# Energetic Particle Stability and Confinement Issues in 3D Configurations

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#### **Energetic Particle Group Working Session**

21st US Transport Taskforce Workshop March 25 - 28, 2008 Three main topics will be discussed (work in progress):

- Discrete shear Alfvén mode structure in 3D configurations
- Coupled Alfvén/sound continua for stellarators
   application to HSX observations
- Self-consistent ripple calculations for ITER and Monte Carlo alpha confinement simulation

# Motivations for 3D Alfvén mode structure calculations (AE3D):

- Collaboration with Y. Todo (NIFS)
- Applications:
  - Alpha/beam/ICRF/ECH driven AE instabilities
  - Stellarators (both high R/<a> and compact)
  - Tokamaks with ripple, RWM coils, etc.
- Useful in development of reduced models
  - Perturbative linear stability models
  - Limited mode/particle nonlinear models
- Improve understanding of instability drive, sideband coupling effects
  - more complex in 3D systems
- Optimization target function
- Useful for interpretation of experimental measurements
  - Mode identification
  - synthetic diagnostics
- Input for particle orbit confinement studies in presence of Alfvén instabilities (collaborations at NIFS with Isobe, Okasabe)
- Other efforts: A. Könies, Phys. Plasmas 7 (2000) 1139.

Reduced MHD description leads to a single shear Alfvén eigenmode equation vs. the 3 coupled equations that characterize full MHD

$$\begin{cases} \vec{\nabla} \cdot \left( \frac{\mu_0 \rho_{ion}}{B} \frac{\partial}{\partial t} \frac{\vec{\nabla} \phi}{B} \right) = \left( \vec{B} \cdot \vec{\nabla} \right) \frac{J_{\parallel}}{B} \\ \frac{\partial \psi}{\partial t} = \frac{\vec{B}}{B} \cdot \vec{\nabla} \phi \end{cases}$$

Setting  $\frac{\partial}{\partial t} = -i\omega$  and combining equations :

$$\omega^2 \vec{\nabla} \cdot \left( v_A^{-2} \vec{\nabla} \phi \right) + \vec{B} \cdot \vec{\nabla} \left[ \frac{1}{B} \nabla^2 \left( \frac{\vec{B}}{B} \cdot \vec{\nabla} \phi \right) \right]$$

Reduced MHD shear Alfvén model [based on "Generalized reduced MHD Equations," by S. E. Kruger, et al., Phys. of Plasmas **5** (1998) 4169]

#### Variational form of 3D shear Alfvén equation:

- Multiply by weighting function  $\tilde{\phi}$  and reduce to surface terms plus symmetric terms:
- Inertial term:

$$\tilde{\phi} \ \omega^2 \ \vec{\nabla} \cdot \left( \mathbf{v}_A^{-2} \ \vec{\nabla} \phi \right) = \omega^2 \ \vec{\nabla} \cdot \left[ \mathbf{v}_A^{-2} \tilde{\phi} \ \vec{\nabla} \phi \right] - \frac{\omega^2}{\mathbf{v}_A^2} \vec{\nabla} \tilde{\phi} \cdot \vec{\nabla} \phi$$

• Bending energy term:

$$\tilde{\phi} \ \vec{B} \cdot \vec{\nabla} \left[ \frac{1}{B} \nabla^2 \left( \frac{\vec{B}}{B} \cdot \vec{\nabla} \phi \right) \right] = \vec{\nabla} \cdot \left[ \frac{\tilde{\phi}}{B} \vec{B} \ \nabla^2 \left( \frac{\vec{B} \cdot \vec{\nabla} \phi}{B} \right) \right] - \vec{\nabla} \cdot \left[ \left( \frac{\vec{B} \cdot \vec{\nabla} \phi}{B} \right) \vec{\nabla} \left( \frac{\vec{B} \cdot \vec{\nabla} \phi}{B} \right) \right] + \vec{\nabla} \left( \frac{\vec{B} \cdot \vec{\nabla} \phi}{B} \right) \cdot \vec{\nabla} \left( \frac{\vec{B} \cdot \vec{\nabla} \phi}{B} \right)$$

 Ignoring terms (indicated in blue) that will lead to surface integrals, this results in the following symmetric, selfadjoint equation:

$$-\frac{\omega^2}{v_A^2}\vec{\nabla}\vec{\phi}\cdot\vec{\nabla}\phi+\vec{\nabla}\left(\frac{\vec{B}\cdot\vec{\nabla}\vec{\phi}}{B}\right)\cdot\vec{\nabla}\left(\frac{\vec{B}\cdot\vec{\nabla}\phi}{B}\right)=0$$

#### **Reduced MHD shear Alfvén model: solution method**

- Solution procedure
- Galerkin method (weighted residuals with trial function = weight function)

 $\phi = \sum_{m,l} f_m(\rho) \cos(m_l \theta - n_l \zeta) \qquad (based on stellarator symmetry)$ 

mod *e* selection rules (e.g., for n = 1 family):

 $n = 1, \quad m = 0, 1, 2, 3, ...$   $n = -1, \quad m = 1, 2, 3, ...$   $n = 1 \pm N_{fp}, \quad m = 1, 2, 3, ...$   $n = -1 \pm N_{fp}, \quad m = 1, 2, 3, ...$ etc.

Up to 200 Fourier modes have been used with 40 - 60 radial elements





# The Alfvén equation is solved in $\rho$ , $\theta$ , $\zeta$ Boozer coordinates:

Inertial term:

$$\frac{\omega^{2}}{v_{A}^{2}}\left[g^{\rho\rho}\frac{\partial\tilde{\phi}}{\partial\rho}\frac{\partial\phi}{\partial\rho}+g^{\theta\theta}\frac{\partial\tilde{\phi}}{\partial\theta}\frac{\partial\phi}{\partial\theta}+g^{\zeta\zeta}\frac{\partial\tilde{\phi}}{\partial\zeta}\frac{\partial\phi}{\partial\zeta}+g^{\rho\theta}\left(\frac{\partial\tilde{\phi}}{\partial\rho}\frac{\partial\phi}{\partial\theta}+\frac{\partial\tilde{\phi}}{\partial\theta}\frac{\partial\phi}{\partial\rho}\right)+g^{\theta\zeta}\left(\frac{\partial\tilde{\phi}}{\partial\zeta}\frac{\partial\phi}{\partial\theta}+\frac{\partial\tilde{\phi}}{\partial\theta}\frac{\partial\phi}{\partial\zeta}\right)\right]$$

Bending energy term: 
$$P\tilde{P}g^{\rho\rho} + Q\tilde{Q}g^{\theta\theta} + R\tilde{R}g^{\zeta\zeta} + (P\tilde{Q} + Q\tilde{P})g^{\rho\theta} + (P\tilde{R} + R\tilde{P})g^{\rho\zeta} + (Q\tilde{R} + R\tilde{Q})g^{\zeta\theta}$$

where 
$$P = \frac{\partial}{\partial \rho} \left[ \left( \frac{\psi'}{B\sqrt{g}} \right) \left( \frac{\partial \phi}{\partial \zeta} + \dot{\tau} \frac{\partial \phi}{\partial \theta} \right) \right]$$
  
 $Q = \frac{\partial}{\partial \theta} \left[ \left( \frac{\psi'}{B\sqrt{g}} \right) \left( \frac{\partial \phi}{\partial \zeta} + \dot{\tau} \frac{\partial \phi}{\partial \theta} \right) \right]$   
 $R = \frac{\partial}{\partial \zeta} \left[ \left( \frac{\psi'}{B\sqrt{g}} \right) \left( \frac{\partial \phi}{\partial \zeta} + \dot{\tau} \frac{\partial \phi}{\partial \theta} \right) \right]$ 

and  $\tilde{P}, \tilde{Q}, \tilde{R} = P, Q, R$  with  $\phi$  replaced by  $\tilde{\phi}$ 

#### **Remaining steps:**

• Final step: integrate inertia and bending energy terms over plasma volume to obtain matrix equations:

$$\iiint \sqrt{g} d\rho d\theta d\zeta (\cdots) \Rightarrow Ax = \omega^2 Bx$$

• Involves convolution integrals of the form:

$$I_{ijk}^{(ccc)} \equiv \int_{-\pi}^{\pi} d\zeta \int_{-\pi}^{\pi} d\theta \cos\left(n_{i}\zeta - m_{i}\theta\right) \cos\left(n_{j}\zeta - m_{j}\theta\right) \cos\left(n_{k}\zeta - m_{k}\theta\right)$$
$$I_{ijk}^{(ssc)} \equiv \int_{-\pi}^{\pi} d\zeta \int_{-\pi}^{\pi} d\theta \sin\left(n_{i}\zeta - m_{i}\theta\right) \sin\left(n_{j}\zeta - m_{j}\theta\right) \cos\left(n_{k}\zeta - m_{k}\theta\right)$$

- these can be done exactly/analytically
- Setting coefficient of highest order ρ derivatives to zero leads to continuum equation:

$$\frac{\omega^2}{v_A^2}\frac{g^{\rho\rho}}{B^2}\frac{\partial\phi}{\partial\rho} + \vec{B}\cdot\vec{\nabla}\left\{\frac{g^{\rho\rho}}{B^2}\left(\vec{B}\cdot\vec{\nabla}\right)\frac{\partial\phi}{\partial\rho}\right\} = 0$$

 [previously studied in D. A. Spong, R. Sanchez, A. Weller, Phys. of Plasmas 10 (3217) 2003]

#### Metric elements (variation on a flux surface shown here for LHD)



- Coefficients that involve these matrix elements are calculated on a  $\theta$ ,  $\zeta$  grid on each surface
- This data is then Fourier transformed

### NCSX (quasi-toroidal stellarator)



### NCSX: Open gap mode





### Closed gap (continuum damped) mode





# LHD Beam-driven AE's

- Alfvén instabilities driven by beam ions have been studied in LHD shot #47645 [M. Osakabe, et al., Nucl. Fusion 46, S911 (2006)]
- Hole-clump pairs in the fast ion distribution are observed with the neutral particle analyzer
- Shown to be consistent with AE3D Alfvén mode structures [Y. Todo, et al., Proc. Of ITC/ISHW2007]



# LHD gap structure



<sup>#</sup> of AE modes/3 kHz frequeuncy bin







6167, omega = 0.12256E+03





6195, omega = 0.10438E+03





## **Topics**

- Discrete shear Alfvén mode structure in 3D configurations
- Coupled Alfvén/sound continua for stellarators
  - Collaborators: B. Breizman, D. Brower, C. B. Deng
  - Coherent modes observed in HSX with Alfvénic scaling
  - Alfvénic activity near 50 kHz (and sometimes 70kHz) is observed in HSX correlated with energetic electron tails
  - 1/q = iota ~ 1.05
  - Frequencies sometimes intersect continua near magnetic axis
  - Frequency insensitive to small iota variations
  - Motivated inclusion of sound continua
- Self-consistent ripple calculations for ITER and effects on alpha confinement

Solving the coupled Alfvén-sound continuum equations (derived by B. Breizman) provides a more complete picture of modes that may be available for energetic particle coupling (STELLGAP)

Shear Alfvén-sound continuum (B. Breizman) Coupled sound-shear Alfvén waves  $-\frac{\omega^2 g}{B^2}\eta - \frac{1}{4\pi\rho}(\mathbf{B}\cdot\nabla)\frac{g}{B^2}(\mathbf{B}\cdot\nabla)\eta = C_s^2[(\mathbf{B}\cdot\nabla)\zeta - \eta G]G \qquad (\text{color scale} \approx \sqrt{\sum_{m,n} \eta_{mn}^2}$ 



Upward frequency shift in coupled sound-shear Alfvén continua near magnetic axis is more consistent with HSX observations:





#### 52.6 kHz eigenmode (at $\rho$ = 0.2 surface) **Alfvénic** - sound and Alfvén component continua meet with similar frequency and mode structure 5.365e-0 -2.728e-07 – 1.387e-09 Max: 2.076 Min: 1.387e-09 -0.01 0.01 -0.00 0.00 Freq (kHz)<sup>60</sup> sound component 30 0.10 0.20 0.30 0.40 0.50 0.60 0.70 0.80 0.90 rho -1.58 -0.53 0.53 1.58

# Topics

- Discrete shear Alfvén mode structure in 3D configurations
- Coupled Alfvén/sound continua for stellarators
- Self-consistent ripple calculations for ITER and effects on alpha confinement
  - Free boundary VMEC equilibria including both updown asymmetry and TF ripple
  - Allows more self-consistent study of finite β, Shafranov shift, diamagnetic/Pfirsh-Schlüter currents on alpha confinement than previously possible
    - Y. Suzuki, Y. Nakamura, K. Kondo, "Finite beta effects on the toroidal field ripple in threedimensional tokamak equilibria," Nucl. Fusion **43** (2003) 406.
    - J. L. Johnson, A. H. Reiman, "Self-consistent, three-dimensional equilibrium effects on tokamak magnetic field ripple," Nucl. Fusion **38** (1988) 1116.
  - DELTA5D + VMEC + MFBE\* codes will be used modeling of alpha transport in ITER
    - \*E. Strumberger, P. Merkel, et al., IPP Report 5/100 (May, 2002)

View of the future: an operating ITER or DEMO will need accurate "cradle-to-grave" modeling of alphas in order to assure protection the first wall/divertor/PFCs



Accurate modeling of the edge region and large numbers of simulation particles are needed to obtain good alpha loss statistics at the 1st wall

#### Codes:

- VMEC 3D rippled equilibria
- MFBE
  - E. Strumberger, P. Merkel, et al., IPP Report 5/100 (May, 2002)
  - Accurate magnetic field (virtual casing method) evaluation in vacuum region between plasma and coils
  - Follow alphas to their demise at the wall -> predict localized wall heating
- DELTA5D-Magcoords
  - Monte Carlo orbits in Boozer coordinates
  - R. H. Fowler, J. A. Rome, J. F. Lyon, Phys. Fl. 28, (1985) 338.
- DELTA5D-CyISVD
  - Monte Carlo orbits in VMEC and cylindrical coordinates using SVD compression methods
  - Can include islands/AE modes
  - "Compression of magnetohydrodynamic simulation data using singular value decomposition,"del-Castillo-Negrete, D., Hirshman, S. P., Spong, D. A., D'Azevedo, E. F., JOURNAL OF COMPUTATIONAL PHYSICS 222, 265 (2007)

#### Recent parallel scaling on Cray XT3



# ITER free boundary VMEC equilibria have been calculated using a filamentary coil model:

Coil model: 18 TF's with 25 filaments each 5 filaments per PF |B| contours on outer flux surface (compressed color map)



- Next step: complete connections between VMEC's external Green's function and ANSYS vacuum field calculation
  - Allows direct incorporation of the effects of ferritic inserts, TBM's, finite volume coil currents, RWM's etc. into VMEC equilibrium

Contours of the main n = 18 ripple harmonic have been benchmarked against vacuum data from ANSYS: Contours of  $B_{n=18}(R,Z)$  where  $B(R,\phi,Z) = \sum_{n} B_n(R,Z) \cos(n\phi)$ ANSYS vacuum n = 18 ripple VMEC  $\beta = 2.4\%^n$  n = 18 ripple



At finite β's ripple contours permeate somewhat further into core (i.e., ripple amplification by diamagnetic currents)

note: edge ripple( $\delta$ ) ~ B<sub>n=18</sub>/5 ~ 0.2 - 1%

Contours of  $|B_{n=18}(\theta,s)|$  for  $<\beta>=0\%$ 















Test problem: follow slowing-down of alpha populations launched at various flux surfaces for  $n(0) = 10^{20} \text{ m}^{-3}$ ,  $T_i(0) = T_e(0) = 20 \text{ keV}$ (more realistic centrally peaked profiles show zero losses) 10,240 alphas followed for 1 second physical time up-down symmetric ITER fixed boundary equilibrium



### Variation of losses with equilibrium $<\beta>$



### Summary

- New code (AE3D) developed for Alfven spectral analysis and mode structure calculations in stellarators
  - Can calculate all or a subset of Alfvén spectrum
    - Up to 8000 eigenmodes have been kept per configuration
  - Mode density function vs. frequency
    - Density minima associated with more open gaps
  - 3D mode structure visualization
- STELLGAP code has been upgraded for Alfvén-sound wave continua
  - Regions identified where Alfvén and sound continua couple
  - Application to HSX shows that minima of coupled continua is somewhat higher in frequency
  - Further experimental information needed to identify observed mode
- ITER rippled equilbria calculated with VMEC and used for alpha loss calculations
  - Self-consistent finite  $\beta$  3D model including ripple
    - future upgrades to include effects of ferritic steel inserts, RWM coils, etc.
  - Coupled to Monte Carlo alpha loss code (DELTA5D)
  - Can be extended to include turbulence/follow alphas to the 1st wall