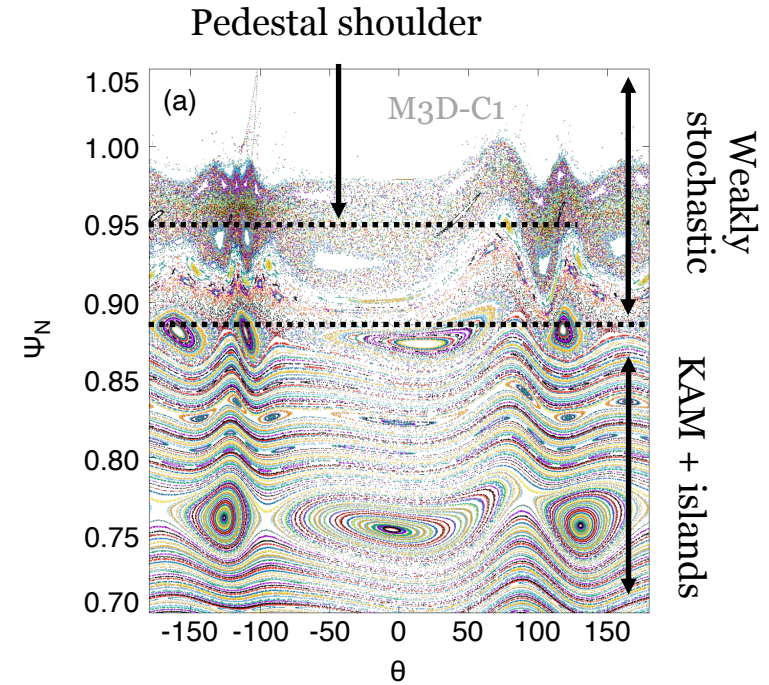
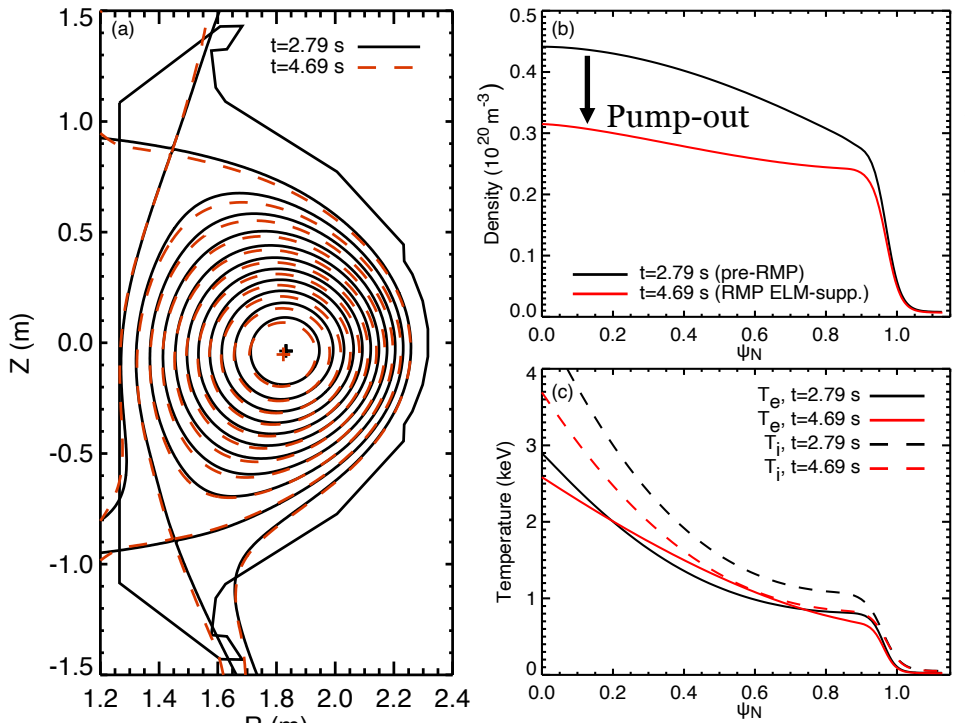


- KSTAR #18451
 - ELMing H-mode at $t=2.79$ s before RMP application
 - ELM mitigation at $I_{RMP}=2$ kA/t
 - ELMs suppressed at $t \geq 4.4$ s
- Study two time slices with M3D-C1+XGC
 - $t=2.79$ s, $I_{RMP}=0$ kA/t \rightarrow Before RMP application
 - $t=4.69$ s, $I_{RMP}=2.69$ kA/t \rightarrow ELM suppression

Pedestal ($\psi_N \geq 0.9$) is Weakly Stochastic \rightarrow Expect Stronger Role of Neoclassical Transport



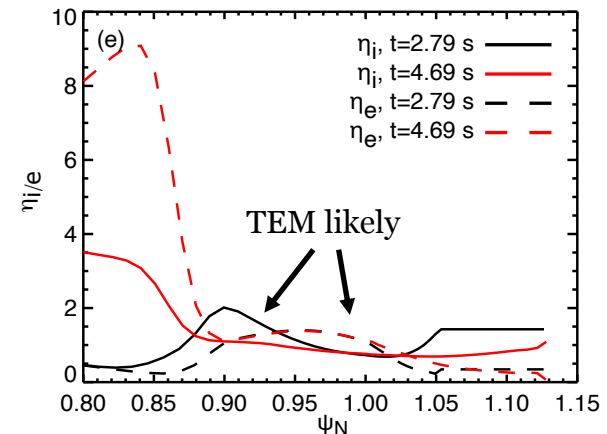
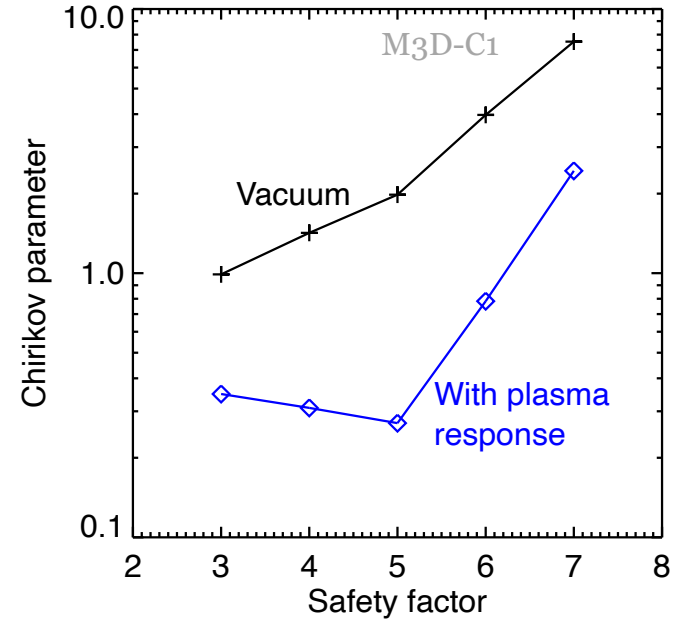
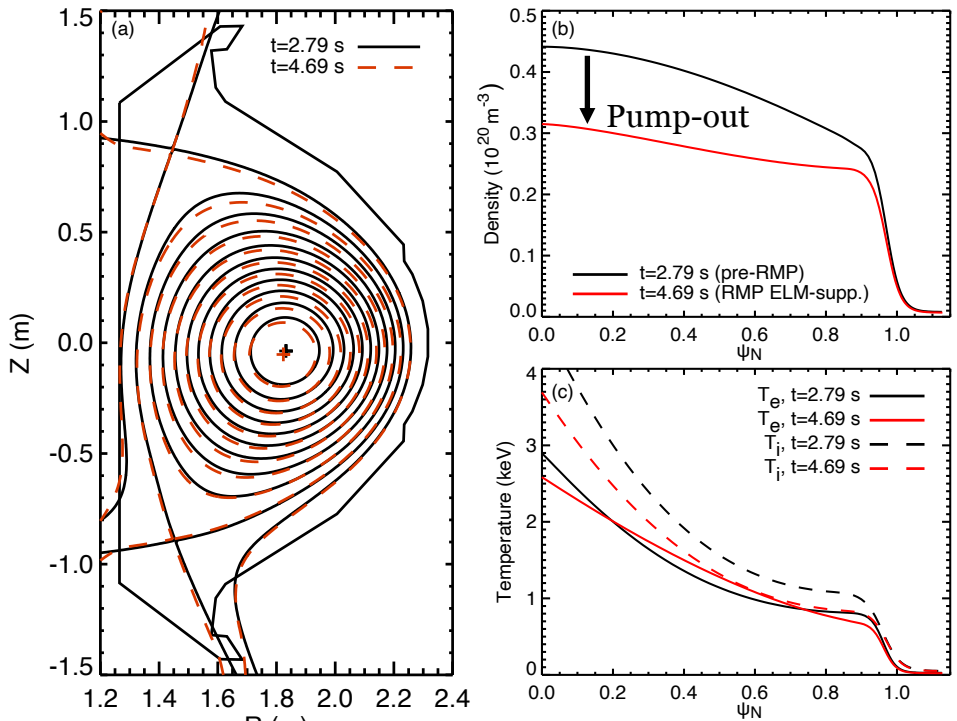
- Plasma shapes and safety factor profile are very similar at $t=2.79$ s and $t=4.69$ s
- Pump-out: central density reduced by $\sim 25\%$
- Perturbed magnetic field weakly stochastic in the pedestal (from the $q=4$ surface outwards)
- Low η_i and $\eta_e \rightarrow$ Trapped electron mode (TEM) likely unstable



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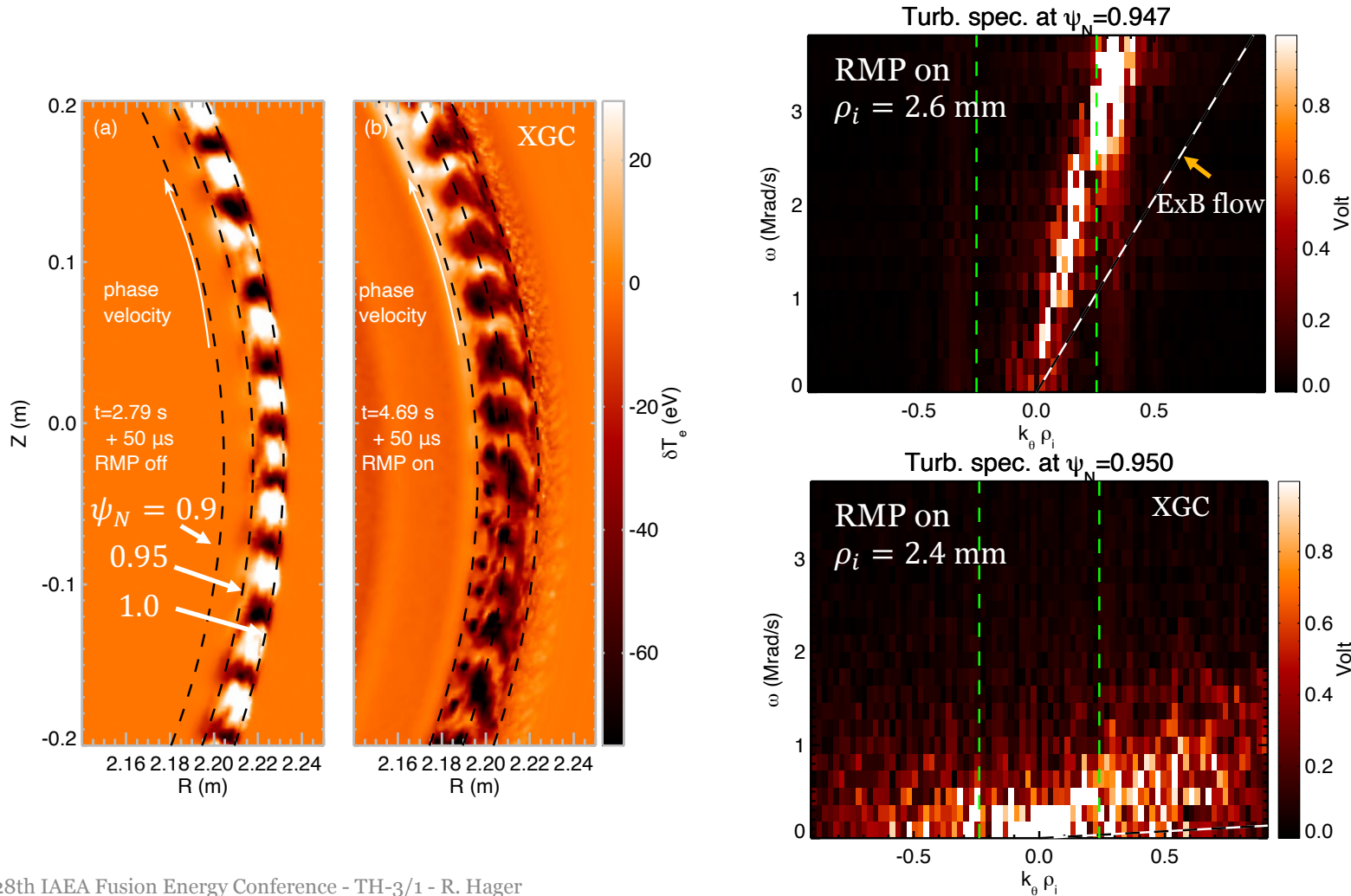
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XGC Simulation: Unstable Trapped-electron Modes (TEMs) extends to pedestal top, covering $0.9 \leq \psi_N \leq 1$



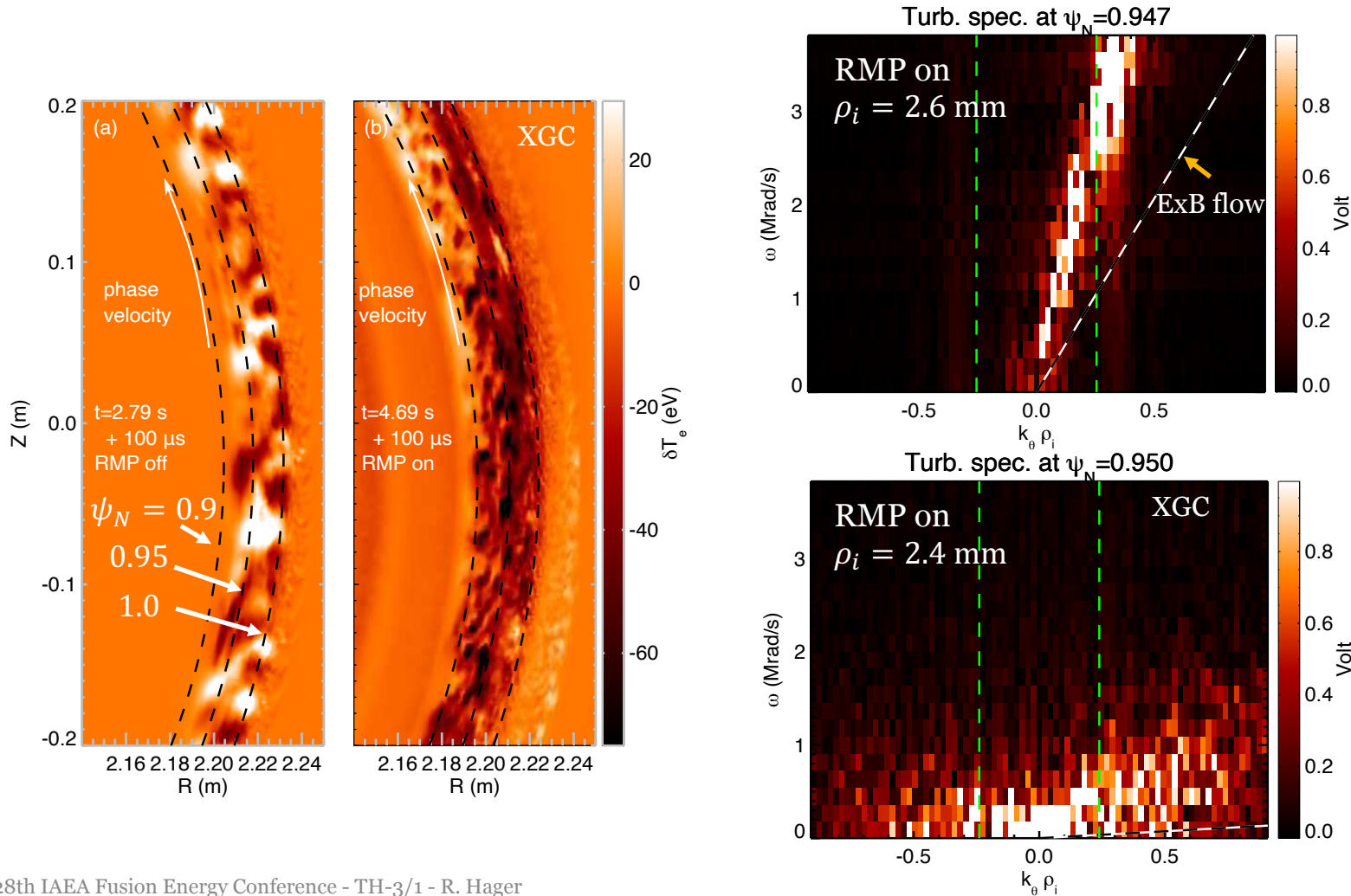
- Mode structure propagates in elec. diamagnetic direction in ExB frame \rightarrow TEM
- Turbulence grows faster around pedestal top with RMPs.
- Need synthetic diagnostic for better comparison with experiment



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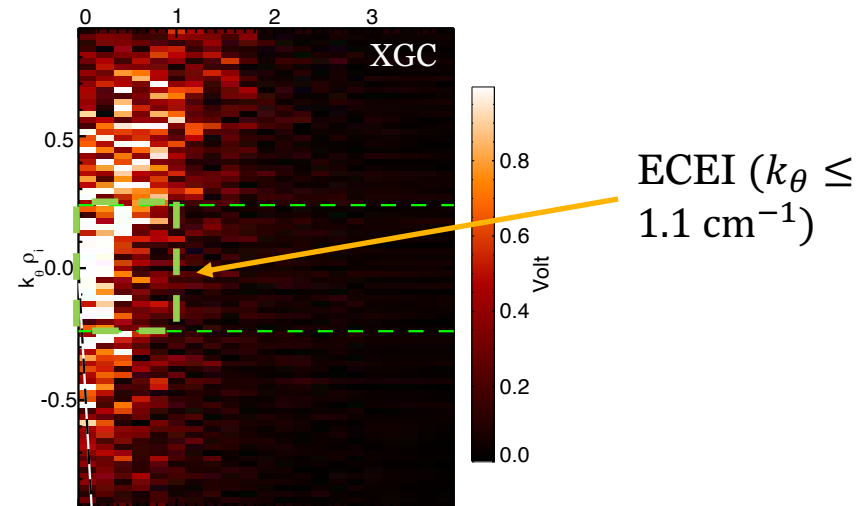
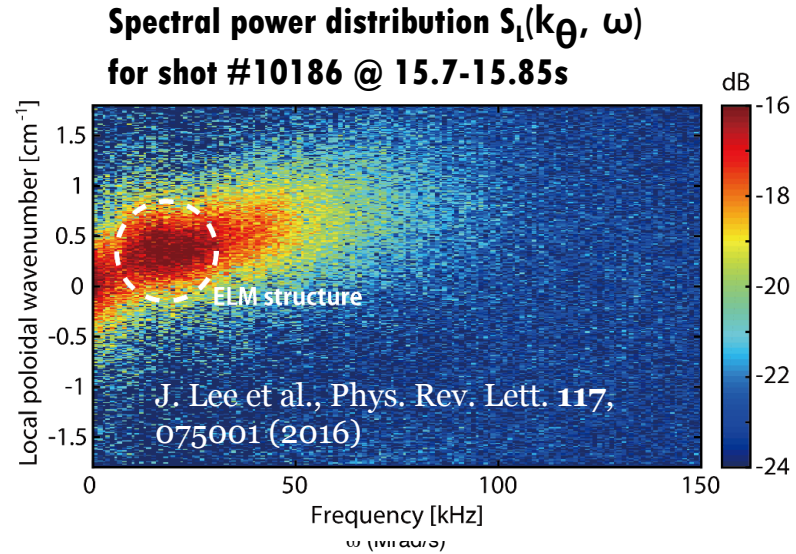
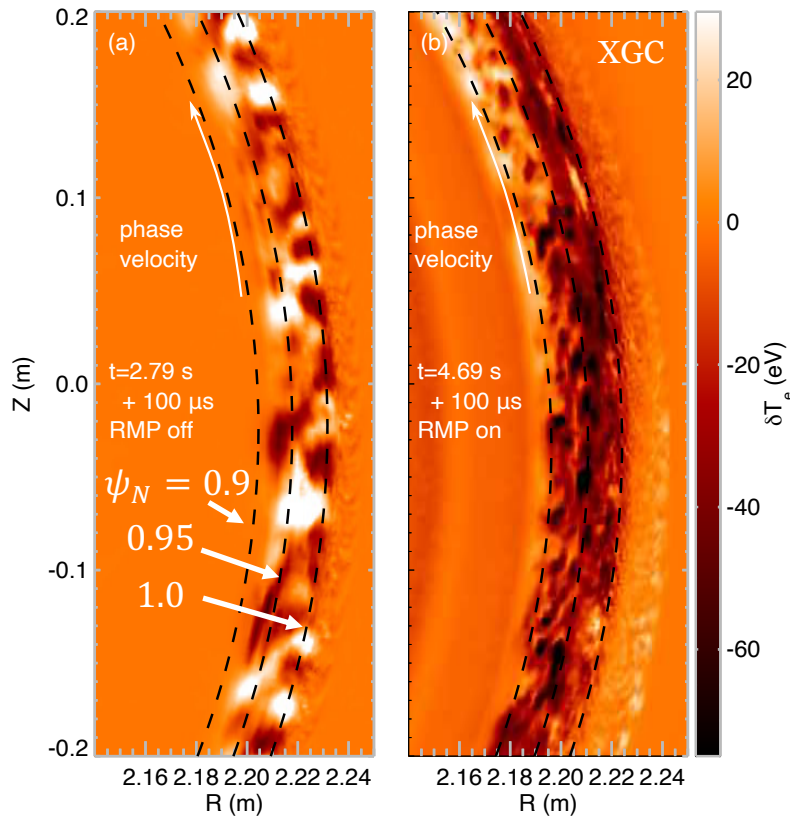
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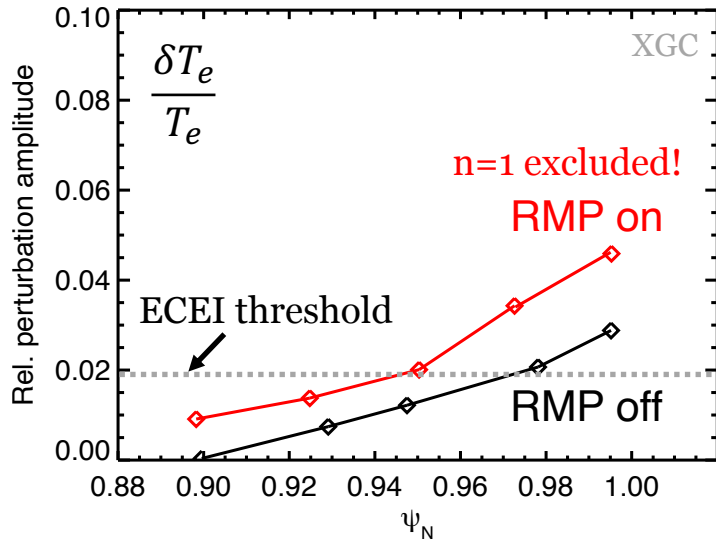
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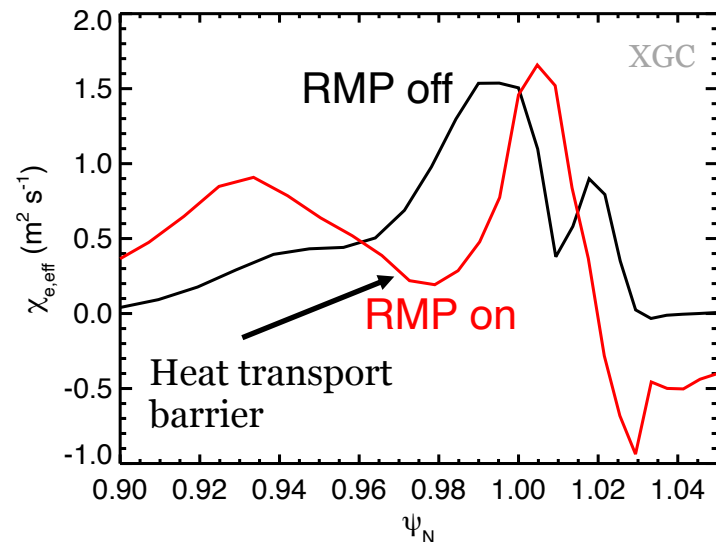
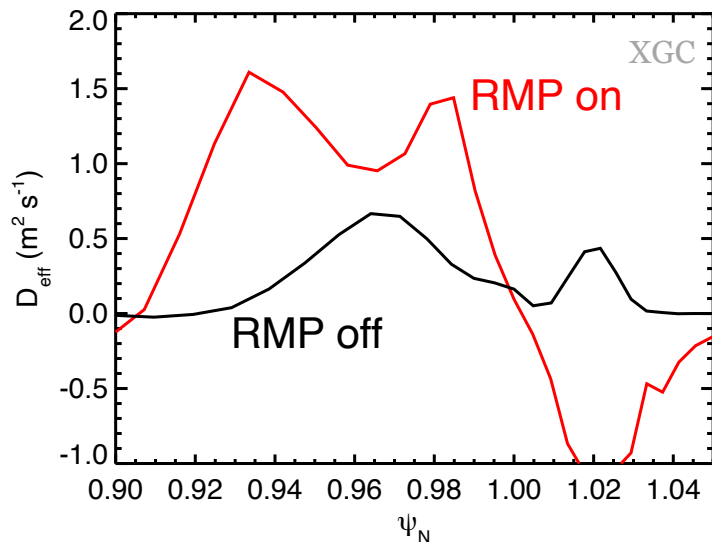
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XGC Simulation Exhibits Similarities to KSTAR, but Electromagnetic Turbulence may be Missing



- Relative perturbation amplitudes $\frac{\delta\phi}{T_e}, \frac{\delta n}{n}, \frac{\delta T_e}{T_e}$ are $\gtrsim 2x$ higher (experiment: 4x)
- Particle flux is enhanced in the pedestal with contribution from turbulence.
- High neoclassical electron heat flux, but heat barrier at $\psi_N \sim 0.98$





- Simulations of DIII-D ($n_{\text{RMP}}=3$) and KSTAR ($n_{\text{RMP}}=1$) exhibit
 - Higher particle flux in the pedestal → **density pump-out**
 - Electron heat transport barrier in the pedestal → **maintains steep T_e gradient**
 - TEM turbulence at $\psi_N \geq 0.94$ (DIII-D) and $\psi_N \geq 0.90$ (KSTAR) are affected by RMPs
 - KSTAR shows higher degree of stochasticity than DIII-D
 - Fundamental result is similar between $n=1$ (KSTAR) and $n=3$ (DIII-D)
- XGC simulations of KSTAR $n=1$ RMP discharge find
 - ~2x higher fluctuation amplitudes with RMPs compared to pre-RMP phase
 - Experiment finds ~4x increase → electromagnetic effects might be missing
- Working on self-consistent RMP penetration in XGC → mitigate uncertainty due to experimental plasma and rotation profiles
- **Electromagnetic XGC** will be used to study effect on ELM-turbulence interaction